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Perspectives on cavity field control

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Low-Level RF Workshop, 2019-10-03



Thoughts from an automatic-control (and linac) perspective

1. Complex-coefficient LTI systems
2. Energy-based parametrization of cavity dynamics
3. Normalized cavity dynamics, sensitivity to disturbances
4. Parasitic modes

Baseband cavity models

Complex SISO representation

$$P_a(s) = \frac{\omega_{1/2}}{s + \omega_{1/2} - i\Delta\omega}$$

Real TITO representation

$$\mathbf{P}_a(s) = \frac{\omega_{1/2}}{(\Delta\omega)^2 + (s + \omega_{1/2})^2} \begin{bmatrix} \omega_{1/2} & -\Delta\omega \\ \Delta\omega & \omega_{1/2} \end{bmatrix}$$

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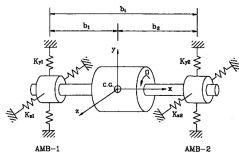
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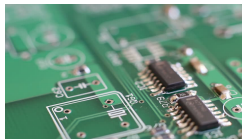
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Other control applications:



Vibration damping of rotating machinery



FB linearization of RF amps.

Relation between real and complex representations

Complex SISO representation

$$G(s) = G_{\text{Re}}(s) + iG_{\text{Im}}(s)$$

Real TITO representation

$$G_{\text{equiv}}(s) = \begin{bmatrix} G_{\text{Re}}(s) & -G_{\text{Im}}(s) \\ G_{\text{Im}}(s) & G_{\text{Re}}(s) \end{bmatrix}$$

Relation between real and complex representations

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Consider eigendecomposition

$$G_{\text{equiv}}(s) = U \begin{bmatrix} G(s) & 0 \\ 0 & G^*(s) \end{bmatrix} U^H, \quad U = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ -i & i \end{bmatrix}.$$

Note: $G^*(i\omega) = \overline{G(-i\omega)}$. Positive and negative frequencies are intertwined in $G_{\text{equiv}}(s)$

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Advantages of the complex-coefficient representation

- Simplifies understanding, calculations, life in general, etc
- Structure is implicit, good for system identification
- More efficient computations

Control theory for complex-coefficient systems

Standard tools and results apply but

- Change A^T to A^H
- Remember to consider negative frequencies

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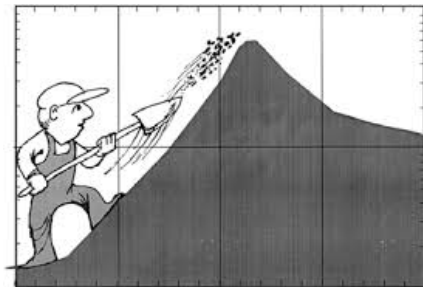


Illustration of Bode's sensitivity integral (the water-bed effect)

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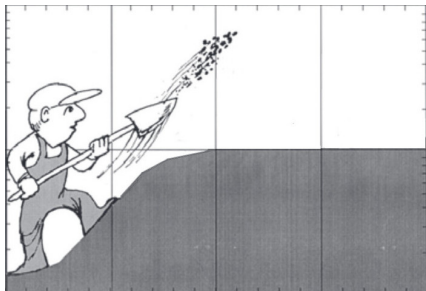


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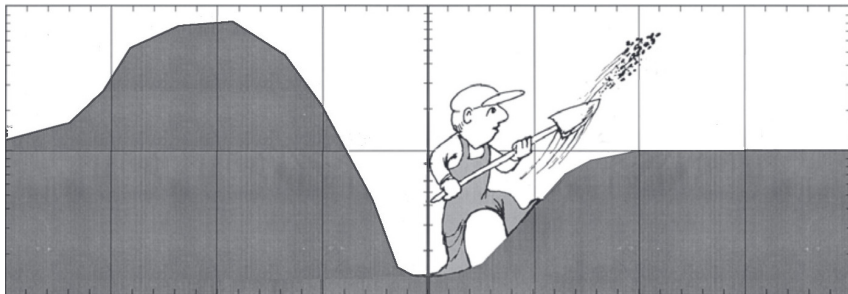
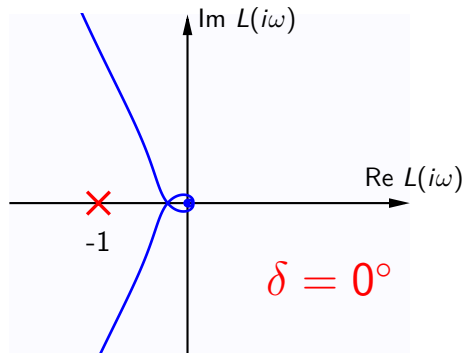
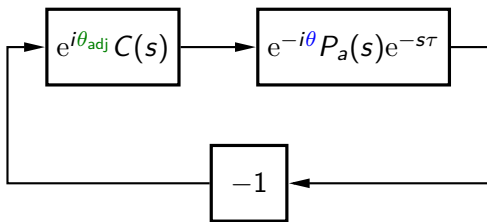


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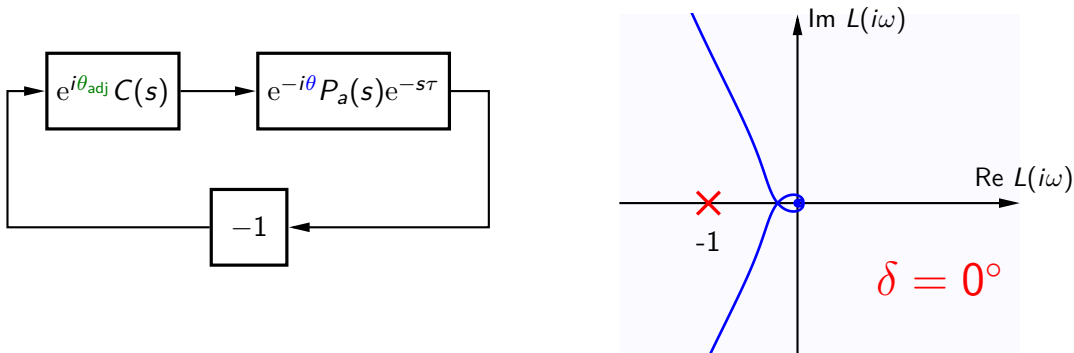
Intuitive understanding of loop-phase adjustment

Open-loop system: $e^{i\theta_{\text{adj}}} C(s) e^{-i\theta} P_a(s) e^{-s\tau} = e^{\delta} L_0(s)$ where $\delta := \theta_{\text{adj}} - \theta$.



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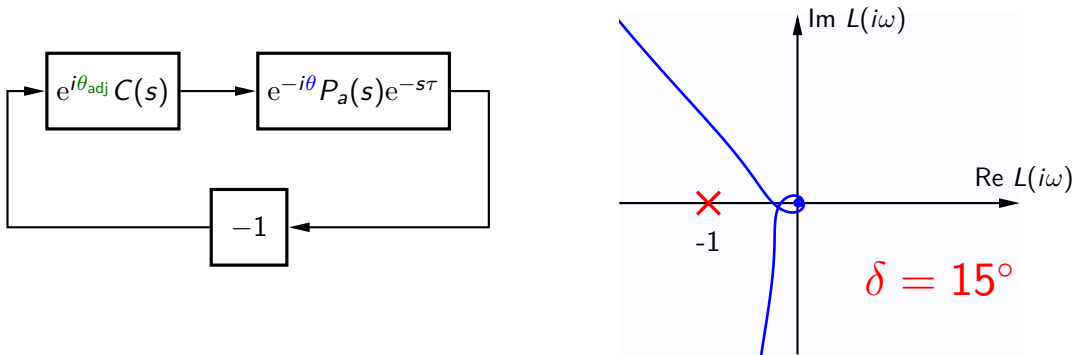
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Loop-phase-adjustment error δ gives corresponding phase-margin reduction!

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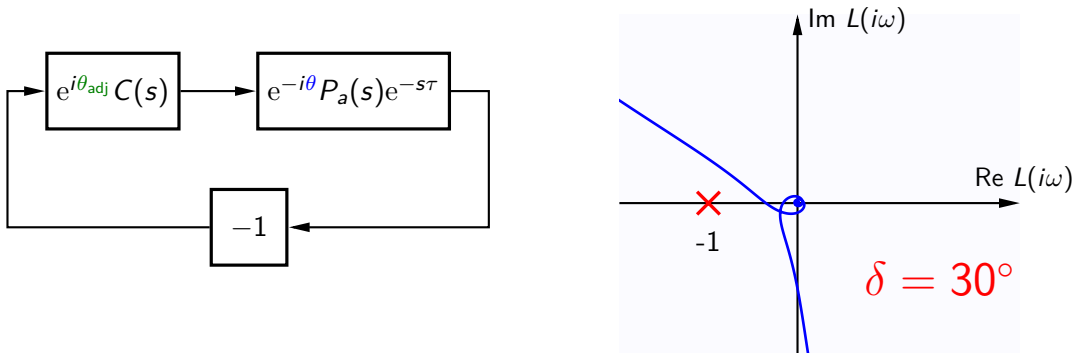
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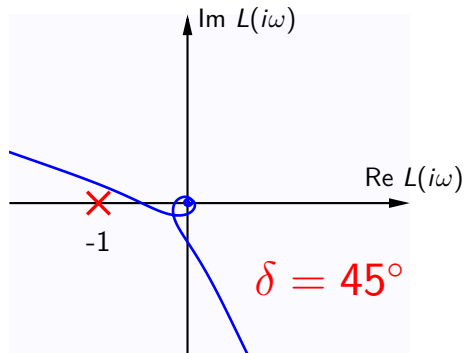
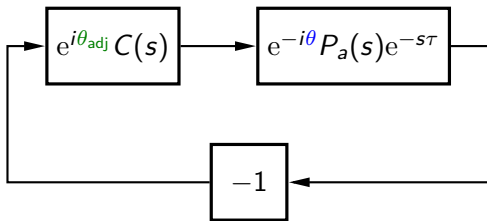
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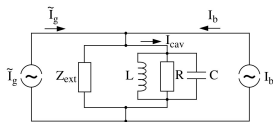
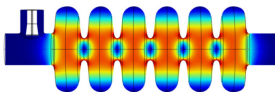
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Energy-based parametrization of cavity dynamics

Accelerator Cavity Modeling (1/2)

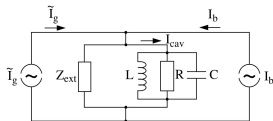
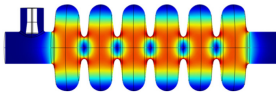


Equivalent-circuit parametrization:

$$\frac{d\mathbf{V}}{dt} = (-\omega_{1/2} + i\Delta\omega)\mathbf{V} + R_L\omega_{1/2}(2\mathbf{I}_g + \mathbf{I}_b)$$

$$P_g = \frac{1}{4} \frac{r}{Q} Q_{\text{ext}} |\mathbf{I}_g|^2$$

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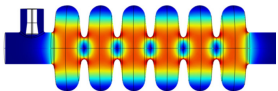
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RF drive is modeled as fictitious generator current \mathbf{I}_g . Problematic.

“One word of caution is required here: /.../ for considerations where Q_{ext} varies /.../ or where (R/Q) varies /.../ the model currents cannot be considered constant; they have to be re-adapted” [Tückmantel (2011)]

Accelerator Cavity Modeling (1/2)



Energy-based parametrization:

$$\frac{d\mathbf{A}}{dt} = (-\gamma + i\Delta\omega)\mathbf{A} + \sqrt{2\gamma_{\text{ext}}}\mathbf{F}_{\mathbf{g}} + \frac{\alpha}{2}\mathbf{I}_{\mathbf{b}}$$

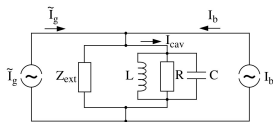
\mathbf{A} – Mode amplitude [$\sqrt{\text{J}}$]

$\mathbf{F}_{\mathbf{g}}$ – Forward wave [$\sqrt{\text{W}}$]

$$\mathbf{V} = \alpha\mathbf{A} \quad \left(\alpha = \sqrt{\omega_a(r/Q)}\right)$$

$$P_g = |\mathbf{F}_{\mathbf{g}}|^2$$

Haus (1984) *Waves and fields in optoelectronics*
plus beam loading

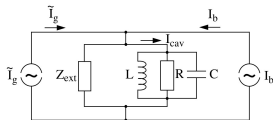
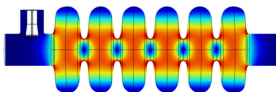


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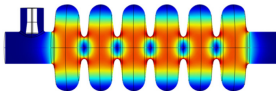
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Advantages of energy-based parameterization:

- Cleaner expressions, e.g., $P_g = |\mathbf{F}_g|^2$
- States and parameters are well defined
- Direct connection to physical quantities of interest

Accelerator Cavity Modeling (1/2)

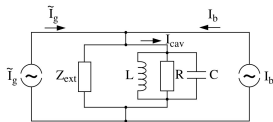


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Helpful to think of γ as both decay rate and bandwidth.

The total decay rate $\gamma = \gamma_0 + \gamma_{\text{ext}}$, not so intuitive if considered as bandwidths
Common for laser cavities



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Optimal coupling and detuning

$$\text{Minimize } |\mathbf{F}_{\mathbf{g}0}|^2 = \frac{1}{2\gamma_{\text{ext}}} \left| (-\gamma_0 - \gamma_{\text{ext}} + i\Delta\omega)\mathbf{A}_0 + \frac{\alpha}{2}\mathbf{I}_{\mathbf{b}0} \right|$$

with respect to $\Delta\omega$ and γ_{ext}

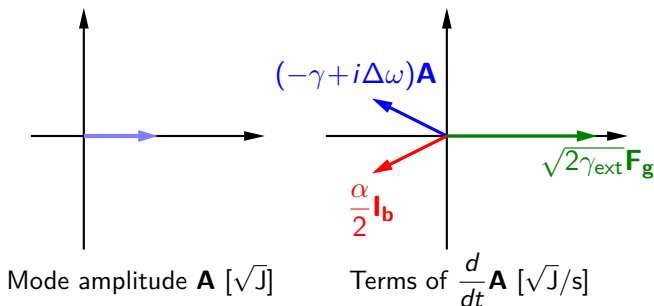
$$\text{Solution: } \Delta\omega = -\frac{1}{A_0} \text{Im} \frac{\alpha}{2} \mathbf{I}_{\mathbf{b}0}, \quad \gamma_{\text{ext}} = \gamma_0 - \frac{1}{A_0} \text{Re} \frac{\alpha}{2} \mathbf{I}_{\mathbf{b}0}$$

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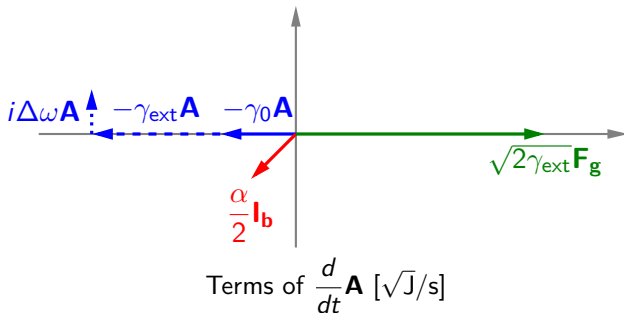


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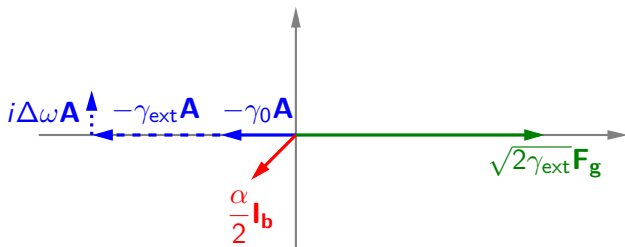


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Terms of $\frac{d}{dt}\mathbf{A} [\sqrt{\text{J/s}}]$

Normalized cavity dynamics

Normalization

$$\mathbf{a} := \frac{1}{A_0} \mathbf{A}, \quad \mathbf{f}_g := \frac{1}{\gamma A_0} \sqrt{2\gamma_{\text{ext}}} \mathbf{F}_g, \quad \mathbf{i}_b := \frac{1}{\gamma A_0} \frac{\alpha \mathbf{l}_b}{2}.$$

Normalized cavity dynamics

$$\dot{\mathbf{a}} = (-\gamma + i\Delta\omega)\mathbf{a} + \gamma(\mathbf{f}_g + \mathbf{i}_b).$$

At nominal operating point, with optimal coupling and tuning,
 $1 \leq \mathbf{f}_{g0} \leq 2$, $0 \leq \text{Re } \mathbf{i}_{b0} \leq 1$.

Relative disturbances \mathbf{d}_b and \mathbf{d}_b give rise to

$$\begin{aligned} \mathbf{f}_g &\approx (1 + \mathbf{d}_b)(\mathbf{f}_{g0} + \tilde{\mathbf{f}}_g) \approx \mathbf{f}_{g0} + \tilde{\mathbf{f}}_g + \mathbf{f}_{g0} \mathbf{d}_b \\ \mathbf{i}_b &= (1 + \mathbf{d}_b) \mathbf{i}_{b0} \end{aligned}$$

Introducing the relative field error $\mathbf{z} = 1 - \mathbf{a}$, we have

$$\dot{\mathbf{z}} = (-\gamma + i\Delta\omega)\mathbf{z} + \gamma \tilde{\mathbf{f}}_g + \gamma \mathbf{f}_{g0} \mathbf{d}_b + \gamma \mathbf{i}_{b0} \mathbf{d}_b$$

Transfer functions around operating point

At nominal operating point

$$\dot{\mathbf{z}} = (-\gamma + i\Delta\omega)\mathbf{z} + \gamma\tilde{\mathbf{f}}_{\mathbf{g}} + \gamma\mathbf{f}_{\mathbf{g}0}\mathbf{d}_{\mathbf{g}} + \gamma\mathbf{i}_{\mathbf{b}0}\mathbf{d}_{\mathbf{b}}$$

The transfer function from control action to the cavity field is given by

$$P_a(s) := \frac{\gamma}{s + \gamma - i\Delta\omega}$$

Transfer functions from relative disturbances to relative field errors are given by

$$P_{d_{\mathbf{g}} \rightarrow \mathbf{z}}(s) = \mathbf{f}_{\mathbf{g}0}P_a(s) \quad (1)$$

$$P_{d_{\mathbf{b}} \rightarrow \mathbf{z}}(s) = \mathbf{i}_{\mathbf{b}0}P_a(s) \quad (2)$$

For optimally tuned and coupled superconducting cavities $\mathbf{f}_{\mathbf{g}0} = 2$.

Additional factor 2 in disturbance sensitivity to relative amplifier variations!

Impact of disturbances

For clarity, assume that $\Delta\omega = \phi_b = 0$, so $\gamma_{\mathbf{i}_{b0}} = \gamma_{\text{beam}}$ and $\gamma_{\mathbf{f}_{g0}} = \gamma_0 + \gamma_{\text{ext}} + \gamma_{\text{beam}}$

Transfer functions from relative disturbances to relative field errors are given by

$$P_{d_g \rightarrow z}(s) = \frac{\gamma_0 + \gamma_{\text{ext}} + \gamma_{\text{beam}}}{s + \gamma_0 + \gamma_{\text{ext}}}, \quad (3)$$

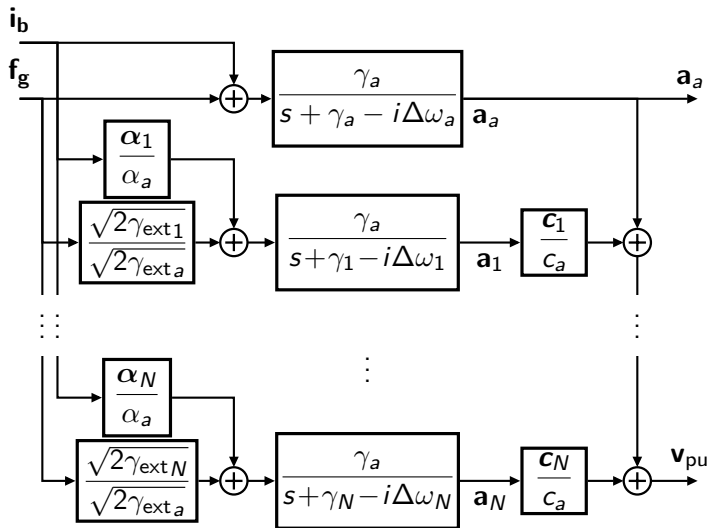
$$P_{d_b \rightarrow z}(s) = \frac{\gamma_{\text{beam}}}{s + \gamma_0 + \gamma_{\text{ext}}}. \quad (4)$$

Sensitivity to amplifier ripple, equation (4), cannot be made smaller than $\frac{\gamma_0 + \gamma_{\text{beam}}}{s + \gamma_0}$

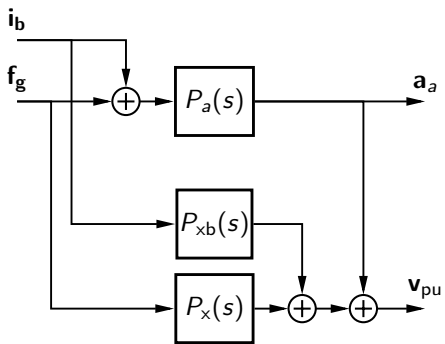
Difficulty of field control is determined by γ_0 and γ_{beam} , but typically $\gamma = 2\gamma_0 + \gamma_{\text{beam}}$

Parasitic modes

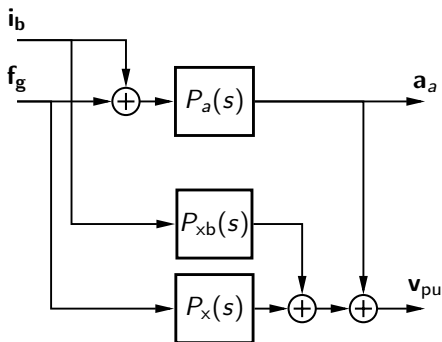
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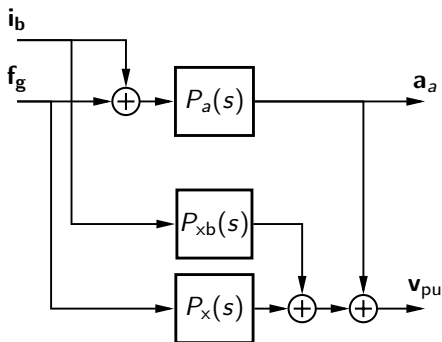
Parasitic modes



A. Calibrate setpoint v_{pu} for $a_a = a_a^*$ with short/low-current beam pulses ($i_{b0} = 0$)

B. Operation with nominal beam current and regulation to the set point

Parasitic modes



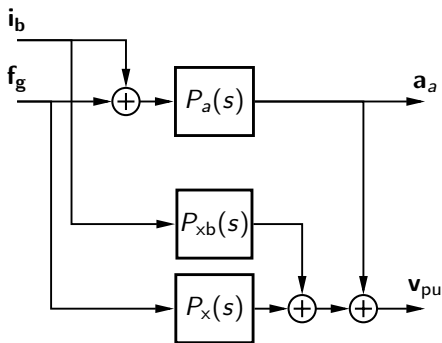
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Gives error steady-state error:

$$\delta = a_a^B - a_a^* = \frac{P_x(0) - P_{xb}(0)}{P_a(0) + P_x(0)} P_a(0) i_{b0} \approx (P_x(0) - P_{xb}(0)) i_{b0}$$

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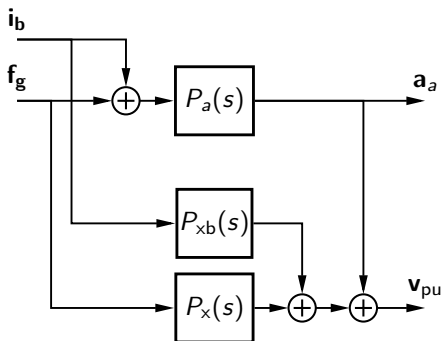
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ESS medium- β cavity: $\delta = P_x(0) \mathbf{i}_{b0} \approx \frac{\gamma_{5\pi/6}}{i\Delta\omega_{5\pi/6}} = \frac{R_5^2 \gamma_\pi}{i\Delta\omega_{5\pi/6}} \approx 0.00187i \leftrightarrow 0.11^\circ$

Parasitic modes



A. Calibrate setpoint v_{pu} for $a_a = a_a^*$ with short/low-current beam pulses ($i_{b0} = 0$)

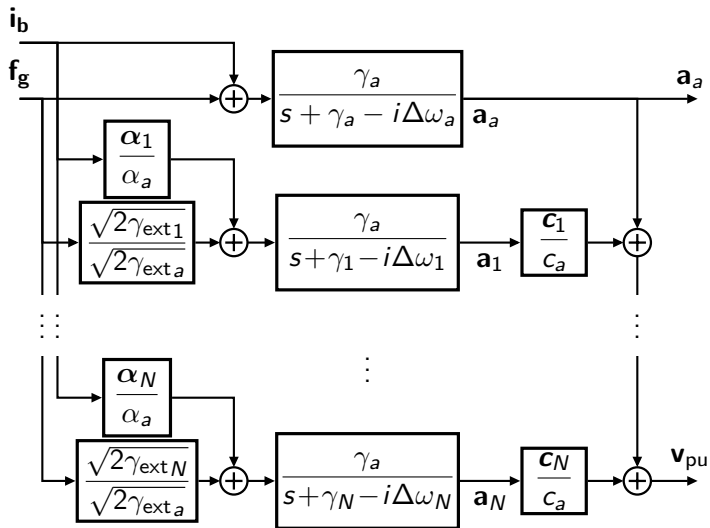
B. Operation with nominal beam current and regulation to the set point

Gives error steady-state error:

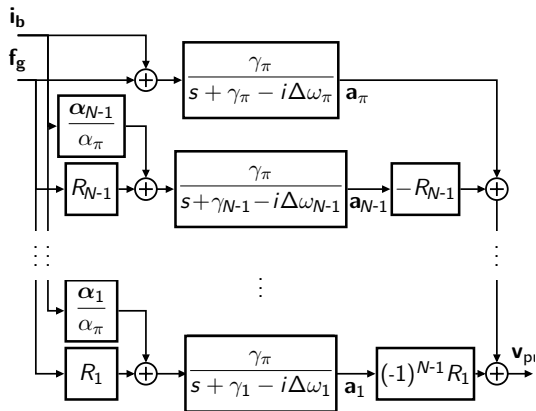
$$\delta = a_a^B - a_a^* = \frac{P_x(0) - P_{xb}(0)}{P_a(0) + P_x(0)} P_a(0) i_{b0} \approx (P_x(0) - P_{xb}(0)) i_{b0}$$

How to handle this? Do nothing, Kalman filter, re-calibrate?

Parasitic modes



Relations between same-order modes

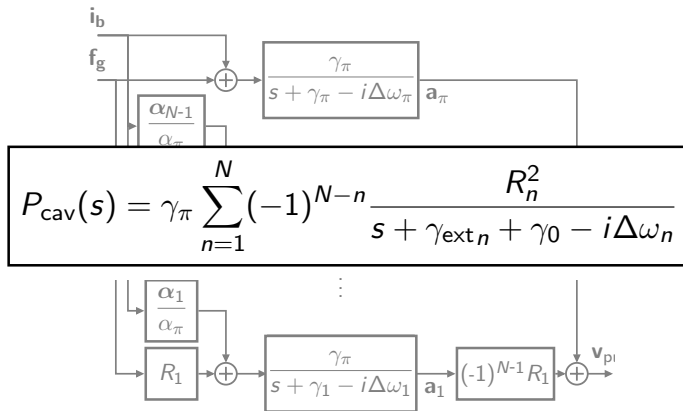


$$\gamma_{\text{ext}_n} = R_n^2 \gamma_{\text{ext}_\pi}$$

$$\Delta\omega_n \approx (R_n^2 - 2) k_{cc} \omega_{\text{cell}}$$

$$\text{where } R_n := \sqrt{2} \sin(n\pi/(2N))$$

Relations between same-order modes

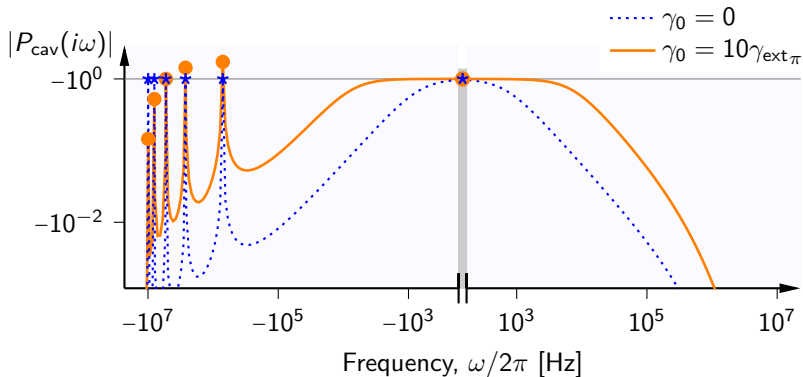


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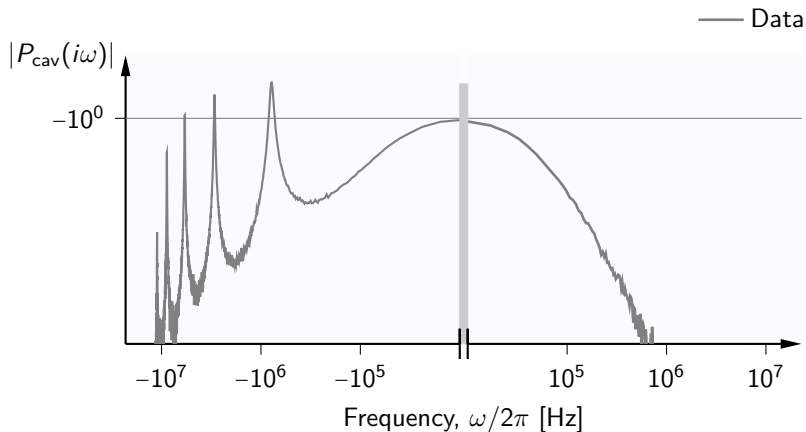
$$\text{where } R_n := \sqrt{2} \sin(n\pi/(2N))$$

Bode magnitude plot for 6-cell cavity



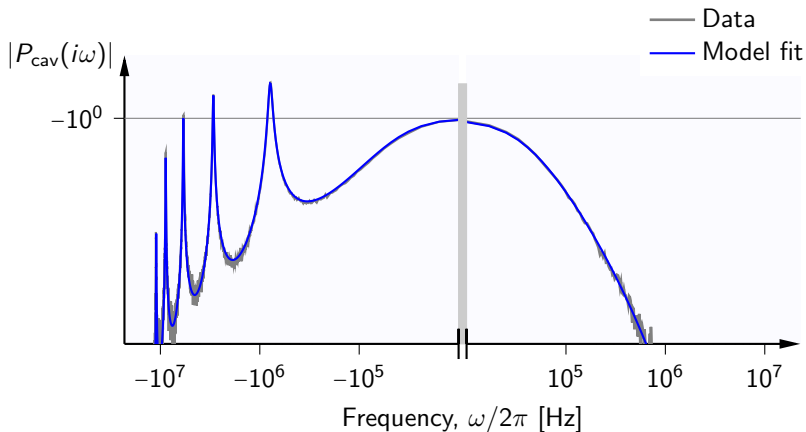
Similar to ESS medium- β cavity, $\gamma_{\text{ext}\pi} = 700$ Hz

Fit to measured data



Measurements by P. Pierini on *warm* 6-cell ESS medium- β cavity

Fit to measured data



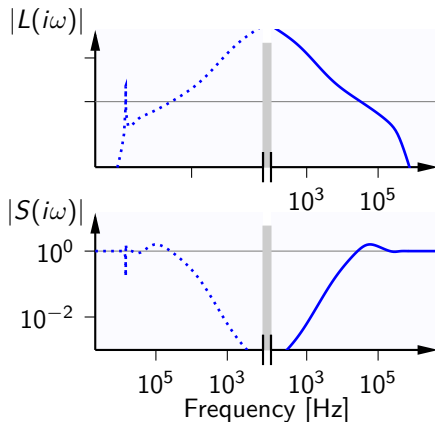
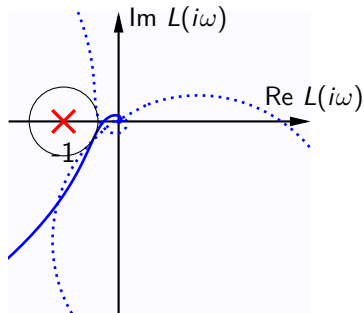
Measurements by P. Pierini on *warm* 6-cell ESS medium- β cavity

Four parameters were fitted. Estimated resistive decay rate, $\gamma_0/2/\pi = 35$ kHz.

Example Control Strategies for Parasitic Modes (1/3)

PI controller + 3rd order filter

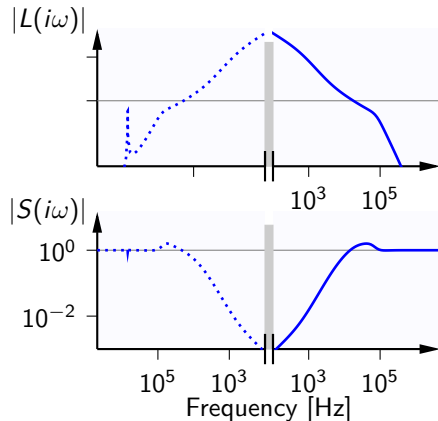
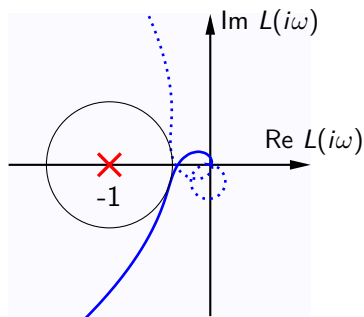
Set controller parameters for good phase of resonant “bubble”



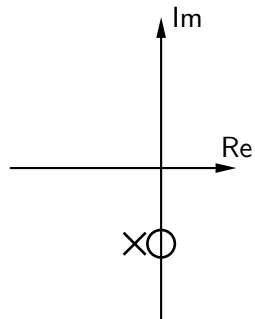
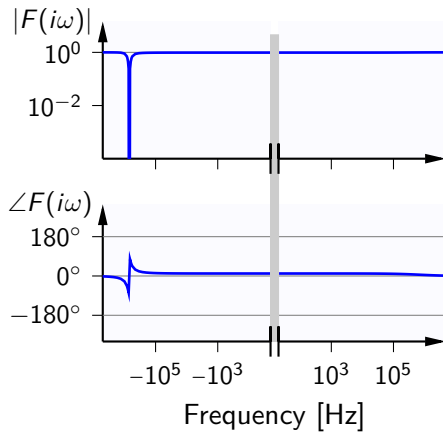
Example Control Strategies for Parasitic Modes (2/3)

PI controller + 2nd order filter

Wide-band suppression of the “bubble”

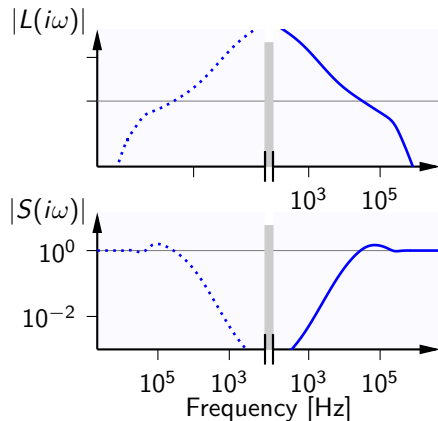
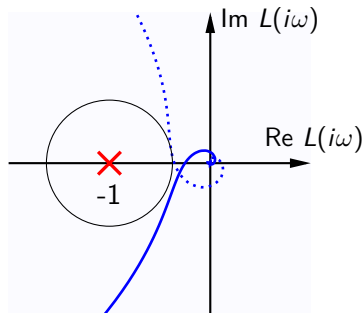


One-Sided Notch Filter



Example Control Strategies for Parasitic Modes (3/3)

PI controller + one-sided notch filter + 2nd order filter
Notch out the “bubble”



Summary

- Analyzing the field control loop as a complex-coefficient system is easier and gives more understanding. Particularly for loop-phase adjustment and parasitic modes.
- Energy-based cavity parametrization is more convenient and fundamental.
- There is a factor ≈ 2 in relative sensitivity to amplifier variations.
- Parasitic modes may give systematic control error since the controlled variable is not measured.

Summary

- Analyzing the field control loop as a complex-coefficient system is easier and gives more understanding. Particularly for loop-phase adjustment and parasitic modes.
- Energy-based cavity parametrization is more convenient and fundamental.
- There is a factor ≈ 2 in relative sensitivity to amplifier variations.
- Parasitic modes may give systematic control error since the controlled variable is not measured.

More details in upcoming PhD thesis.

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Bo Bernhardsson
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Anders J Johansson (EIT)
Project Leader
co-supervisor



Rolf Johansson
co-supervisor

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