

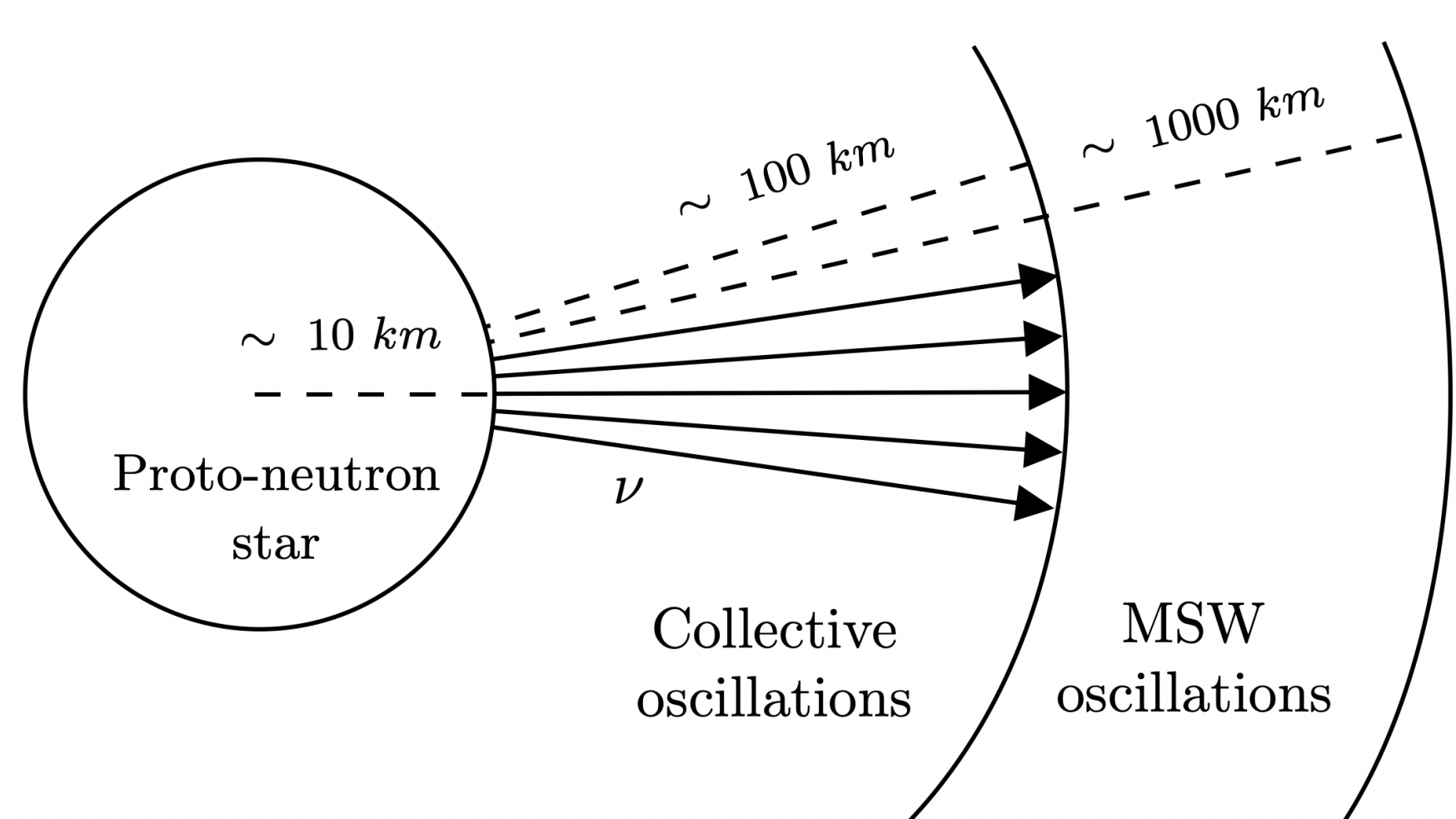
# QUANTUM COMPUTING SIMULATION FOR COLLECTIVE NEUTRINO OSCILLATIONS [1]

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## MOTIVATIONS

- **Core-collapse supernovae** of massive stars  $M \gtrsim 8M_\odot$  emit a huge number of neutrinos ( $\sim 10^{58}$ ).
- The physics of matter under extreme conditions is strongly **flavor-dependent** (nucleosynthesis, neutron-proton ratio, spectrum splits...).
- Interesting quantum many-body problem governed by the weak interaction.
- Describing the full dynamic is very complicated due to the **collective neutrino oscillations** that make the equation non linear.



- We want to simulate the real time evolution:

$$|\Psi(t)\rangle = U(t)|\Psi_0\rangle, \quad U(t) = e^{-iHt}. \quad (1)$$

## PHYSICAL DESCRIPTION

- Two-flavors approximation ( $SU(2)$  model) to encode the flavor state in a qubit state:

$$|\nu_e\rangle \mapsto |0\rangle, \quad |\nu_x\rangle \mapsto |1\rangle \quad (2)$$

- $N$  neutrinos encoded into  $N$  qubits.

- The flavor Hamiltonian of  $N$  neutrinos is:

$$H = \sum_{i=1}^N \mathbf{b} \cdot \boldsymbol{\sigma}_i + \sum_{i < j}^N J_{ij} \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j \quad (3)$$

- **1-body term:** vacuum mixing

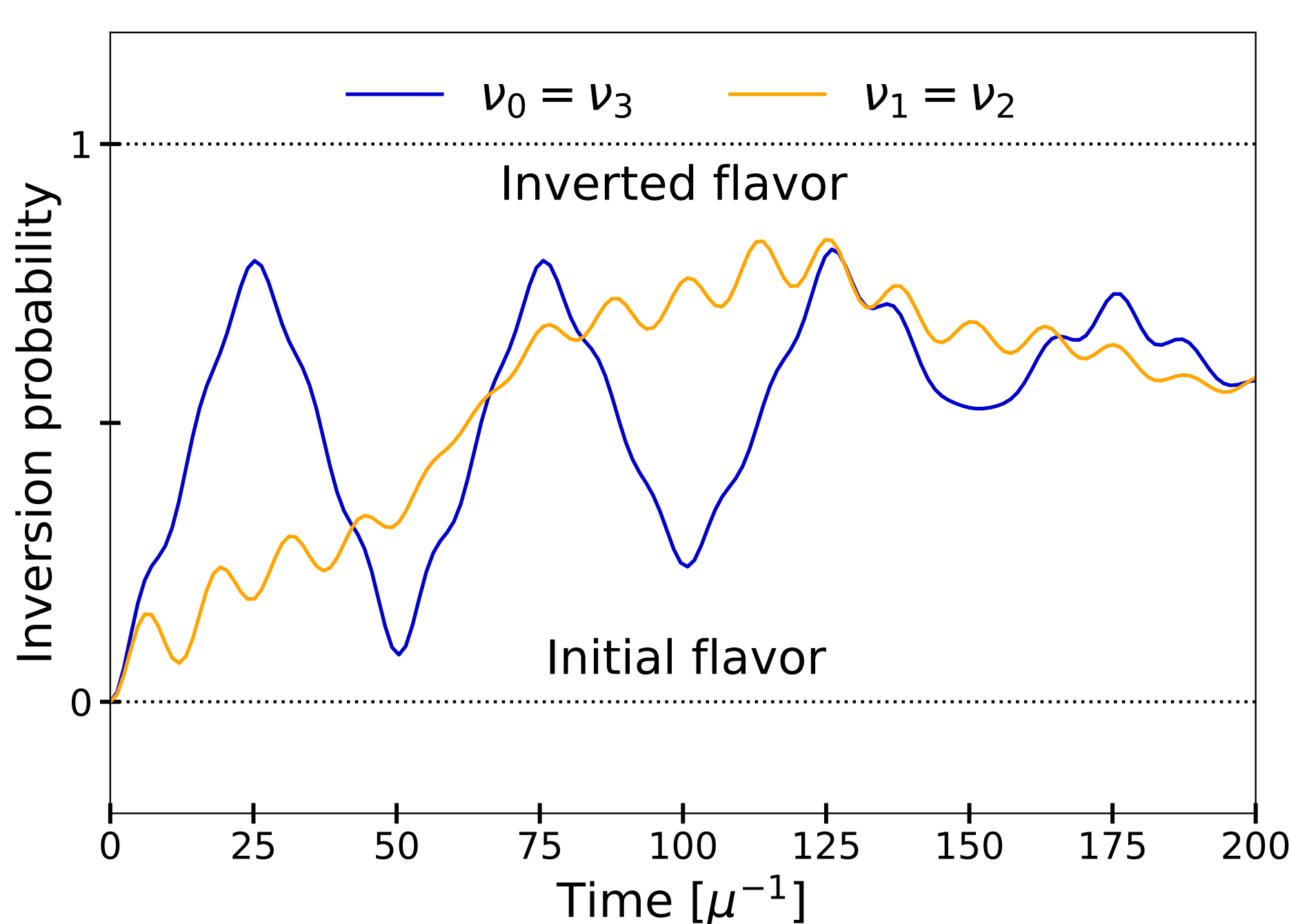
$$\mathbf{b} = \frac{\delta m^2}{4E_\nu} (\sin(2\theta_\nu), 0, -\cos(2\theta_\nu)).$$

- **2-body term:**  $\nu\nu$ -interaction

$$J_{ij} := \frac{\mu}{N} (1 - \cos(\theta_{ij})), \quad \mu = \sqrt{2} G_F \rho_\nu.$$

- Initial state for  $N = 4$ :  $|\Psi_0\rangle = |0011\rangle$ .

- We can look at the inversion probability:



- Symmetry under particle exchange:

$$0 \leftrightarrow 3, \quad 1 \leftrightarrow 2. \quad (4)$$

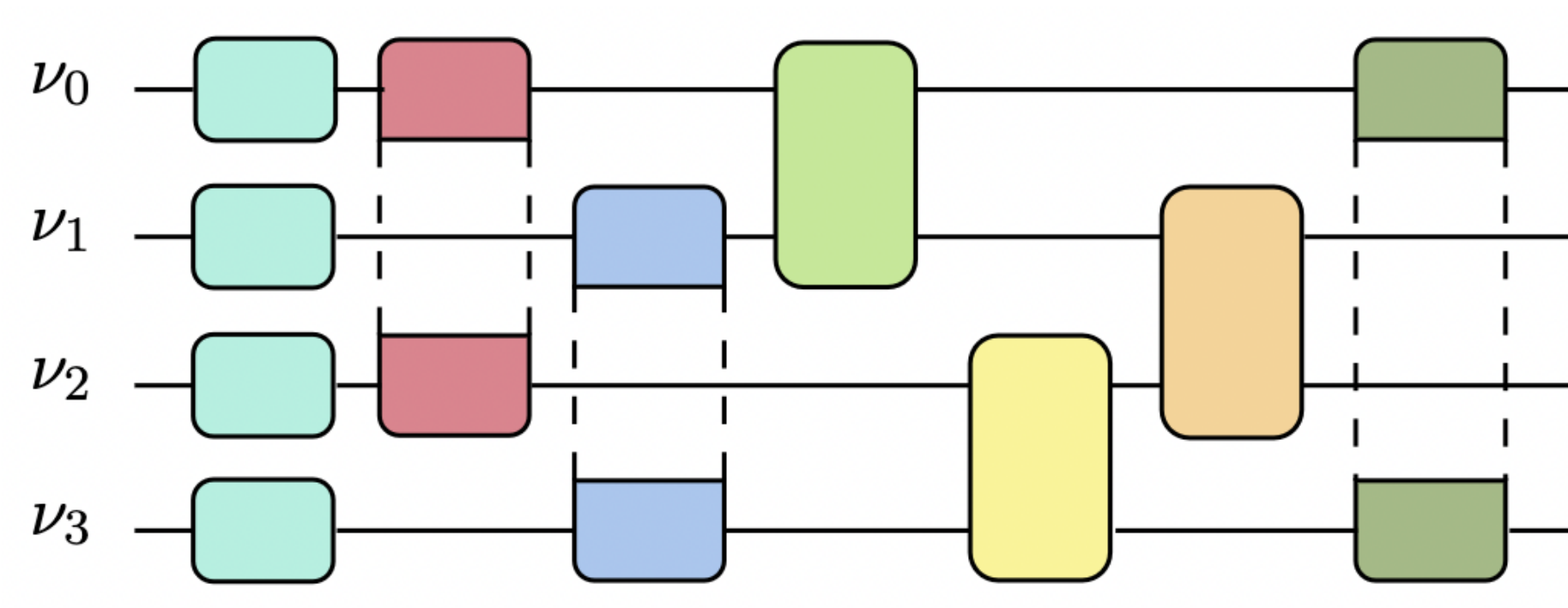
## UNITARY IMPLEMENTATION

- To perform the quantum simulation we need a **quantum gate decomposition** of the  $U(t)$  operator ( $2^N \times 2^N$  unitary matrix on the flavor basis):

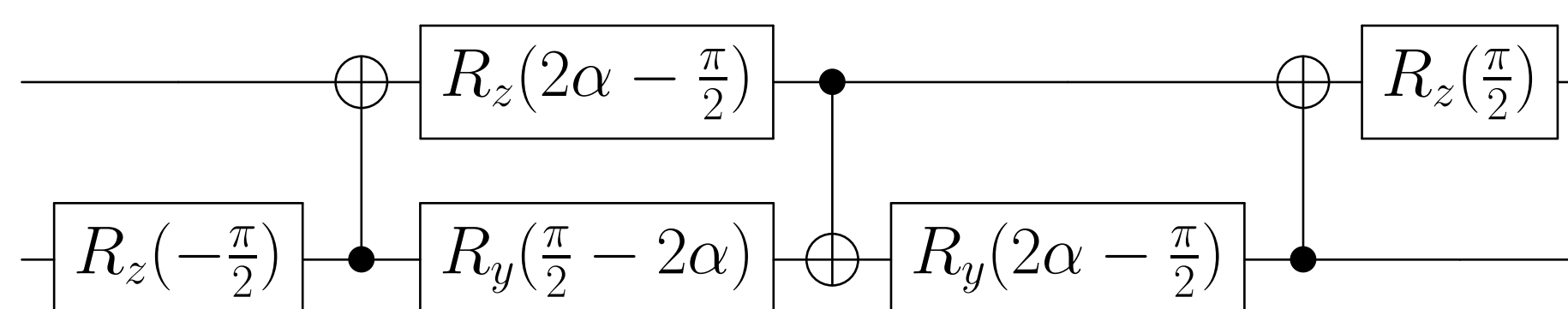
- Divide 1-body and 2-body parts that commute:

$$U(t) = U_2(t)U_1(t). \quad (5)$$

- Approximate the 2-body part as a product of pair interactions.



- Each 2-qubit gate  $u_{ij} = e^{i\alpha(X \otimes X + Y \otimes Y + Z \otimes Z)}$  can be implemented as [3]:



- The order in which the pairs interact changes the error due to the commutators.

- The **swap network** proposed in Ref. [2] implements the interaction on a chain of linearly connected qubits.

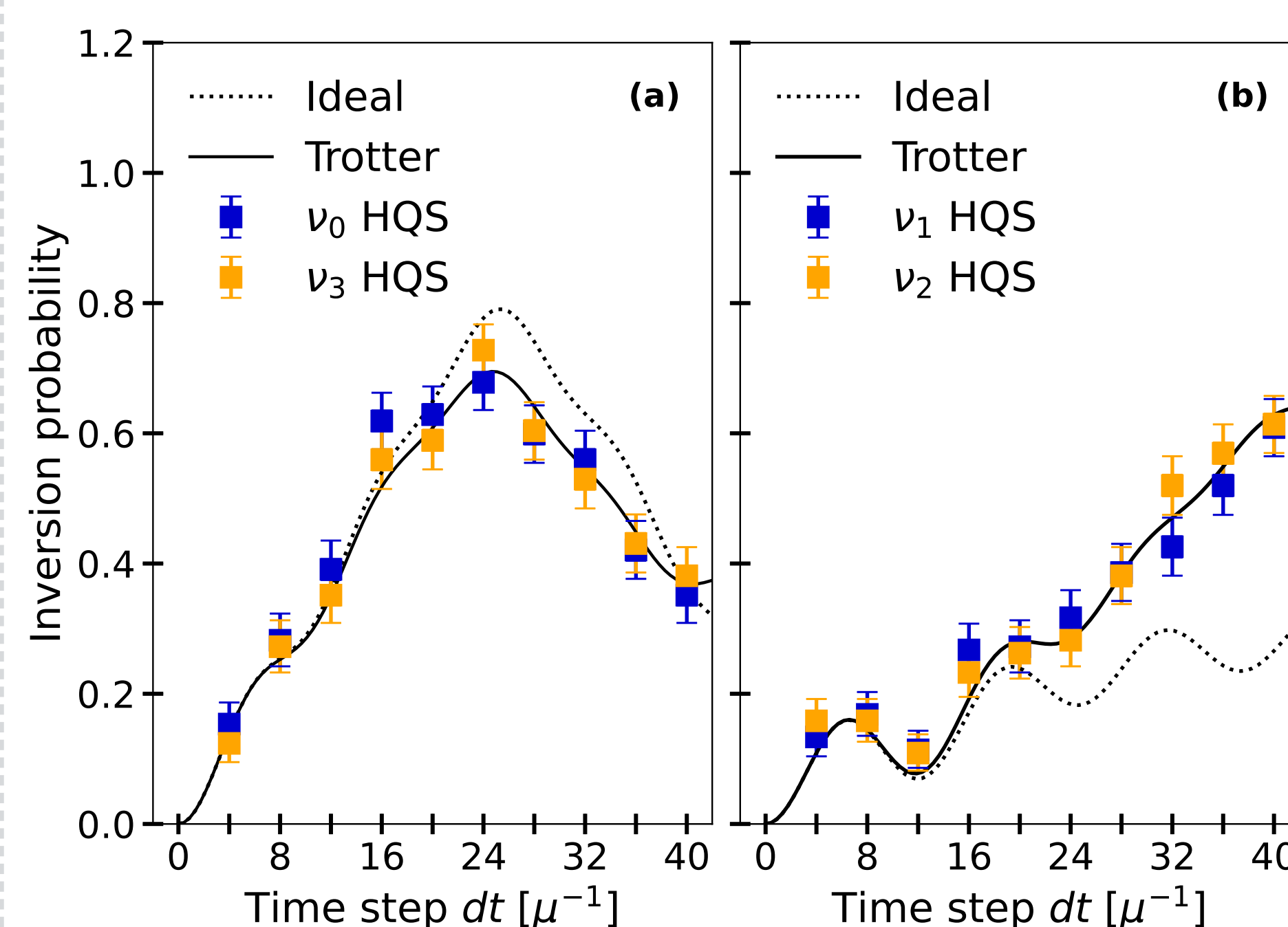
- **All-to-all connectivity** allows for best ordering and lower circuit complexity.

- Machine-aware compilation.

## SINGLE TROTTER STEP PROPAGATION

- The propagator is applied to the same initial state for different Trotter steps  $dt = 4, 8, \dots, 40 \mu^{-1}$ .

- Results obtained from the simulations on Quantinuum's trapped-ion device:



## References

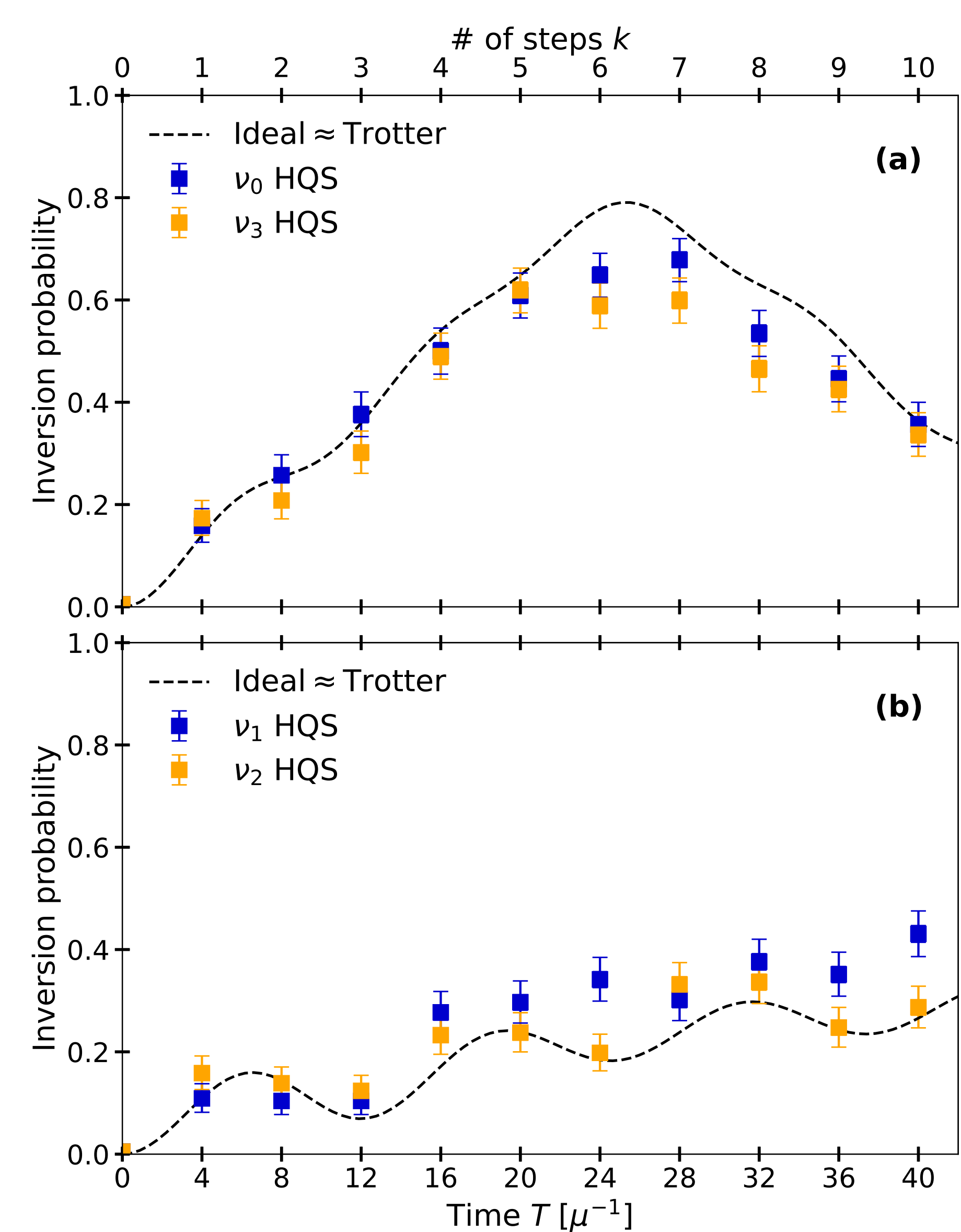
- [1] V. Amitrano et al. "Trapped-Ion Quantum Simulation of Collective Neutrino Oscillations". In: *arXiv:2207.03189* (2022).
- [2] B. Hall et al. "Simulation of Collective Neutrino Oscillations on a Quantum Computer". In: *Phys. Rev. D* 104, 063009 (2021).
- [3] F. Vatan et al. "Optimal Quantum Circuits for General Two-Qubit Gates". In: *Phys. Rev. A* 69, 032315 (2004).

## MULTIPLE STEPS EVOLUTION

- Evolve the system until  $T$  applying  $k = T/dt$  Trotter steps:

$$|\Psi(T)\rangle = U_2(dt)^k U_1(dt)^k |\Psi_0\rangle. \quad (6)$$

- Real quantum machine results:



- Very long circuits with a huge number of gates:

$k$	1	2	3	4	5	6	7	8	9	10
# ZZ	18	36	54	72	90	108	126	144	162	180
# $SU(2)$	36	68	100	132	164	196	228	260	292	324

## COMPLEXITY ALGORITHM SCALING

- The number of 2-qubit gates needed to evolve up to  $T$  with an error  $< \epsilon$  scales polynomially with  $N$ :

Decomposition type	Circuit complexity
First order Trotter	$\mathcal{O}\left(\frac{T^2 \mu^2 N^3}{\epsilon}\right)$
Second order trotter	$\mathcal{O}\left(\frac{T^3 \mu^3 N^{5/2}}{\sqrt{\epsilon}}\right)$
Qubitization	$\mathcal{O}(T \mu N^3 + \log(1/\epsilon) N^2)$

## CONCLUSION

- Fully connected qubits allow for more freedom in gate decomposition.
- Quantum circuit optimization is a crucial step in order to perform simulations on a near-term quantum device.
- The complexity of the implementation of time evolution scales polynomially with the number of neutrinos.

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