Elastic and Inelastic Scattering at Residual Gas in IOTA

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<u>Measurements</u>

- At small intensity
 - Beam lifetime is set by gas scattering (no Touschek)
 - Vertical emittance is set by
 - Elastic gas scattering (Coulomb, Rutherford) and
 - by quantum nature of SR due to coupling of vertical and horizontal motion
- Comparison of beam lifetime and emittances showed large discrepancy if only Coulomb scattering is accounted ⇒Bremsstrahlung has to be added

<u>Analysis</u>

Usage standard formulas showed that bremsstrahlung can add only 5-10% and does not allow to match ends Inelastic (Bremsstrahlung) Scattering Lifetime

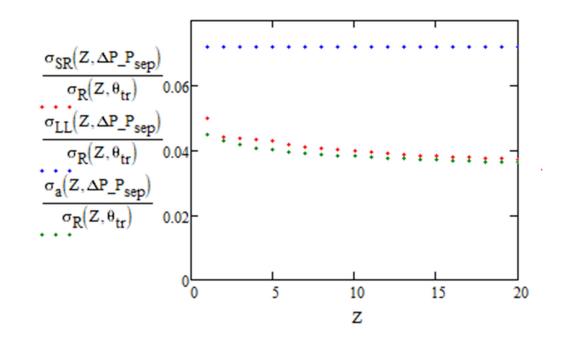


$$\frac{1}{\tau_{brem}(h)} = 2.42 \times 10^{-3} L\left(\mathcal{S}_{acc}^{E}\right) \ln\left(\frac{183}{\sqrt[3]{Z}}\right) \frac{P[nTorr]}{T[K]} \sum_{i}^{n} Z_{i}^{2} N_{i} f_{i}$$

Where, $L\left(\delta_{acc}^{E}\right) = \frac{4}{3}\left(\ln\frac{1}{\delta_{acc}^{E}} - \frac{5}{8}\right)$

and $\delta_{acc}^{\mathcal{E}}$ is energy acceptance (%) of the ring. The Bremsstrahlung lifetime has similar dependence on residual gas properties (Z and N), but a weak dependence on relative energy acceptance.

Scattering at atomic electrons is also too weak to help



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<u>Analysis (2)</u>

- The problem is that the standard formula is not applicable to IOTA
 - It looks at probability to generate a single photon with given energy
- While we need a probability to lose energy above certain value
- Classical estimate allows us to match ends
 - Its usage is justified by small energy of photons: ω < m_e

$$\sigma_{cl}(Z,\theta_{s}) \coloneqq \pi \cdot r_{e}^{2} \cdot \left(\frac{\pi}{4} \cdot \frac{\gamma}{\theta_{s}}\right)^{\frac{2}{3}} \cdot \left(\frac{4}{Z+Z^{\frac{3}{3}}}\right)$$

