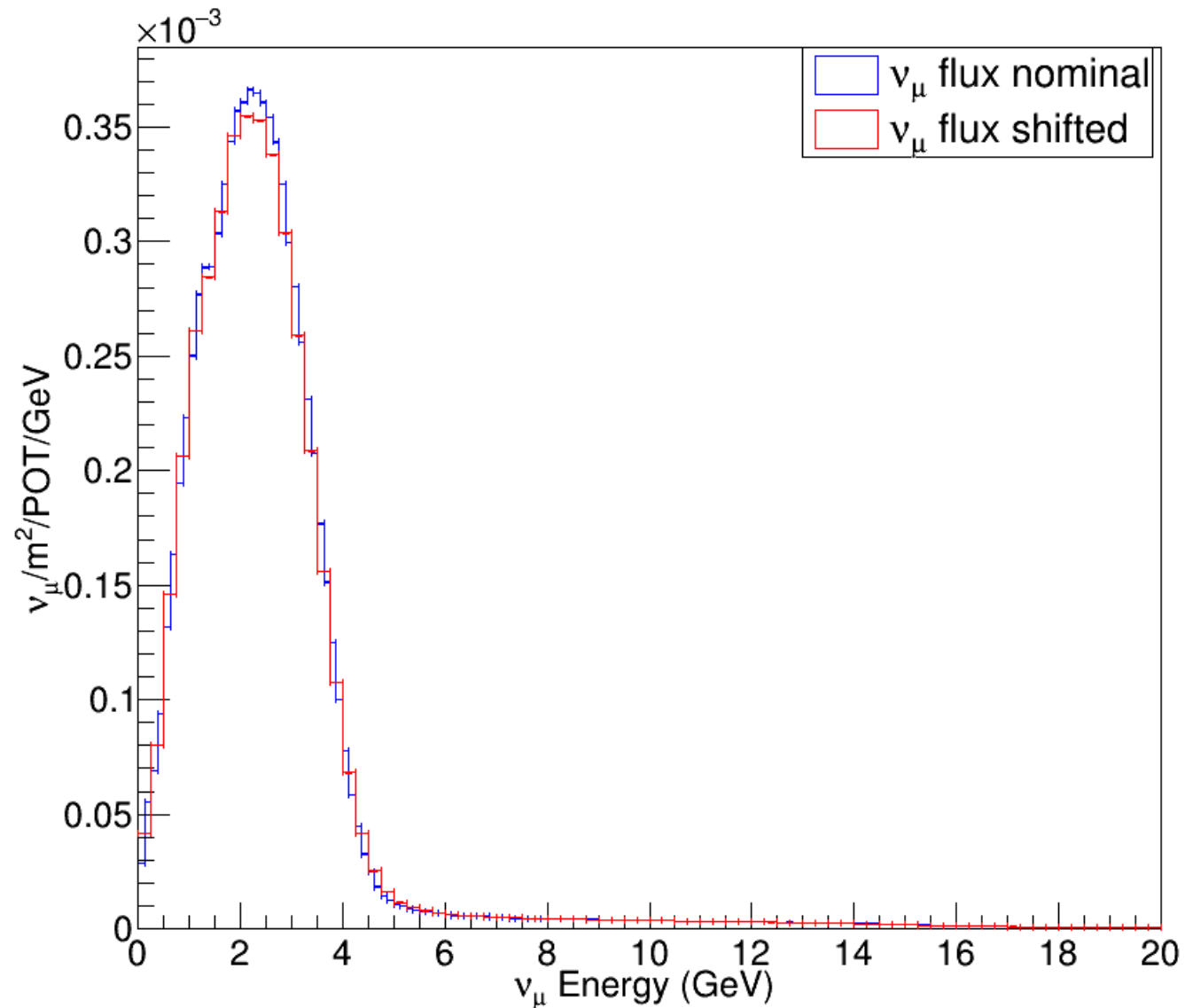


BEAM MONITORING WITH SAND

- My thesis studied the beam monitoring capabilities of the SAND detector, via the comparison of the distribution of an observable sensitive to beam anomalies
- **Observable:** reconstructed muon momentum
- **First systematic:** horn-1 0.5 mm Y shift
- **Test statistics:** $T \sim \chi^2$

SIMULATED SAMPLES

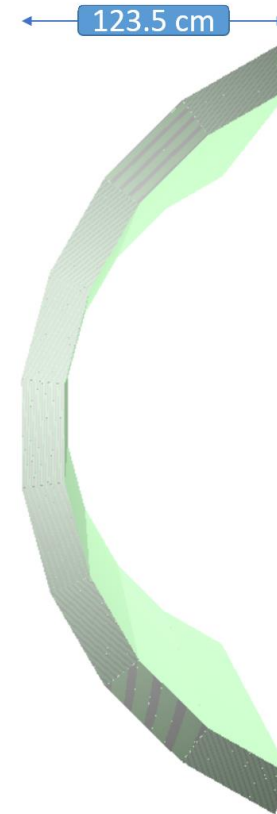
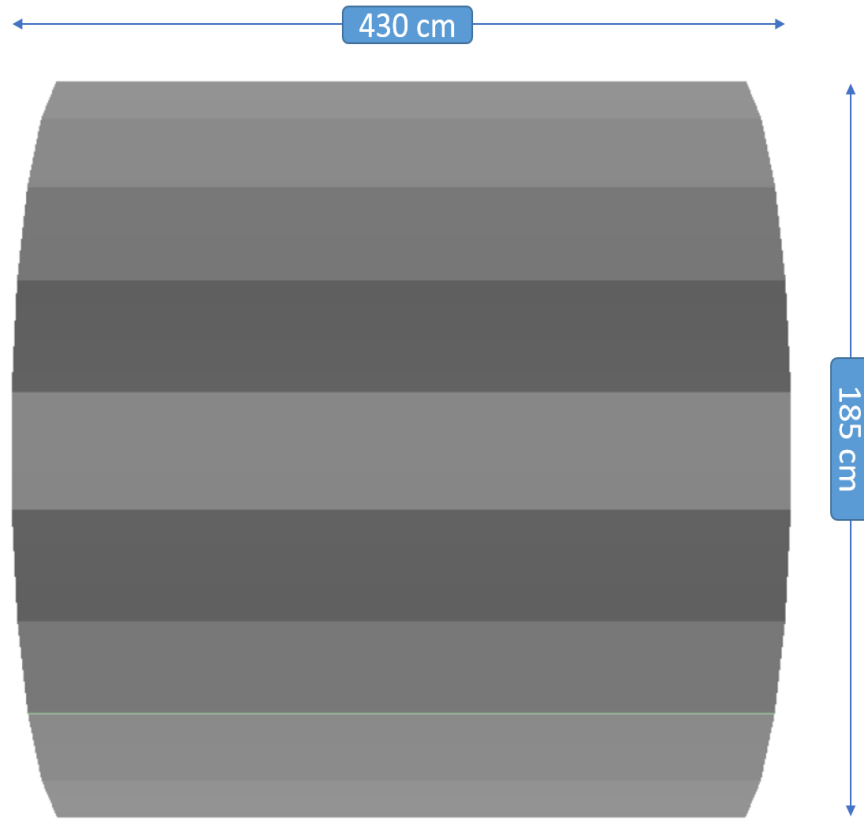


Simulated samples from :

- **Nominal** neutrino flux
- **Shifted** neutrino flux
(Y +0.5mm in the first beam horn)

Note : Histograms retrieved from
<https://home.fnal.gov/~ljf26/DUNEFluxes>

TARGET: FRONT CALORIMETERS

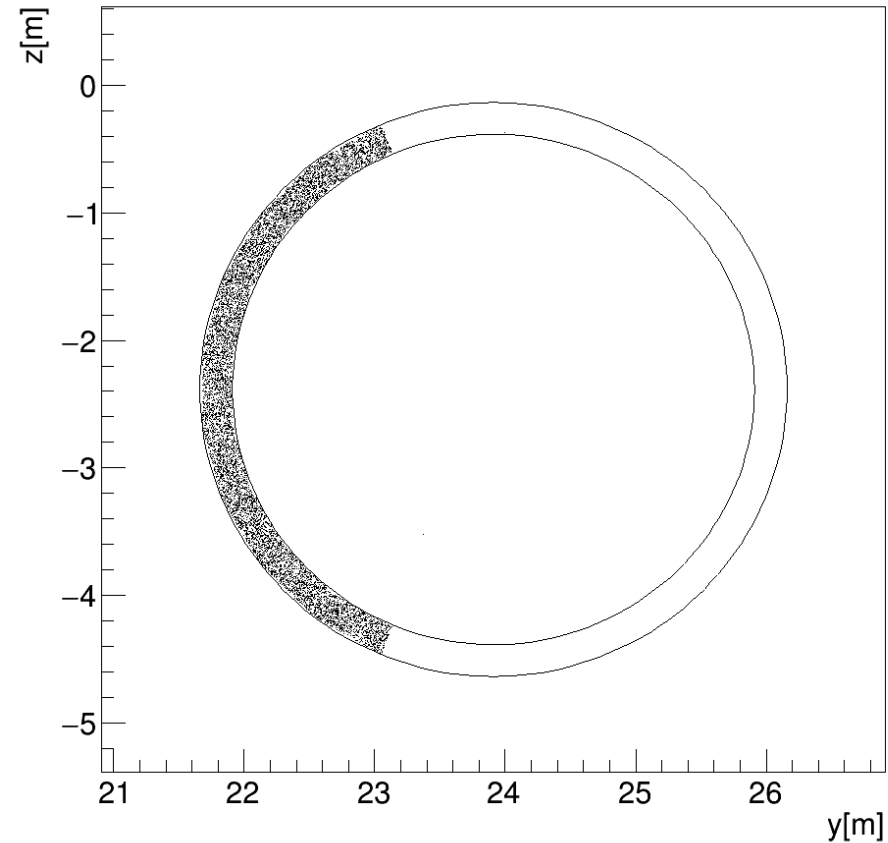
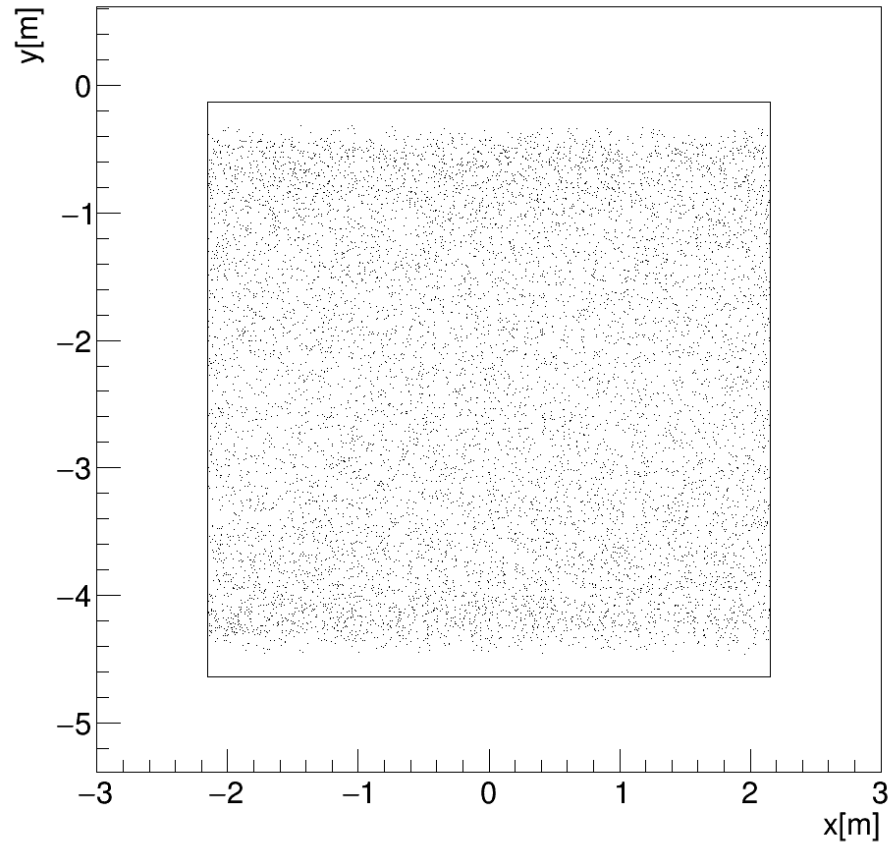


- *Dunendggd* simulated geometry
- Target: 9 front calorimeter barrel modules

$$\begin{aligned} N_{week} &= r_{CC} \times f_{CC}^{-1} \times m_{mod} \times n_{mod} \\ &= (4.865 \times 10^5 \text{ ev}[v_{\mu}(CC)]/\text{ton/week}) \times (2/3)^{-1} \times (2.75 \text{ ton}) \times (9) \\ &\simeq 1.75 \times 10^6 \end{aligned}$$

One week's worth of statistics

TARGET: FRONT CALORIMETERS

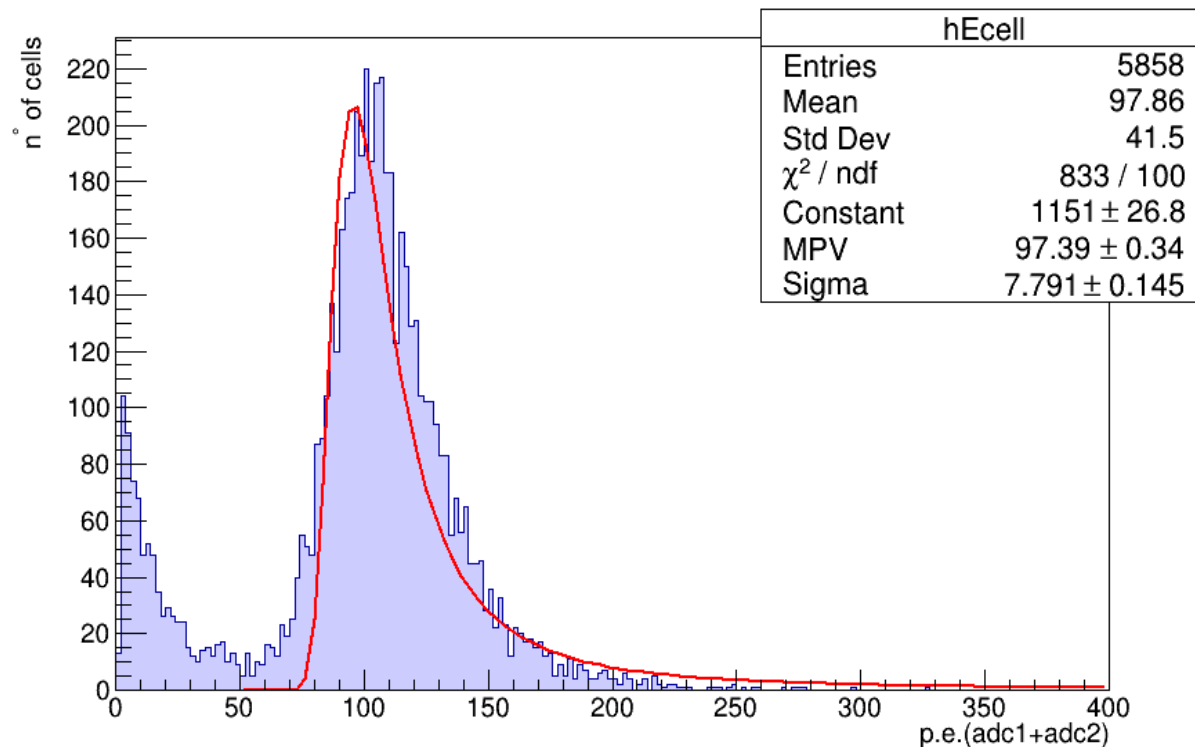
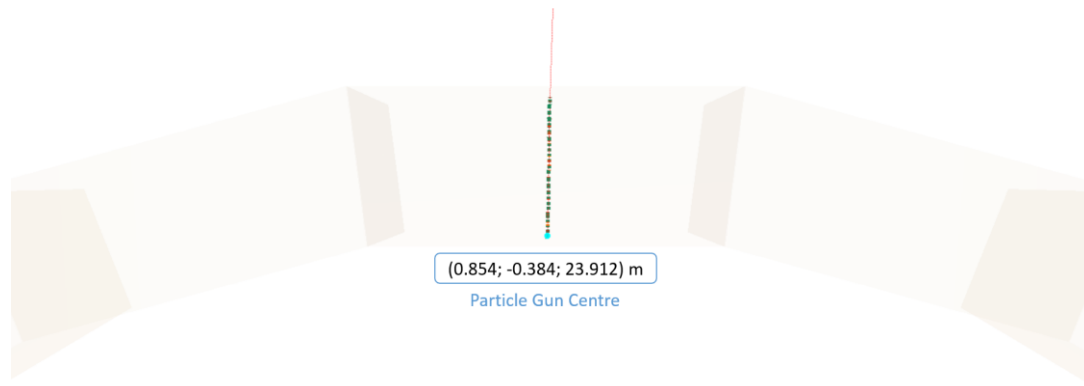


- Interaction vertex distributions on the xy and yz planes

$$\begin{aligned} N_{week} &= r_{CC} \times f_{CC}^{-1} \times m_{mod} \times n_{mod} \\ &= (4.865 \times 10^5 \text{ ev}[v_{\mu}(CC)]/\text{ton}/\text{week}) \times (2/3)^{-1} \times (2.75 \text{ ton}) \times (9) \\ &\simeq 1.75 \times 10^6 \end{aligned}$$

One week's worth of statistics

PRELIMINARY MEASUREMENT: LIGHT YIELD



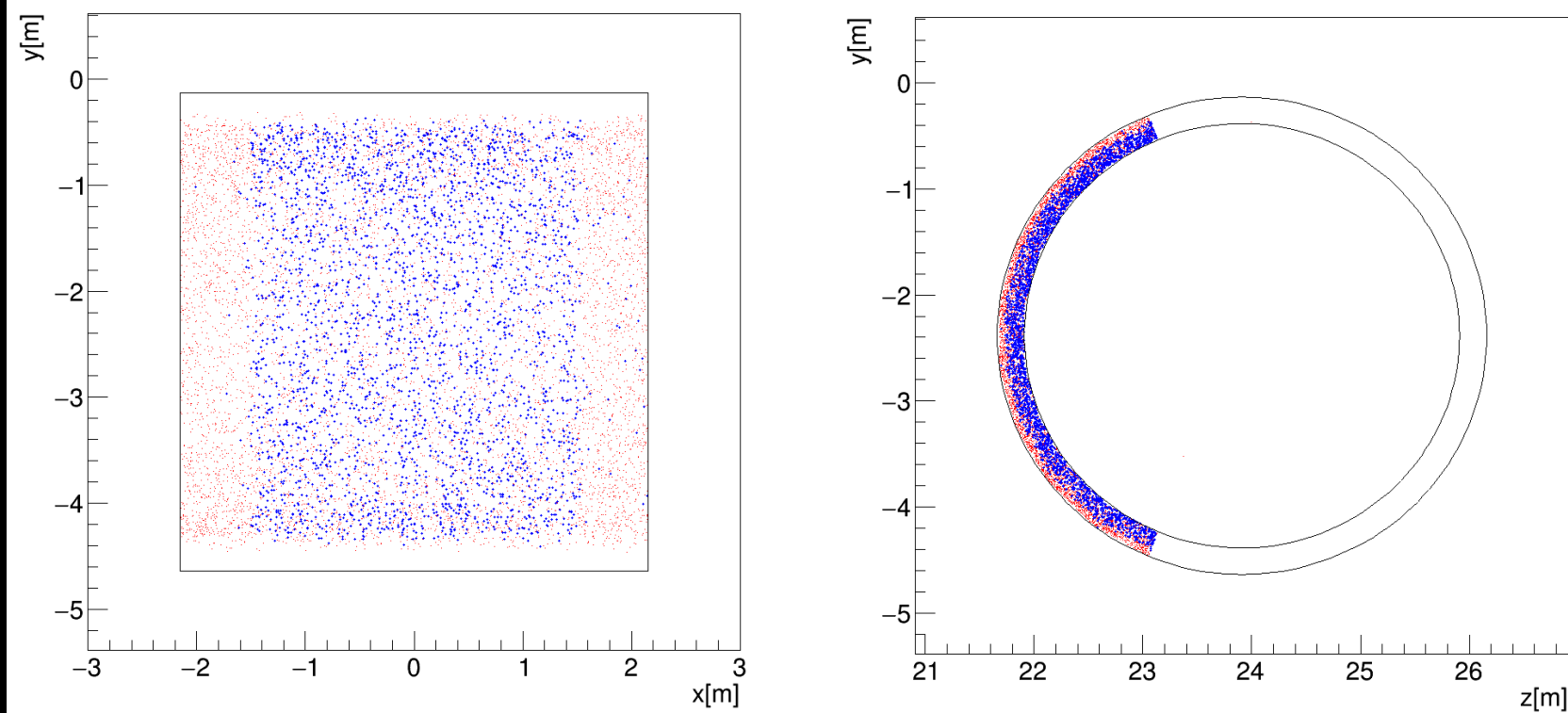
- Simulated 1000 muons at 10 GeV passing through an ECAL barrel module

$$\Delta E_{\text{cell}} \simeq \left(\frac{dE}{dx} \right)^{\text{MIP}} \rho_{\text{Pb}} \Delta x_{\text{Pb}} + \left(\frac{dE}{dx} \right)^{\text{MIP}} \rho_{\text{Sc}} \Delta x_{\text{Sc}} \simeq 42.22 \text{ MeV}$$

$$N_{\text{p.e.}}^{\text{cell}} = (97.4 \pm 0.3) \text{ p.e.}$$

$$c = \frac{N_{\text{p.e.}}^{\text{cell}}}{\Delta E_{\text{cell}}} \simeq 2.31 \text{ [p.e./MeV]}$$

FIDUCIAL CUTS



Spatial distribution in ND hall global coordinates of the true neutrino interaction vertexes of the events that survive the fiducial cut (blue) and those that don't (red).

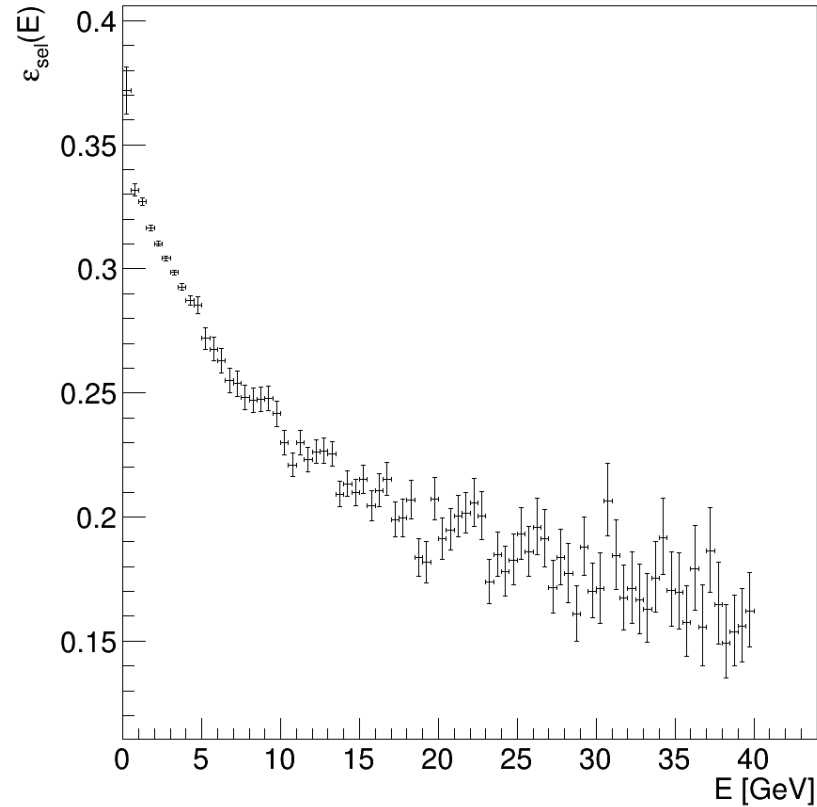
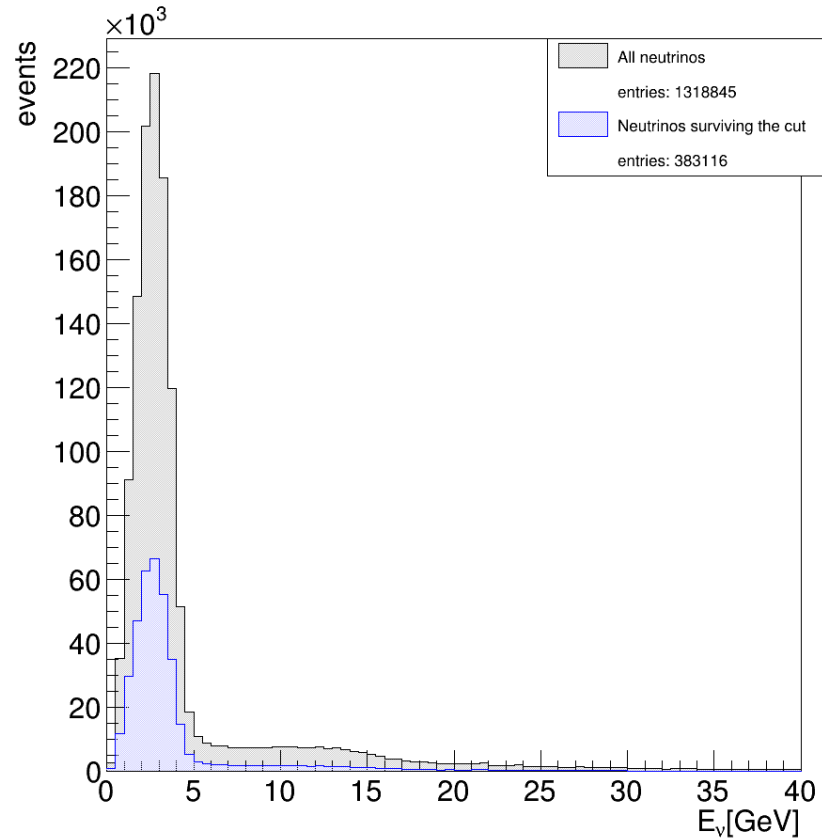
- Threshold on the energy deposition on the outer layer ($E < 15\text{MeV}$):

$$N_{p.e.}^{th} = c \times \Delta E_{th} \simeq 35 \text{ p.e.}$$

- X vertex, estimated as a weighted average on the energy deposition on the cells, is selected as:

$$|x_V| \leq 1.5 \text{ m}$$

FIDUCIAL CUT



- Selection efficiency:

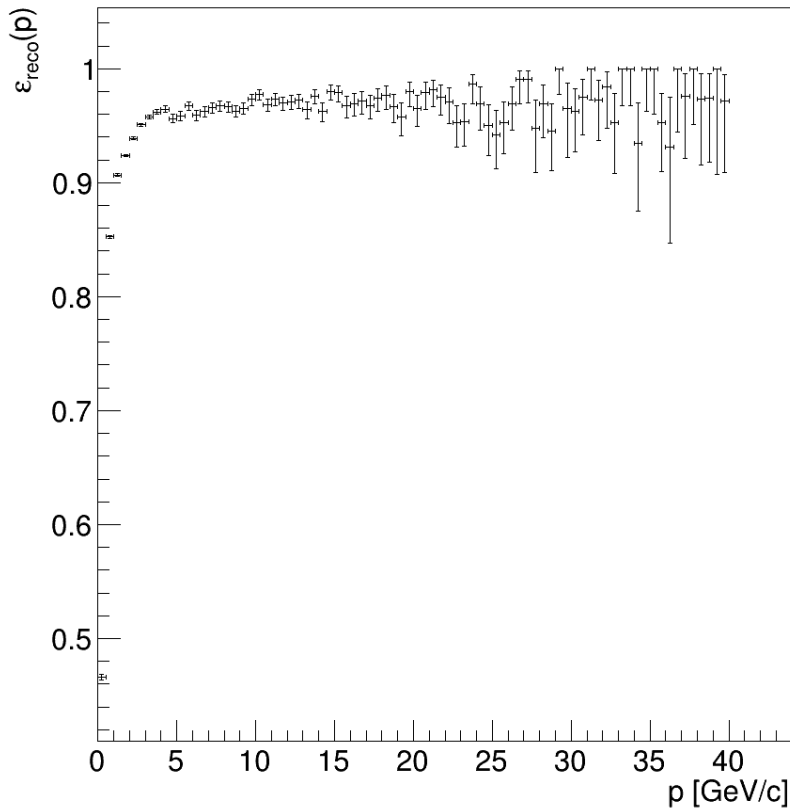
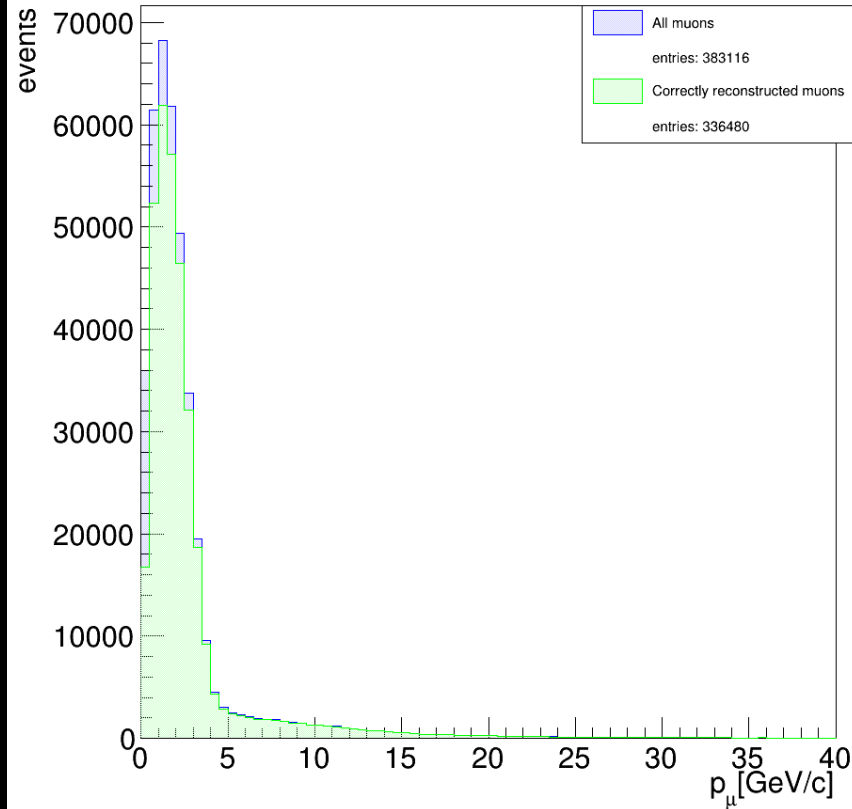
$$\epsilon_{cut} = \frac{N_{fid}}{N_{CC}} = 0.2905 \pm 0.0004$$

Note: efficiency decreases at higher energy; might be due to nuclei fragmentation in DIS

(Left) Energy (Monte Carlo truth) distribution of neutrinos from the CC nominal sample (grey); distribution surviving the fiducial cut (blue).

(Right) Selection efficiency as a function of neutrino energy from the Monte Carlo truth.

MOMENTUM RECONSTRUCTION SELECTION



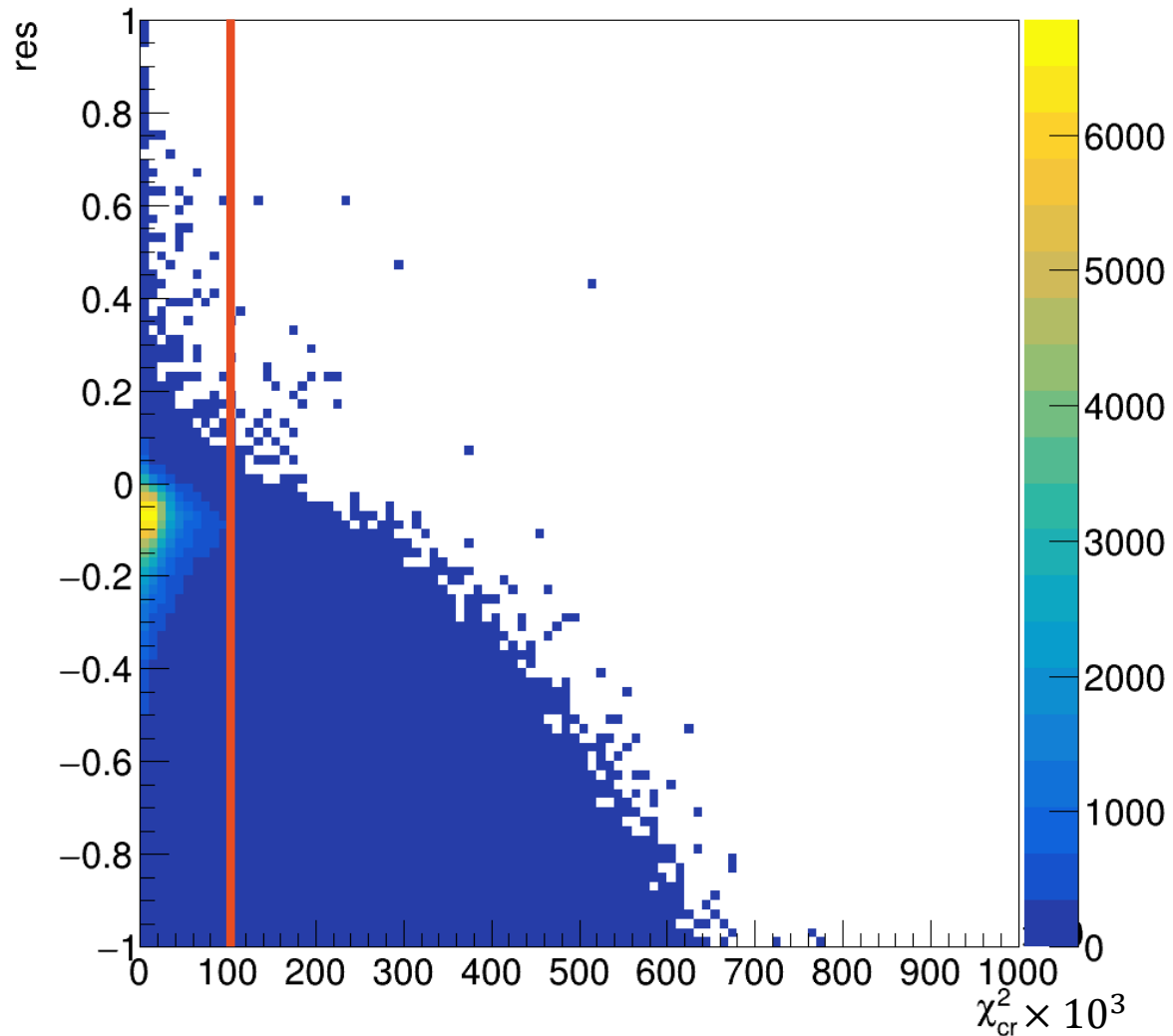
- Momentum reconstruction efficiency:

$$\varepsilon_{reco} = \frac{N_{reco}}{N_{fid}} = 0.9168 \pm 0.0004$$

(Left) Distributions of the true Monte Carlo momenta of the muons from the fiducial sample (blue) and only the ones correctly reconstructed (green).

(Right) Reconstruction algorithm efficiency as a function of the true Monte Carlo muon momentum

MOMENTUM RECONSTRUCTION QUALITY SELECTION



$$\chi_{cr}^2 = \frac{1}{N_{hits}} \sum_{i=1}^{N_{hits}} |(y_i - y_C)^2 + (z_i - z_C)^2 - R^2|$$

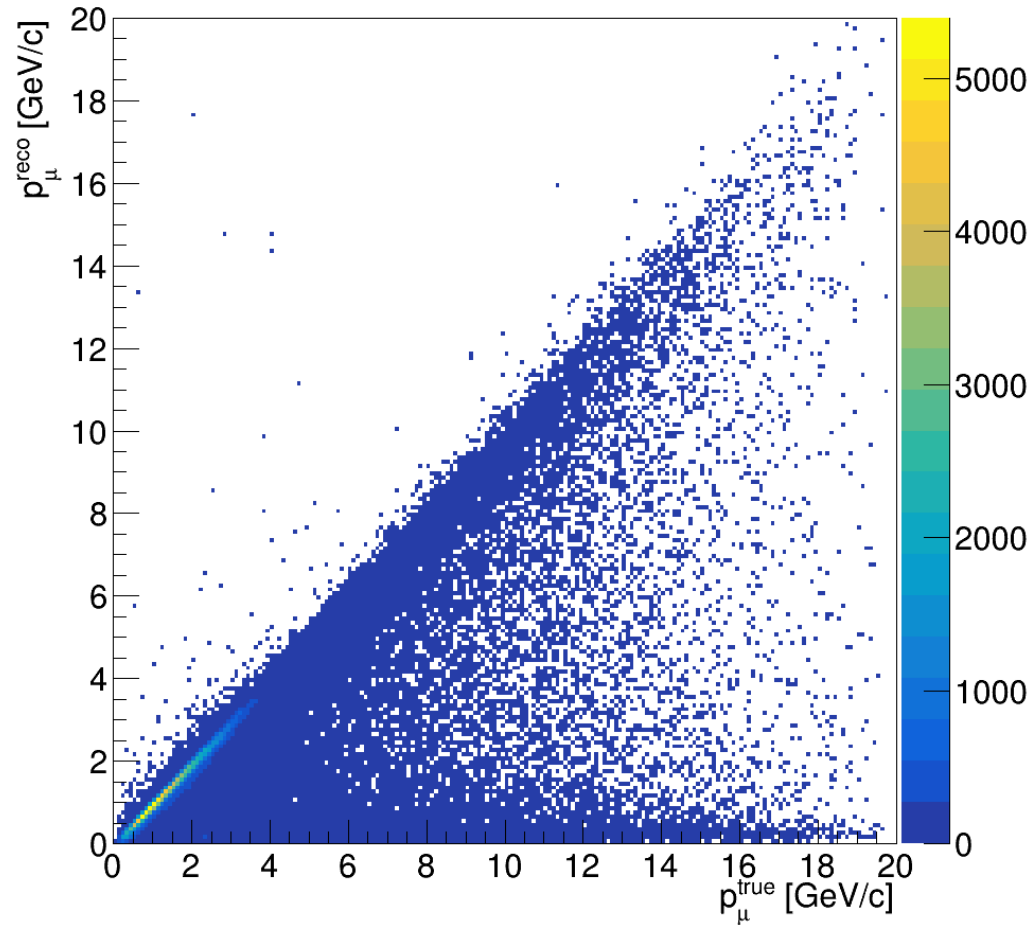
$$res = 1 - p_{\mu}^{true} / p_{\mu}^{reco}$$



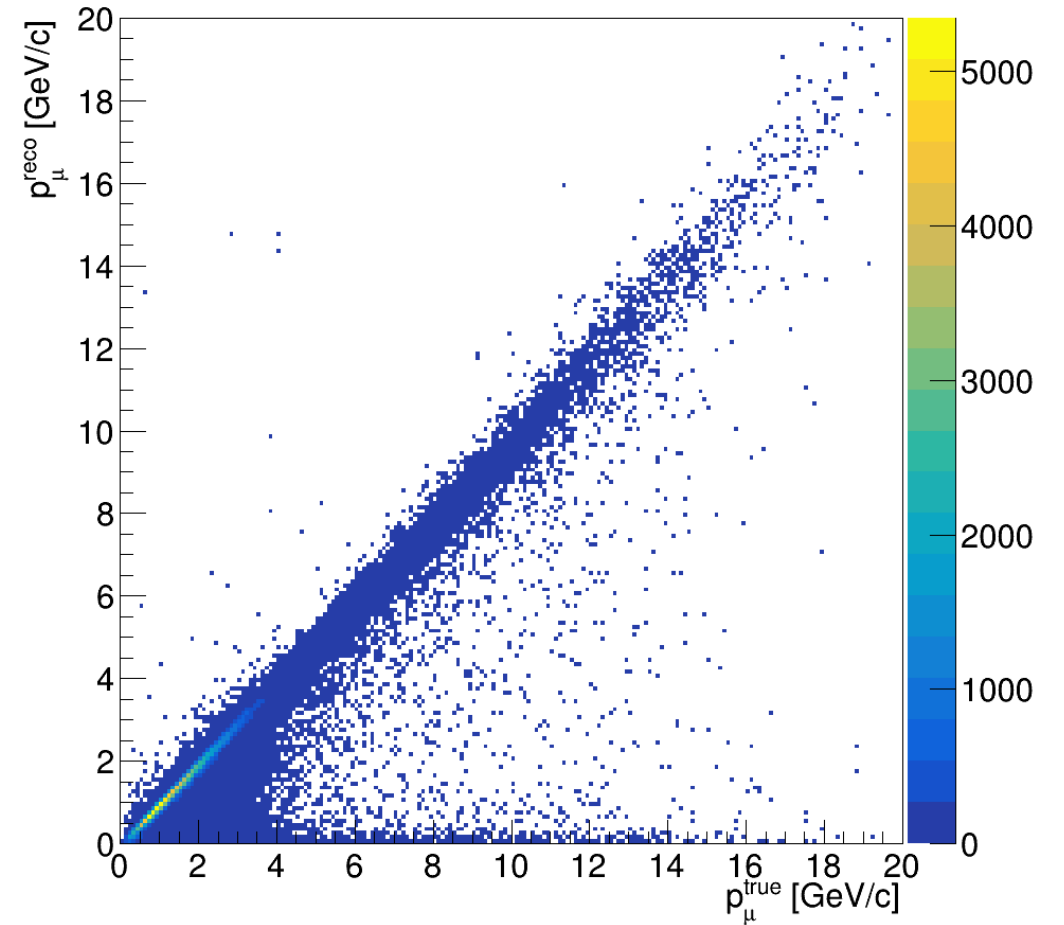
$$\chi_{cr}^2 < 10^5$$

$$\varepsilon_{qual} = \frac{N_{qual}}{N_{reco}} = 0.9290 \pm 0.0004$$

MOMENTUM RECONSTRUCTION QUALITY SELECTION



BEFORE CUT



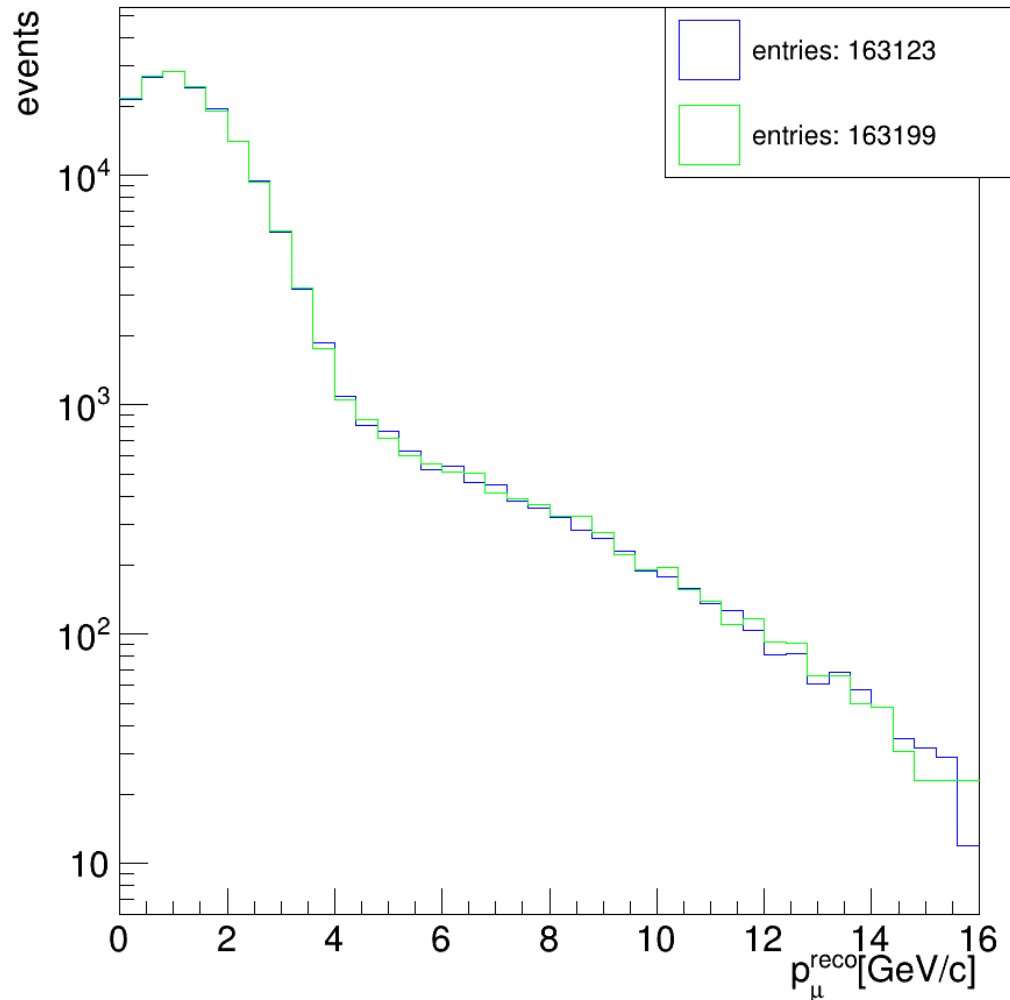
AFTER CUT

CHI-SQUARED TWO-SAMPLE TEST STATISTICS

$$T = \sum_{i=1}^k \frac{(u_i - v_i)^2}{u_i + v_i}$$

- Where k is the number of bins in the histograms and u_i and v_i are their contents
- T approximately follows a chi-squared distribution

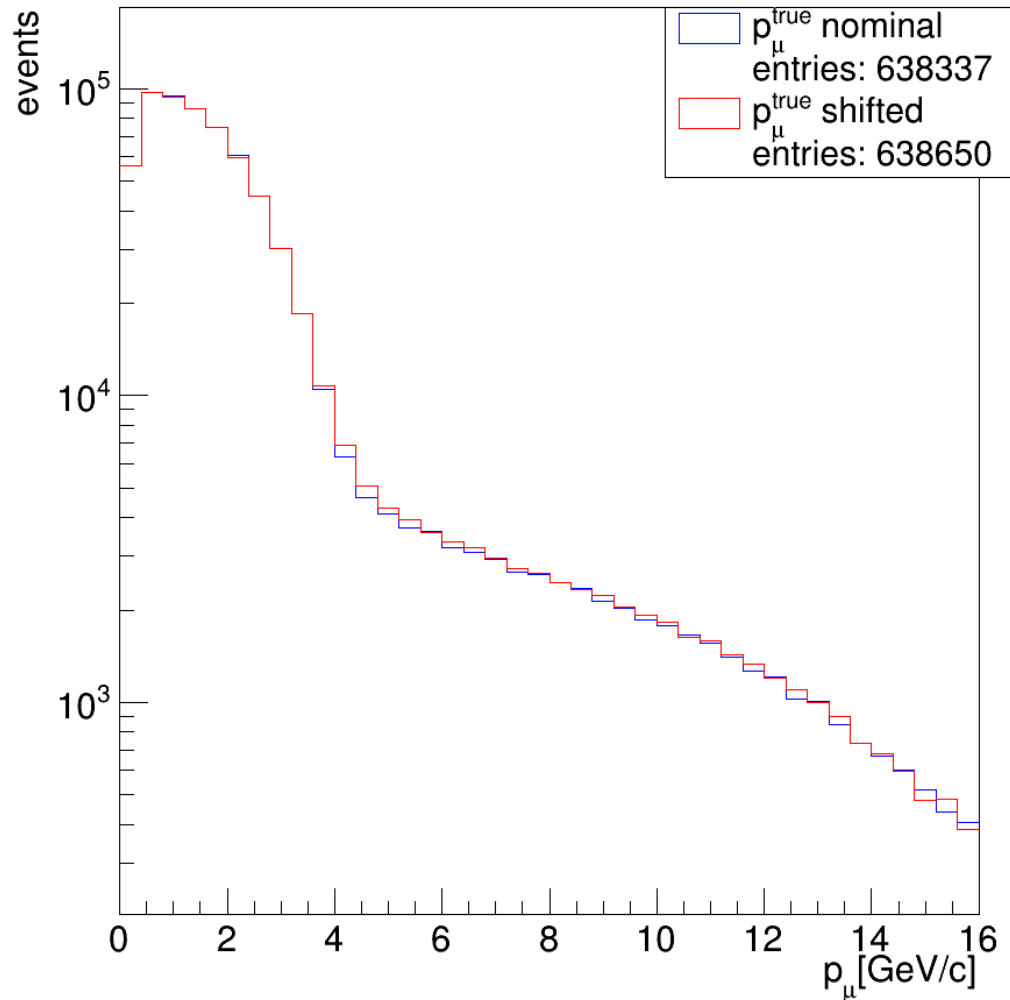
CONSISTENCY CHECK: TWO NOMINAL SAMPLES



- We apply the T two equally large nominal samples ($\sim 0.8 \times 10^6$ events):

$$p_{control} = 0.527; \quad \sigma_{control} = 0.633 \text{ } (\chi^2)$$

CONSISTENCY CHECK: IDEAL DETECTOR

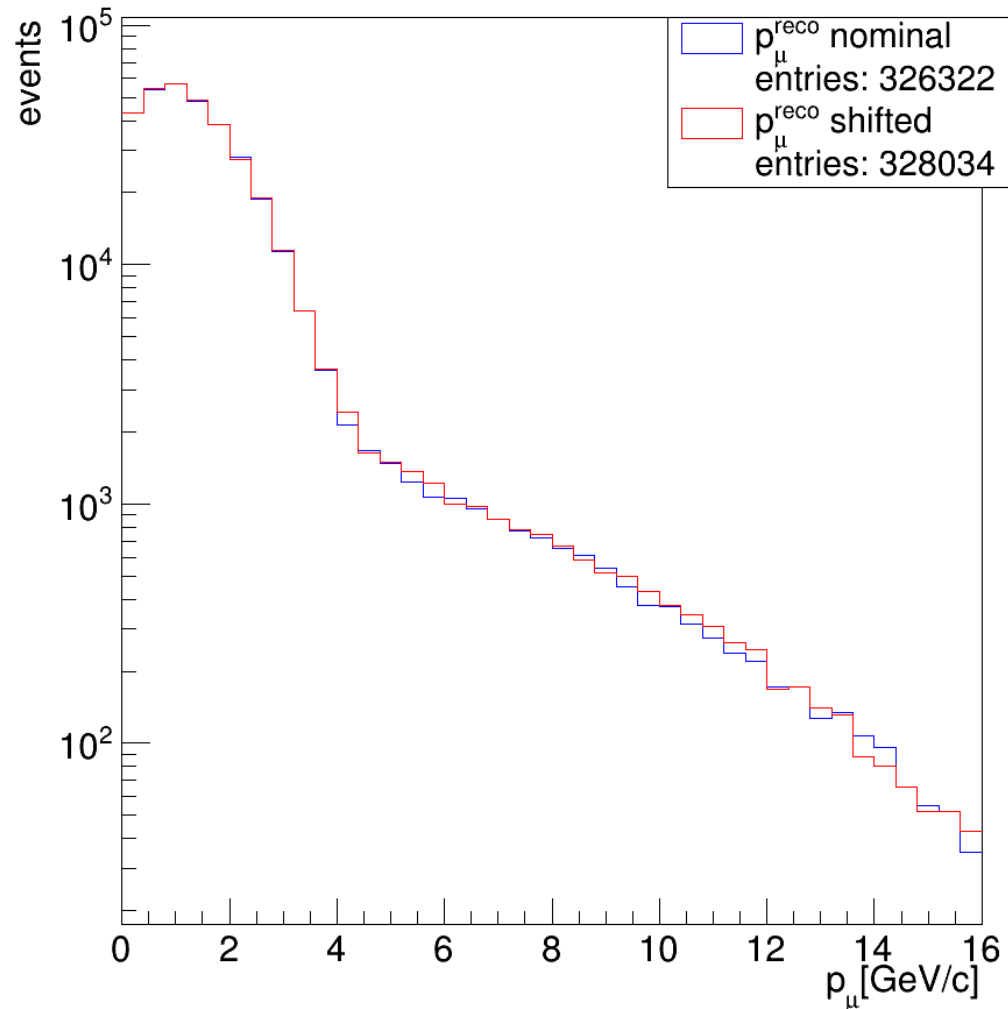


- We applied T the test to the true Monte Carlo momenta from the nominal and shifted samples.
- A fiducial cut was applied by selecting events whose true interaction position was in the fiducial volume
- This was done in order to gauge what the best possible p-value (i.e. the smallest and most decisive) might be:

$$p_{\text{truth}} = 5.15 \times 10^{-7}; \quad \sigma_{\text{truth}} = 5.02 (\chi^2)$$

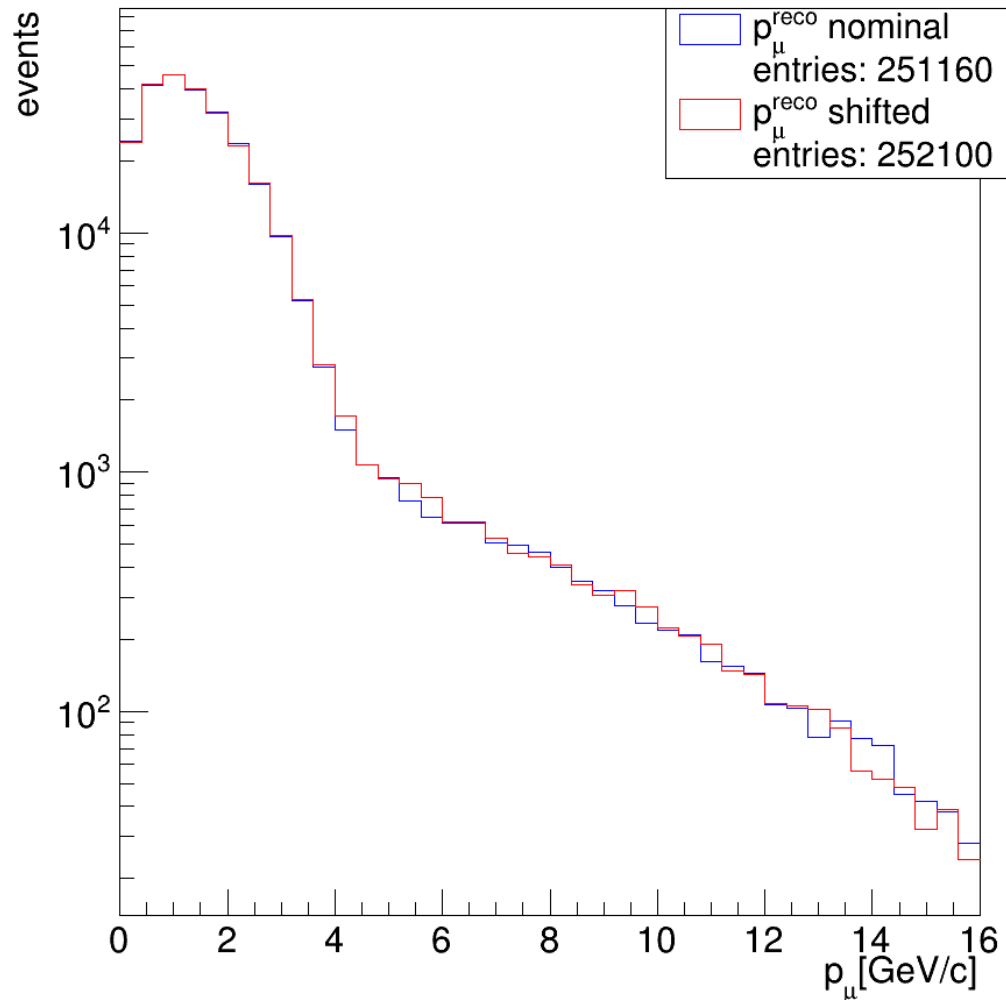
Note: in principle no p-value should be smaller than p_{truth}

RESULTS: FIDUCIAL + RECONSTRUCTION



$$p_{reco} = 1.55 \times 10^{-3};$$
$$\sigma_{reco} = 3.17$$

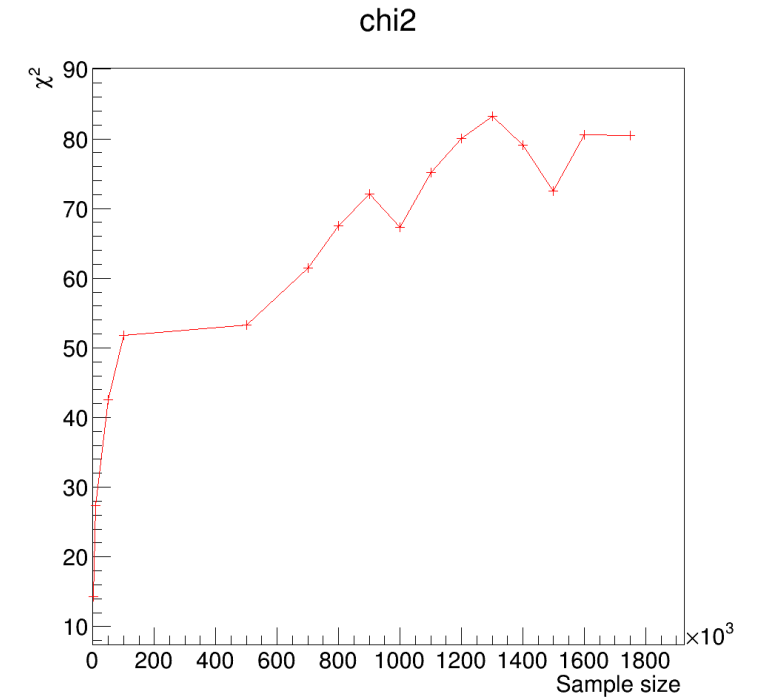
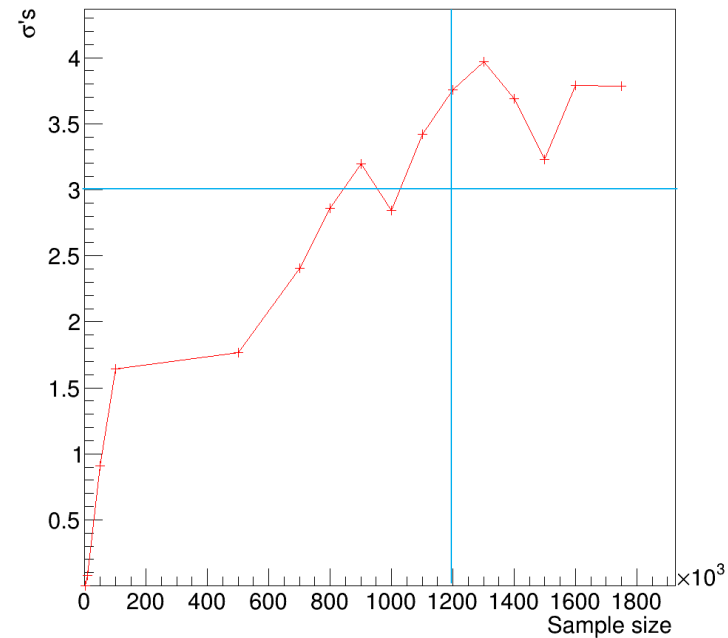
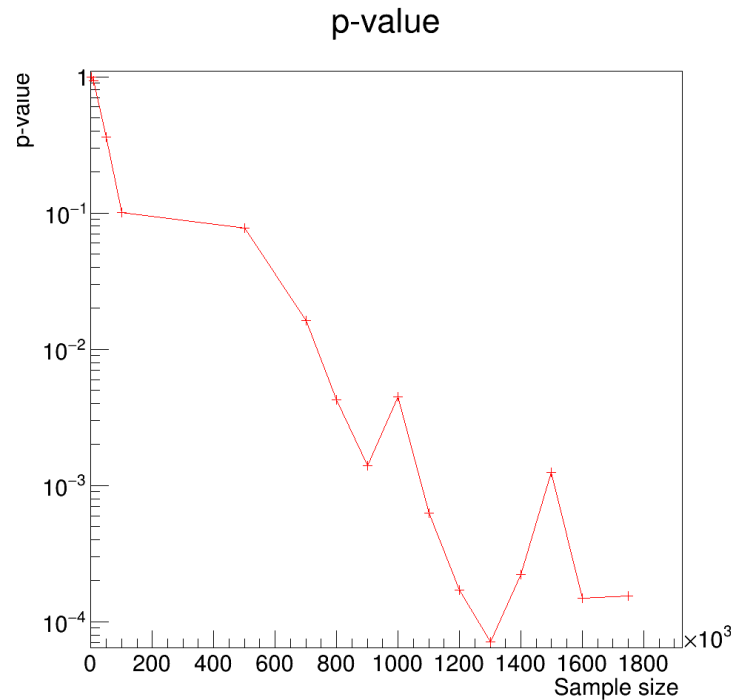
RESULTS: FIDUCIAL + RECONSTRUCTION + QUALITY



We repeat the procedure by applying the quality cut to both samples:

$$p_{reco} = 1.55 \times 10^{-4};$$
$$\sigma_{reco} = 3.78$$

P-VALUE EVOLUTION WITH THE SAMPLE



- As we should expect, as the samples become larger the χ^2 and number of σ 's increase, while the p-value decreases.
- With a sample greater than 1.2 million events, comparable to (even if smaller) that the one expected during a week of data taking, is possible to identify the beam anomaly with a confidence level corresponding to more than 3σ

CONCLUSIONS

- I studied the **beam monitoring capabilities** of the SAND detector, comparing on a weekly basis the muon momentum distribution, from a sample in standard conditions (**nominal**) and one where an anomaly was introduced (**shifted**)
- We observed that in the case of a perfect detector with a perfect reconstruction, (i.e. using the Monte Carlo “truth”), the significance of the difference among the nominal and shifted samples is around **5σ**
- Using reconstructed quantities, fiducial volume and quality selection we reach **$\sim 3.8\sigma$ confidence level**
- The result can be improved: for example include neutrinos with interactions in the STT, or consider the reconstructed position on the xy plane of the interaction vertexes, both in the front calorimeters and the STT.

EXTRA

ECAL: DIGITIZATION AND RECONSTRUCTION

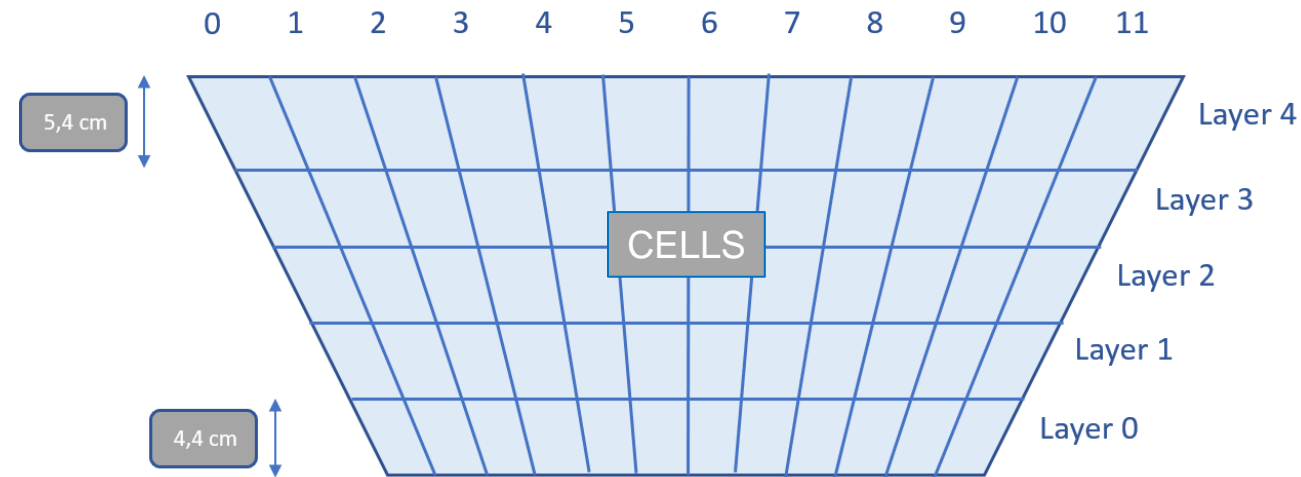
DIGITIZATION

- Calorimeter read-out is segmented in **cells** ($\sim 4.5 \times 4.5 \text{ cm}^2$) and hits are assigned to them
- Hits are converted into number of photo-electrons and their arrival time at PMT on each side

$$N_{p.e.} = 25 \times E_A \times dE$$

$$t_{p.e.} = t_{part} + t_{decay} + d \cdot u_{ph} + Gauss(1ns)$$

- For each PMT:
 - **ADC**: signal is proportional to $N_{p.e.}$
 - **TDC**: t_{TDC} is evaluated as constant fraction (15%) discriminator



RECONSTRUCTION

- The particle arrival time t and its coordinate x along the fiber direction are evaluated as:

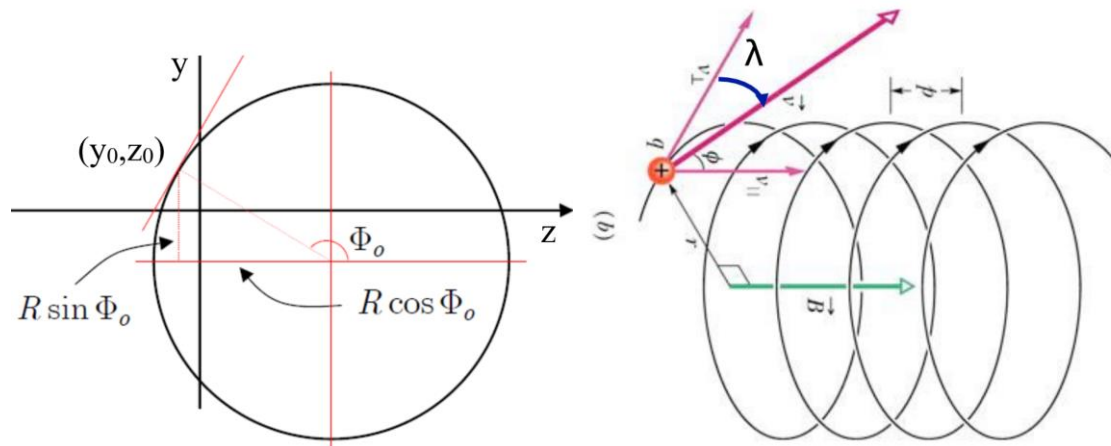
$$x = \frac{t_{TDC1} - t_{TDC2}}{2u_{p.e.}} + x_{cell};$$

$$t = (t_{TDC1} + t_{TDC2} - u_{p.e.} \times L);$$

STT: DIGITIZATION AND RECONSTRUCTION

DIGITIZATION

- **Straw tubes** provide:
 - Transversal position and time of the passage of the particles
 - Their energy deposit



RECONSTRUCTION

- Elicoidal motion due to the magnetic field is split into:
 - **Circular**: fit on the yz plane perpendicular to the magnetic field to obtain radius R and angle Φ_0
 - **Linear**: fit on $x\rho$ plane to find dip angle λ :

$$x = x_0 + \rho \tan \lambda$$

- Transverse momentum:

$$p_T [\text{GeV}] = 0.3 B [\text{T}] R [\text{m}]$$

- Momentum components:

$$p_x = p_T \tan \lambda$$

$$p_y = p_T \cos \Phi_0$$

$$p_z = p_T \sin \Phi_0$$