

# Prospects on theory inputs for the first-row CKM unitarity tests

***Chien-Yeah Seng***

**Helmholtz-Institut für Strahlen- und Kernphysik  
and**

**Bethe Center for Theoretical Physics,  
Universität Bonn**

[cseng@hiskp.uni-bonn.de](mailto:cseng@hiskp.uni-bonn.de)

Snowmass Community Summer Study Workshop, University of Washington, Seattle

19 July, 2022

# Anomalies in beta decays

Beta decays place **one of the most stringent tests of SM** through precision measurements of the **first-row CKM matrix elements  $V_{ud}$  and  $V_{us}$**

$V_{ud}$

	$ V_{ud} $
Superaligned nuclear decays ( $0^+ \rightarrow 0^+$ )	0.97373(31)
Free $n$ decay	0.97377(90)
Mirror nuclei decays	0.9739(10)
Pion semileptonic decay ( $\pi_{e3}$ )	0.9740(28)

$V_{us}$

	$ V_{us} $
Kaon semileptonic decays ( $K_{\ell 3}$ )	0.22308(55)
Tau decays	0.2207(14)
Hyperon decays	0.2250(27)

HFLAV,  
2206.07501

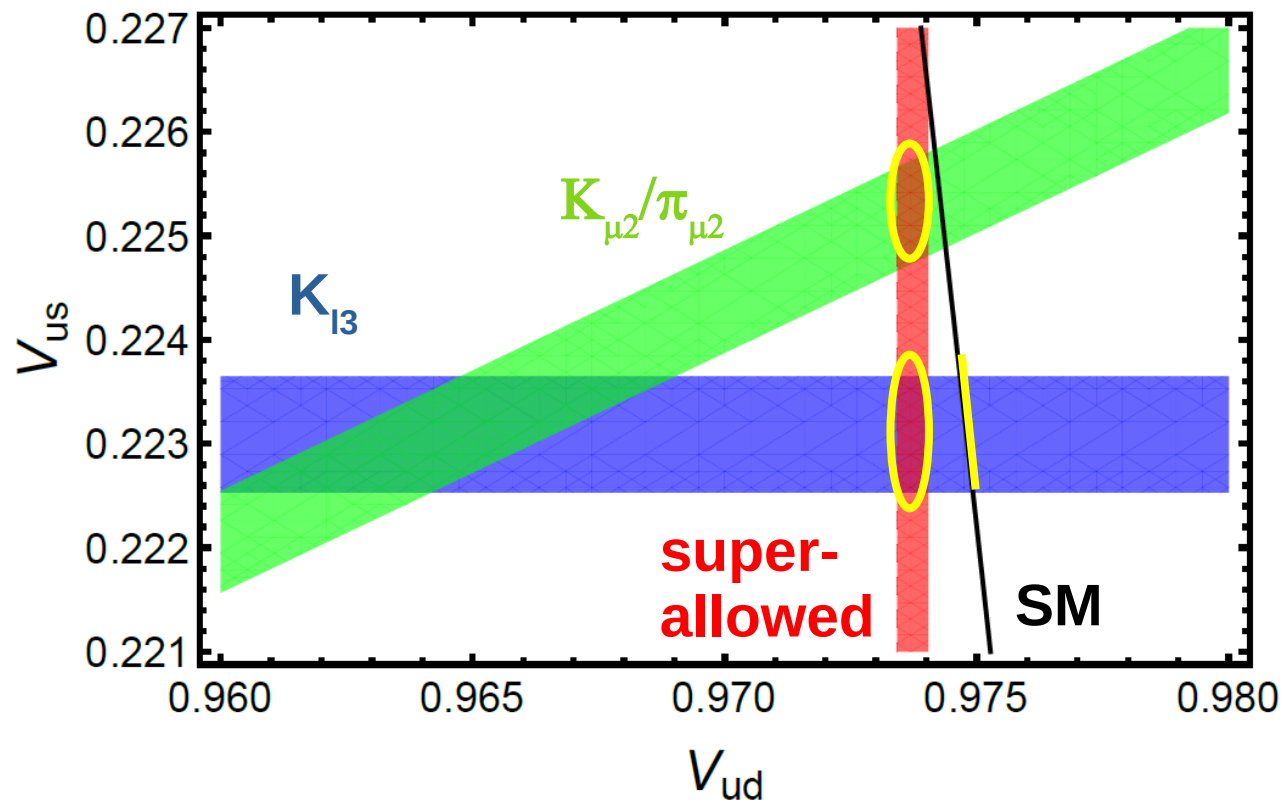
$V_{us}/V_{ud}$

	$ V_{us}/V_{ud} $
$K/\pi$ leptonic decays ( $K_{\mu 2}/\pi_{\mu 2}$ )	0.23131(51)
$K/\pi$ semileptonic decays ( $K_{\ell 3}/\pi_{e 3}$ )	0.22908(87)

# Anomalies in beta decays

Several **anomalies** are recently observed in the **first-row CKM matrix elements**!

**SM prediction:**  $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$



**"Cabibbo Angle Anomaly (CAA)"  $\sim 3\sigma$**

# A glance at the error analysis

**An example:** First-row CKM unitarity with  $|V_{ud}|$  from  $0^+$  beta decay  
and  $|V_{us}|$  from  $K_{\ell 3}$  decay

$$|V_{ud}|_{0^+}^2 + |V_{us}|_{K_{\ell 3}}^2 + \cancel{|V_{ub}|^2} - 1 = -0.0021(7)$$

## SOURCES OF UNCERTAINTY:

$ V_{ud} _{0^+}^2 +  V_{us} _{K_{\ell 3}}^2 - 1$	$-2.1 \times 10^{-3}$
$\delta V_{ud} _{0^+}^2, \text{ exp}$	$2.1 \times 10^{-4}$
$\delta V_{ud} _{0^+}^2, \text{ RC}$	$1.8 \times 10^{-4}$
$\delta V_{ud} _{0^+}^2, \text{ NS}$	$5.3 \times 10^{-4}$
$\delta V_{us} _{K_{\ell 3}}^2, \text{ exp+th}$	$1.8 \times 10^{-4}$
$\delta V_{us} _{K_{\ell 3}}^2, \text{ lat}$	$1.7 \times 10^{-4}$
Total uncertainty	$6.5 \times 10^{-4}$
Significance level	$3.2\sigma$

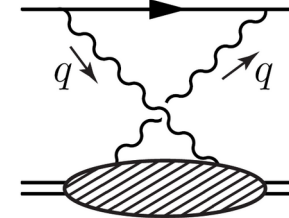
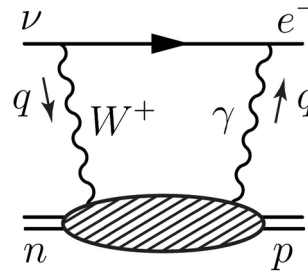
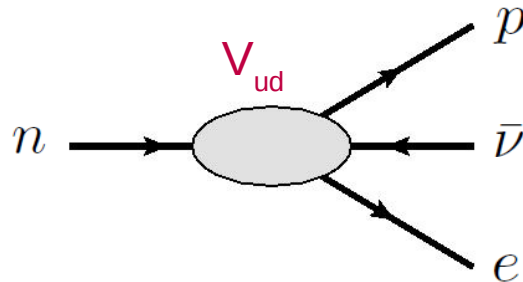
*CYS, Galviz, Marciano and Meißner, 2022 PRD*

$V_{ud}$ : Theory errors dominate

$V_{us}$ : Theory/experimental errors comparable

# Single-nucleon radiative corrections (RC)

Primary source of uncertainty: the “single-nucleon axial  $\gamma W$ -box diagram”

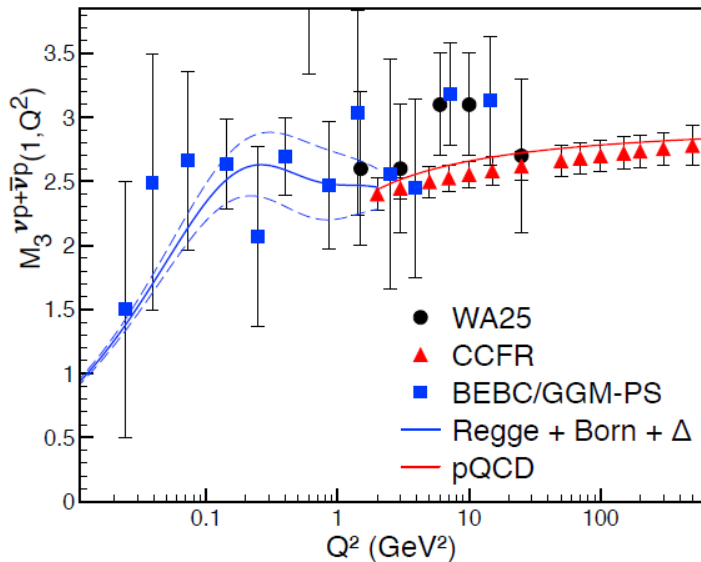


$$Q^2 = -q^2$$

**Dispersive representation:**

*CYS, Gorchtein, Patel and Ramsey-Musolf, 2018 PRL*

$$\square_{\gamma W}^V = \frac{\alpha_{em}}{\pi g_V} \int_0^\infty \frac{dQ^2}{Q^2} \frac{M_W^2}{M_W^2 + Q^2} \int_0^1 dx \frac{1 + 2r}{(1 + r)^2} F_3^{(0)}(x, Q^2)$$



Data input: **Parity-odd structure function  $F_3$**   
from **neutrino-nucleus scattering**

**$|V_{ud}|$ : 0.97420(21)  $\rightarrow$  0.97370(14)**

*Pre-2018*

*2018*

Confirmation by independent studies:

*Czarnecki, Marciano and Sirlin, 2019 PRD*

*CYS, Feng, Gorchtein and Jin, 2020 PRD*

*Hayen, 2021 PRD*

*Shiells, Blunden and Melnitchouk, 2021 PRD*

# Single-nucleon radiative corrections (RC)

**Major limiting factor** of the DR treatment: **low quality of the neutrino data** in the most interesting region:  $Q^2 \sim 1\text{GeV}^2$

**Ongoing program:** Calculate the box diagram directly with **lattice QCD**

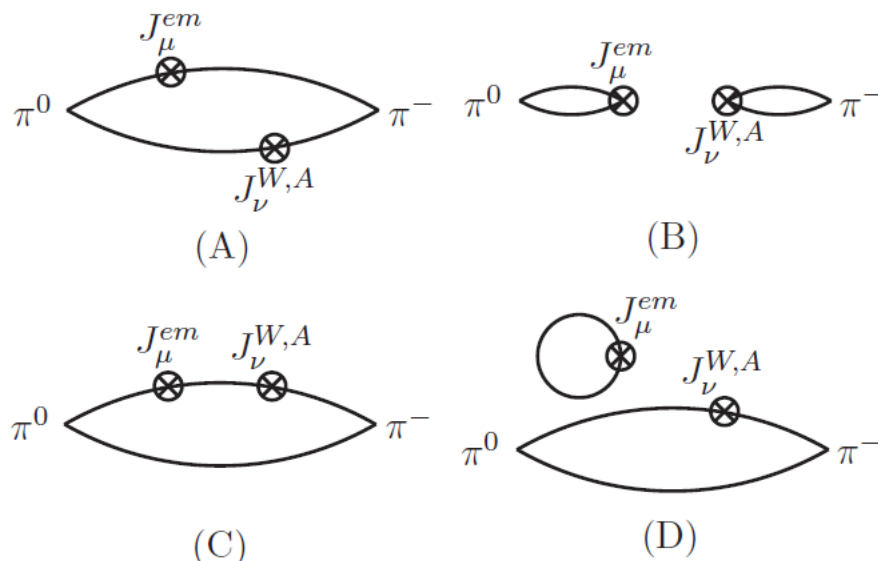
**Year 2020:** First realistic lattice QCD calculation of the simpler **pion axial  $\gamma W$ -box diagram**

*Feng, Gorchtein, Jin, Ma and CYS, 2020 PRL*

## Consequences:

- Significant reduction of the theory uncertainty in **pion semileptonic decay ( $\pi_{e3}$ )**
- Indirect implications on the **free-neutron** axial  $\gamma W$ -box diagram

*CYS, Feng, Gorchtein and Jin, 2020 PRD*



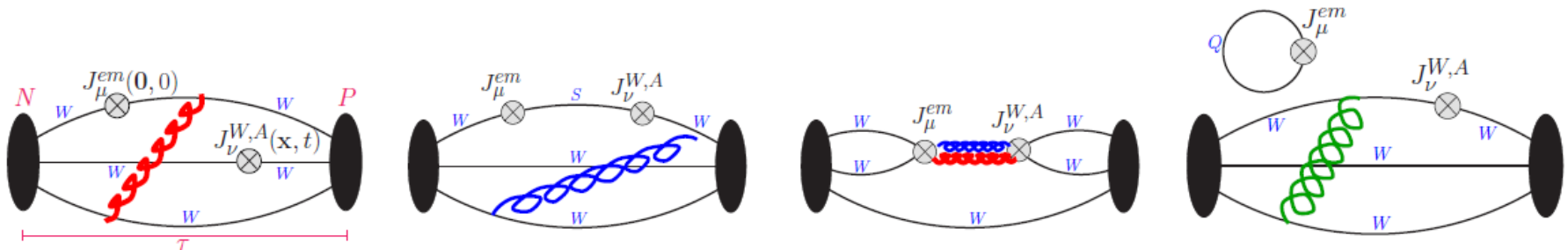
**More in Luchang's talk!**

# Single-nucleon radiative corrections (RC)

**Major limiting factor** of the DR treatment: **low quality of the neutrino data** in the most interesting region:  $Q^2 \sim 1\text{GeV}^2$

**Ongoing program:** Calculate the box diagram directly with **lattice QCD**

**Neutron axial  $\gamma W$ -box diagram** is more complicated, but on the way.



*(R. Gupta, Rare Processes and Precision Frontier Townhall Meeting, 2020)*

Possible alternative approach using **Feynman-Hellmann theorem (FHT)**

*CYS and Meißner, 2019 PRL*

# Nuclear Structure corrections

Superaligned  $0^+ \rightarrow 0^+$  nuclear beta decays provides the best measurement of  $V_{ud}$

Master formula:

$$|V_{ud}|^2 = \frac{2984.43 \text{ s}}{\mathcal{F}t (1 + \Delta_R^V)}.$$

Corrected ft (half-life\*statistical function)-value:

$$\mathcal{F}t = ft (1 + \delta'_R) (1 + \delta_{NS} - \delta_C)$$

Measured ft-value

Nucleus-dependent "outer corrections" (under control)

Nuclear structure effects in inner RC

Isospin-breaking corrections

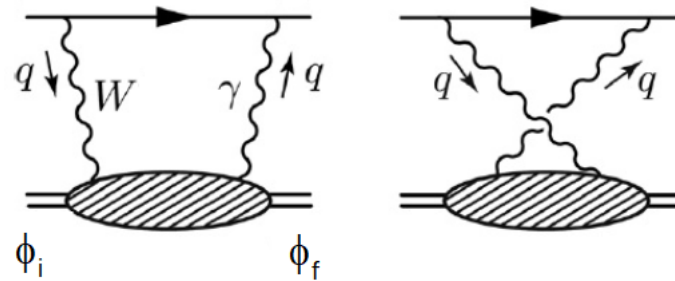
Best-measured decays:

$T_Z = -1$
$^{10}\text{C} \rightarrow ^{10}\text{B}$
$^{14}\text{O} \rightarrow ^{14}\text{N}$
$^{22}\text{Mg} \rightarrow ^{22}\text{Na}$
$^{26}\text{Si} \rightarrow ^{26}\text{Al}$
$^{34}\text{Ar} \rightarrow ^{34}\text{Cl}$
$^{38}\text{Ca} \rightarrow ^{38}\text{K}$
$T_Z = 0$
$^{26m}\text{Al} \rightarrow ^{26}\text{Mg}$
$^{34}\text{Cl} \rightarrow ^{34}\text{S}$
$^{38m}\text{K} \rightarrow ^{38}\text{Ar}$
$^{42}\text{Sc} \rightarrow ^{42}\text{Ca}$
$^{46}\text{V} \rightarrow ^{46}\text{Ti}$
$^{50}\text{Mn} \rightarrow ^{50}\text{Cr}$
$^{54}\text{Co} \rightarrow ^{54}\text{Fe}$
$^{62}\text{Ga} \rightarrow ^{62}\text{Zn}$
$^{74}\text{Rb} \rightarrow ^{74}\text{Kr}$



# Nuclear Structure corrections

$\delta_{\text{NS}}$ : nuclear modifications of the free-nucleon  $\gamma W$  box diagram



**LARGEST** source of uncertainty in  $V_{ud}$ !

$$\square_{\gamma W}^{\text{nucl.}} = \square_{\gamma W}^n + \underbrace{[\square_{\gamma W}^{\text{nucl.}} - \square_{\gamma W}^n]}$$

**Ongoing program:** Nuclear structure functions at low  $Q^2$  with **ab-initio** methods

$$\begin{aligned} & \frac{1}{8\pi} \int d^4x e^{iq \cdot x} \langle \phi_f(p) | [J_{\text{em}}^\mu(x), J_W^{\dagger\nu}(0)] | \phi_i(p) \rangle \\ &= \left( -g^{\mu\nu} + \frac{q^\mu q^\nu}{q^2} \right) \underline{F_1} + \frac{\hat{p}^\mu \hat{p}^\nu}{p \cdot q} \underline{F_2} - \frac{i\varepsilon^{\mu\nu\alpha\beta} q_\alpha p_\beta}{2p \cdot q} \underline{F_3} \end{aligned}$$

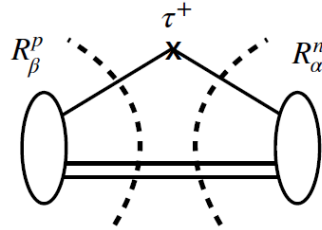
Inputs to the dispersion integral

**Light nuclei:** Quantum Monte Carlo, No-Core Shell Model, ...

**Medium-size nuclei:** Coupled-Cluster method, ....

# Nuclear Structure corrections

$\delta_C$ : **isospin-breaking (ISB) corrections to nuclear wavefunctions**



Mainly due to **Coulomb interaction** between protons

$$V_C = V_C^{(0)} + V_C^{(1)}$$

Computed systematically within **shell model** (Hardy-Towner), but were questioned in several aspects:

1. Theoretical inconsistencies *Miller and Schwenk, 2008 PRC, 2009 PRC; Condren and Miller, 2201.10651*
2. Cannot be reproduced by other nuclear theory calculations, which generically predict **smaller**  $\delta_C$  *Hartree-Fock, DFT, RPA, isovector monopole resonance...*

**New Thoughts:** Possible relation to the **neutron skin** of the stable nuclei?

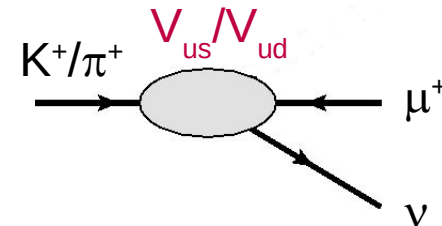
$$\delta_C \propto \frac{1}{3} \sum_a \frac{|\langle a; 0 || V_C^{(1)} || g; 1 \rangle|^2}{(E_{a,0} - E_g)^2} + \frac{1}{2} \sum_{a \neq g} \frac{|\langle a; 1 || V_C^{(1)} || g; 1 \rangle|^2}{(E_{a,1} - E_g)^2} - \frac{5}{6} \sum_a \frac{|\langle a; 2 || V_C^{(1)} || g; 1 \rangle|^2}{(E_{a,2} - E_g)^2}$$

$$(R_n - R_p)_{V_C^{(1)}} \propto \sum_{a \neq g} \frac{|\langle a; 1 || V_C^{(1)} || g; 1 \rangle|^2}{E_{a,1} - E_g} + \sum_a \frac{|\langle a; 2 || V_C^{(1)} || g; 1 \rangle|^2}{E_{a,2} - E_g}$$

# Kaon/pion leptonic decay ( $K_{\mu 2}/\pi_{\mu 2}$ )

$$\frac{|V_{us}|f_{K^+}}{|V_{ud}|f_{\pi^+}} = \underbrace{\left[ \frac{\Gamma_{K_{\mu 2}} M_{\pi^+}}{\Gamma_{\pi_{\mu 2}} M_{K^+}} \right]^{1/2}}_{\text{"axial ratio" } R_A} \frac{1 - m_\mu^2/M_{\pi^+}^2}{1 - m_\mu^2/M_{K^+}^2} (1 - \delta_{\text{EM}}/2)$$

**"axial ratio"  $R_A$**  *Marciano, 2004 PRL; Cirigliano and Neufeld, 2011 PLB*



**Lattice QCD inputs:**  $K^+/\pi^+$  decay constants

$$N_f = 2 + 1 + 1 \quad : \quad f_{K^+}/f_{\pi^+} = 1.1932(21)$$

$$N_f = 2 + 1 \quad : \quad f_{K^+}/f_{\pi^+} = 1.1917(37)$$

$$N_f = 2 \quad : \quad f_{K^+}/f_{\pi^+} = 1.205(18)$$

*FLAG 2021*

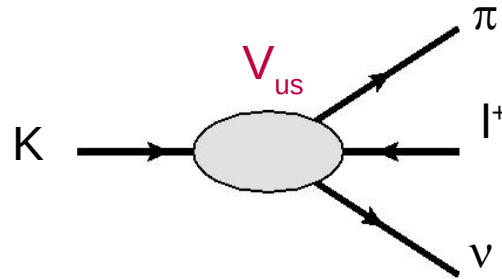
**Electromagnetic RC in ChPT:**  $\delta_{\text{EM}} = \delta_{\text{EM}}^K - \delta_{\text{EM}}^\pi = -0.0069(17)$  *Knecht et al., 2000 EPJC; Cirigliano and Neufeld, 2011 PLB*

Advantage: **LECs cancel in the ratio**

**Direct lattice QCD calculation** of the EMRC+isospin breaking correction (contained in the physical  $K^+/\pi^+$  decay constants) consistent with ChPT result, with slightly lower uncertainty *Giusti et al, 2018 PRL*

**Total:**  $|V_{us}/V_{ud}| = 0.23131(41)_{\text{lat}}(24)_{\text{exp}}(19)_{\text{RC}}$

# Kaon semileptonic decays ( $K_{l3}$ )



$$\Gamma_{K_{\ell 3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} S_{\text{EW}} \underbrace{|f_+^{K^0 \pi^-}(0)|^2}_{\text{purple}} \underbrace{I_{K\ell}^{(0)}}_{\text{blue}} \left( 1 + \underbrace{\delta_{\text{EM}}^{K\ell}}_{\text{red}} + \underbrace{\delta_{\text{SU}(2)}^{K\pi}}_{\text{green}} \right)$$

Measurements of **branching ratio** exist in all channels (most recent:  $K_{\mu 3}^S$ )

Theory Inputs:

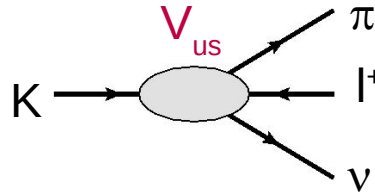
**K $\pi$  form factor at t=0**

**Phase-space factor**

**Long-distance electromagnetic RC**

**ISB correction**

# Kaon semileptonic decays ( $K_{l3}$ )



Master formula:

$$\Gamma_{K_{l3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} S_{EW} |f_+^{K^0\pi^-}(0)|^2 I_{K\ell}^{(0)} \left(1 + \delta_{EM}^{K\ell} + \delta_{SU(2)}^{K\pi}\right)$$

**$K\pi$  form factor at  $t=0$ :**  $\langle \pi^-(p') | J_W^\mu | K^0(p) \rangle = f_+^{K^0\pi^-}(t)(p+p')^\mu + f_-^{K^0\pi^-}(t)(p-p')^\mu$

Lattice QCD inputs:

$$N_f = 2 + 1 + 1 : f_+(0) = 0.9698(17)$$

$$N_f = 2 + 1 : f_+(0) = 0.9677(27)$$

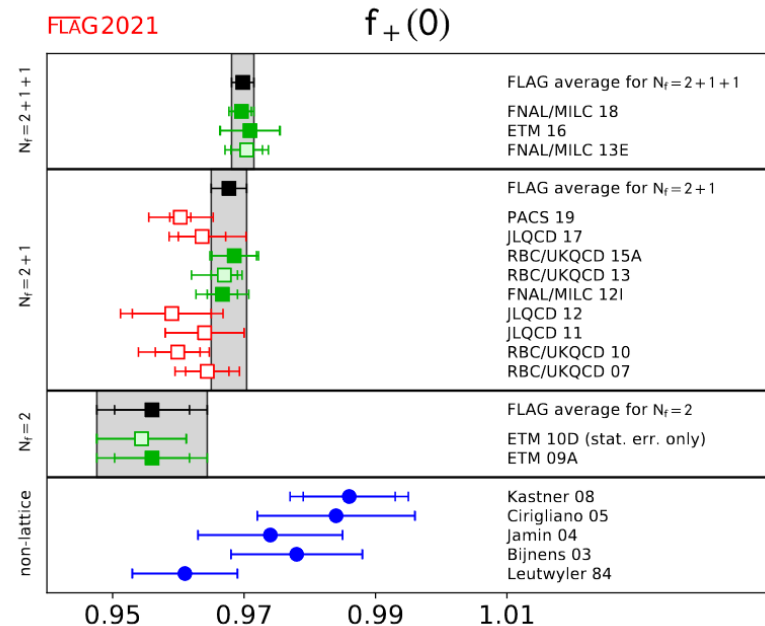
$$N_f = 2 : f_+(0) = 0.9560(57)(62)$$

FLAG 2021

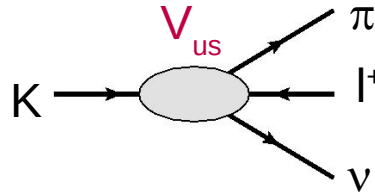
New result from PACS ( $N_f=2+1$ ):

$$f_+(0) = 0.9615(10)(^{+47}_{-2})(5)$$

*Ishikawa et al, 2206.08654*



# Kaon semileptonic decays ( $K_{l3}$ )



Master formula:

$$\Gamma_{K_{\ell 3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} S_{\text{EW}} |f_+^{K^0 \pi^-}(0)|^2 \underbrace{I_{K\ell}^{(0)}}_{\text{Phase-space factor}} \left(1 + \delta_{\text{EM}}^{K\ell} + \delta_{\text{SU}(2)}^{K\pi}\right)$$

**Phase-space factor:** 
$$I_{K\ell}^{(0)} = \int_{m_\ell^2}^{(M_K - M_\pi)^2} \frac{dt}{M_K^8} \bar{\lambda}^{3/2} \left(1 + \frac{m_\ell^2}{2t}\right) \left(1 - \frac{m_\ell^2}{t}\right)^2 \left[ \bar{f}_+^2(t) + \frac{3m_\ell^2 \Delta_{K\pi}^2}{(2t + m_\ell^2) \bar{\lambda}} \bar{f}_0^2(t) \right]$$

probes the **t-dependence** of the  $K\pi$  form factors.

Obtained by fitting to the  $K_{l3}$  **Dalitz** plot with **specific parameterizations**

**of  $f(t)$**  (Taylor expansion, z-expansion, dispersive parameterization, pole parameterization ...)

Rescaled  
 $K\pi$  form factors

The **dispersive parameterization** currently quotes the smallest uncertainty:

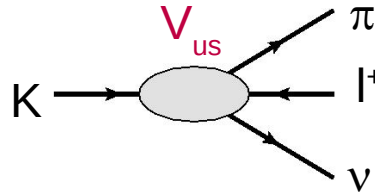
*Future: Direct lattice calculation of the t-dependence?*

Mode	Update
$K_{e3}^0$	<b>0.15470(15)</b>
$K_{e3}^+$	<b>0.15915(15)</b>
$K_{\mu 3}^0$	<b>0.10247(15)</b>
$K_{\mu 3}^+$	<b>0.10553(16)</b>

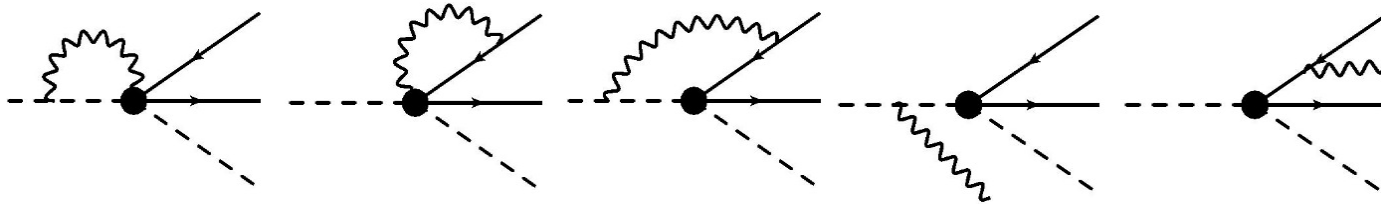
*M. Moulson,  
in the 11<sup>th</sup>  
International  
Workshop on the  
CKM Unitarity  
Triangle, 2021*

# Kaon semileptonic decays ( $K_{l3}$ )

Master formula:



$$\Gamma_{K_{l3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} S_{EW} |f_+^{K^0\pi^-}(0)|^2 I_{K\ell}^{(0)} \left( 1 + \delta_{EM}^{K\ell} + \delta_{SU(2)}^{K\pi} \right)$$

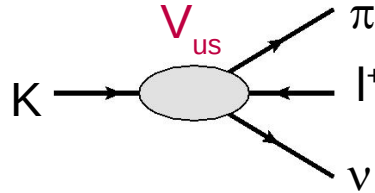


Long-distance electromagnetic RC ( $10^{-3}$ )

	$\delta_{EM}^{K\ell}$ "Sirlin's representation"	ChPT
$K^0 e$	$11.6(2)_{\text{inel}}(1)_{\text{lat}}(1)_{\text{NF}}(2)_{e^2 p^4}$	$9.9(1.9)_{e^2 p^4}(1.1)_{\text{LEC}}$
$K^+ e$	$2.1(2)_{\text{inel}}(1)_{\text{lat}}(4)_{\text{NF}}(1)_{e^2 p^4}$	$1.0(1.9)_{e^2 p^4}(1.6)_{\text{LEC}}$
$K^0 \mu$	$15.4(2)_{\text{inel}}(1)_{\text{lat}}(1)_{\text{NF}}(2)_{\text{LEC}}(2)_{e^2 p^4}$	$14.0(1.9)_{e^2 p^4}(1.1)_{\text{LEC}}$
$K^+ \mu$	$0.5(2)_{\text{inel}}(1)_{\text{lat}}(4)_{\text{NF}}(2)_{\text{LEC}}(2)_{e^2 p^4}$	$0.2(1.9)_{e^2 p^4}(1.6)_{\text{LEC}}$

Lattice inputs:  
see Luchang's  
talk!

# Kaon semileptonic decays ( $K_{l3}$ )



Master formula:

$$\Gamma_{K_{l3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} S_{EW} |f_+^{K^0\pi^-}(0)|^2 I_{K\ell}^{(0)} \left( 1 + \delta_{EM}^{K\ell} + \delta_{SU(2)}^{K\pi} \right)$$

**ISB correction:** presents only in the  $K^+$  channel by construction.

$$\delta_{SU(2)}^{K^+\pi^0} \equiv \left( \frac{f_+^{K^+\pi^0}(0)}{f_+^{K^0\pi^-}(0)} \right)^2 - 1 = \frac{3}{2} \frac{1}{Q^2} \left[ \frac{\hat{M}_K^2}{\hat{M}_\pi^2} + \frac{\chi_{p^4}}{2} \left( 1 + \frac{m_s}{\hat{m}} \right) \right] \quad (\text{neglecting small EM contributions})$$

$$Q^2 = (m_s^2 - \hat{m}^2)/(m_d^2 - m_u^2)$$

**Most recent lattice QCD inputs:** FLAG 2021

$$Q = 23.3(5) , \quad m_s/\hat{m} = 27.42(12) \quad N_f = 2 + 1$$

$$\text{returns: } \delta_{SU(2)}^{K^+\pi^0} = 0.0457(20)$$

Phenomenological inputs from  $\eta \rightarrow 3\pi$  returns a somewhat larger value:

$$\delta_{SU(2)}^{K^+\pi^0} = 0.0522(34)$$



# Summary

---

- Several **anomalies** at the level  $\sim 3\sigma$  have been observed in the measurements of the **first-row CKM matrix elements**  $V_{ud}$  and  $V_{us}$  in **beta decay processes**.
- **SM theory inputs** that require further improvements are:
  - $V_{ud}$  **sector**: RC in single-nucleon and nuclear systems, ISB corrections in nuclear wavefunctions
  - $V_{us}$  **sector**: Lattice inputs of Kaon/pion decay constants and  $K\pi$  form factor, RC in leptonic and semileptonic kaon decays,  $K_{l3}$  phase-space factor, ISB corrections in  $K^\pm$  semileptonic decays
- Successful reduction of theory uncertainties above could increase the significance of the anomalies to more than  $5\sigma$
- Desirable future **experimental improvements**:  $K_{l3}$  and  $\pi_{e3}$  branching ratios, neutron lifetime and  $g_A$ , ...

*Thanks for your attention!*