

# DUNE Expanded Scope Session: $0\nu\beta\beta$

Julietta Gruszko

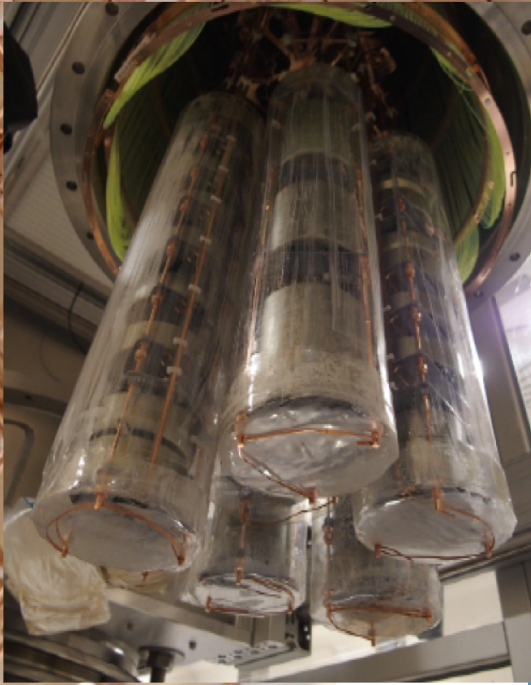
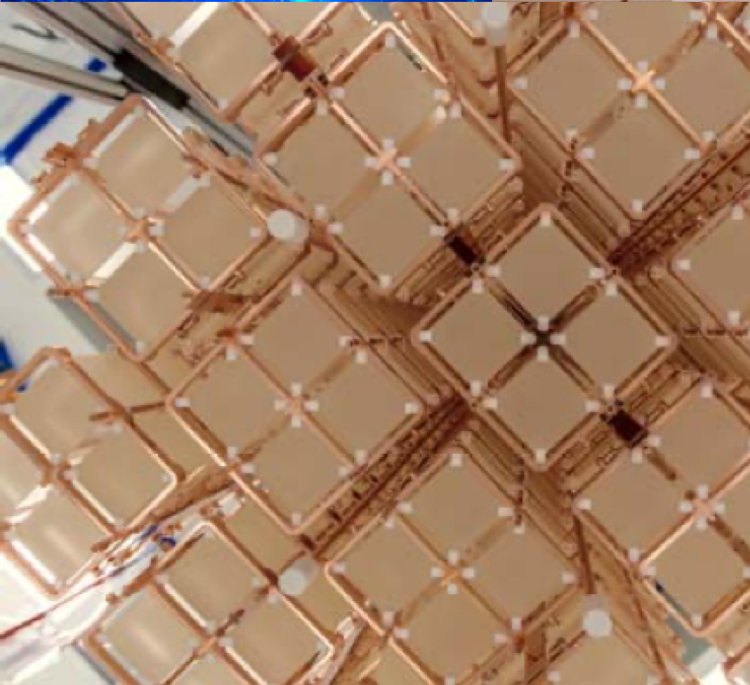
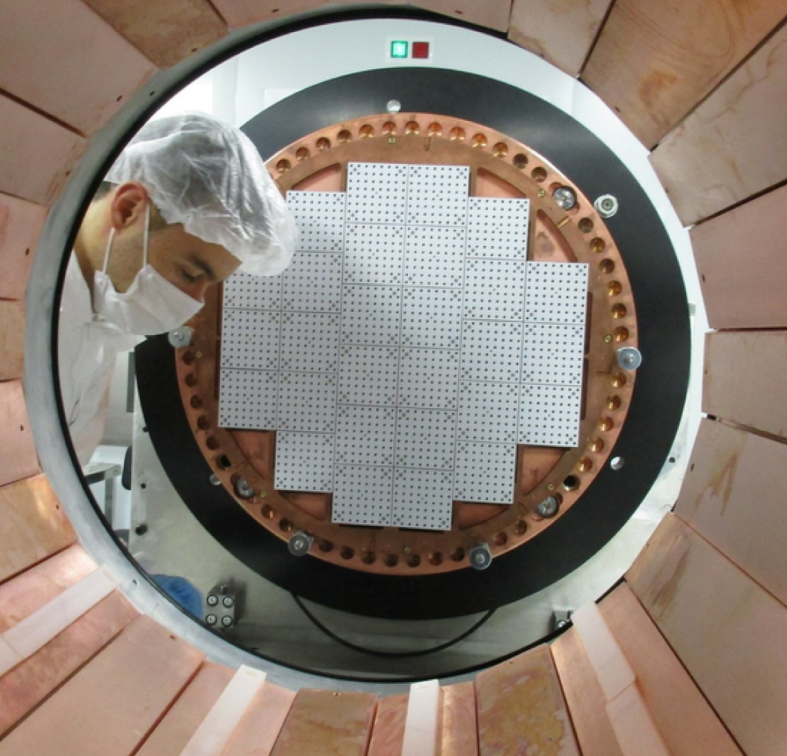
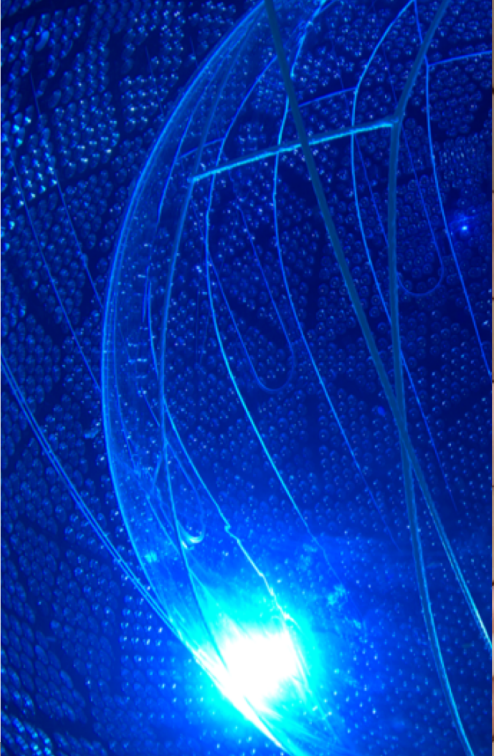
Snowmass Community Summer Study

July 20, 2022



THE UNIVERSITY  
of NORTH CAROLINA  
at CHAPEL HILL





## Outline

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- Why look for  $0\nu\beta\beta$ ?
- $0\nu\beta\beta$  sensitivity and discovery
- Current status and near-term future

Why look for  $0\nu\beta\beta$ ?

# Motivating BSM Physics

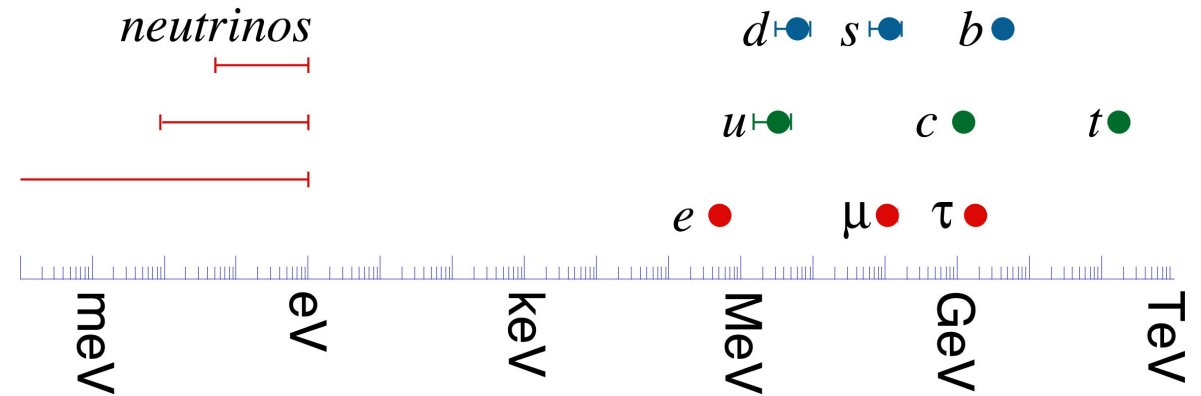
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- Need answers to:
  - What is dark matter?
  - What is dark energy/the mechanism behind the cosmological constant?
  - What generates neutrino mass?
  - What created the matter/antimatter asymmetry?
- Would like answers to:
  - Naturalness problems
  - CP conservation in QCD
  - Unification, flavor, etc...



# The Surprising Neutrino Mass

- A reminder: neutrino mass is not in the Standard Model!
- This is one of the few observations we have of Beyond-the-Standard Model physics
- Another surprise: neutrino mass is very small



# The Surprising Neutrino Mass

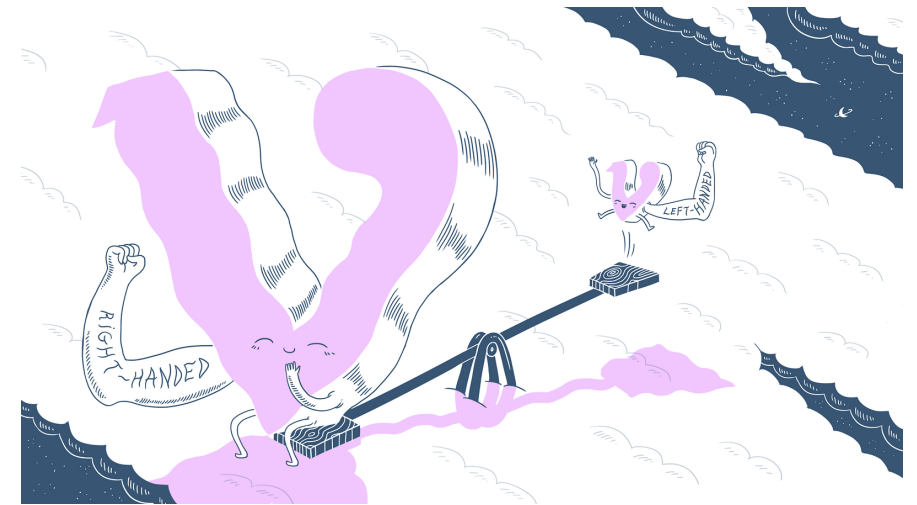
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Two options for neutrino mass terms:

- Dirac mass:
  - Requires two **non-interacting** new fields,  $\nu_R$  and  $\bar{\nu}_L$
  - Leads to hierarchy problem
- Majorana mass:
  - No new fields required;  $\bar{\nu}_R = \nu_R$  and  $\nu_L = \bar{\nu}_L$
  - Can be generated by new physics at TeV - GUT scale
- Both may be present; any non-zero Majorana mass makes the neutrino a Majorana fermion
- Majorana neutrino masses can be generated by a range of models

# The Type I See-Saw Mechanism

- Including both Majorana and Dirac mass terms can generate two light neutrinos,  $\nu$  and  $\bar{\nu}$ , and two heavy neutrinos,  $N$  and  $\bar{N}$
- If the Majorana mass term is of the GUT scale ( $\sim 10^{14}$  GeV) and Dirac mass term is of EW scale ( $\sim 100$  GeV):
  - $m_\nu \sim 0.1$  eV
  - $m_N \sim 10^{14}$  GeV
- This gives a “natural” neutrino of the correct mass by introducing a new GUT-scale particle



# The Matter Asymmetry Problem

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- Today, all the structure we see in the universe is made up of matter, with no significant quantity of antimatter
- Baryon asymmetry measurements give  $\eta \equiv \frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 6 \times 10^{-10}$
- We believe this asymmetry has to have been **generated dynamically**, not as an initial condition

# Making an Asymmetry: The Sakharov Conditions

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In 1967, Sakharov proposed 3 conditions required for baryon-generating interactions that would generate an asymmetry:

1. Baryon number violation From SM at high temperature
2. Interactions out of thermal equilibrium ( ) Majorana neutrinos can do this in many models
3. C and CP violation: need more than the CP violation observed in the SM (even if  $\delta_{CP}$  is maximal) ( ) Majorana neutrinos can do this in many models

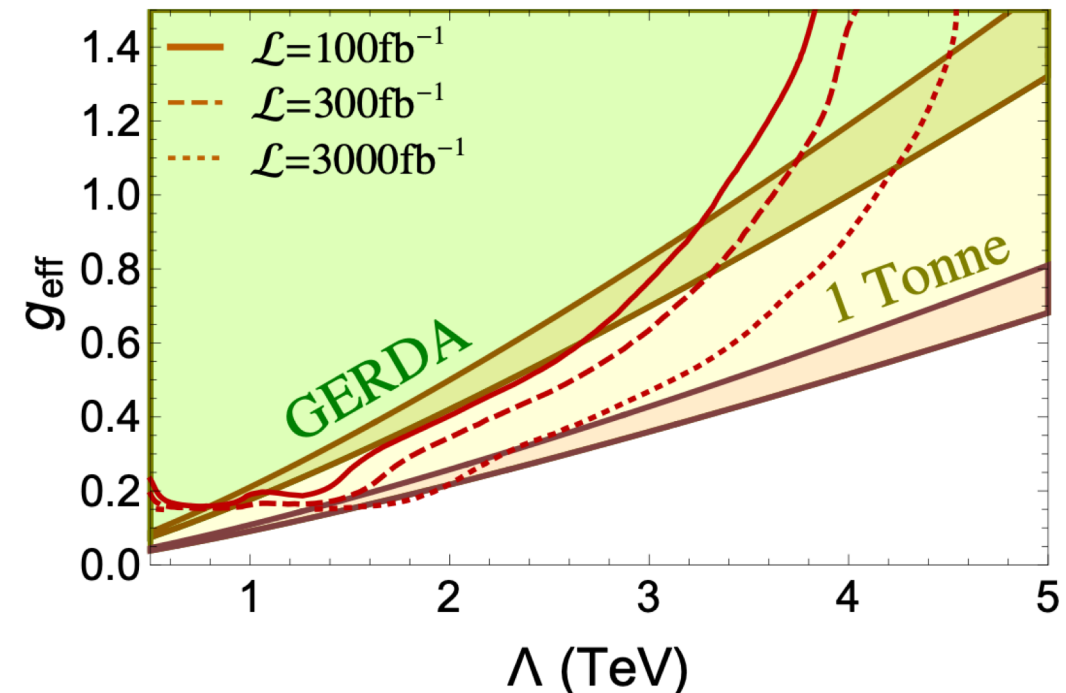
Majorana neutrinos could be a low-energy signature of the high-energy physics that generated baryon asymmetry



## Other Majorana Mass Mechanisms and Model-Building

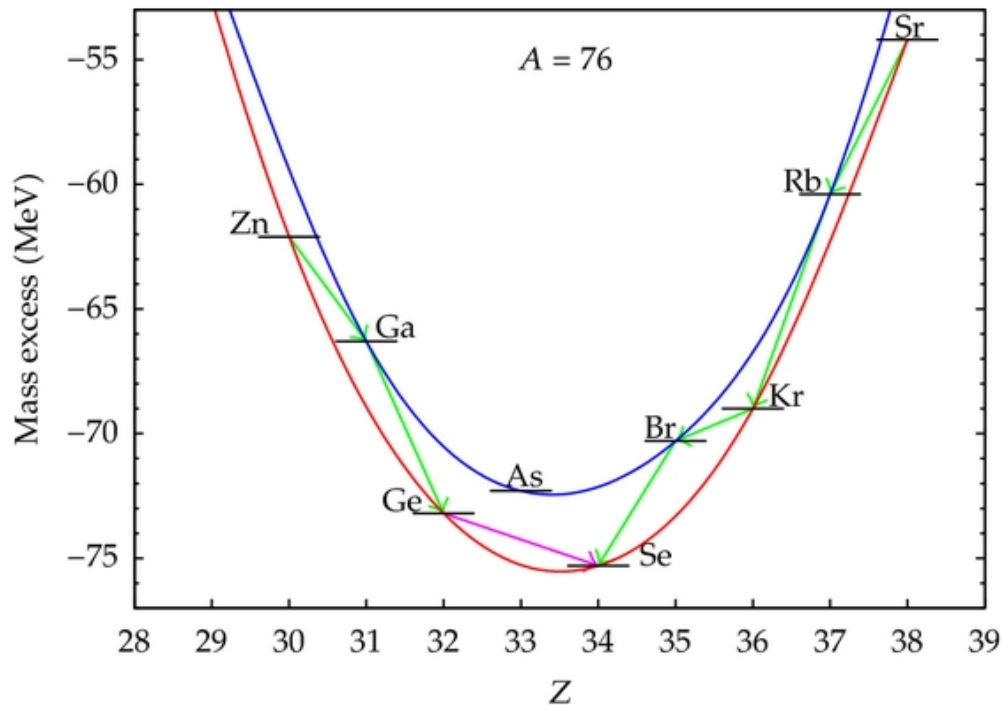
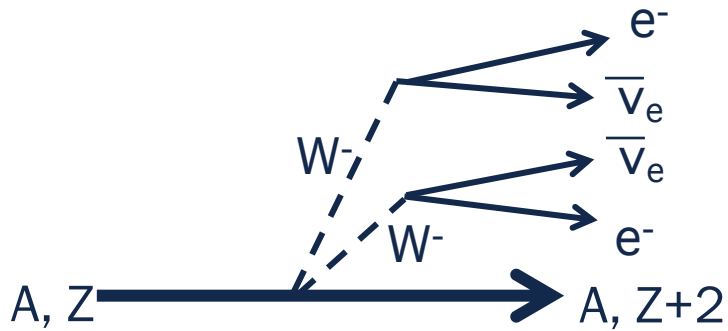
- There are many mechanisms beyond Type I see-saw that would generate neutrino mass
- Some generate the baryon asymmetry or dark matter candidate particles
- Many of these also predict new particles that could be observed at accelerators ( $O(1\text{-}10\text{'s of TeV})$ )
- Many models of flavor predict Majorana neutrinos with specific Majorana phases

Comparing LHC and  $0\nu\beta\beta$  limits on TeV-scale Lepton number violation



Peng, Ramsey-Musolf, and Winslow  
Phys. Rev. D **93**, 093002 (2016)

## Standard Model: $2\nu\beta\beta$



## Double-Beta Decay

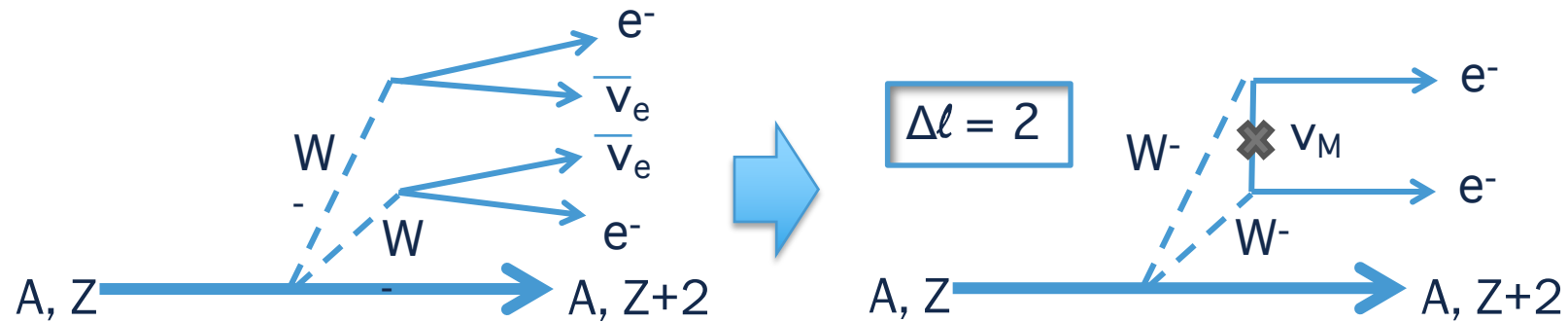
- For certain even-even nuclei, single beta decay is disallowed b/c of energy or momentum
- Instead, they double-beta decay
- Second-order weak process  
 $T_{1/2} \sim 10^{19} - 10^{21}$  years
- Electron capture variant is longest-lifetime process ever observed

# Double-Beta Decay Isotopes

- 35 naturally-occurring isotopes are capable of double-beta decay; we've observed it in 14 of these
- These 14 “golden nuclei” are particularly well-suited to experiments:
  - High Q-values
  - High abundance or ability to enrich (with some exceptions)
  - Other abundant isotopes of the element not highly radioactive

Double-beta candidate	Q-value (MeV)	Phase space $G_{01}(\text{y}^{-1})$	Isotopic abundance (%)	Enrichable by centrifugation
$^{48}\text{Ca}$	4.27226 (404)	$6.05 \times 10^{-14}$	0.187	No
$^{76}\text{Ge}$	2.03904 (16)	$5.77 \times 10^{-15}$	7.8	Yes
$^{82}\text{Se}$	2.99512 (201)	$2.48 \times 10^{-14}$	9.2	Yes
$^{96}\text{Zr}$	3.35037 (289)	$5.02 \times 10^{-14}$	2.8	No
$^{100}\text{Mo}$	3.03440 (17)	$3.89 \times 10^{-14}$	9.6	Yes
$^{116}\text{Cd}$	2.81350 (13)	$4.08 \times 10^{-14}$	7.5	Yes
$^{130}\text{Te}$	2.52697 (23)	$3.47 \times 10^{-14}$	33.8	Yes
$^{136}\text{Xe}$	2.45783 (37)	$3.56 \times 10^{-14}$	8.9	Yes
$^{150}\text{Nd}$	3.37138 (20)	$1.54 \times 10^{-13}$	5.6	No

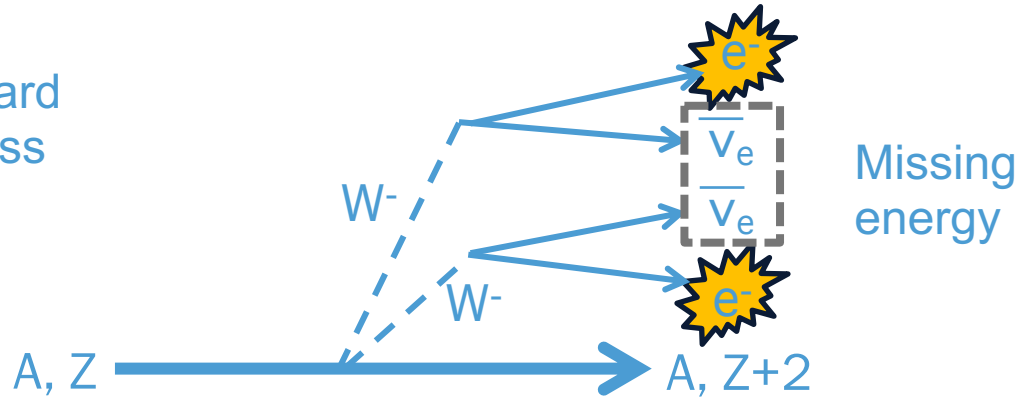
# Neutrinoless Double-Beta Decay



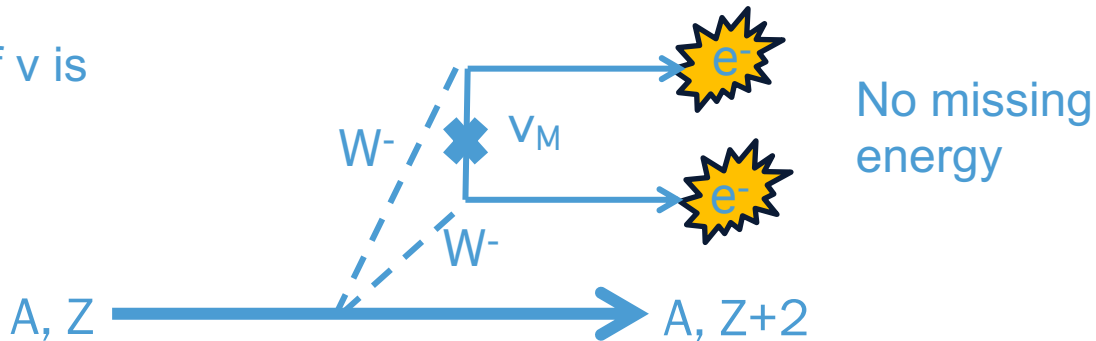
- If neutrinos are Majorana,  $0\nu\beta\beta$  could occur
- Lepton number conservation is violated by 2 units
- In this case, I've drawn the exchange of a light neutrino, but you can think of that "x" as a contracted diagram of any sort (with new physics in it)

# The Decay Signature

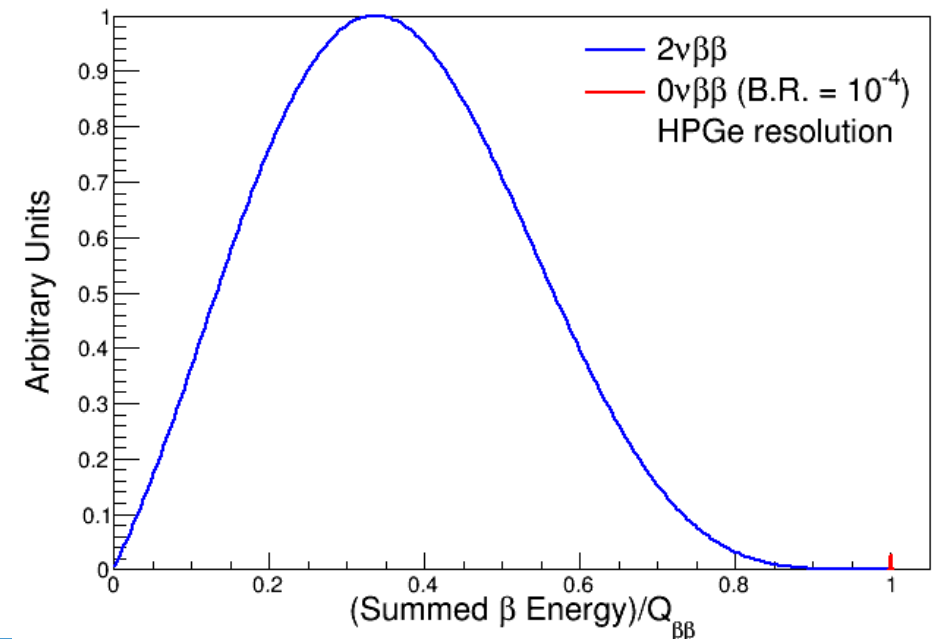
$2\nu\beta\beta$ : Standard Model process



$0\nu\beta\beta$ : Only if  $\nu$  is Majorana

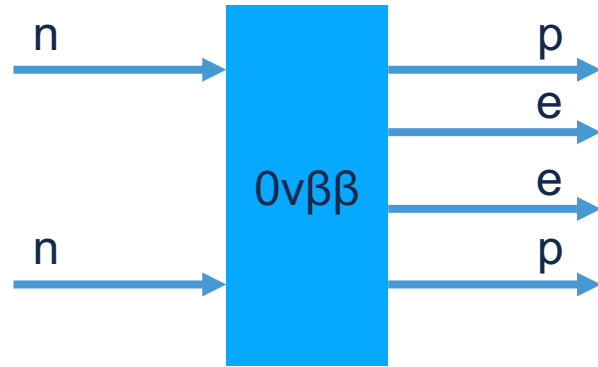


- In non-gaseous detectors, looks like a single-site event (modulo Cherenkov light)

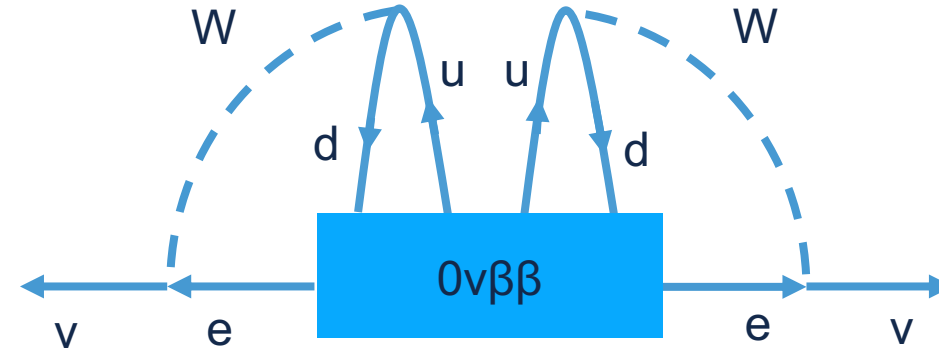




# Majorana Neutrinos and $0\nu\beta\beta$



$$(A, Z) \rightarrow (A, Z+2) + 2e^-$$



$$(A, Z) \rightarrow (A, Z+2) + 2e^-$$

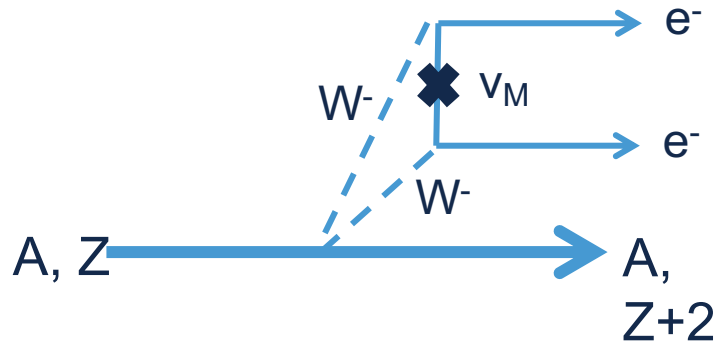
Model-independent implications of  $0\nu\beta\beta$ :

- Lepton number violation
- Neutrino-antineutrino oscillation, implying a non-zero Majorana mass term

The mechanism of  $0\nu\beta\beta$  determines the rate along with the parameters of the model

# $0\nu\beta\beta$ Sensitivity and Discovery

# The $0\nu\beta\beta$ Rate for Light Majorana Neutrino Exchange



$$(T_{1/2}^{0\nu})^{-1} = G^{0\nu} |M_{0\nu}|^2 \left( \frac{\langle m_{\beta\beta} \rangle}{m_e} \right)^2$$

Effective Majorana mass for light neutrino exchange:

$$\langle m_{\beta\beta} \rangle = \left| \sum_{i=1}^3 U_{ei}^2 m_i \right|$$

$$U = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{bmatrix}$$

$$= \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{bmatrix} \begin{bmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$c_{ij} = \cos \theta_{ij}$ ,  $s_{ij} = \sin \theta_{ij}$ ,  $\delta$  = Dirac CP violation,  $\alpha_i$  = Majorana CP violation

Even under simple assumptions, the  $0\nu\beta\beta$  rate depends on:

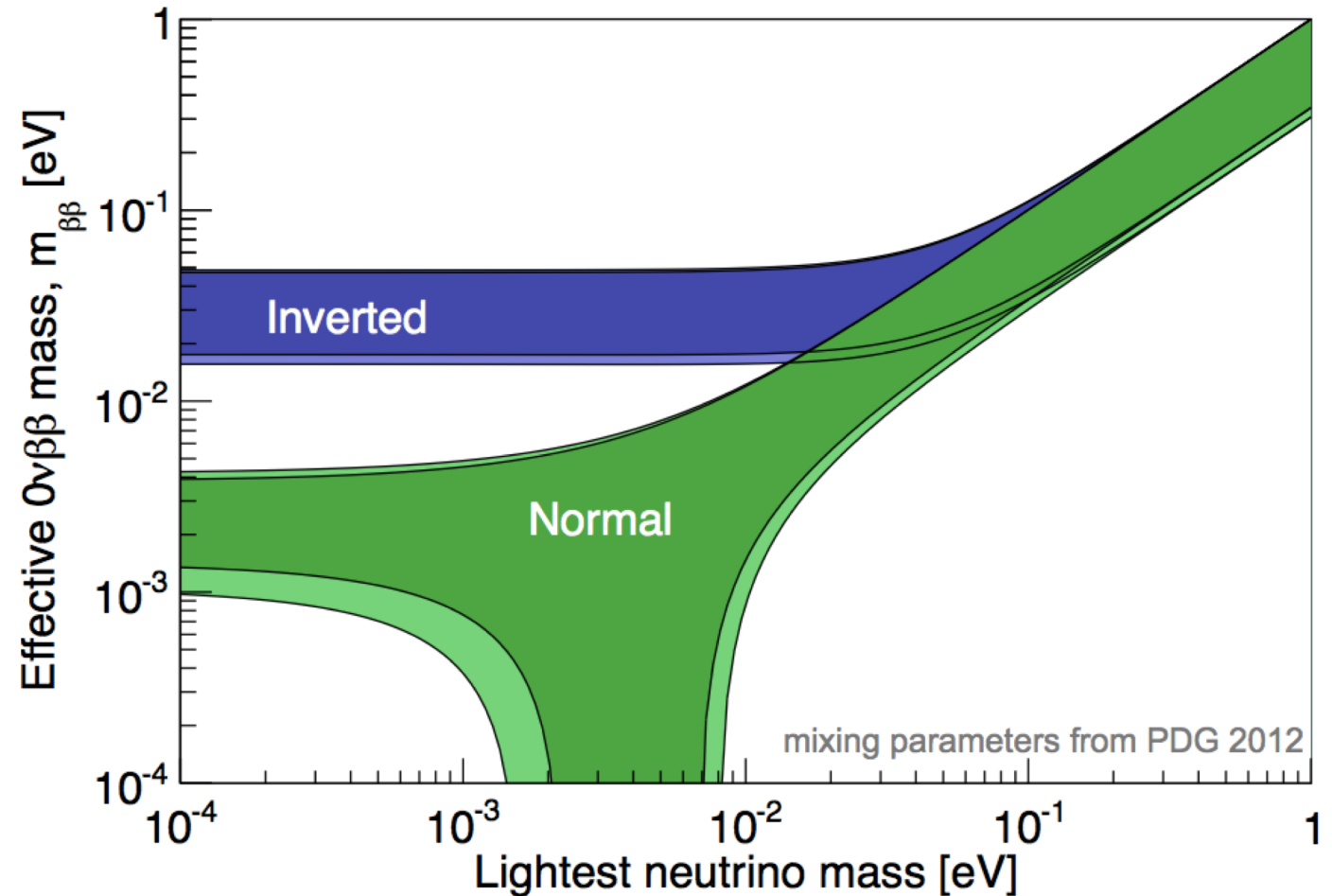
- $\nu$  mixing angles
- $\nu$  masses
- mass hierarchy
- 2 totally unknown phases

# Interpretation of Half-Life Sensitivity

$$(T_{1/2}^{0\nu})^{-1} = G^{0\nu} |M_{0\nu}|^2 \left( \frac{\langle m_{\beta\beta} \rangle}{m_e} \right)^2$$

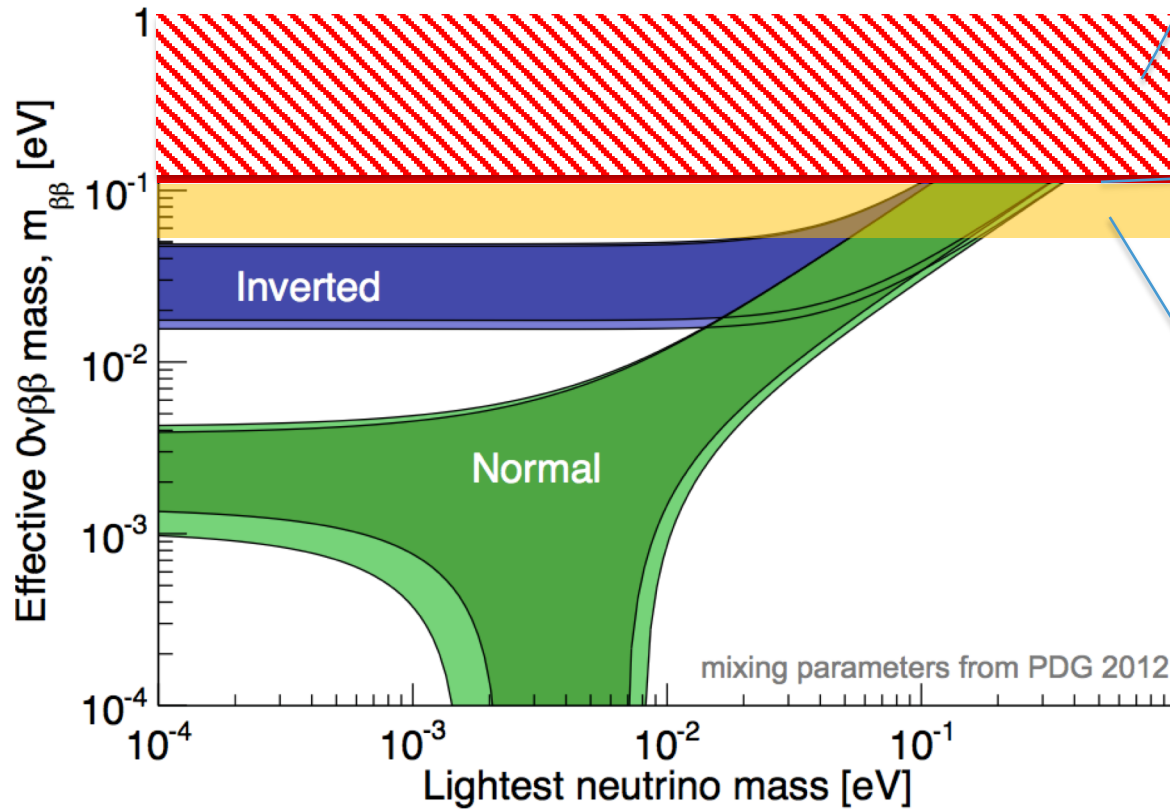
$$\langle m_{\beta\beta} \rangle = \left| \sum_{i=1}^3 U_{ei}^2 m_i \right|$$

- Light Majorana neutrino exchange: assumes new physics is at GUT scale,  $0\nu\beta\beta$  mediated by dim. 5 operator
- Used to compare and set goals for future experiments



# Translating Half-Life to $m_{\beta\beta}$

- Need to use a particular model, the phase space factor and a nuclear matrix element to turn half-life into  $m_{\beta\beta}$
- Results are generally reported for the full set of NMEs, so the upper limit in  $m_{\beta\beta}$  has a range



This area excluded

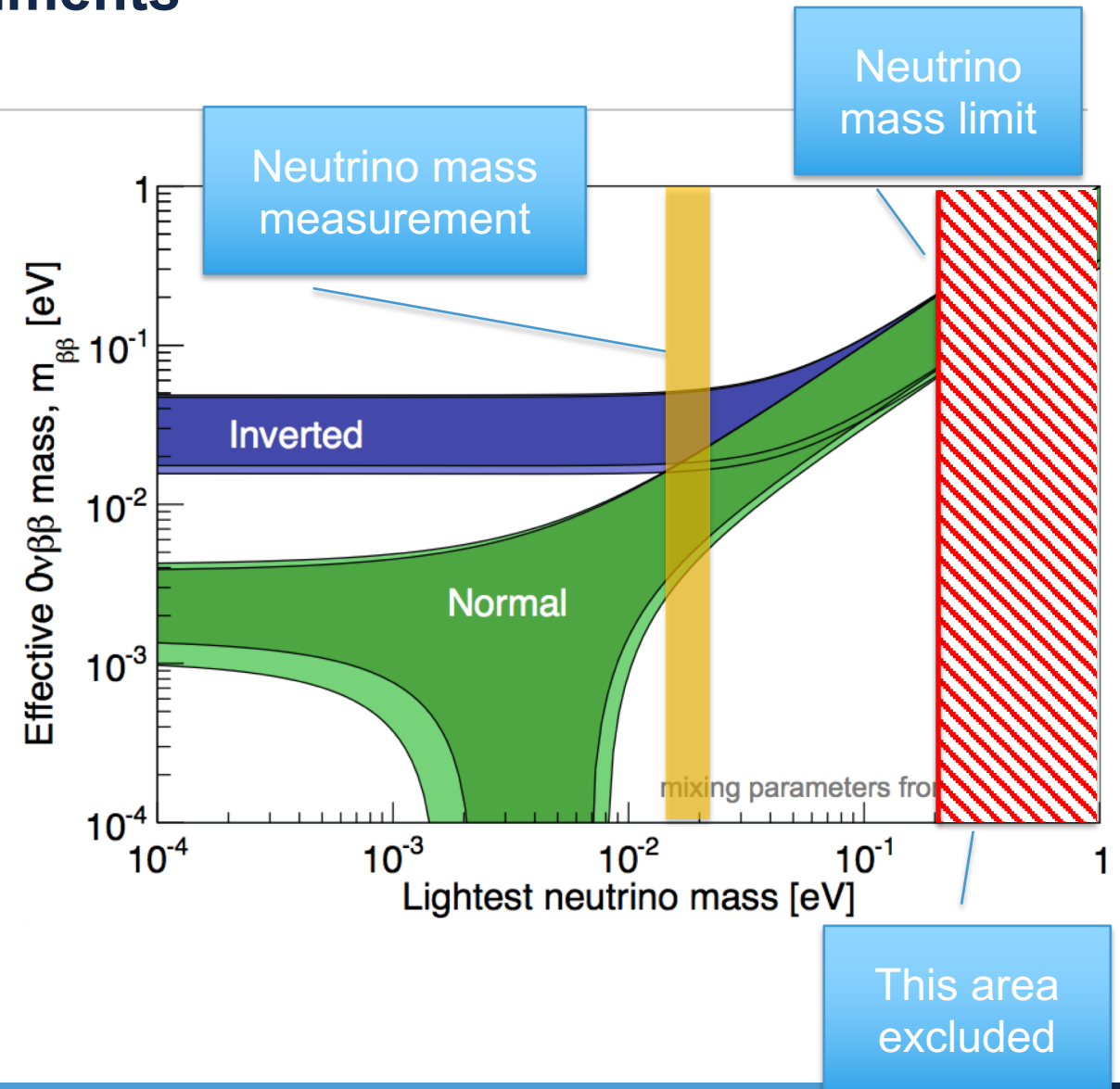
Upper limit for a single NME

Excluded for larger NME



## Information from Other Neutrino Experiments

- Light-colored edges are  $3\sigma$  uncertainty on neutrino mixing and mass splittings
- Measuring hierarchy would tell us which branch we need to look in
- Mass measurement would tell us which vertical band to look in



## Discovery and Sensitivity

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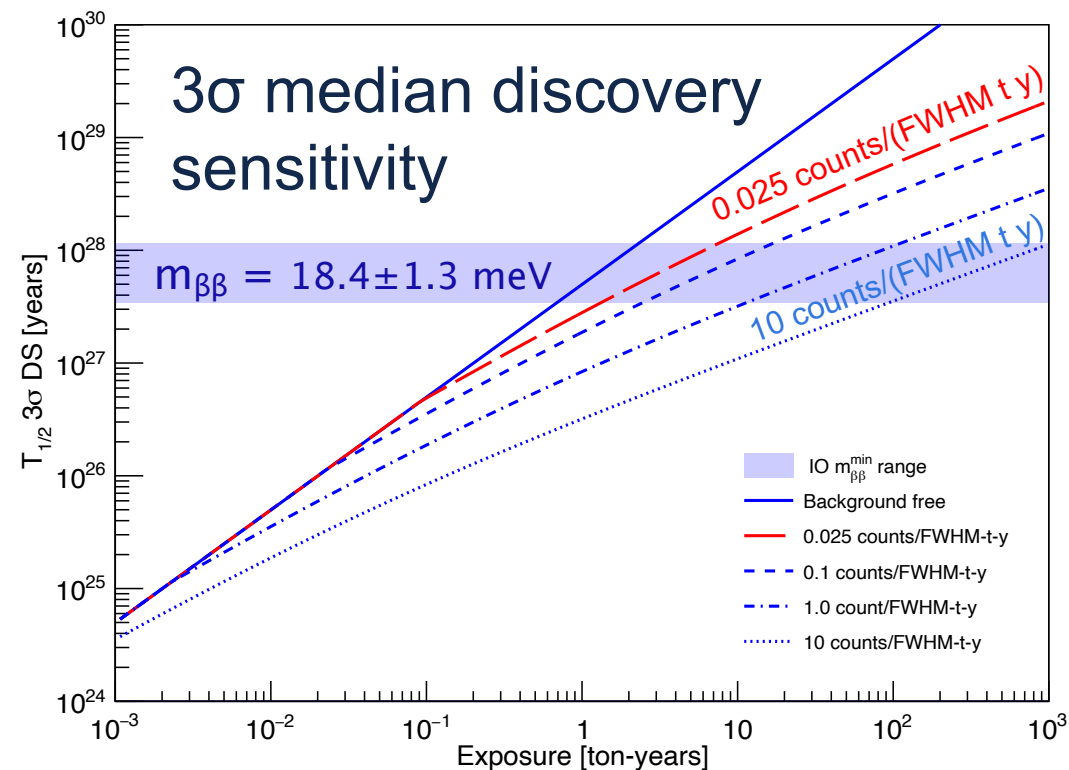
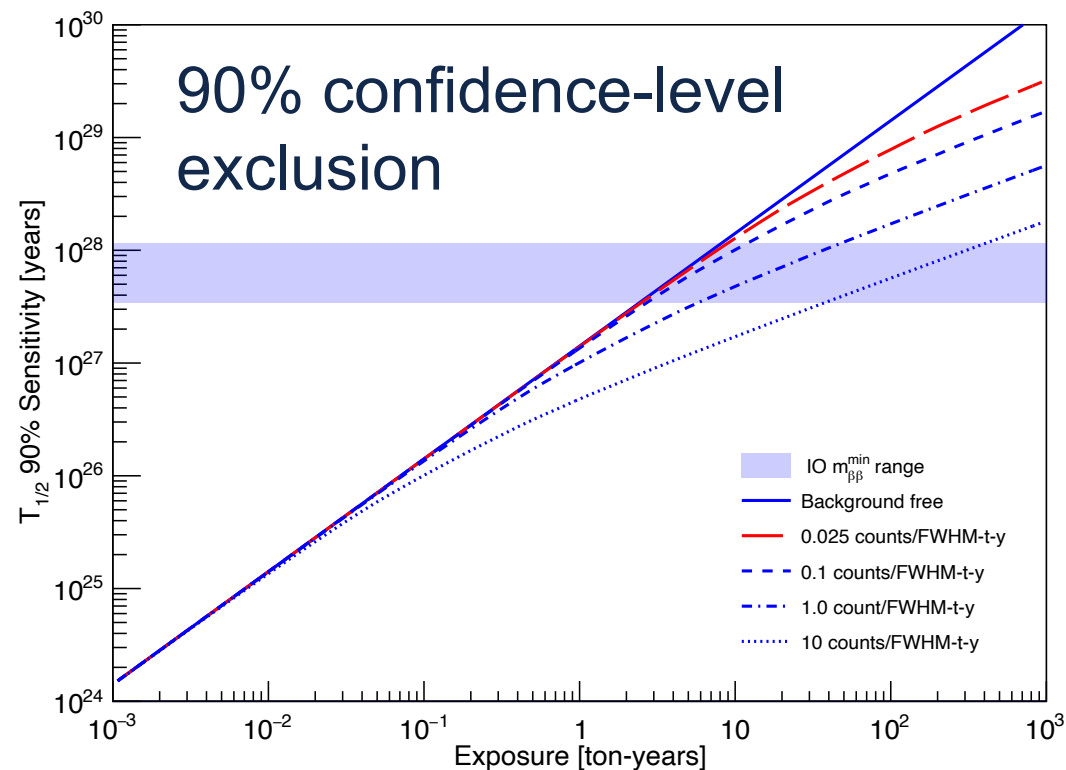
After you run a  $0\nu\beta\beta$  search...

- You either see an excess at the Q value, and fit a peak with some rate to it.
- Or you don't see an excess. In that case, you set a lower limit on half-life:

The diagram illustrates the equation for the lower limit on half-life,  $T_{1/2}^{0\nu} > \ln(2) \frac{N_a T \epsilon}{S}$ , with four blue boxes connected to its terms by lines:

- Number of  $\beta\beta$  atoms** is connected to  $N_a$ .
- Lifetime** is connected to  $T$ .
- Efficiency** is connected to  $\epsilon$ .
- Upper limit on the number of signal counts** is connected to  $S$ .

# Sensitivity vs. Discovery



Background demands are more stringent if you want to make a discovery

## Reaching Ultra-Long Half-Life

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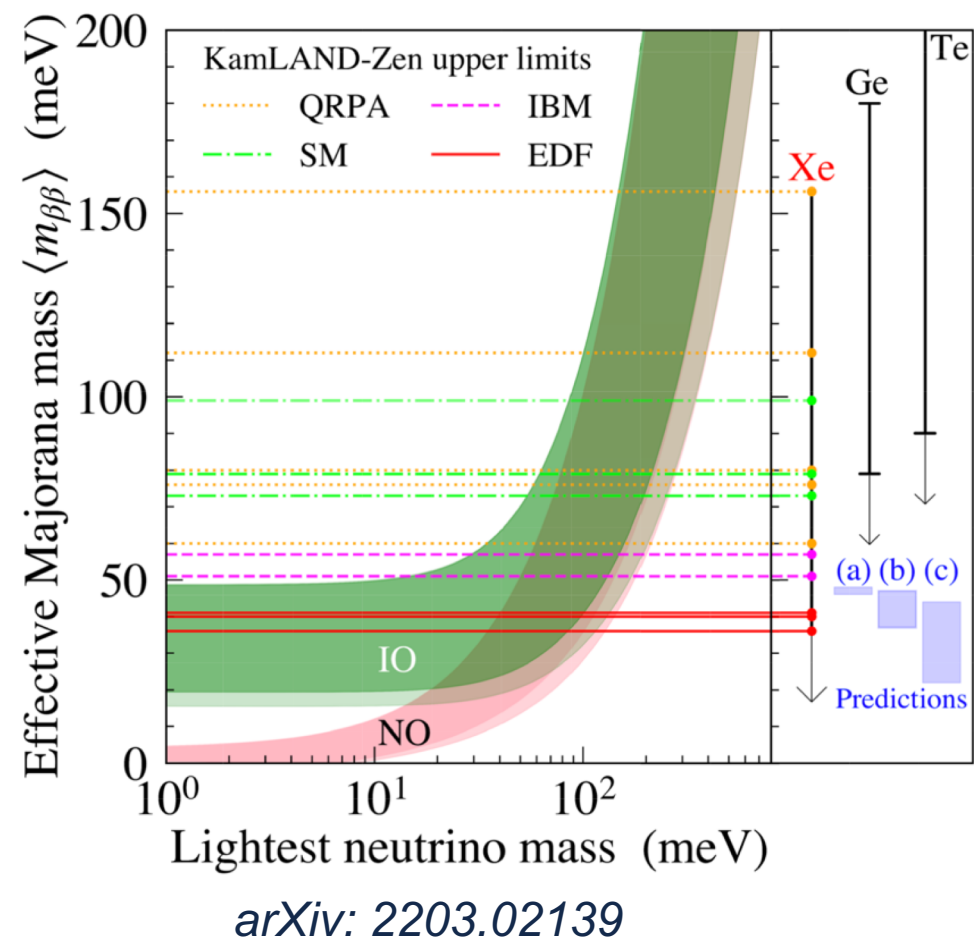
- Best-case scenario: quasi-background-free experiment,  $3\sigma = 3$  counts
- Long half-lives mean you need large exposures. For 3-4 counts of  $0\nu\beta\beta$  at...
  - $10^{26}$  years: 100 kg-years
  - $10^{27}$  years: 1 ton-year
  - $10^{28}$  years: 10 ton-years
- Goal of the next generation of experiments: cover the bottom of the IO region in discovery mode for most nuclear matrix elements
- Implies required discovery sensitivities of  $10^{27}$  to  $10^{28}$  years

For higher backgrounds,  
required exposure  
increases accordingly

Current status and near-term future



# Current Best Limits on $0\nu\beta\beta$



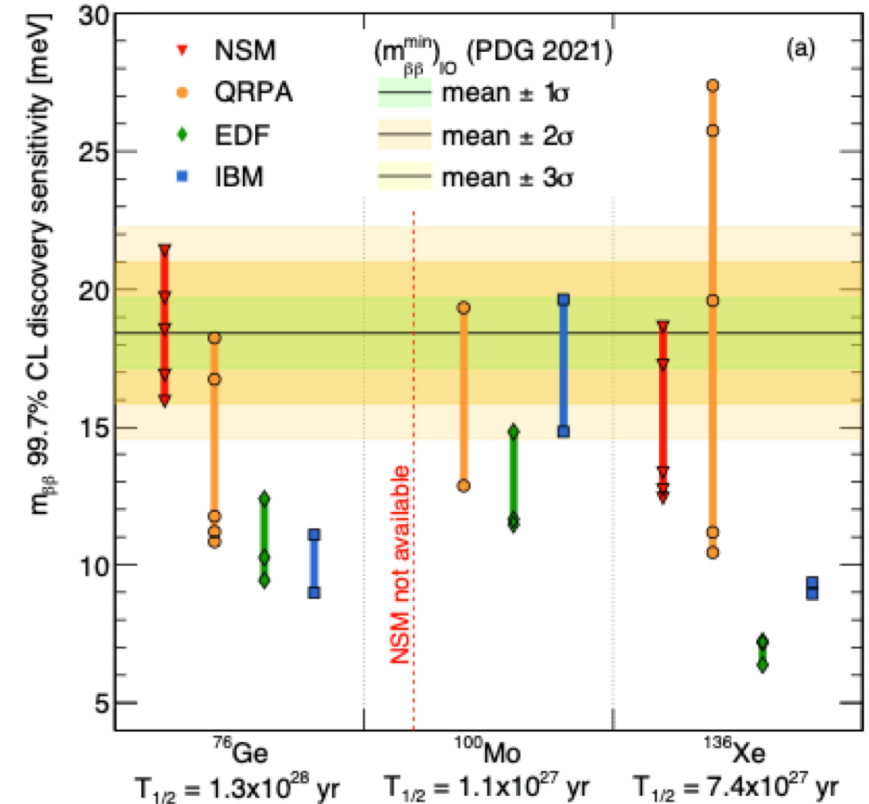
Experiment	Isotope	Exposure [kg yr]	$T_{1/2}^{0\nu}$ [ $10^{25}$ yr]	$m_{\beta\beta}$ [meV]
Gerda	$^{76}\text{Ge}$	127.2	18	79-180
Majorana	$^{76}\text{Ge}$	26	8.3	113-269
KamLAND-Zen	$^{136}\text{Xe}$	970	23	36-156
EXO-200	$^{136}\text{Xe}$	234.1	3.5	93-286
CUORE	$^{130}\text{Te}$	1038.4	2.2	90-305

NSAC recommendation: quote a range of  $m_{\beta\beta}$  using the largest and smallest available NME from the 4 main calculation methods;  $g_A=1.27$ ; no contribution from the contact term

# The Ton-Scale Generation

- Covering the IO in discovery mode requires  $O(1 \text{ ton})$  of isotope
- 3 candidate experiments with US participation, in addition to other ongoing efforts: LEGEND, nEXO, and CUPID
- All 3 experiments cover the IO for some matrix elements, and miss for others
- All 3 were evaluated by the DOE in Summer 2021. DOE-NP is seeking international support to pursue all 3 experiments.

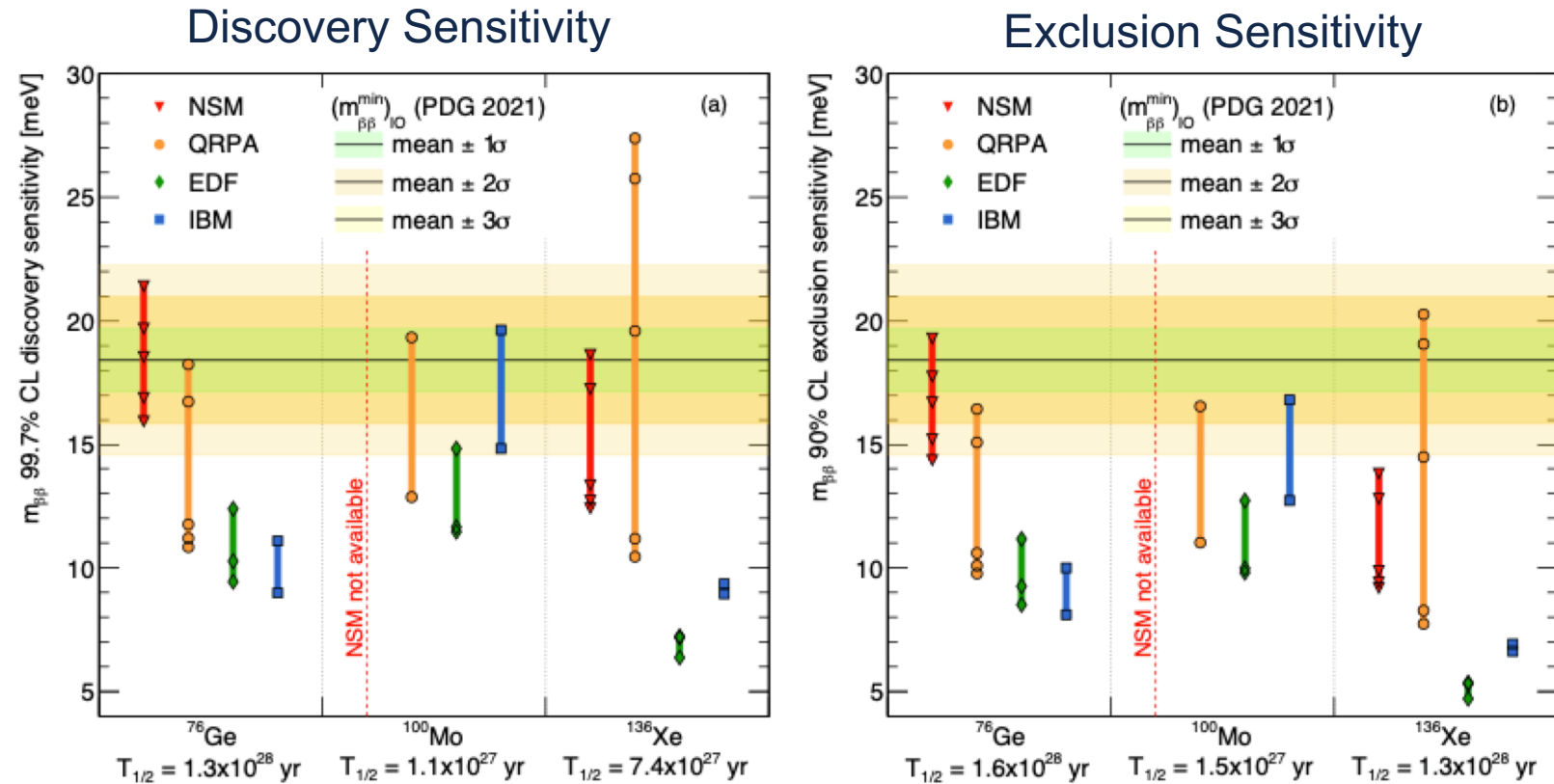
Discovery Sensitivity for the “Big 3”



10.1103/PhysRevC.104.L042501

# Discovery and Sensitivity

- Larger background = more difference between discovery and exclusion
- Liquid scintillator experiments will have competitive sensitivity, but generally don't publish discovery projections:
  - KL2Z:  $T_{1/2} > 2 \times 10^{27}$
  - SNO+ with increased Te loading:  $T_{1/2} > 1 \times 10^{27}$



From Agostini et al., PRC **104**, L042501 (2021)

## Timeline for Ton-Scale Experiments

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- Depends on funding availability
- The 3 ton-scale experiments are moving towards CD-1, hoping for projected construction start in ~2024-2025
- Construction estimate: ~5 to 10 years\*
- All plan for 10 years of running, full-exposure results in ~2040

\* Just my guess,  $0\nu\beta\beta$  collaborations may disagree

# Conclusion

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- $0\nu\beta\beta$  is some of the most exciting physics we can look for! It could provide insight into...
  - The origin of neutrino mass
  - The mechanism that drove baryogenesis
  - The origin of flavor/particle generations, dark matter, etc...
- Regardless of the mechanism,  $0\nu\beta\beta$  would be a direct observation of lepton number violation and prove that neutrinos have Majorana mass
- The coming generation of experiments is exploring very rich parameter space and (hopefully) beginning very soon
- For a competitive  $0\nu\beta\beta$  search, an experiment would need to have high  $\beta\beta$  isotope mass and ultra-low backgrounds at  $Q_{\beta\beta}$