

# Low-Threshold Detectors for Neutrinos: The Long View

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Neutrino Physics Frontier

Blue Sky/Very Long Term Ideas Session

# Overview

Low-Threshold Detectors = detectors with meV to eV scale thresholds for scattering of weakly interacting particles (neutrinos, DM)

## Outline

DM motivation for low-threshold detectors (“quantum calorimeters”)

Potential science interest of low-threshold detectors for neutrino physics

Backups: Quick summary of architectures under development

## Caveat

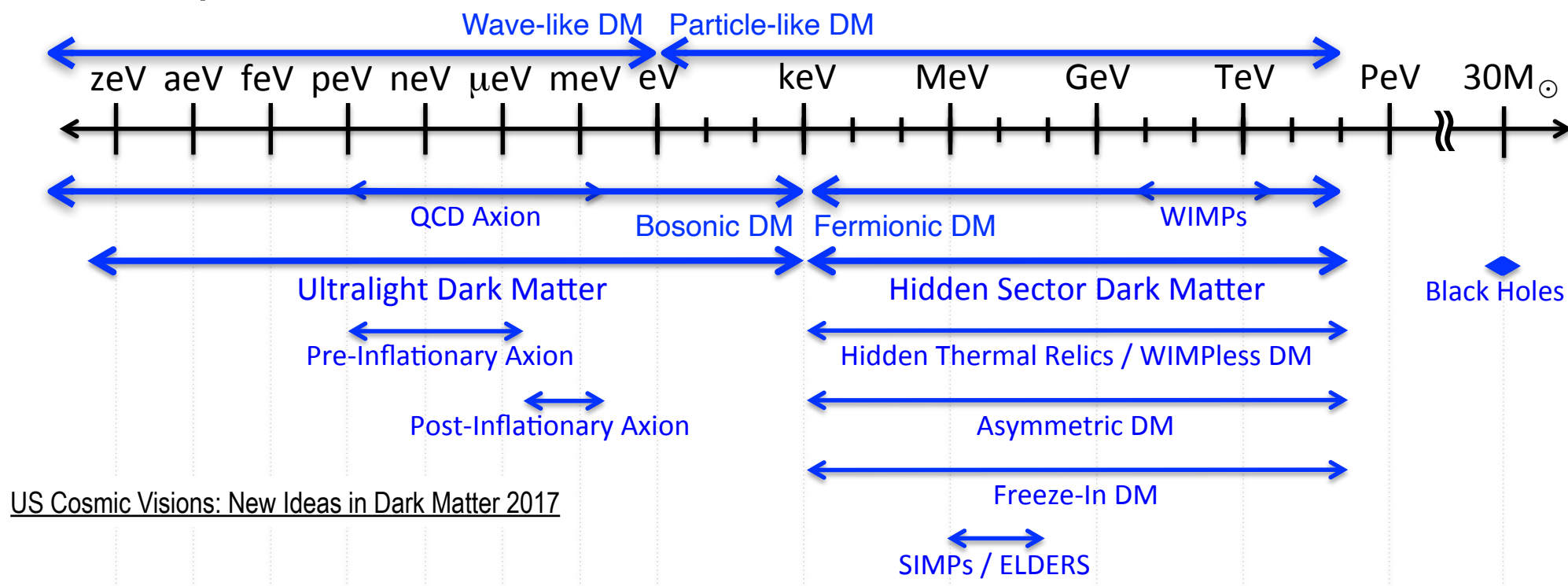
I’m not a neutrino physicist!

But hopefully some developments in DM detection could impact neutrino studies on the 1-2 decade timescale.

For more detail, see Golwala and Figueroa-Feliciano, *Ann Rev Nucl Part Sci*, in press.

# DM Motivation for Low-Threshold Detectors

A new picture for DM:



Why?

No DM in spite of 30 years of dedicated searches in labs, the sky, and accelerators

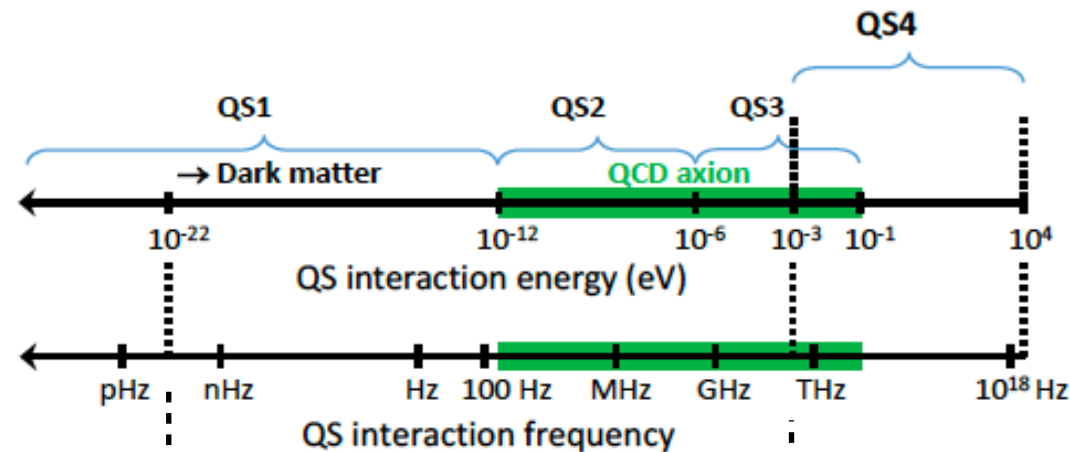
⇒ reconsideration of:

theoretical bias in favor of axions and WIMPs

production mechanisms

cosmological constraints

# Dark Matter → Quantum Sensors



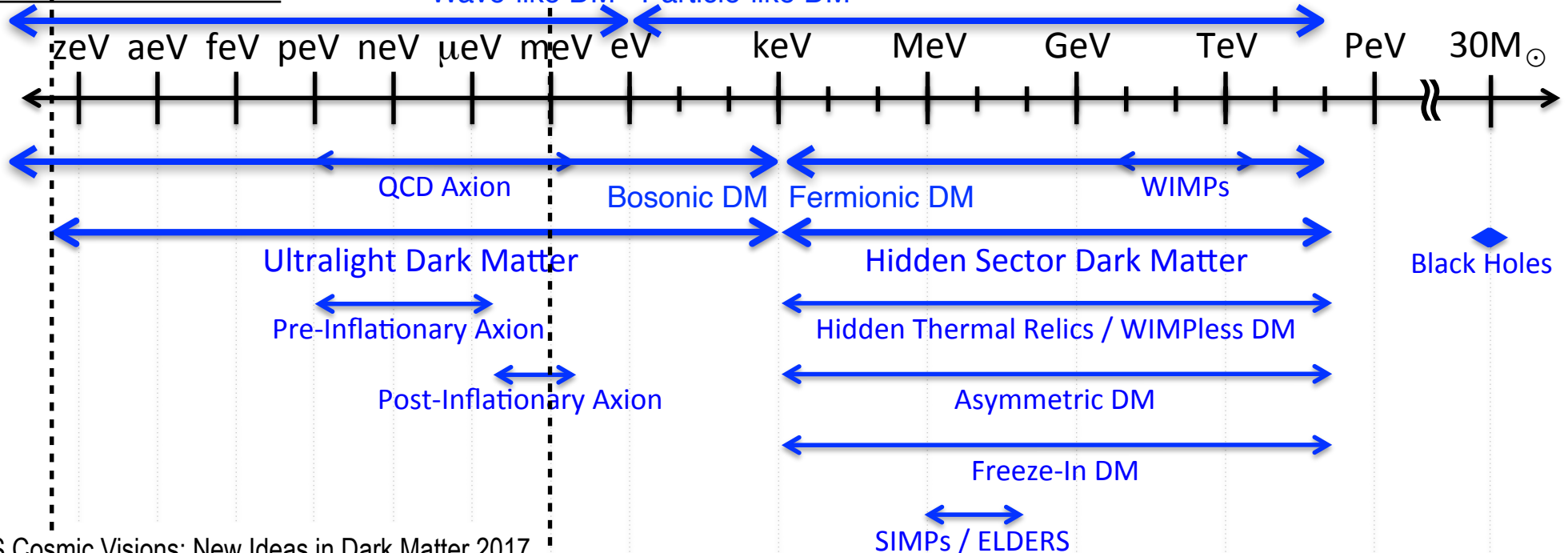
QS for particle physics driven in large part by 2017 US Cosmic Visions *New Ideas in DM* white paper

Seeking to cover  $E_{\text{dep}} \lesssim \text{keV}$  inaccessible to classic particle DM searches

QS regime includes both wave-like (~classical field) and particle-like DM

DOE BRN for HEP Detector R&D

Wave-like DM Particle-like DM

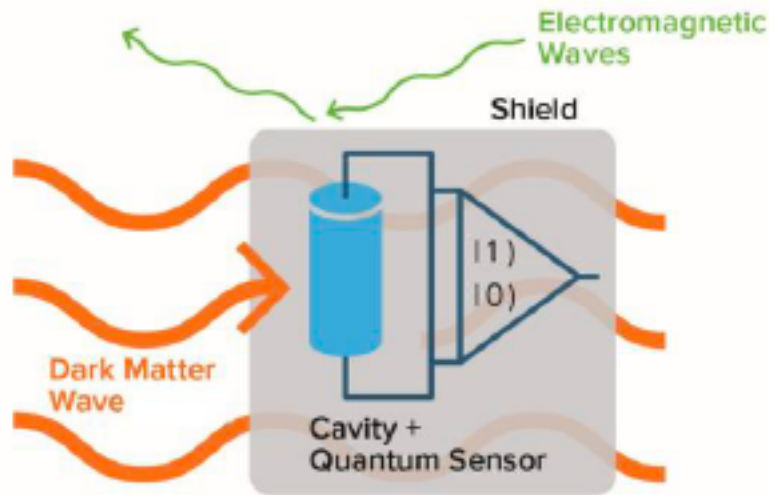


US Cosmic Visions: New Ideas in Dark Matter 2017



# Dark Matter → Quantum Sensors

## Wave-like Dark Matter



Occupation number  $\gg 1$ : difficult to identify individual DM particles, sense “Collective force from macroscopic numbers of particles”

New light scalar (axion-like particle (ALP) or vector (dark photon (DP))

Quantum zero-point fluctuations are the primary obstacle: a quantum problem

Fundamentally quantum approaches needed to circumvent

## Particle-like Dark Matter

relevant for this talk



Occupation number  $\ll 1$ : individual particles can be seen\*

Scattering of fermionic DM

Absorption of heavier ALPs, DPs

“Quantum” effects:

Direct creation of single quantum (not via identifiable particle recoil)

Use of quantum sensors to see the very small energy depositions

eV  $\rightarrow$  meV

\*Occupation number  $> 1$  for bosonic DM does not imply particle-like absorption is not possible, but it makes classical field detection possible.

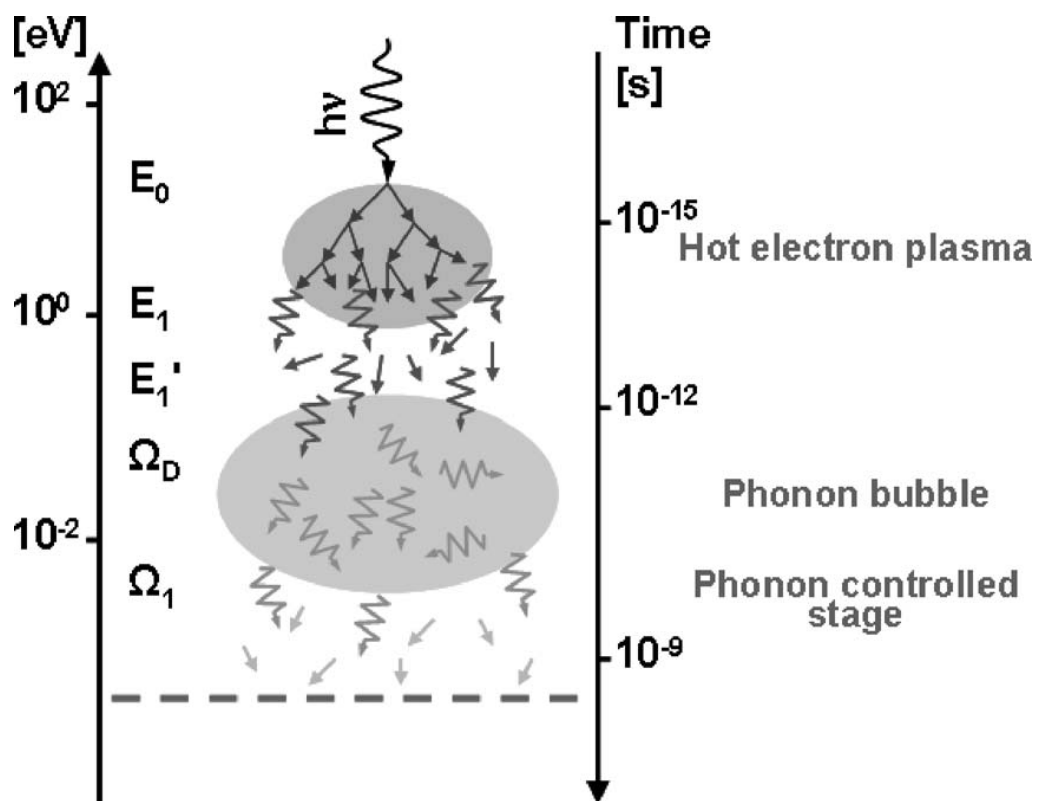
# Quantum Calorimeters for Particle-Like DM

“Logical conclusion” of existing modalities for sensing phonons, ionization, and scintillation

## Semi-classical picture

Enough quanta produced to treat as continuous variable

Quantization not apparent in  $E_{\text{observed}}$ .



Kozorezov et al PRB: 75, 094513 (2007)

## Quantum picture

Production of single charges or scintillation photons (eV quantization)

Arguable whether sensing single charges or photons is “quantum” (e.g., PMTs, SiPMs)

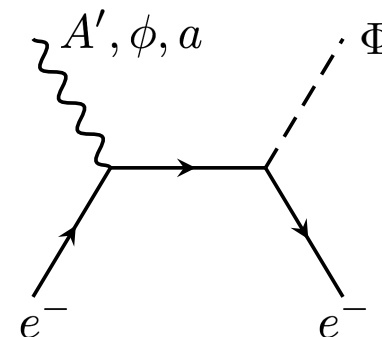
Innovations considered “quantum”:

Detection of single photons with nearly 100% efficiency and very low dark count rates  
vs. PMT, SiPMs

Direct production of single condensed-matter quanta like phonons, plasmons, etc.

Energy depositions so small that “quantum” techniques must be used to detect

e.g. direct optical phonon production mediated by DM-e coupling



# Quantum Calorimeters: Interaction vs. Creation vs. Sensing

## Interaction modalities

### DM

#### Nucleus scattering

single-site

multi-site

#### Electron scattering

single-site

multi-site

(dielectric absorption function)

### Neutrinos

#### Nucleus scattering

single-site

multi-site

### Photons

#### Electron scattering

single-site

multi-site

(dielectric absorption function)

#### Cavity or antenna mode creation

## Creation modalities

Single e-h pair creation

Single scintillation photon  
creation

Single optical phonon creation

Single acoustic phonon  
creation

Single Cooper pair breaking

Single plasmon creation

Exotic quanta

Dirac materials, magnons, ...

Single conversion-photon  
creation

## Sensing Modalities: mostly cryogenic

Transition-Edge Sensors  
(TESs)

Kinetic Inductance Detectors  
(KIDs)

Skipper CCDs

Superconducting Nanowire  
Single Photon Detectors  
(SNSPDs)

Parametric amplifiers  
for squeezing

Quantum non-demolition  
(QND) photon, phonon  
detectors

# Quantum Calorimeters:

## DM Kinematics and Scattering Modes

Energy-momentum transfer of galactic DM  
scattering kinematically limited by max  $v_{\text{DM}}$

Interactions with nuclei ( $\sim \text{CE}\nu\text{NS}$ )

Up to  $E_{\text{nuc}}^{\text{bind}} \sim 10\text{-}20\text{ eV}$ , nucleus cannot leave ionic site:  
direct production of multiple phonons

Lines truncated when single nucleus scattering invalid,  
 $q < (\text{lattice constant})^{-1}$ . Instead, direct single  
phonon creation via coherent multi-site interaction

Interactions with electrons (EC, NC  $e^-$  detection)

FEG = free electron gas, valid for  $E > E_{\text{elec}}^{\text{bind}}$

$E_{\text{elec}}^{\text{bind}} = 0$  for metals, superconductor (SC) in transition

$\sigma_E \propto T^\alpha \sqrt{M}$  because fermionic excitations with no gap  
 $\Rightarrow$  poor  $\sigma_E$  for large target mass for DM

$E_{\text{elec}}^{\text{bind}} \sim \text{eV}$  for semiconductor/insulator, meV for SC

Gapped excitation spectrum suppresses  $\sigma_E$

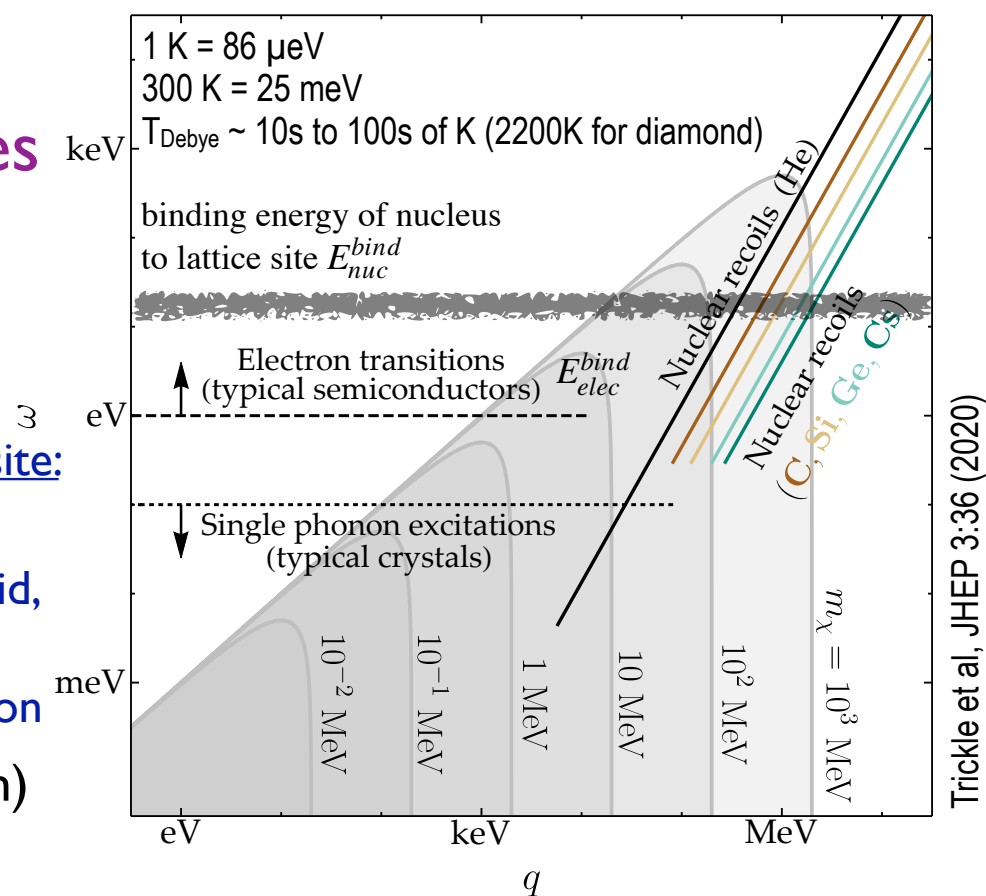
No electronic excitations for  $E < E_{\text{elec}}^{\text{bind}}$ ; instead:

Blue shaded region: plasmons

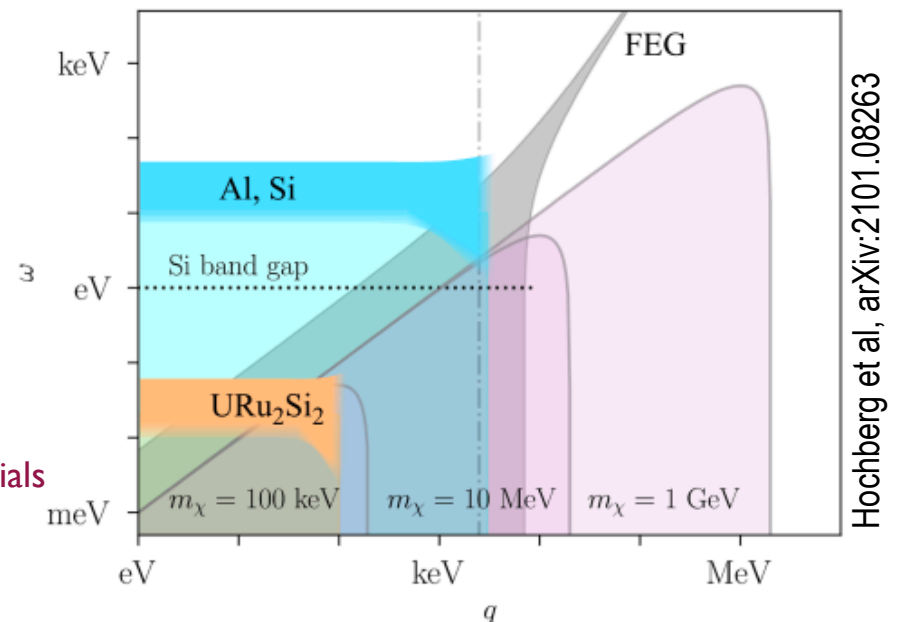
Most phase space inaccessible except for heavy fermion materials

Optical phonon creation by light/dark photon absorption

Coupling to unit cell dipole moment in polar materials



Trickle et al, JHEP 3:36 (2020)



Hochberg et al, arXiv:2101.08263

# What about neutrinos?

eV-scale thresholds interesting for  $\text{CE}\nu\text{NS}$ :  
not quite quantum, but benefits from QS efforts

Solar neutrinos: good potential

e.g. SuperCDMS long-term future:

0.5 eV threshold requires 25 gm/Si detector operating at 0V  
(phonons-only), can deploy 72 kg in SuperCDMS SNOLAB  
( $N \sim 3000$  detectors)

Expect about 250 pp neutrinos per year

Sub-10% *statistical* precision on  $pp$  flux

Or search for deviations from SM with sub-10% *statistical* precision  
if LXe measures  $pp$  flux via  $\nu$ -e scattering

Requires good control of  $^{136}\text{Xe}$  background

Reactor neutrinos: much better potential for studying  $\text{CE}\nu\text{NS}$

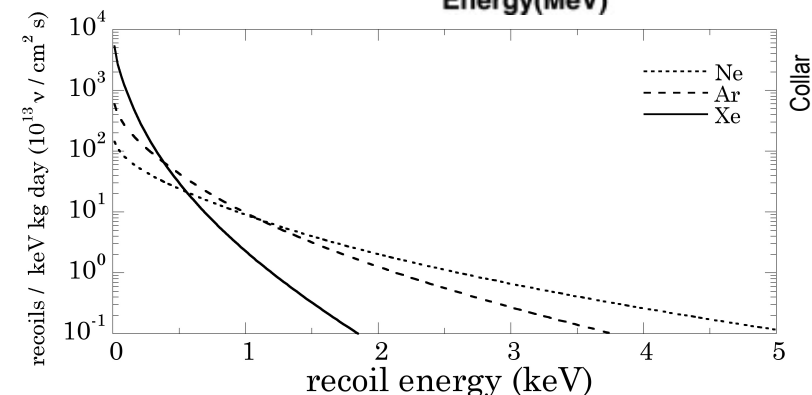
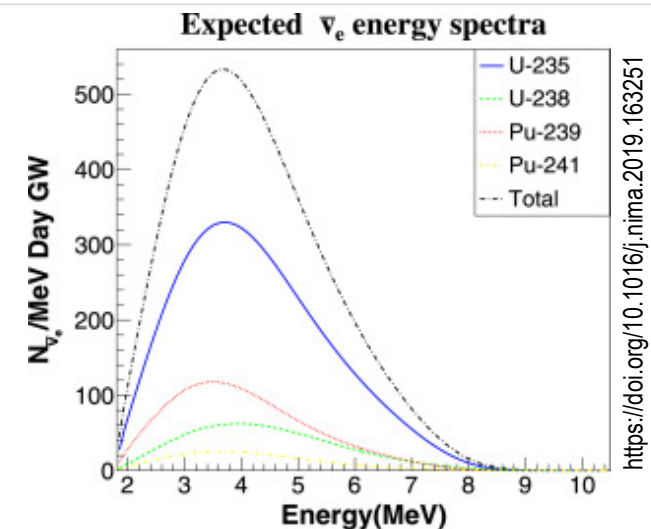
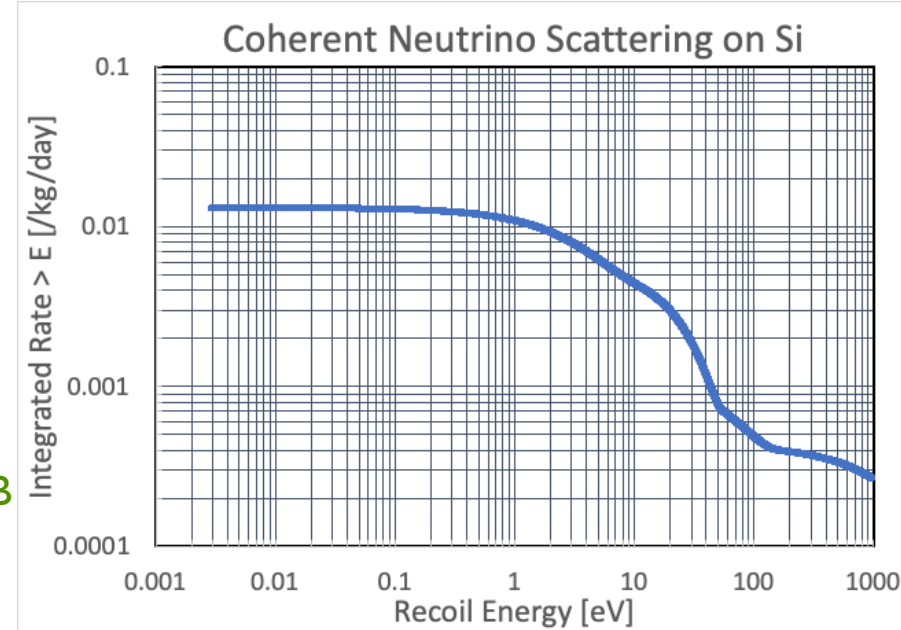
MeV antineutrino energies  $\Rightarrow \langle E \rangle \approx 100$  eV  $\text{CE}\nu\text{NS}$  in Si:  
eV threshold captures full spectrum @  $\sim 50$  events/kg-d

Large event rates!

comparable statistical precision to  $pp$  neutrinos with 5000x  
smaller exposure ( $N \sim 10$  detectors for 100 days)

1% *statistical* precision with intermediate mass  
( $N \sim 100$  detectors for 100 days)

Both require control of 10-100 eV scale low-energy  
backgrounds now becoming visible in DM expts



# What about neutrinos?

$\mathcal{O}(10)$  meV resolution on electron detection possible with quantum calorimeters

Relic cosmic neutrino background

PTOLEMY would use neutrino capture (IBD) on tritium producing 1-10 eV electrons

Requires  $\sigma_E < 0.05$  eV

Small mass (can focus electrons on detector) permits “semi-classical” sensor,  $E_{elec}^{bind} = 0$

Deposit electron energy in free-electron gas directly

Still, benefits from low-threshold/high-resolution goals for sensing modalities (TESs)

See following talk by A. Tan

Neutrino mass via electron capture endpoint in  $^{163}\text{Ho}$  (2.828 keV)

HOLMES embeds isotope in a calorimeter with high resolution (~low threshold)

Requires:

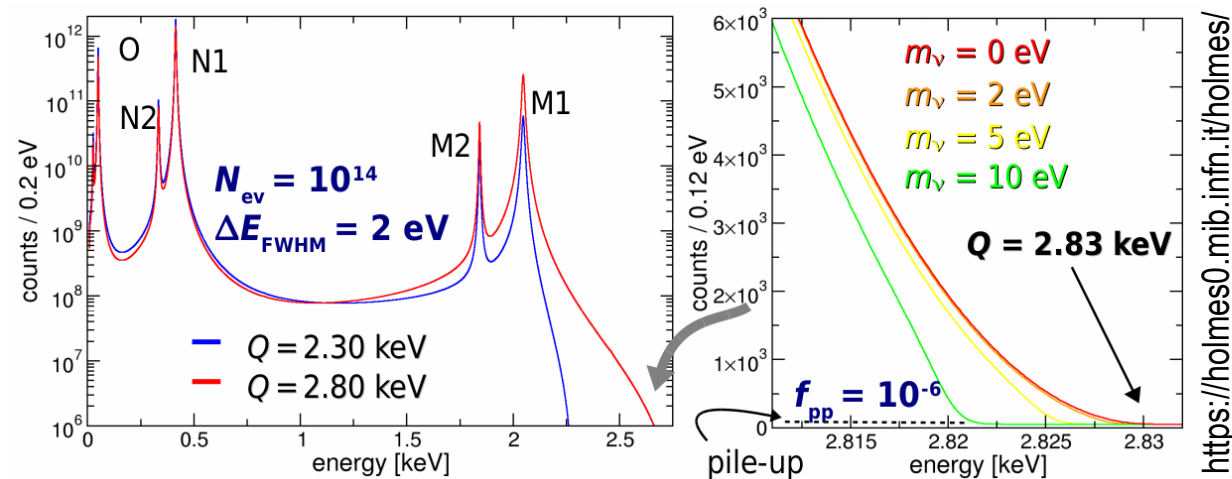
$\lesssim$  eV resolution to compete with tritium endpoint measurements

Characterization of nonlinearity at  $E \approx 10^4 \sigma_E$

Requires “semi-classical” sensor,  $E_{elec}^{bind} = 0$ , to avoid being limited by counting statistics

$E/\sigma_E \gtrsim 10^4$  needed  $\Rightarrow N_{quanta} \gtrsim 10^8$

Again, benefits from low threshold goals for sensing modalities (TESs)





# Conclusions

An expansion in the mass range of interest for DM has led to the development of low-threshold detectors based on quantum sensors

Today: eV in single detectors, 10 eV in full experimental deployment

Long-term goal: meV thresholds

Large payloads of such detectors could be useful for neutrino physics in the longer-term:

Low-threshold detectors with 1-100 kg target mass

Measurements of  $pp$  neutrinos or using them to look for non-SM  $CE\nu NS$

Search for non-SM  $CE\nu NS$  at reactors

High-resolution detectors

Detection of low-energy electrons to measure the  $C\nu B$  via neutrino capture

EC spectrum endpoint measurements for neutrino mass

# Backups: Architectures Under Development



# Quantum Calorimeters Today

## Calorimeters based on Transition-Edge Sensors

TESs provide very sharp resistance vs.  $T$  curve

Electothermal feedback can be used to stabilize

Electrical signal measures received energy

Coupled to calorimetric substrate via athermal phonon collectors

Now approaching the quantum regime

Large bias voltage provides amplification via Neganov-Trofimov-Luke phonons: enables single e-h pair detection

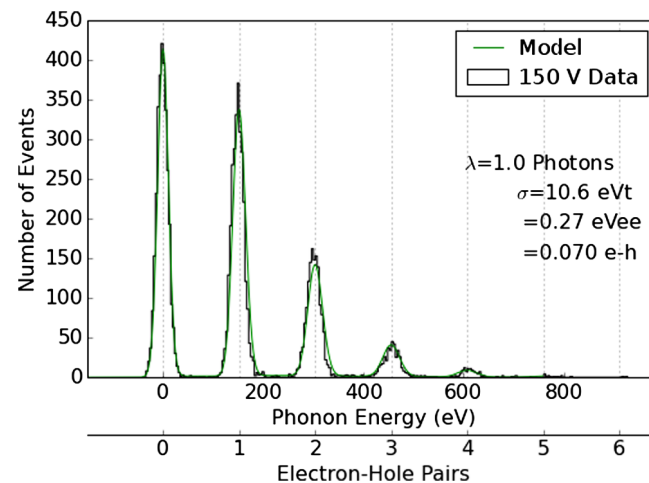
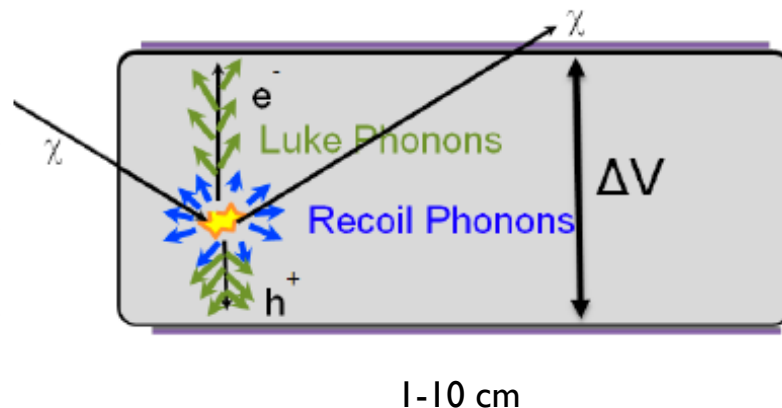
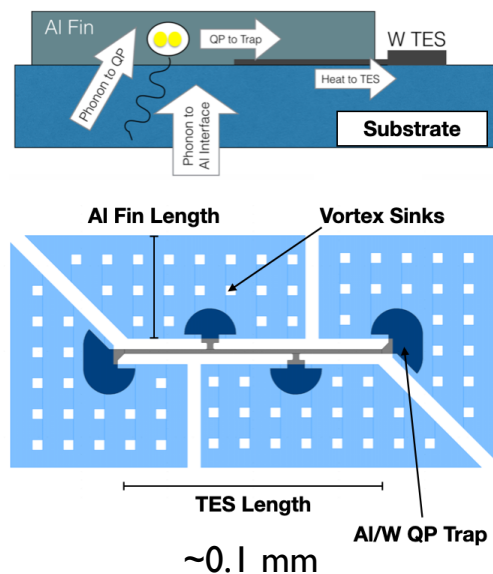
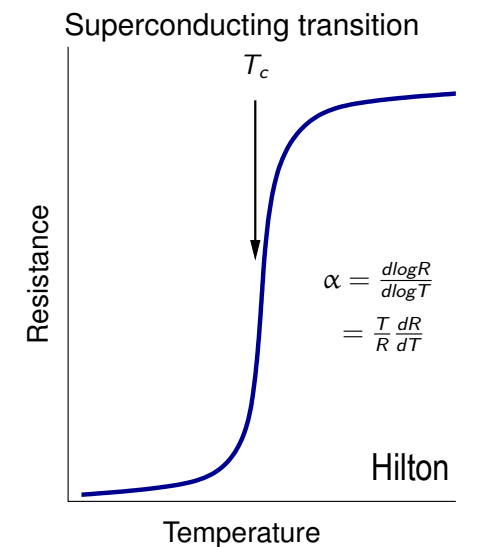
But subject to significant leakage currents: single eh pair detection unreliable

Native resolution  $\sim 3$  eV achieved

Reduced  $T_c$ , improved design expected to yield sub-eV resolutions, eventually approaching 10 meV

Single eh pair detection without leakage

Single phonon detection



PRL 121: 0513401 (2018)

# Quantum Calorimeters Today

## Kinetic Inductance Detectors (KIDs)

Superconductors have an AC inductance due to inertia of Cooper pairs

KID = superconducting film incorporated into LC resonator to sense change in  $L$

Energy resolution:  
sub-eV  $\rightarrow$  meV  
thresholds w/o HV

Direct sensitivity to  
pair-breaking phonons

Large resonators obviate  
phonon collectors

Gapped density of states

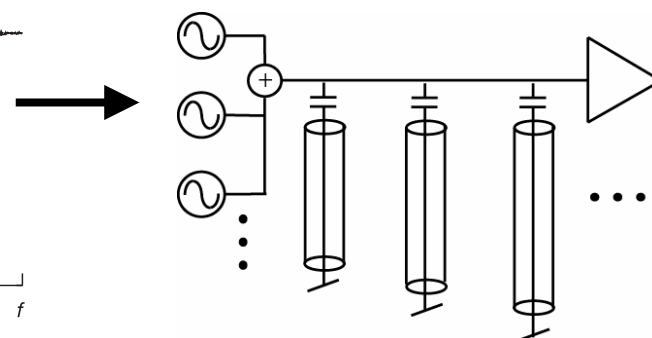
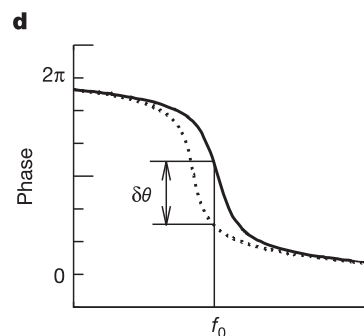
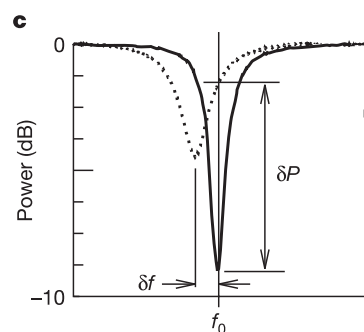
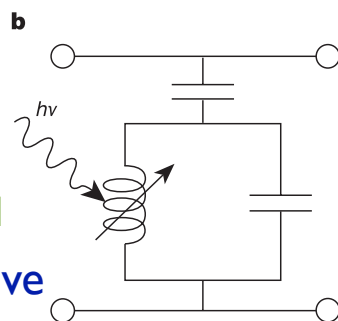
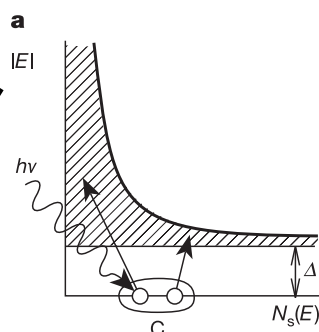
Thermal quasiparticles  
exponentially suppressed

Fundamentally non-dissipative

Amenable to QIS techniques  
(e.g. squeezing, QND)

Noise is limited by

quasiparticle population  
fluctuations  
amplifier noise



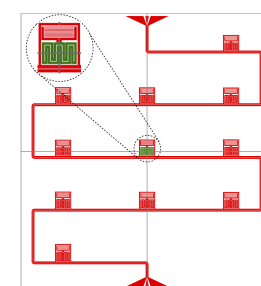
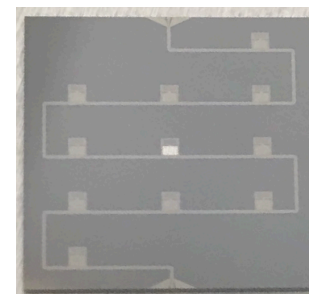
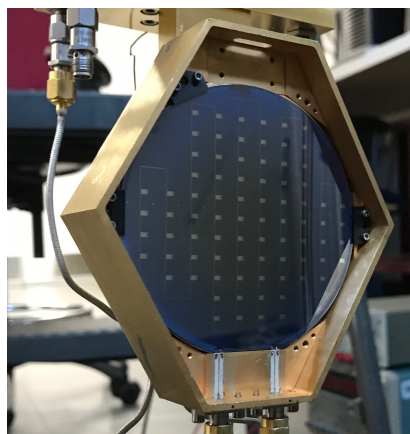
Multiplexing:

KIDs are  $Q > 10^5$  resonators

$\rightarrow$  Readout many with one  
cryo line/amplifier; most  
electronics at 300K

$\rightarrow$  Highly position-resolved  
phonon detection

75 mm  
x 1 mm



22 mm  
x 1 mm

# Quantum Calorimeters Today

## Skipper CCDs

### CCD with two readout innovations

High-frequency differencing to reduce impact of  $1/f$  amplifier noise

Non-destructive multiple read cycles to reduce electronics noise by  $\sqrt{N}$

Provides similar single-eh pair sensitivity

Currently being applied for DM searches, low-light-level astronomy

## Superconducting Nanowire Single Photon Detectors (SNSPDs)

### Threshold detector for single photons

Very narrow ( $\sim 100$  nm) superconducting meander biased close to transition

Absorption of photon drives normal

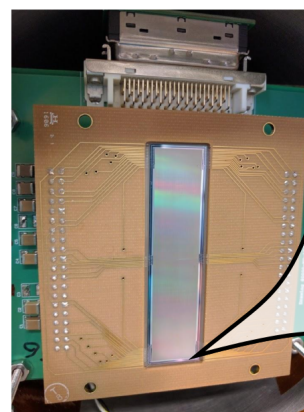
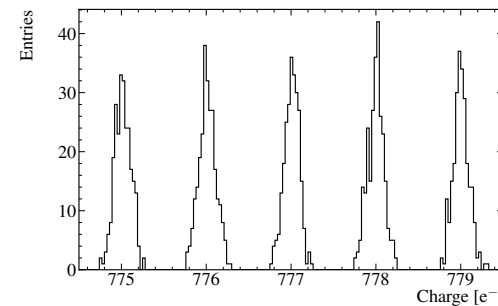
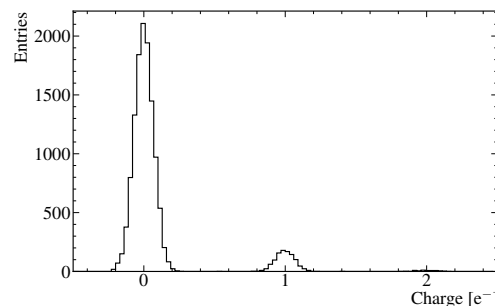
ps timing resolution

Provides high-efficiency, high-fidelity photon counter for QIS applications

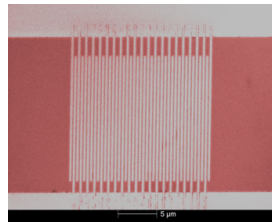
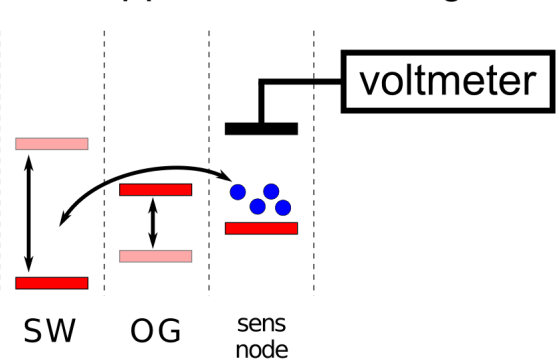
WSi demonstrated with 100 meV threshold

Very low dark count rate demonstrated, applicable for DM searches

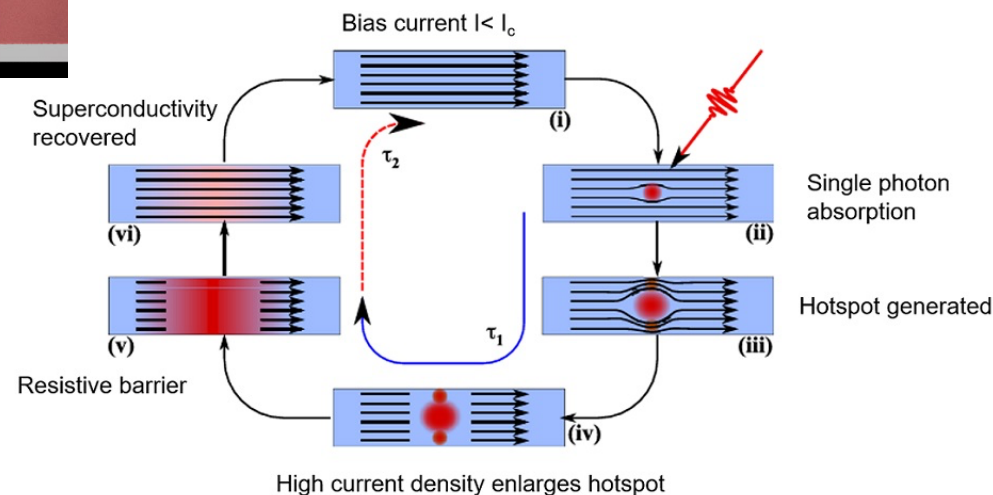
But very small volume



### Skipper read out stage



<https://sensei-skipper.github.io/#SkipperCCD>



<https://singlequantum.com/technology/snsdp/>

# Quantum Calorimeters in the Future

## Single optical phonon creation/detection

### Polar materials:

- > 1 atom/unit cell  $\rightarrow$  optical phonons (10s of meV)
- polar  $\rightarrow$  unit cell EDM can couple to dark photons

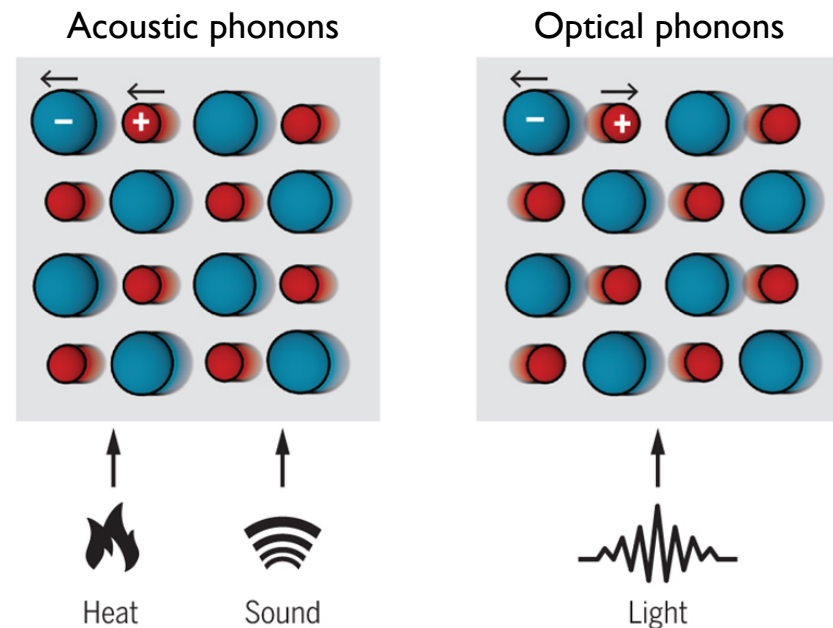
Optical phonon creation provides access to electron scattering down to few keV  $m_{DM}$

## Single acoustic phonon creation/detection

DM interacting with many nuclei coherently

Requires meV-100 meV thresholds

Eventual reach down to few keV  $m_{DM}$



<https://images.app.goo.gl/vvBWneR6noCnseCXA>

