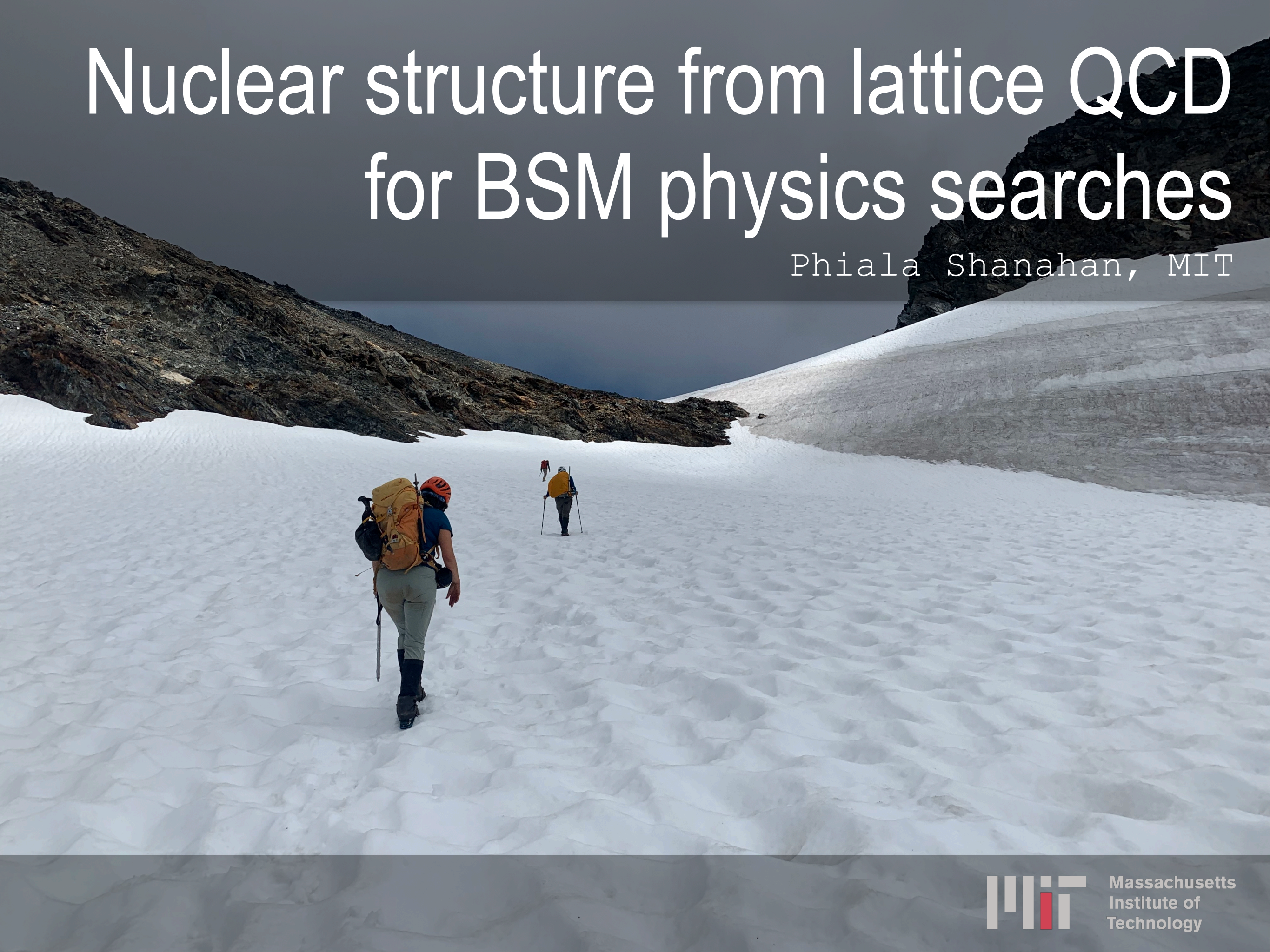


Nuclear structure from lattice QCD for BSM physics searches

Phiala Shanahan, MIT

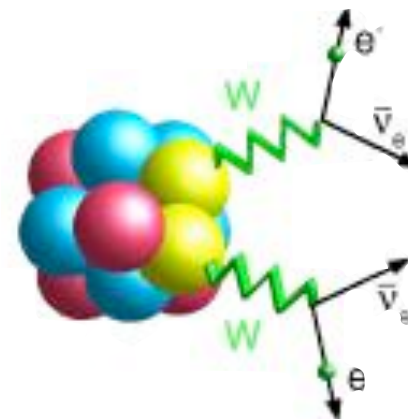


Massachusetts
Institute of
Technology

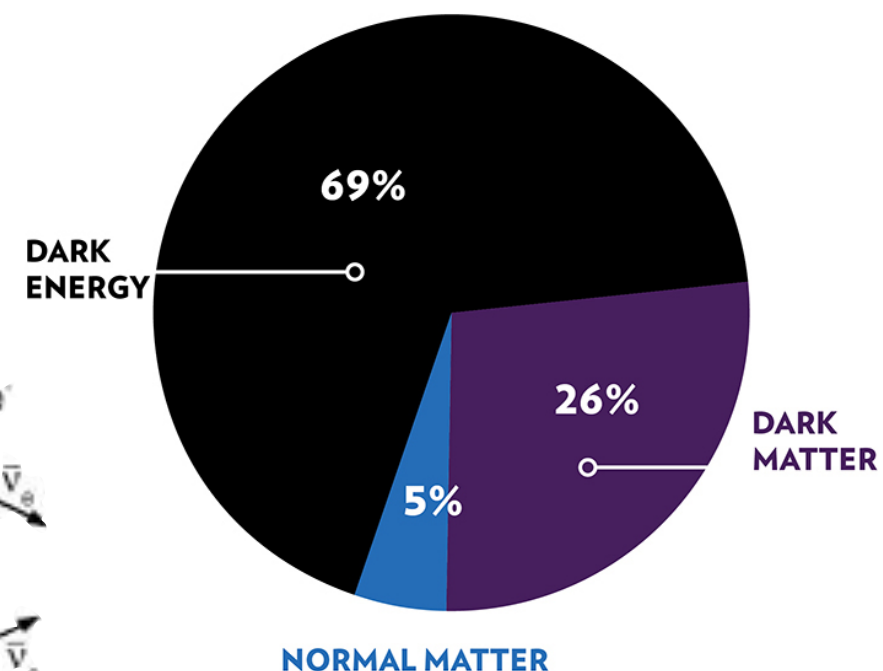
The search for new physics

Precise experiments seek new physics at the “Intensity Frontier”

- Sensitivity to reveal small beyond—Standard-Model effects
- Magnetic moments
- Dark matter direct detection
- Neutrino physics
- Charged lepton flavour violation, $\beta\beta$ -decay, proton decay, neutron-antineutron oscillations...



ENERGY DISTRIBUTION OF THE UNIVERSE

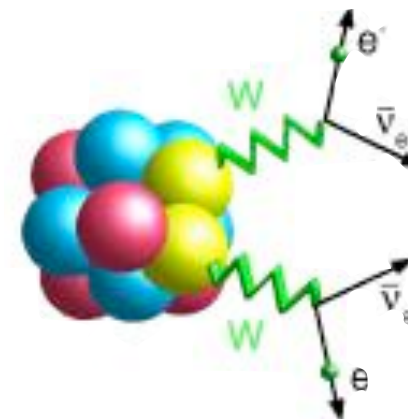


The search for new physics

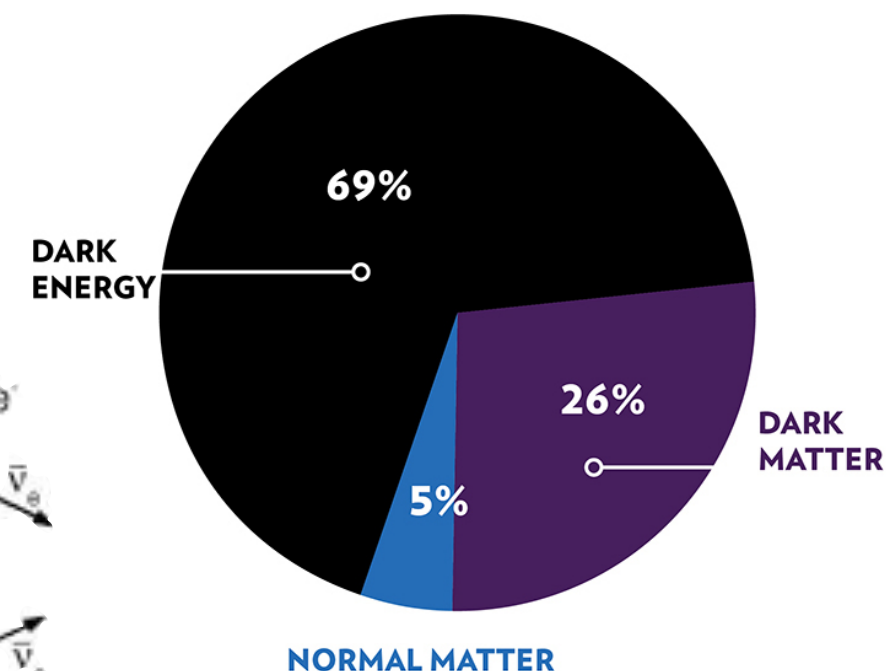
Need to understand the Standard Model
physics of nucleons and nuclei

Interpretation of intensity-frontier experiments

- Axial form factors of Argon
A=40 DUNE long-baseline
neutrino experiment
- Double-beta decay rates of
Calcium A=48
- Scalar matrix elements in A=131
XENON1T dark matter direct
detection search



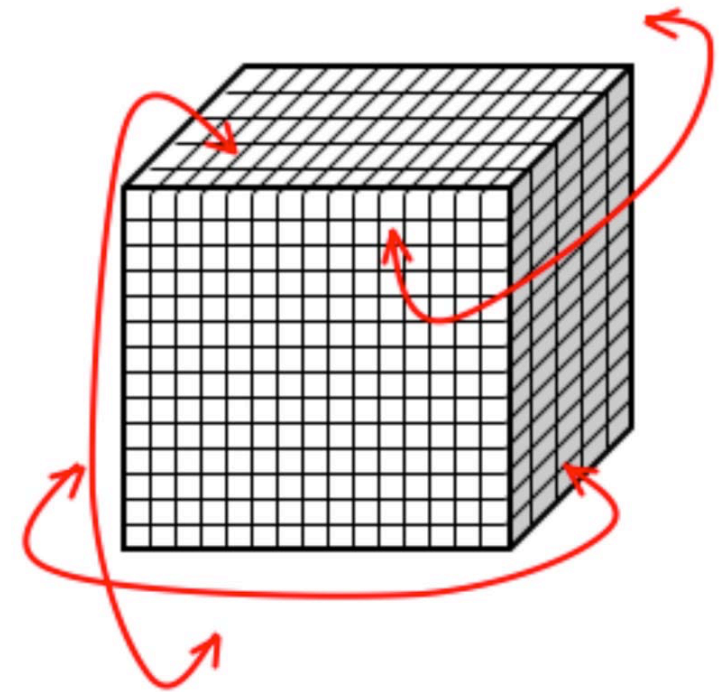
ENERGY DISTRIBUTION
OF THE UNIVERSE



Lattice QCD

Numerical first-principles approach to non-perturbative QCD

- Euclidean space-time
 - Non-zero lattice spacing
 - Finite volume
- Some calculations use larger-than-physical quark masses (cheaper)



Calculate the QCD path integral by **Monte Carlo**

$$\langle \mathcal{O} \rangle = \frac{1}{Z} \int \mathcal{D}A \mathcal{D}\bar{\psi} \mathcal{D}\psi \mathcal{O}[A, \bar{\psi}\psi] e^{-S[A, \bar{\psi}\psi]} \rightarrow \langle \mathcal{O} \rangle \simeq \frac{1}{N_{\text{conf}}} \sum_i^{N_{\text{conf}}} \mathcal{O}([U^i])$$

with field configurations U^i distributed according to $e^{-S[U]}$

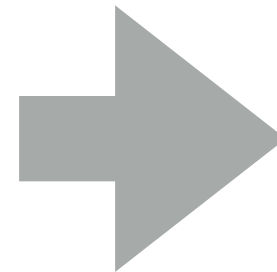
Lattice QCD

Numerical first-principles approach to
non-perturbative QCD

INPUT

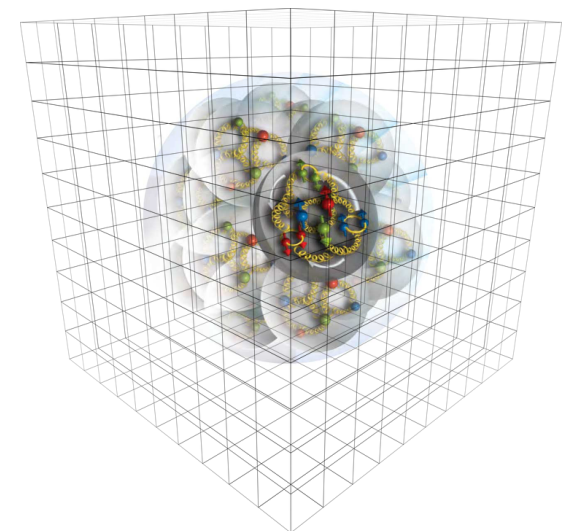
Lattice QCD action has same free parameters as QCD: quark masses, α_S

- Fix quark masses by matching to measured hadron masses, e.g., π, K, D_s, B_s for u, d, s, c, b
- One experimental input to fix lattice spacing in GeV (and also α_S), e.g., $2S-1S$ splitting in Y , or f_π or Ω mass



OUTPUT

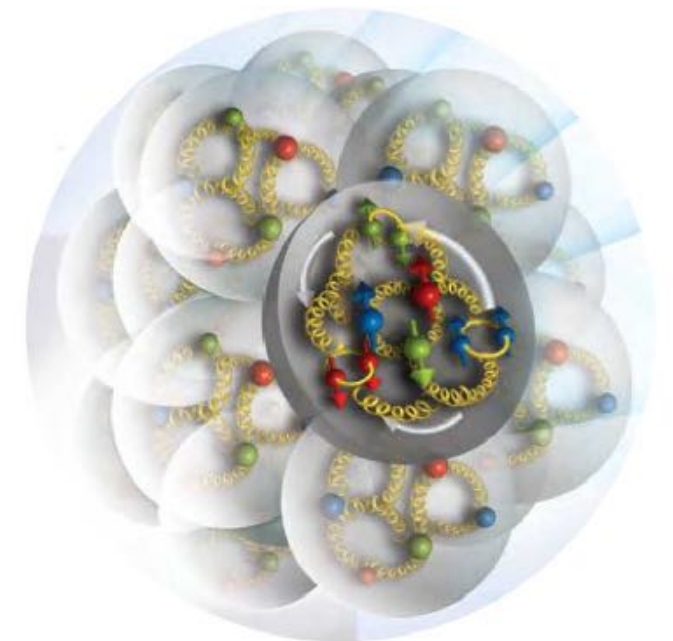
Calculations of all other quantities are QCD predictions



Nuclear physics from lattice QCD

Nuclei on the lattice are HARD

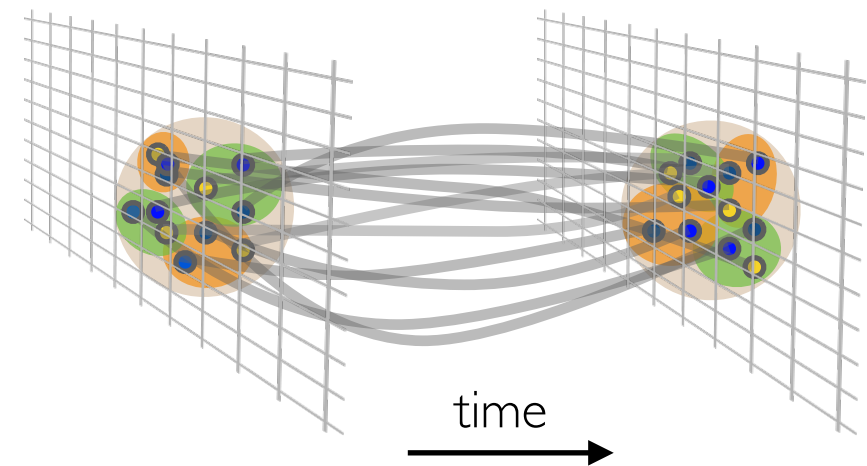
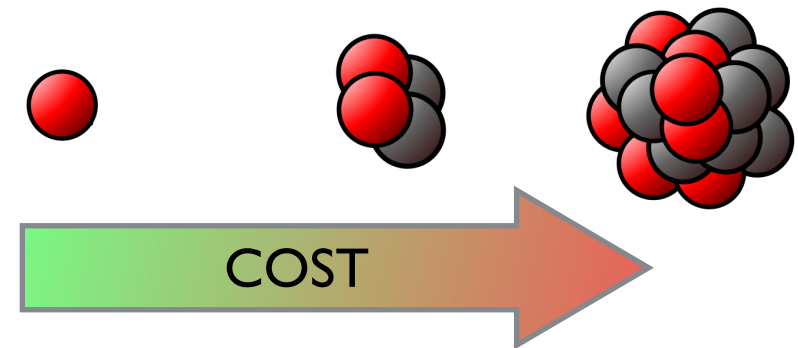
- Calculations of matrix elements of currents in light nuclei just beginning:
 - Controlled calculations of spectrum of light nuclei yet to be achieved
 - First exploratory calculations of matrix elements taking place now
- With sufficient computing resources, calculations are in principle possible:
 - Deeply bound nuclei: same techniques as for single hadron matrix elements
 - Near threshold states: need to be careful with volume effects



Nuclear physics from lattice QCD

Nuclei on the lattice are HARD

- **Noise:**
Statistical uncertainty grows exponentially with number of nucleons
- **Complexity:**
Number of contractions grows factorially

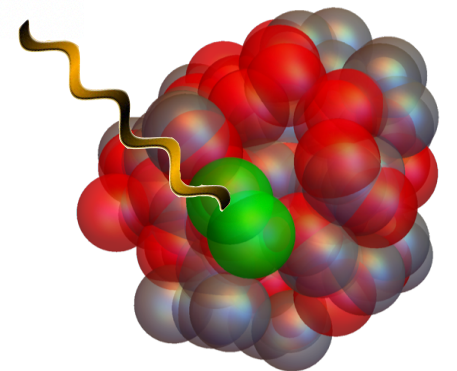
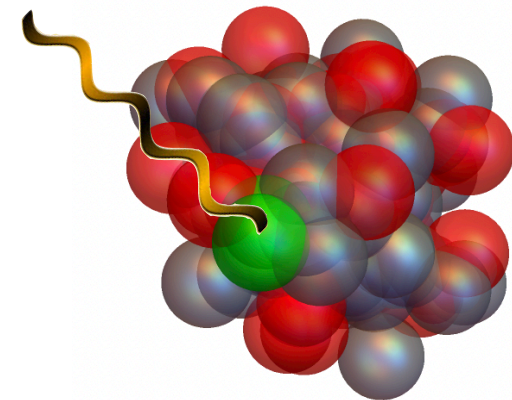


Calculations possible for $A < 5$

Larger nuclei

What about larger
(phenomenologically-relevant) nuclei?

- Nuclear effective field theory:
 - 1-body currents are dominant
 - 2-body currents are sub-leading
but non-negligible
- Determine one body contributions from single nucleon
- Determine few-body contributions from $A=2,3,4\dots$
- Match effective theory and many body methods to lattice results to make predictions for larger nuclei
- Can reproduce axial matrix elements for large nuclei



Nuclear physics from lattice QCD

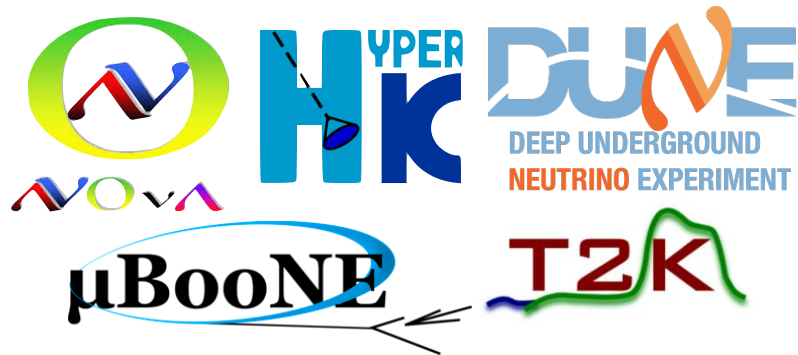
Nuclear matrix elements from lattice QCD studied only by the NPLQCD Collaboration to date

- Proton-proton fusion and tritium β -decay [PRL 119, 062002 (2017)]
- Double β -decay [PRL 119, 062003 (2017), PRD 96, 054505 (2017)]
- Gluon structure of light nuclei [PRD 96 094512 (2017)]
- Scalar, axial, tensor MEs [PRL 120 152002 (2018), PRD 103, 074511 (2021)]
- Baryon-baryon interactions, including QED [PRD 103, 054504 (2021), PRD 103, 054508 (2021)]
- EMC-type effects in light nuclei [PRD 96 094512 (2017), PRL 126, 202001 (2021)]

Many other collaborations are studying nuclei from lattice QCD

- **PACS-CS**
e.g., Yamazaki et al, PRD 92 (2015);
- **Callatt**
e.g., E Berkowitz et al, PLB 765 (2017); Hörz et al, PRC 103 (2021)
- **Mainz**
e.g., A. Francis et al, PRD 99 (2019); Green et al, PRL 127 (2021)
- **HALQCD**
e.g., Ishii et al, PRL 99 (2007)
(potential approach)

Neutrino oscillation experiments

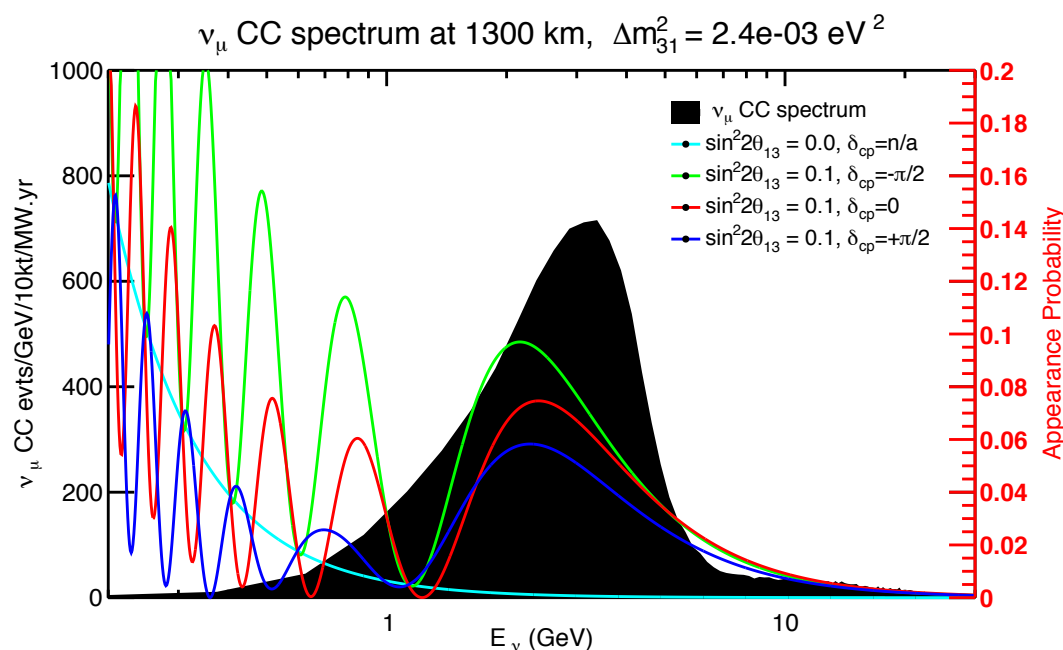


Seek to determine neutrino mass hierarchy, mixing parameters, CP violating phase

To differentiate between mixing & CP parameter scenarios



Need neutrino energy reconstruction from final state to better than 100 MeV

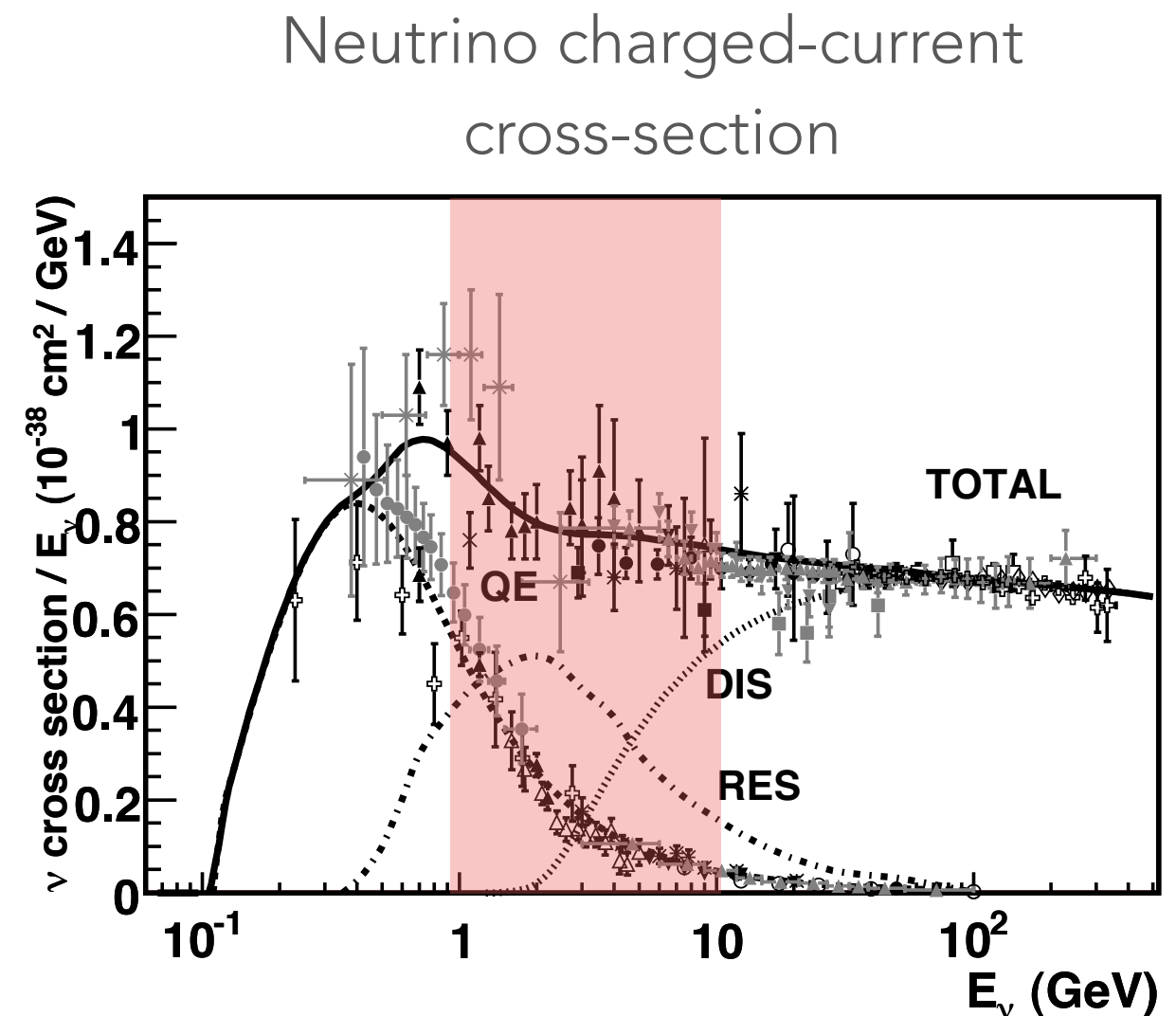


Need robust understanding of relevant nucleon and nuclear level amplitudes

e.g., axial and pseudo-scalar form factors in **quasi-elastic** region

Constraining ν -nucleus interactions

- For DUNE neutrino energy distributions peak at 1-10 GeV
- Challenging region: several processes contribute
 - Quasielastic lepton scattering
 - Deep inelastic scattering
 - Resonances
- Lattice QCD can provide direct non-perturbative QCD predictions of nucleon and **nuclear** matrix elements



J.A. Formaggio, G.P. Zeller, Rev. Mod. Phys. 84 (2012) 1307

Quasi-elastic scattering

Cross-section for quasi-elastic neutrino-nucleon scattering

$$\frac{d\sigma}{dQ^2} = \frac{G_f^2 M^2 \cos^2 \theta_C}{8\pi E_\nu^2} \left[A \mp \frac{(s-u)}{M^2} B + \frac{(s-u)^2}{M^4} C \right]$$

$$A = \frac{(m^2 + Q^2)}{M^2} [(1 + \tau) G_A^2 - (1 - \tau) F_1^2 + \tau(1 - \tau) F_2^2 + 4\tau F_1 F_2 - \frac{m^2}{4M^2} \left((F_1 + F_2)^2 + (G_A + 2G_P)^2 - \left(\frac{Q^2}{M^2} + 4 \right) G_P^2 \right)]$$

$$B = \frac{Q^2}{M^2} G_A (F_1 + F_2)$$

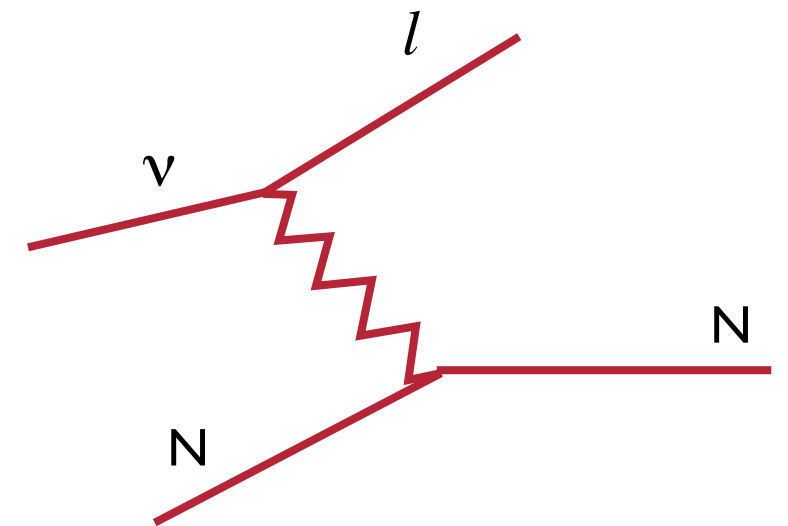
$$C = \frac{1}{4} (G_A^2 + F_1^2 + \tau F_2^2)$$

G_A

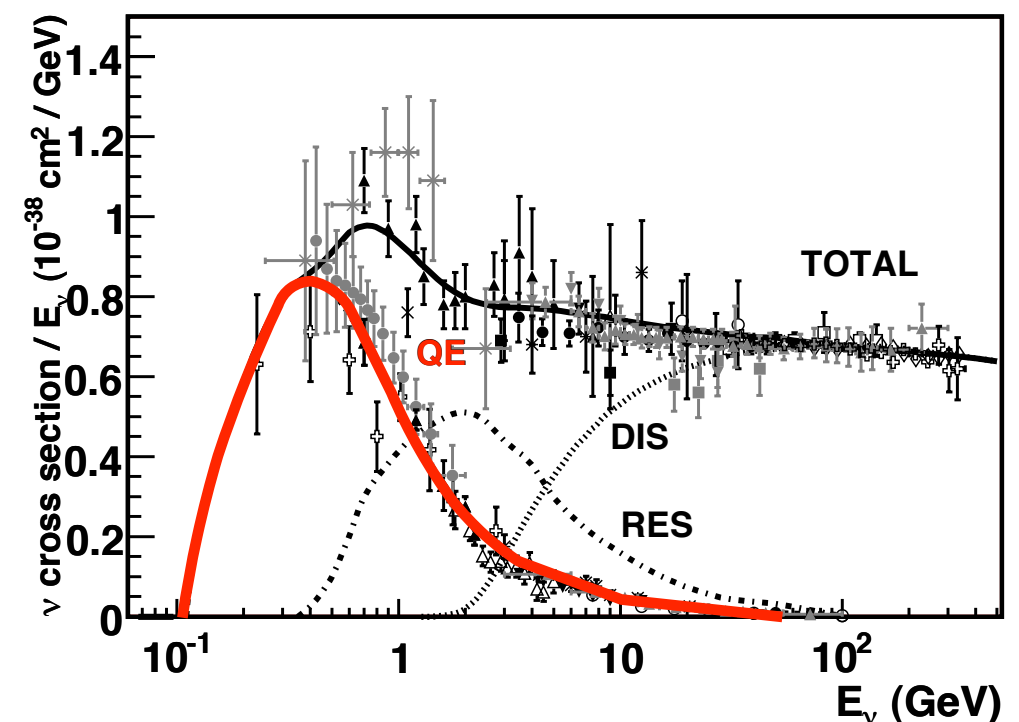
- dominant contribution
- largest uncertainty

$F_{1,2}$ Well-determined from electron scattering expts

G_P can be related to G_A by pion pole dominance

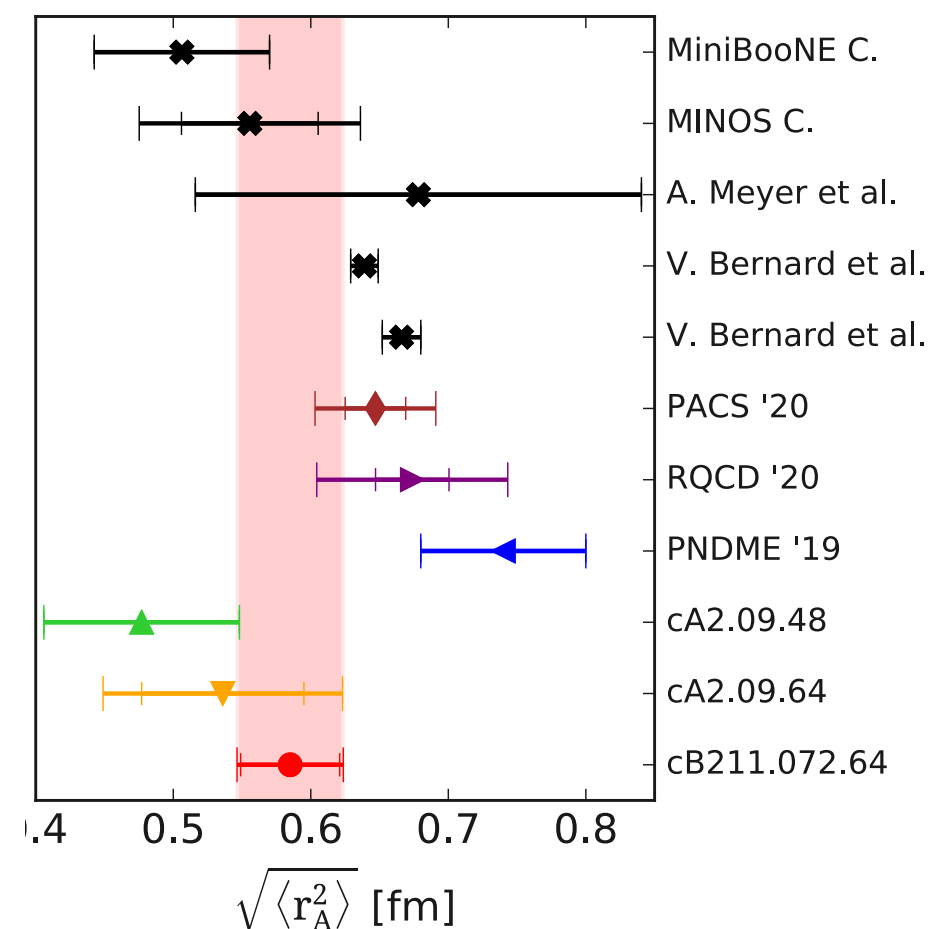
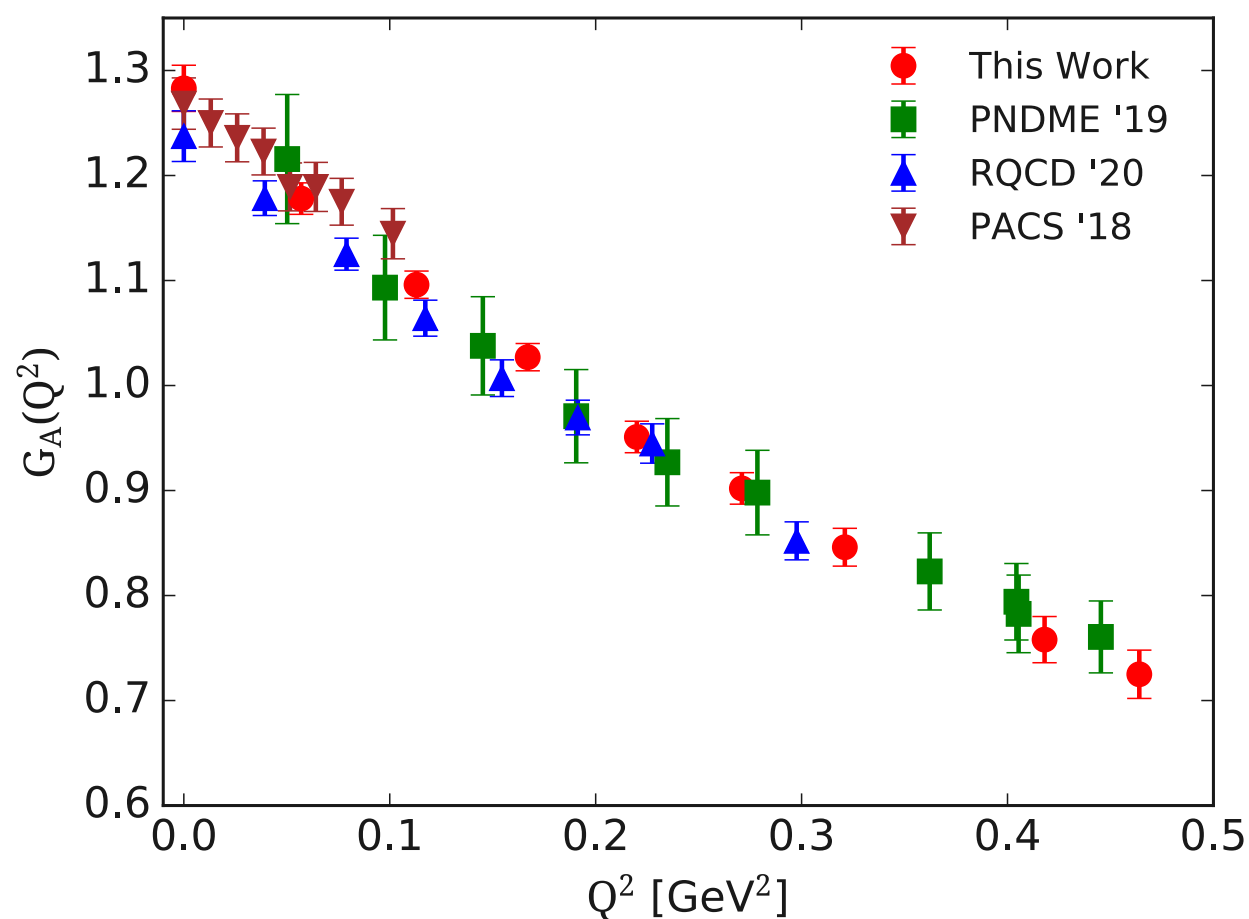


ν charged-current cross-section



Nucleon axial form factors

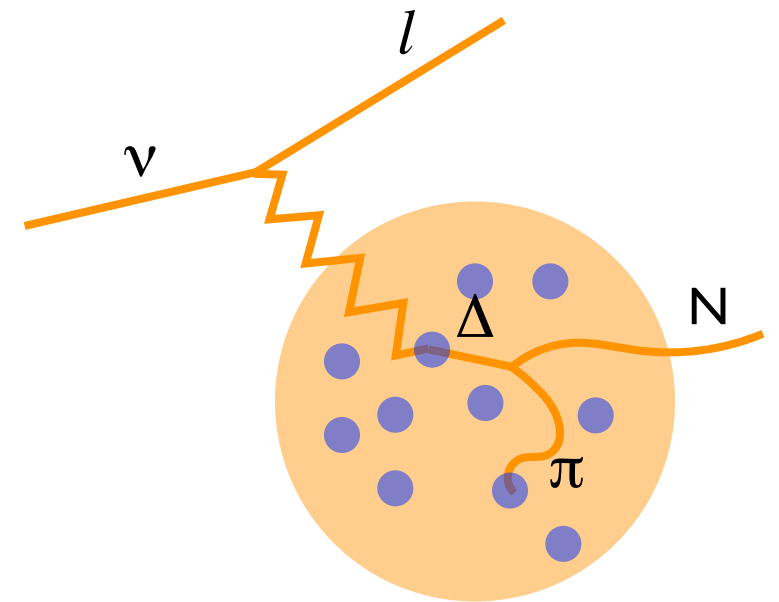
- Recent calculations of nucleon form factors including axial in agreement with experiment with fully-controlled uncertainties
- Q^2 -dependence well-determined in LQCD: competitive with experiment



[Alexandrou et al., Phys. Rev. D 103, 034509 (2021)]

Nuclear effects in matrix elements

- Targets are nuclei (C, Fe, Ar, Pb, H₂O) so how relevant are nucleon FFs?
 - Nuclear effects (EMC effect)
 - Suppression of g_A in Gamow-Teller transitions
- Experimental investigations: MINERvA



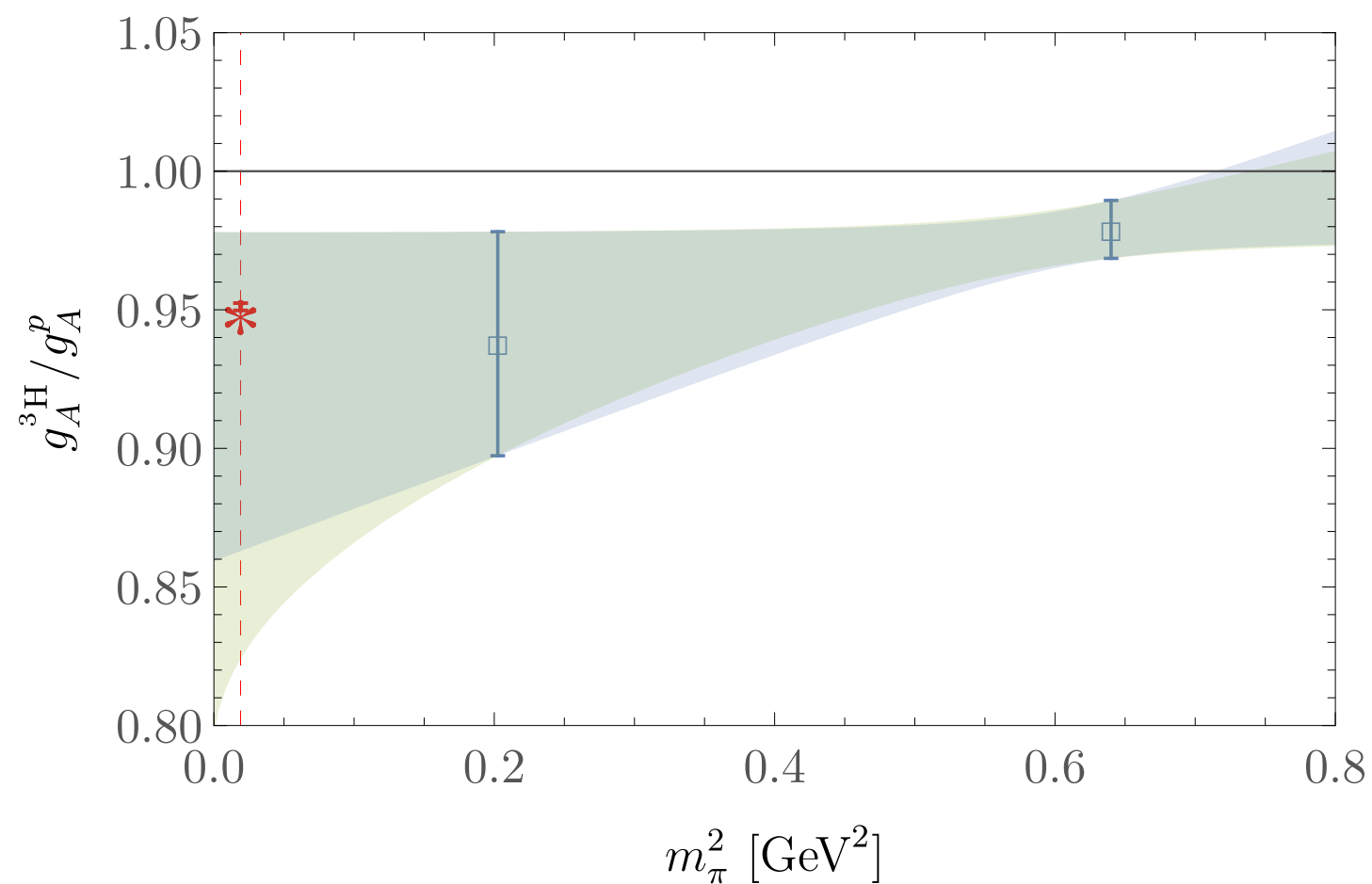
Calculate matrix elements in light nuclei from lattice QCD

➡ EFT to reach heavy nuclear targets relevant to experiment
e.g., First calculations of axial charge of light nuclei, EMC effect in light nuclei



Axial charge of the triton

- Axial charge of He: first extrapolation to the physical quark masses last year
- No axial form factors of nuclei from lattice QCD yet (coming soon!)



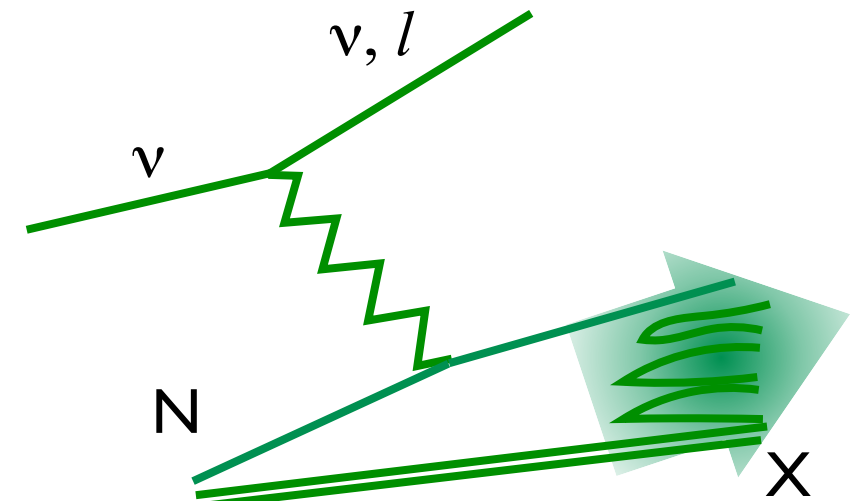
[NPLQCD Phys.Rev.D 103 (2021) 074511]

Inelastic region

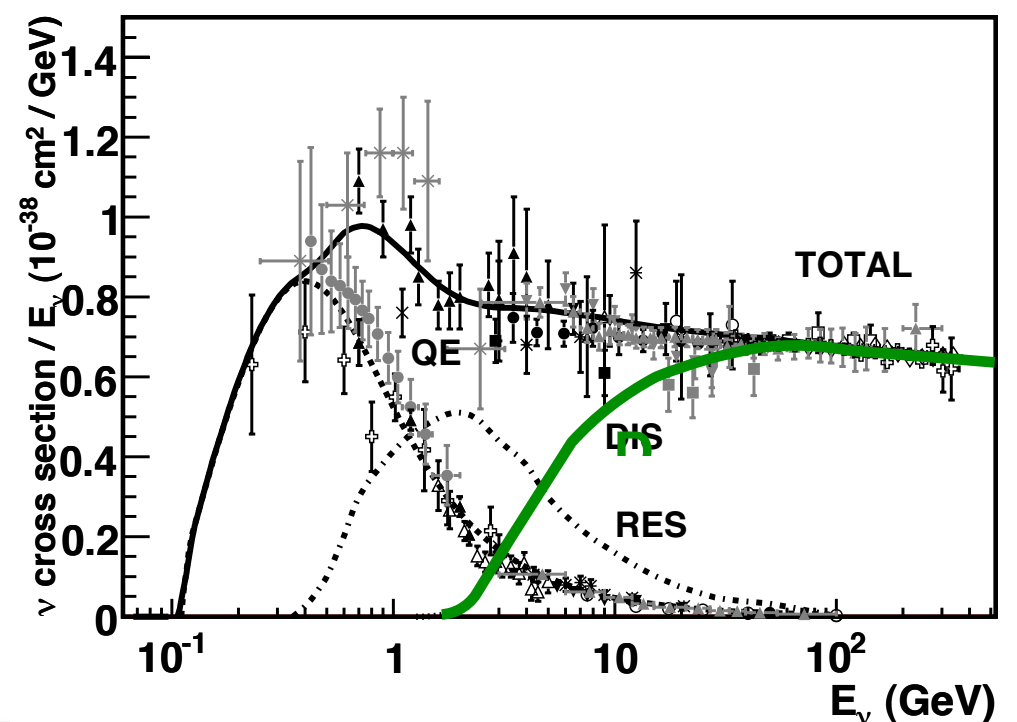
- In inelastic regime, quark PDFs of the nucleon control scattering cross-section
- Both resonances and DIS are important
- Multi-meson channels may become important
- Nuclear effects are different in νA vs. eA
- DIS structure functions accessible in LQCD
 - Low moments of structure functions controlled

$$M_n = \int_{-1}^1 x^n f(x) dx, \quad n \lesssim 4$$

- **x-dependence: systematics challenging, but rapid and exciting progress!**



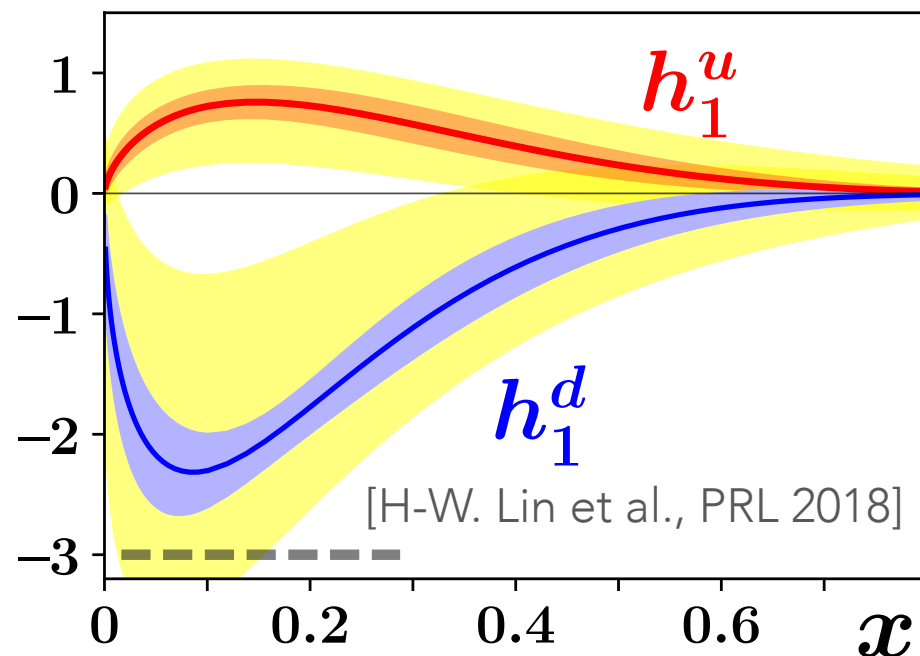
ν charged-current cross-section



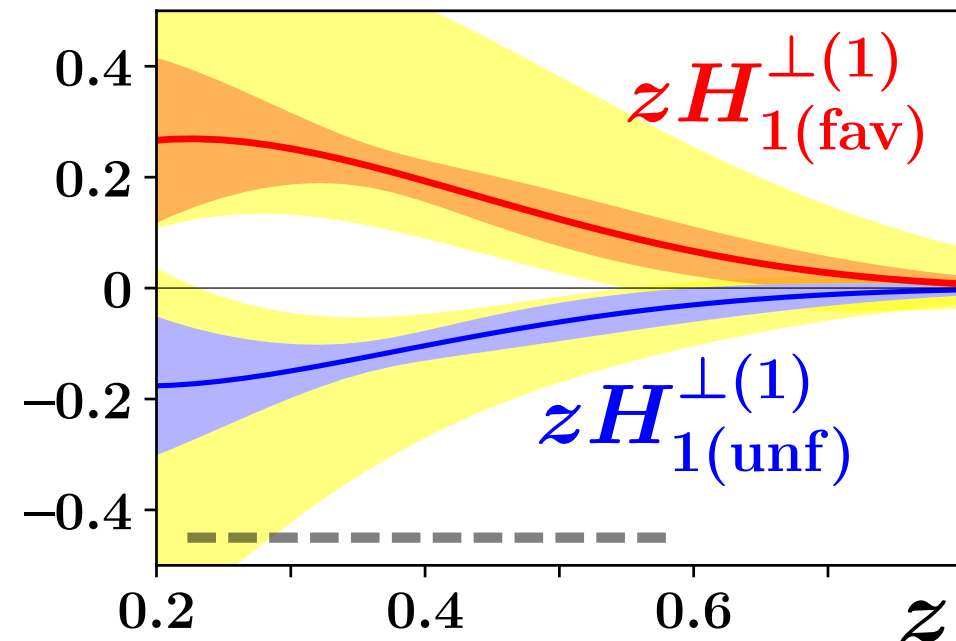
Constraints on global PDF fits

- Including lattice QCD results for moments in global parton distribution function fits can yield significant improvements
- Community white paper (LQCD + phenomenologists) assessed potential impacts [Lin et al., Prog. Part. Nucl. Phys 100 (2018), 107]

Transversity PDFs



Collins fragmentation functions



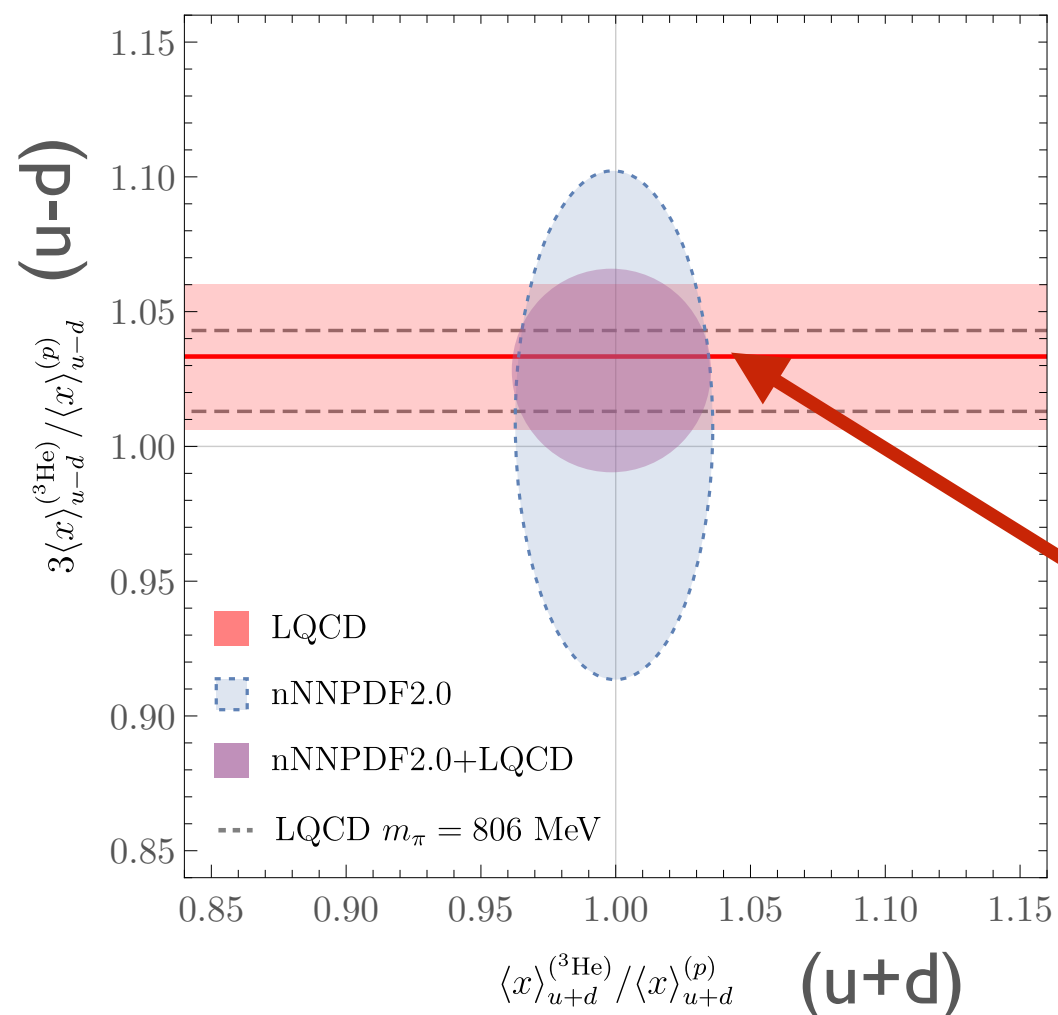
Yellow: SIDIS data only: direct constraints in region indicated by dashes
Blue/Red: SIDIS + lattice QCD for tensor charge (zeroth moment)



Momentum fraction of ^3He

Study nuclear effects in the breakdown of momentum carried by quarks in nuclei

Ratio of ^3He to proton momentum fractions



- Match isovector (u-d quark combination) momentum fraction to low-energy constants of effective field theory, extrapolate to physical quark masses
- Include into nNNPDF global fits of experimental lepton-nucleus scattering data

Blue \rightarrow Purple:
Improvement using theory constraints

[NPLQCD PRL 126, 202001 (2021) [2009.05522]]

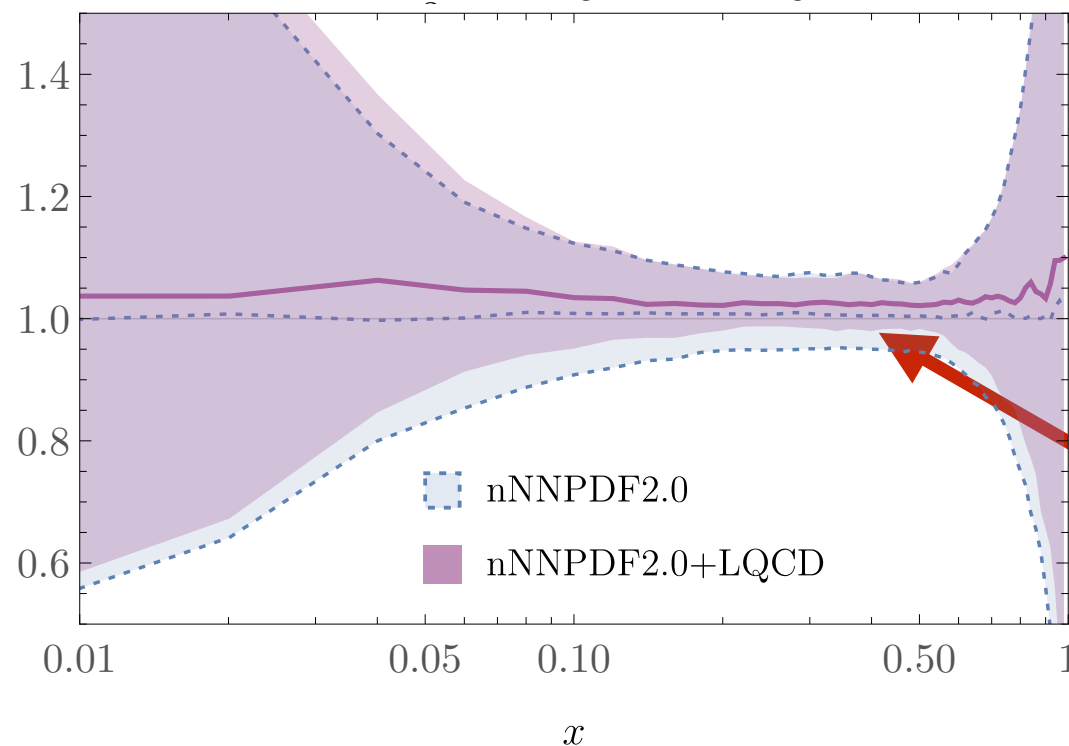


Momentum fraction of ^3He

Study nuclear effects in the breakdown of momentum carried by quarks in nuclei

Ratio of ^3He to proton parton distributions

$$R^{(^3\text{He})}(x) = 3q_3^{(^3\text{He})}(x)/q_3^{(p)}(x)$$



- Match isovector (u-d quark combination) momentum fraction to low-energy constants of effective field theory, extrapolate to physical quark masses
- Include into nNNPDF global fits of experimental lepton-nucleus scattering data

Blue \rightarrow Purple:
Improvement using theory constraints

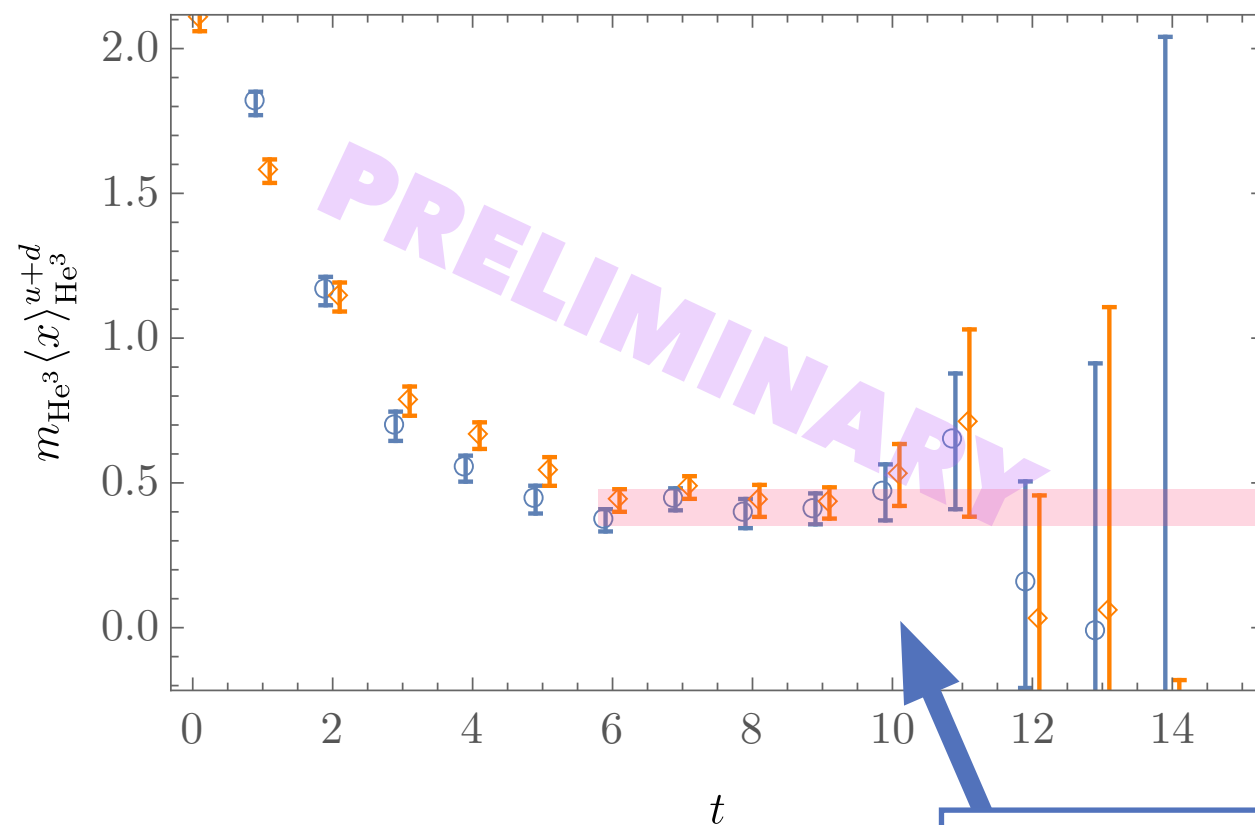
[NPLQCD PRL 126, 202001 (2021) [2009.05522]]



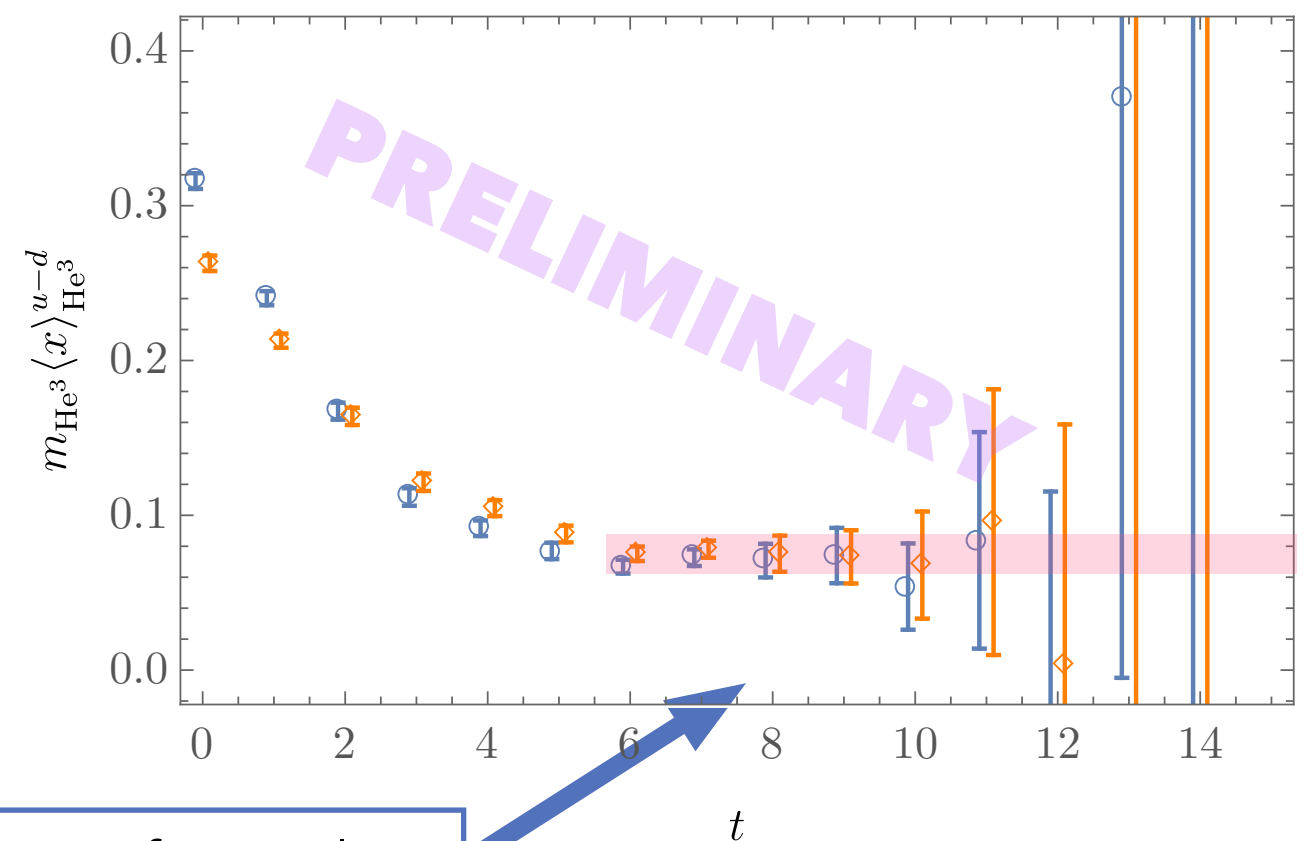
Momentum fraction of ^3He

- Work in progress at close-to-physical values of the quark masses

Isoscalar



Isovector



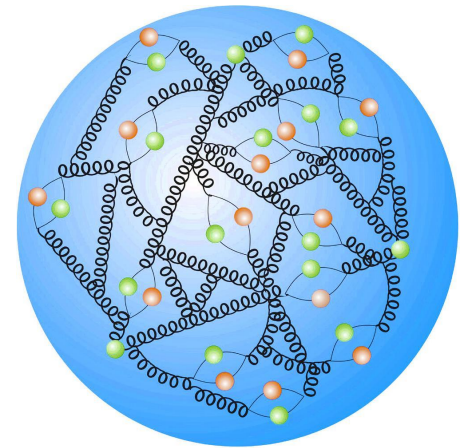
First evidence of signals in physical-point data

- Polarised PDFs, gluon PDFs, also accessible from moments

LQCD input for ν -nucleus interactions

1. Directly access QCD single-nucleon form factors without nuclear corrections

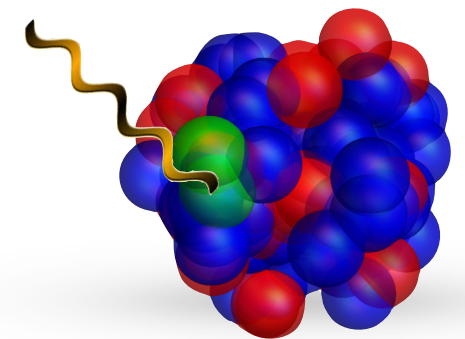
Reliable calculations with fully-controlled uncertainties



2. Calculate matrix elements in light nuclei from first principles

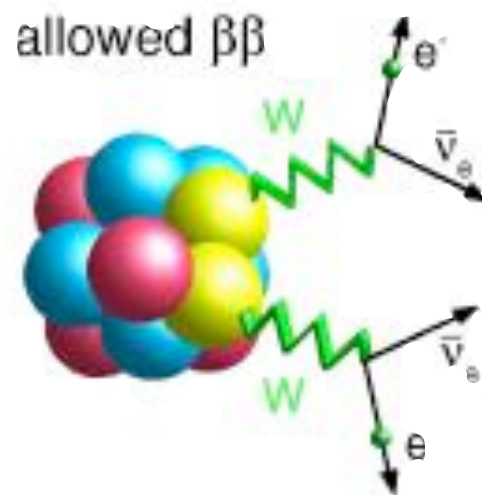
➡ EFT to reach heavy nuclear targets relevant to experiment

e.g., First calculations of axial charge of light nuclei, EMC effect in light nuclei

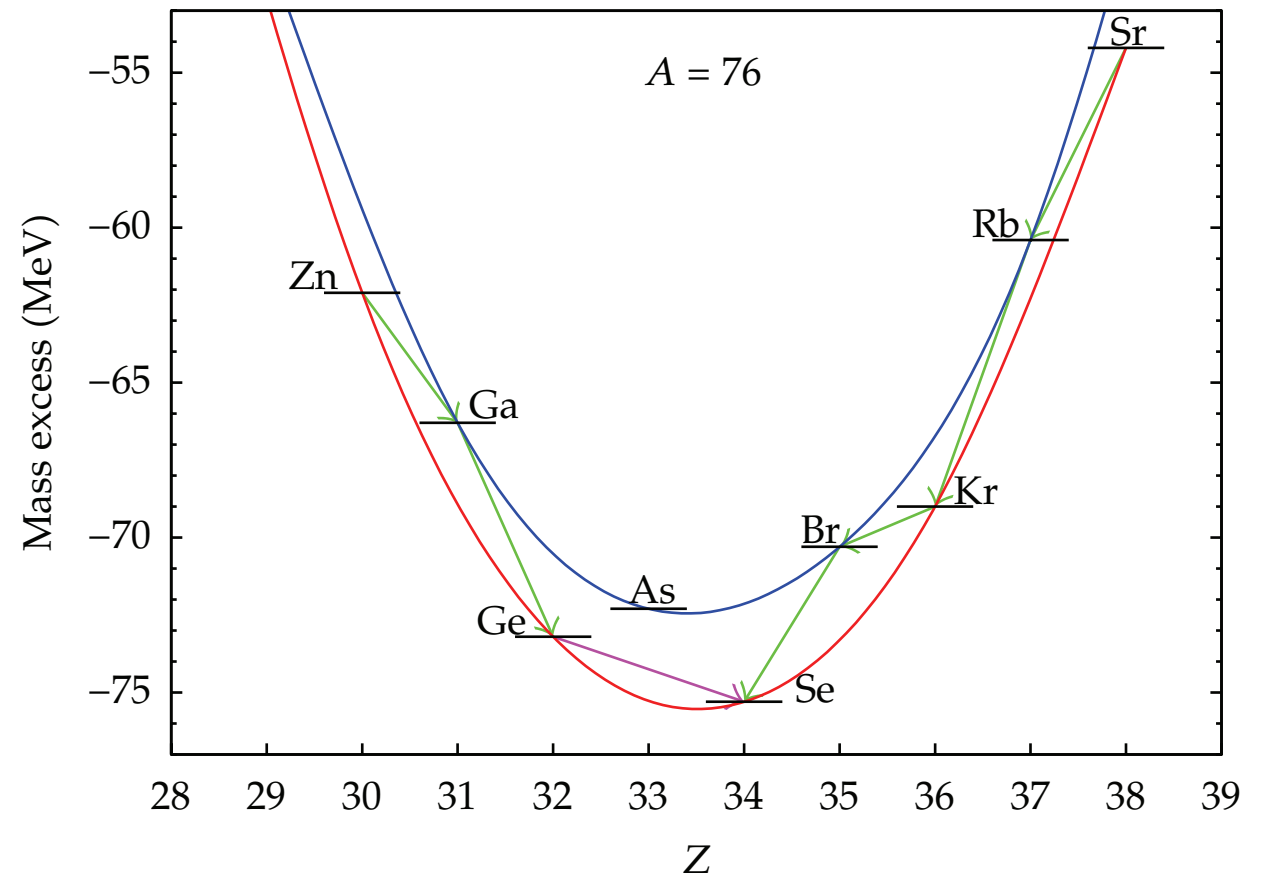


Double-beta decay

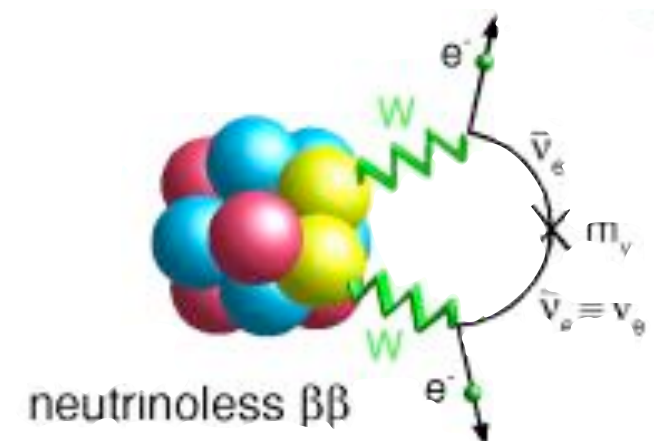
- Certain nuclei allow observable $\beta\beta$ decay



$$T_{1/2}^{2\nu\beta\beta} \gtrsim 10^{19} \text{ y}$$



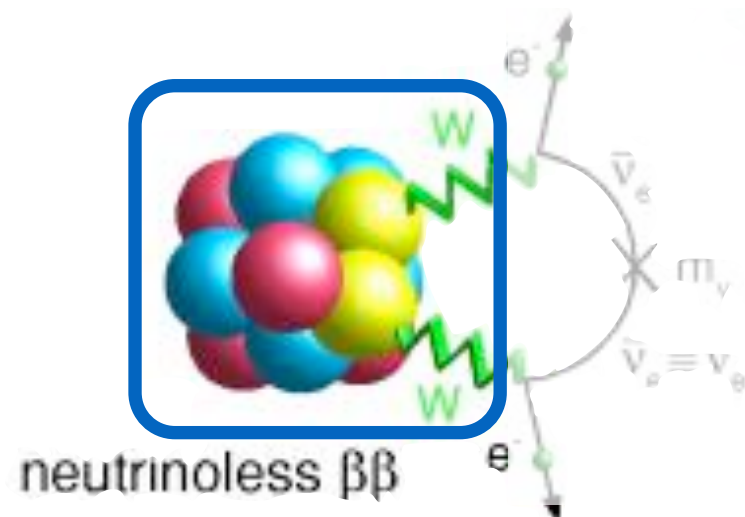
- If neutrinos are massive Majorana fermions $0\nu\beta\beta$ decay is possible
- In addition to light Majorana neutrino exchange, short-distance contributions to $0\nu\beta\beta$ can arise from BSM physics resulting in dim-9 operators in SMEFT



$$T_{1/2}^{0\nu\beta\beta} > 10^{25} \text{ y}$$

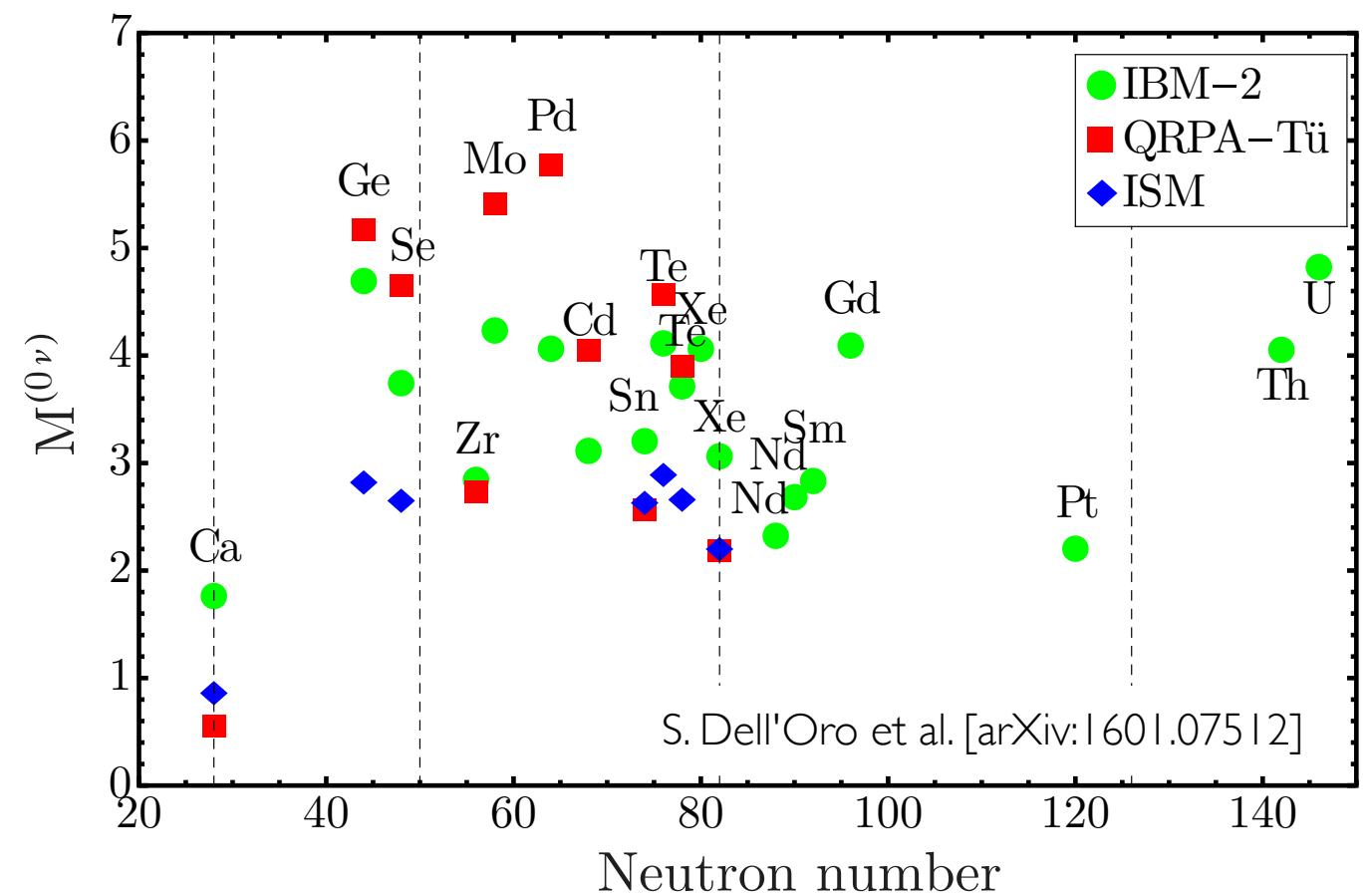
Double-beta decay

Want to understand $2\nu\beta\beta$ and $0\nu\beta\beta$ decay from theory



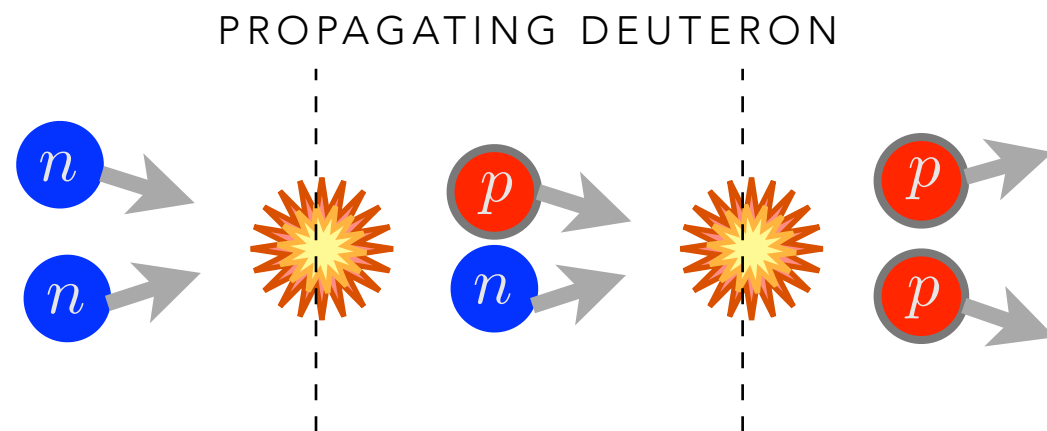
Calculate two-current
nuclear matrix elements
→ dictate half-life

Model calculations have large uncertainties

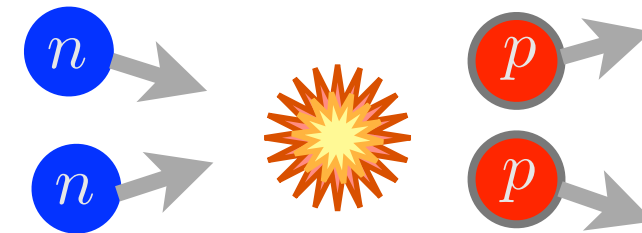


Neutrinoless double-beta decay

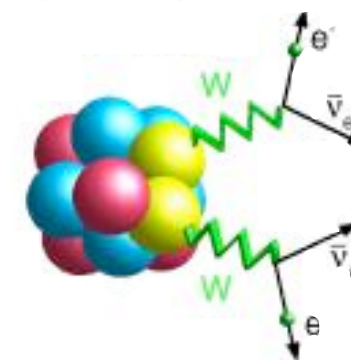
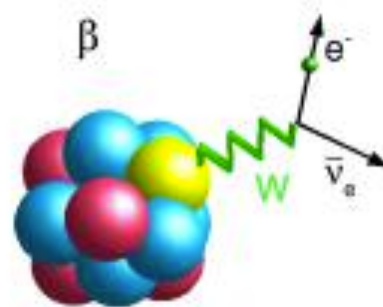
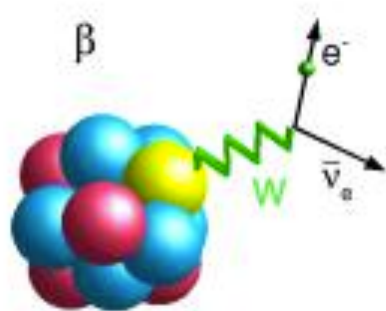
Lattice QCD: Calculate $nn \rightarrow pp$ transition matrix element



Two single-beta decays



Two-body effect



[NPLQCD PRL 119, 062003 (2017), PRD 96, 054505 (2017)]



Neutrino double-beta decay

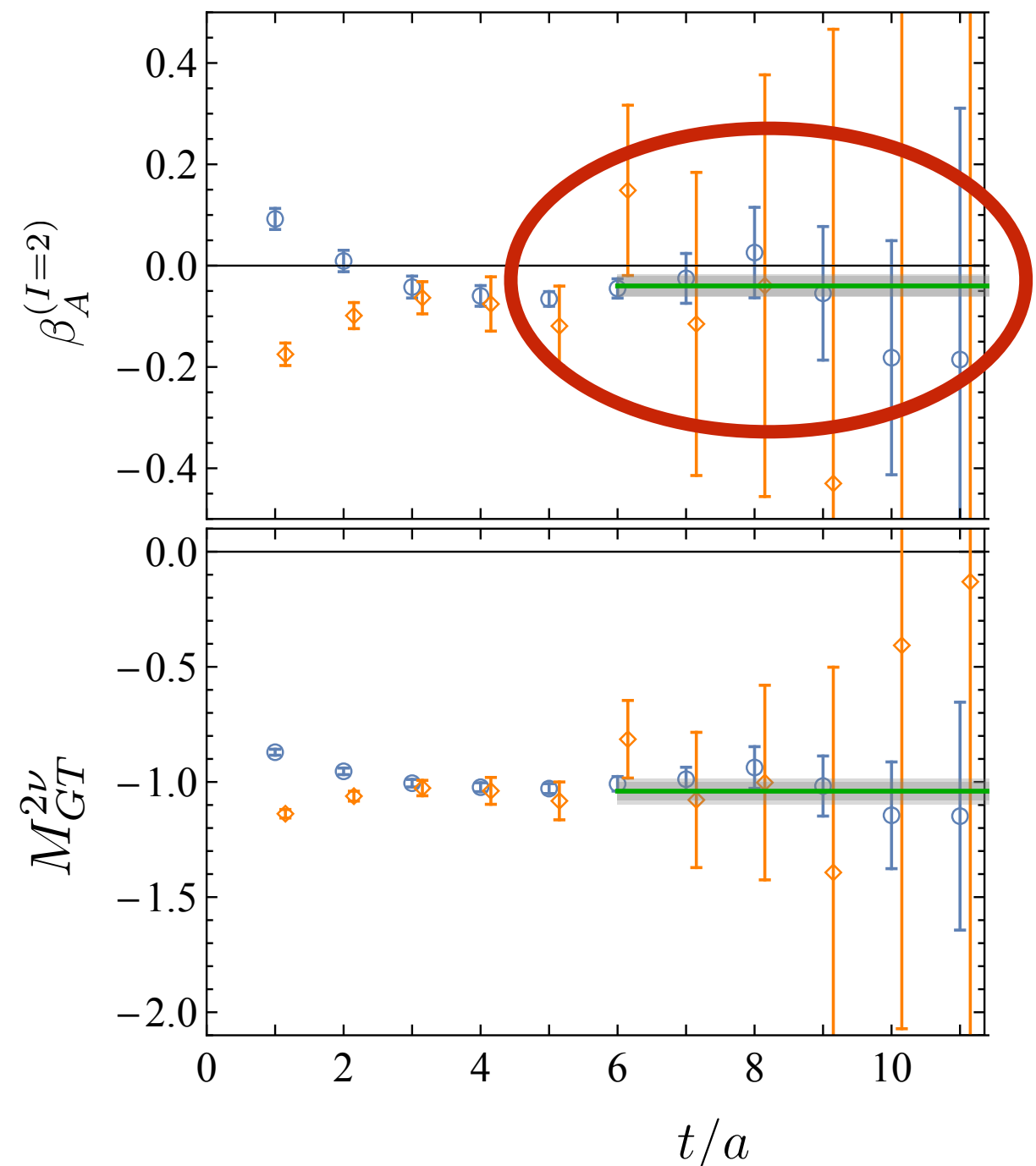
- Non-negligible deviation from deuteron intermediate state contribution

$$M_{GT}^{2\nu} = -\frac{|M_{pp \rightarrow d}|^2}{E_{pp} - E_d} + \beta_A^{(I=2)}$$

Isotensor axial polarisability

➔ Multi-body effects can't be neglected!

- TBD: connect to models / effective field theory for larger systems



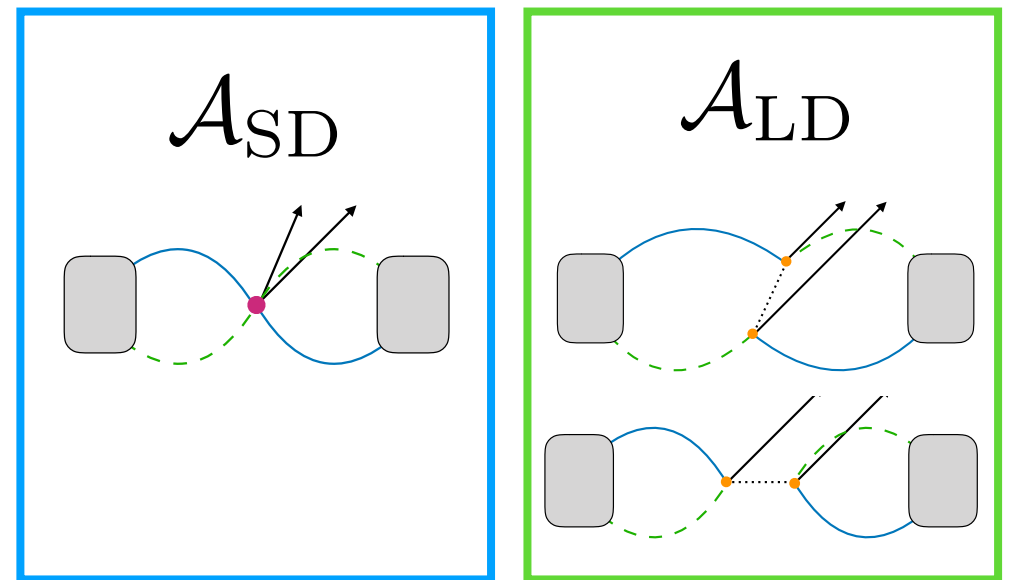
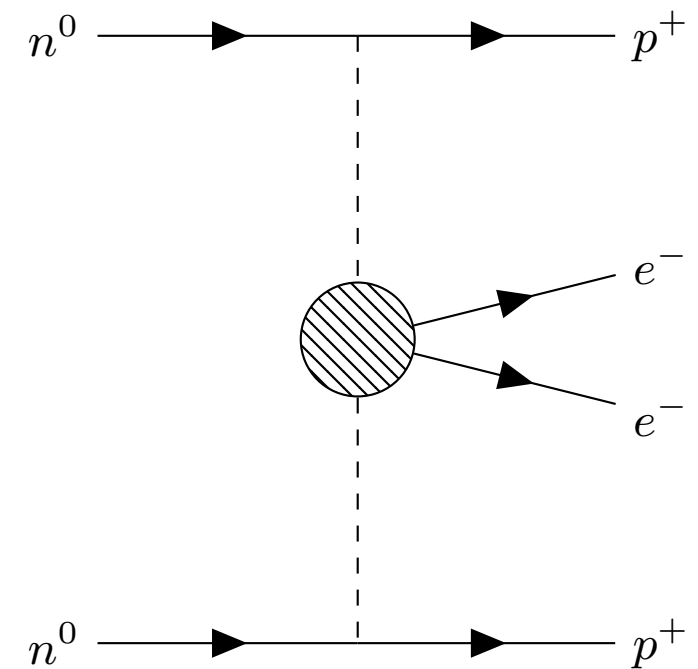
[NPLQCD PRL 119, 062003 (2017), PRD 96, 054505 (2017)]

Neutrinoless double-beta decay

- Calculations of pion matrix elements (statistically clean) give access to key pion exchange contributions to neutrinoless decay process
- New calculations compare short (SD) and long-distance (LD) contributions
- With estimates of BSM Wilson coefficients for minimal left-right symmetric model with heavy right-handed neutrinos:

[Cirigliano et al, JHEP 97 (2018)]:

$$\frac{\mathcal{A}_{\text{SD}}}{\mathcal{A}_{\text{LD}}} = \frac{1}{v^2} \frac{\sum_k |c_k \langle \pi | \mathcal{O}_k | \pi \rangle|}{|M^{0\nu}|} \sim 10^{-4}$$

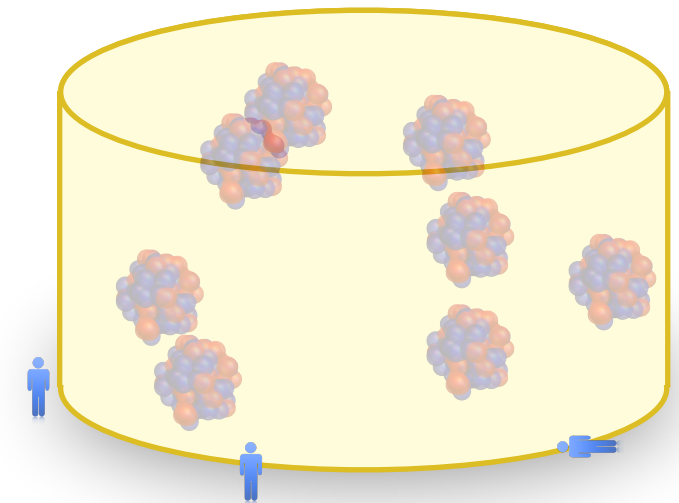


Dark matter direct detection

Look for scattering of WIMP dark matter on nuclear target

Detection rate depends on

- Dark matter properties
- Probability of interaction with nucleus
i.e., *nuclear* effects are important



Low-energy limit of a generic
spin-independent interaction
is scalar



Determine **nucleon and
nuclear scalar matrix
elements** from lattice QCD

Other e.g., spin-dependent couplings can also be constrained

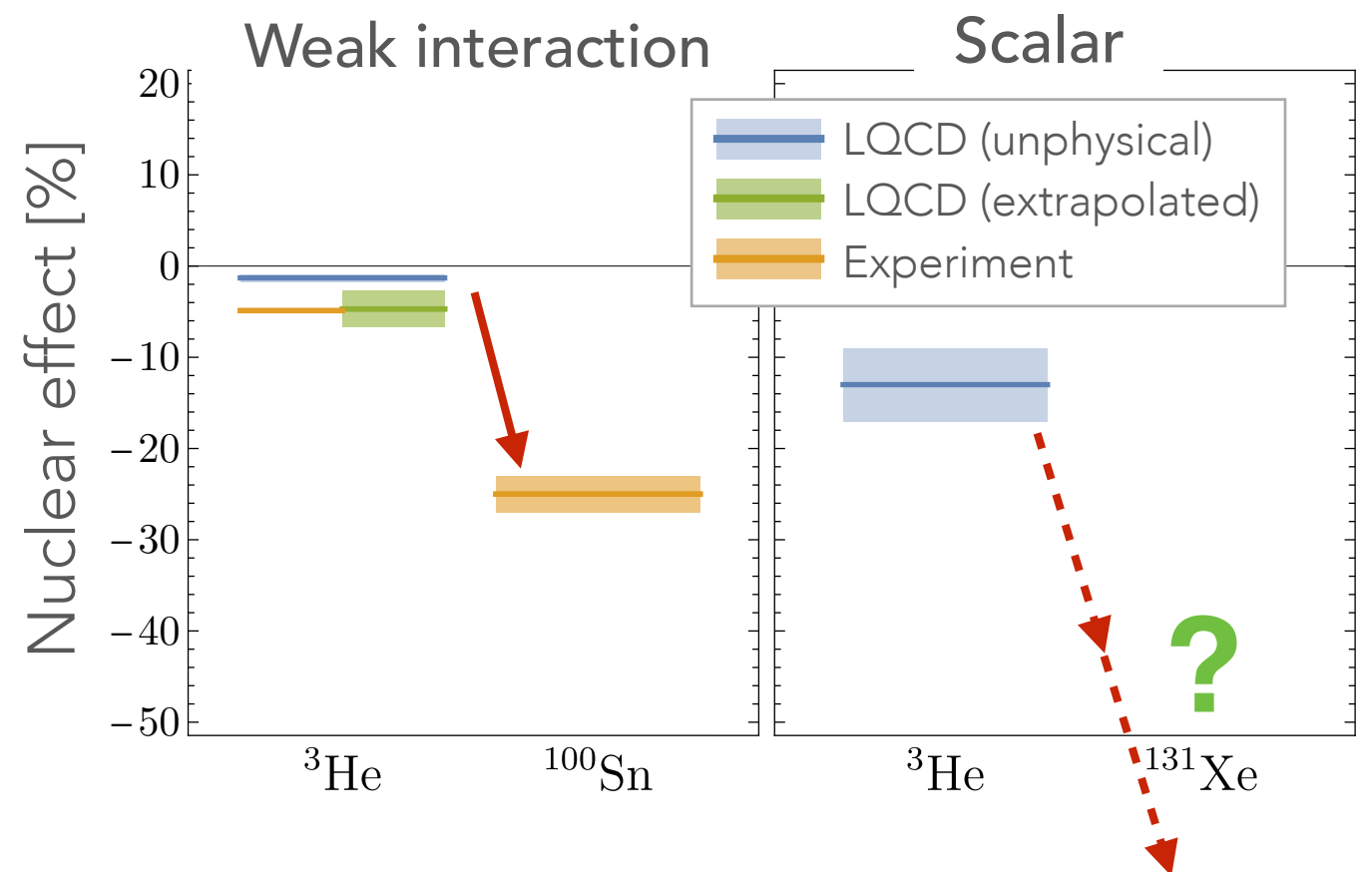
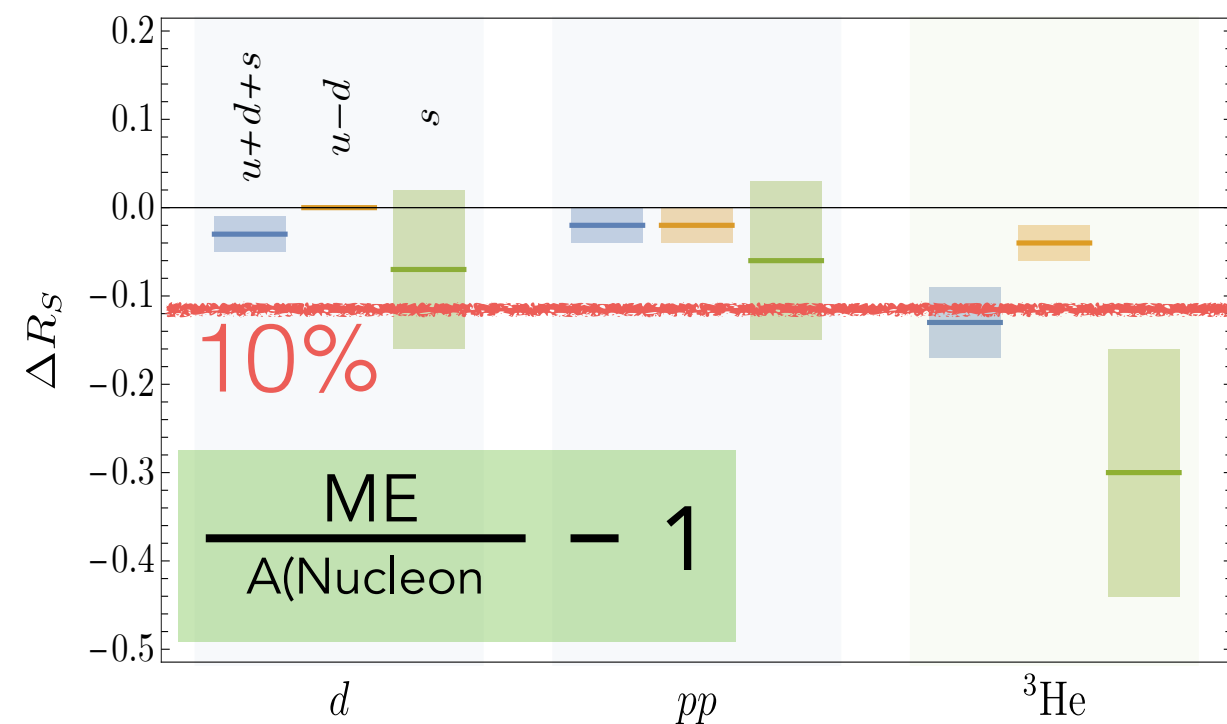
e.g., [Hoferichter et al., arXiv:1503.04811], [Hill et al., arXiv:1409.8290], [Fitzpatrick et al., arXiv:1203.3542]



Scalar matrix elements of nuclei

- Lattice QCD calculation with $m_\pi \sim 800$ MeV shows 10% nuclear effects in scalar MEs of ^3He → potentially very significant effects in larger nuclei e.g., Xenon
- Calculations in progress with \sim physical quark masses


Scalar MEs: deviation from naive scaling with A



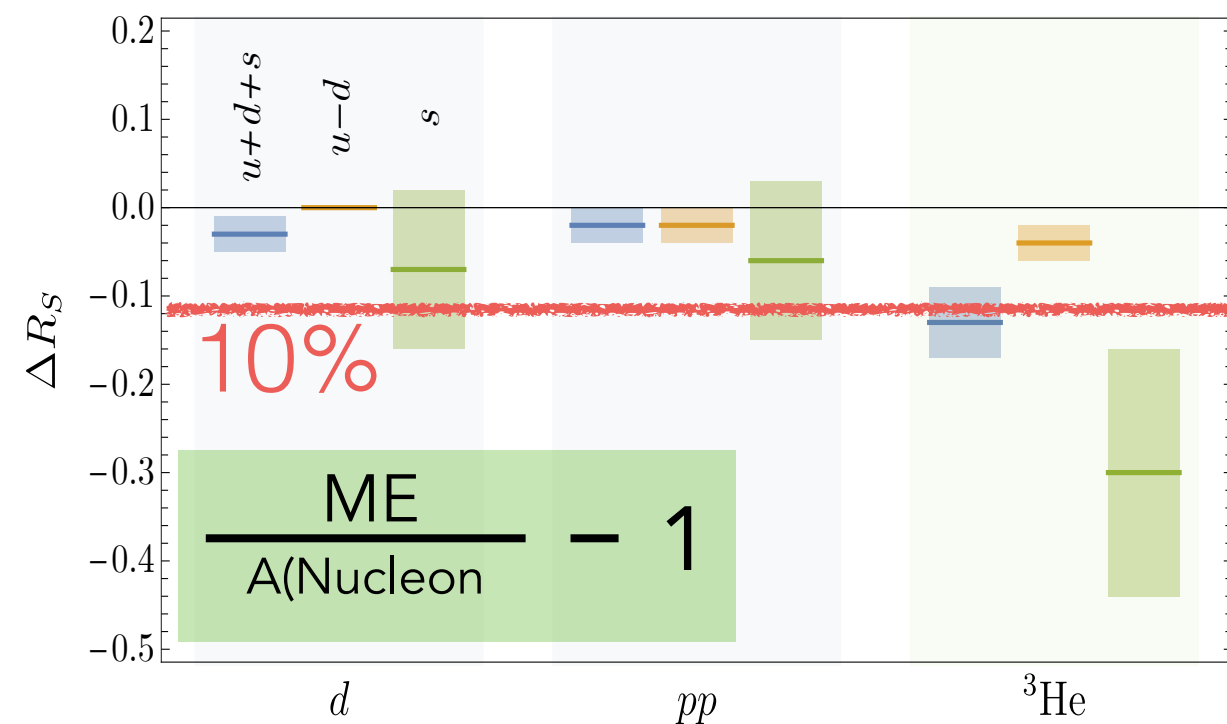
[NPLQCD PRL120 (2018), 152002]



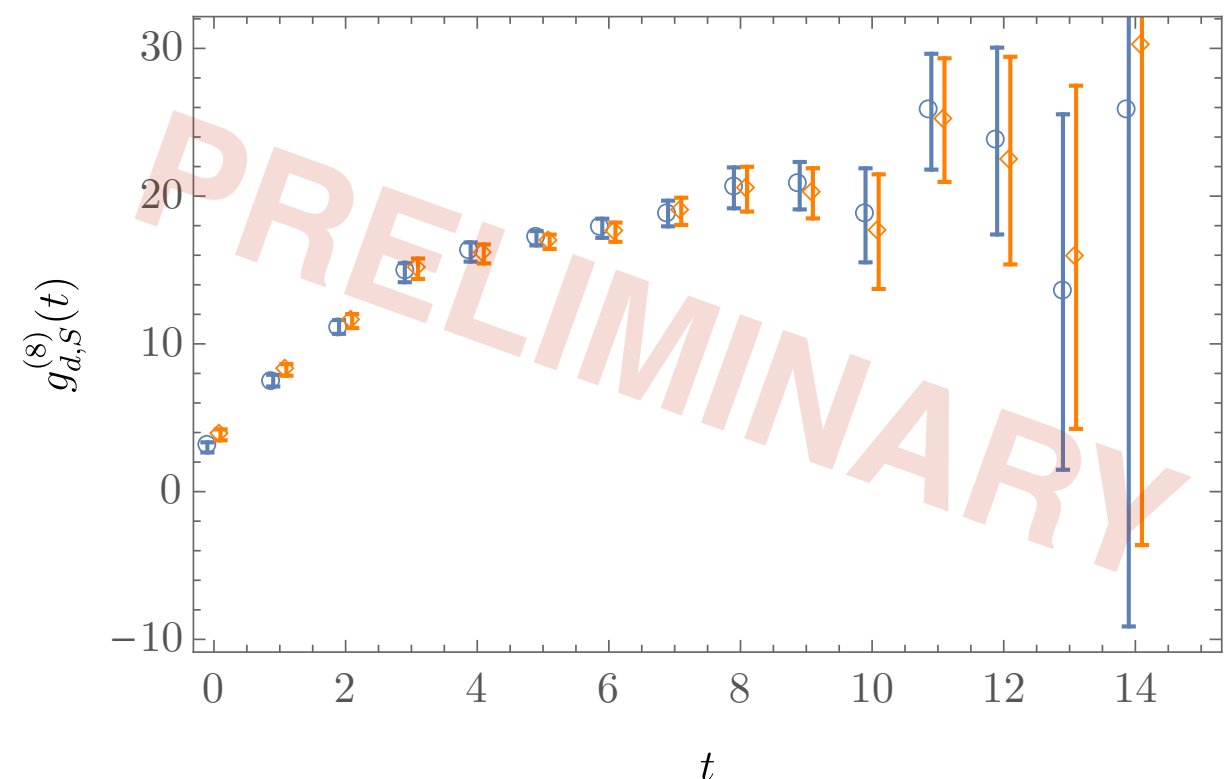
Scalar matrix elements of nuclei

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- Calculations in progress with \sim physical quark masses

Scalar MEs: deviation from naive scaling with A



Signal for physical scalar ME of the deuteron



[NPLQCD PRL120 (2018), 152002]

Nuclear physics from lattice QCD

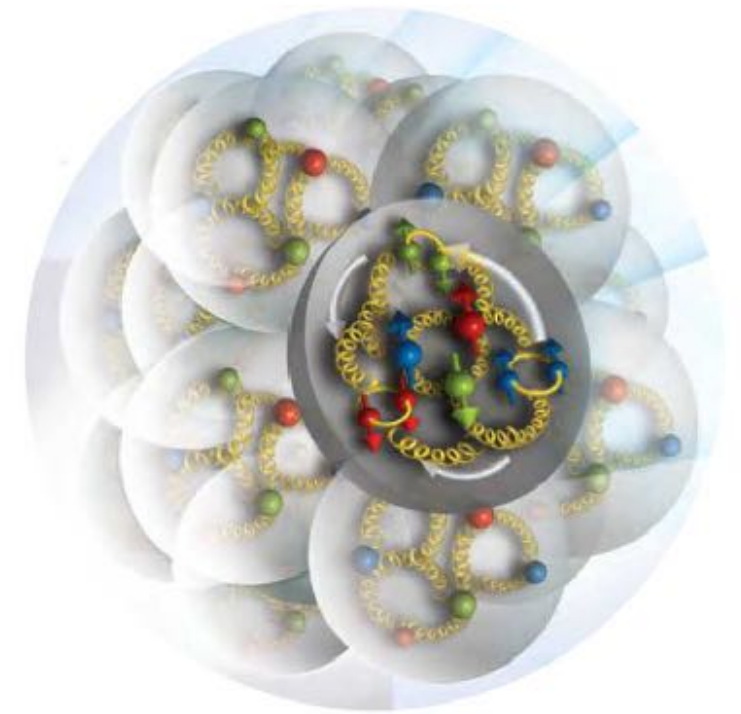
Calculations of nuclear MEs are HARD

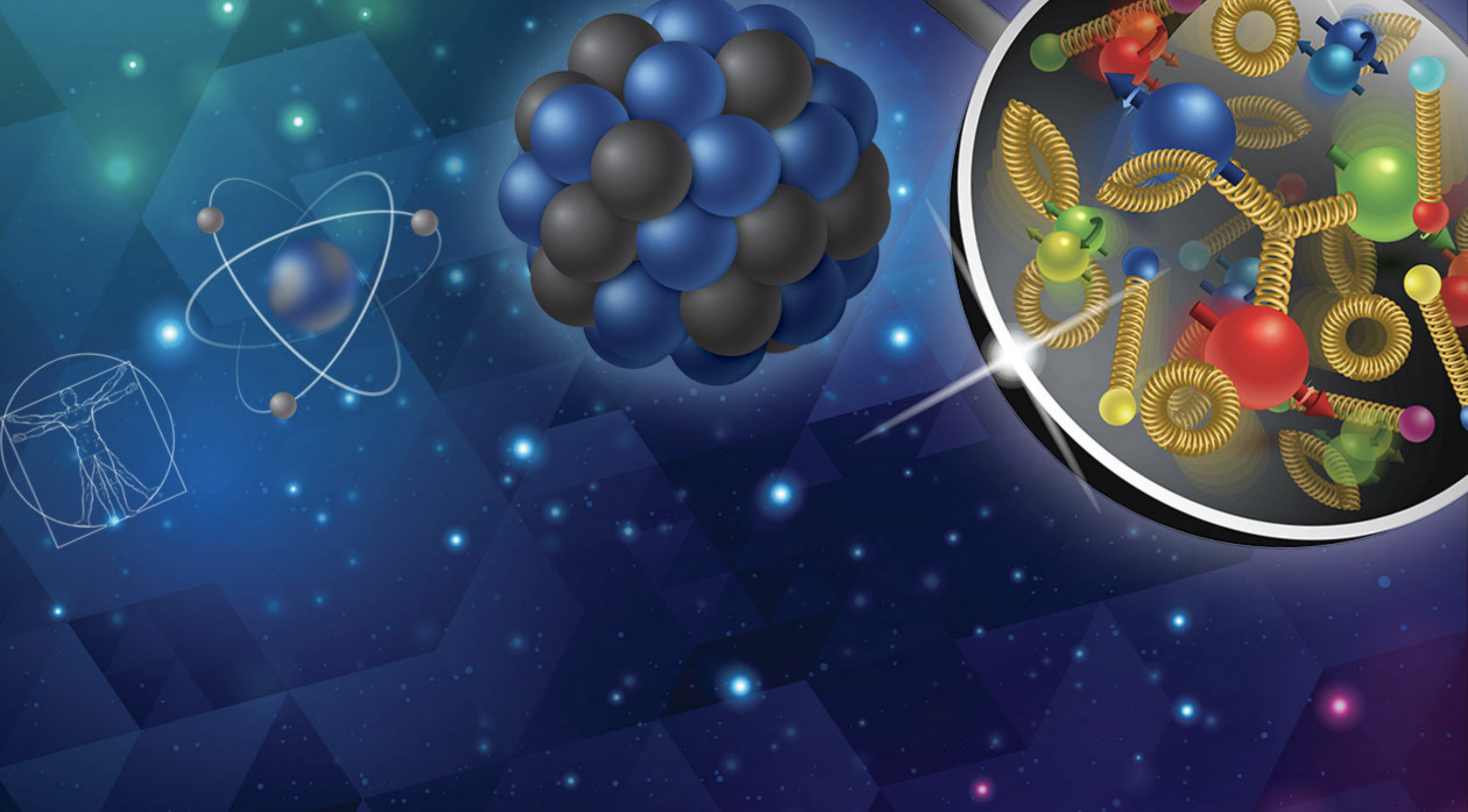
Constraints on nuclear matrix elements are possible:

- Pipeline well-defined and tested
- Still quite far from controlled calculations

Controlled calculations achievable for nuclei with $A < 5$ with ~ 10 - 20% uncertainty in 10-year timeframe:

- Axial MEs, including form factors
- Scalar MEs relevant for e.g., dark matter direct detection
- Double beta-decay matrix elements
- Constraints on PDFs, GPDs of nuclei via moments
- ...





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Technology

