

Boosted $hh \rightarrow bbbb$: a New Topology in Searches for TeV-Scale Resonances at the LHC

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Introduction

Higgs has been discovered with couplings consistent with the Standard Model

Many new physics models predict enhancement of Higgs pair production at high invariant mass

- New resonances

 - KK graviton

- Extended Higgs sectors

 - 2HDM

 - Single Higgs extensions

 - Composite Higgs models

Dominant Higgs decay is $H \rightarrow b\bar{b}$ ($\text{Br} \sim 57\%$)

Motivates searches in $X \rightarrow hh \rightarrow b\bar{b}\bar{b}\bar{b}$

- As yet unexamined search topology

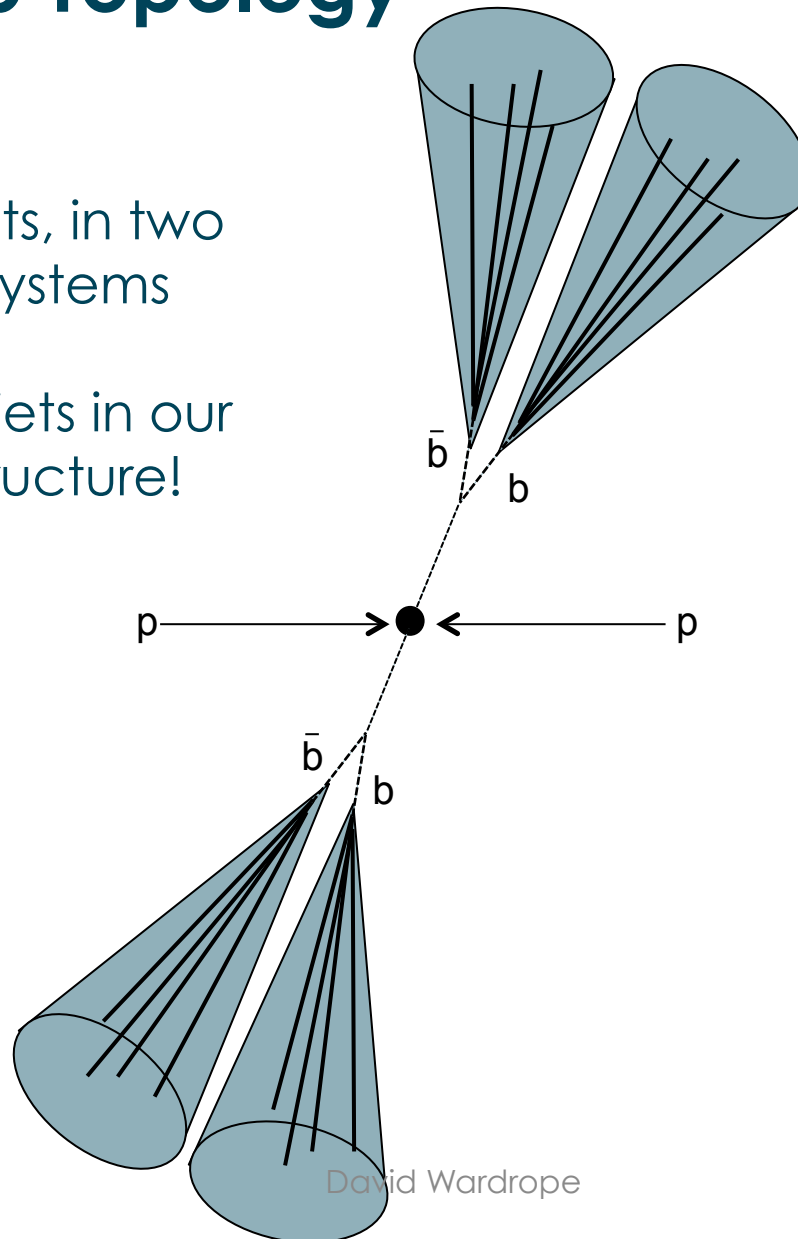
- But is it feasible?

Our short paper assesses the potential: [arXiv:1307.0407](https://arxiv.org/abs/1307.0407)

Boosted 4b Topology

Four b-tagged jets, in two boosted dijet systems

We use resolved jets in our study, no substructure!



Anti- K_T $R=0.4$
b-tagged jet

Boosted 4b Topology

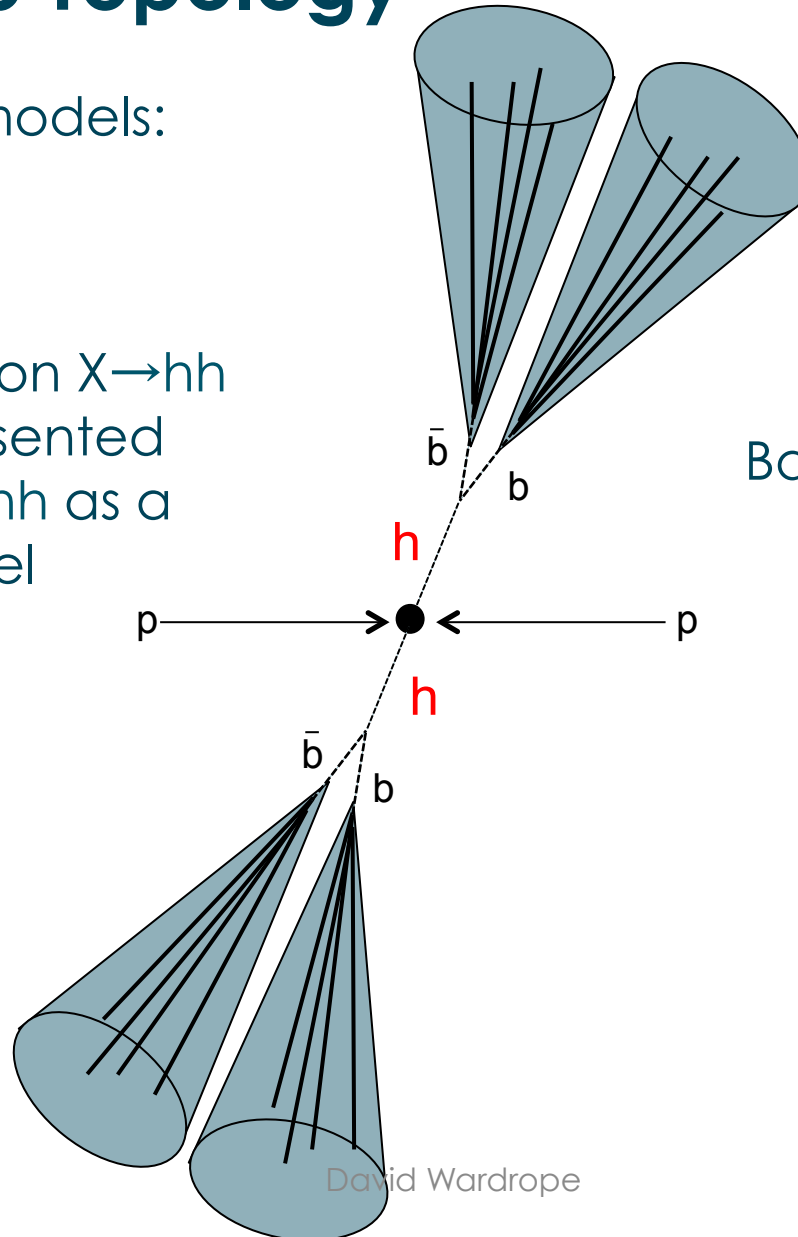
Potential Signal models:

$$G_{kk} \rightarrow hh$$

$$H \rightarrow hh$$

$$VV \rightarrow hh$$

We concentrate on $X \rightarrow hh$ in the studies presented here, using $G_{kk} \rightarrow hh$ as a benchmark model



Anti- K_T $R=0.4$
b-tagged jet

Backgrounds:
QCD multijet
 $t\bar{t}$ bar

Boosted 4b Topology

The topology also allows
for $X \rightarrow ZZ$

Potential signal models:

$$G_{kk} \rightarrow ZZ$$

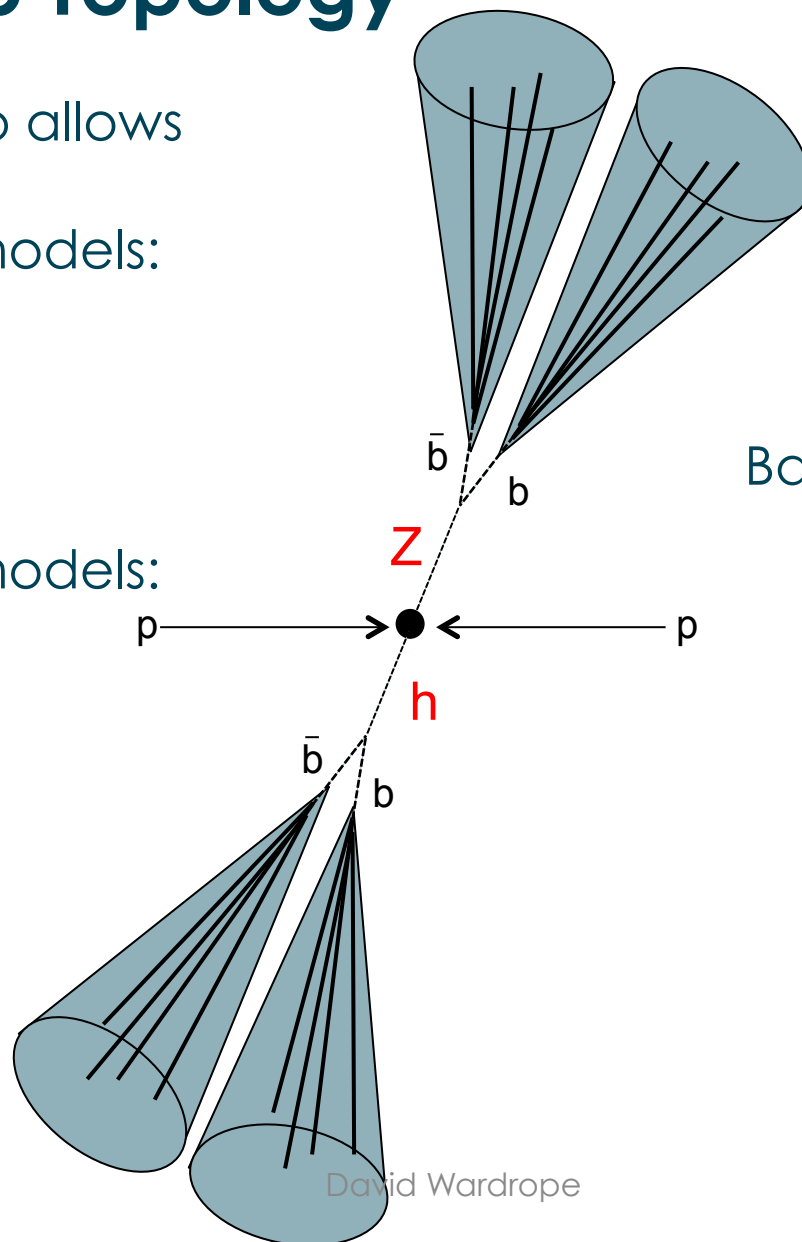
$$H \rightarrow ZZ$$

$$VV \rightarrow ZZ$$

Or even $X \rightarrow hZ$

Potential signal models:

$$A \rightarrow hZ$$



Anti- K_T $R=0.4$
b-tagged jet

Backgrounds:

QCD multijet

$t\bar{t}$ bar

$Z \rightarrow b\bar{b}$ + jets

Diboson ZZ

Pros and Cons of Boosted 4b Channel

Advantages

Double dijet topology is distinctive – reduces background

Resonances with large Higgs couplings benefit from $\text{Br}(h \rightarrow b\bar{b})$

This benefit is squared!

Resonances with large Z couplings benefit from $\text{Br}(Z \rightarrow b\bar{b})$

$\text{BR}(ZZ \rightarrow b\bar{b}b\bar{b}) / \text{BR}(ZZ \rightarrow l\bar{l}l\bar{l}) \sim 5$ (where $l = e, \mu$)

Multiple high p_T b-jets make efficient triggering possible

multijet triggers at first level

b-jet triggers at higher levels

Boost removes ambiguity in assigning jets to parent

Disadvantages

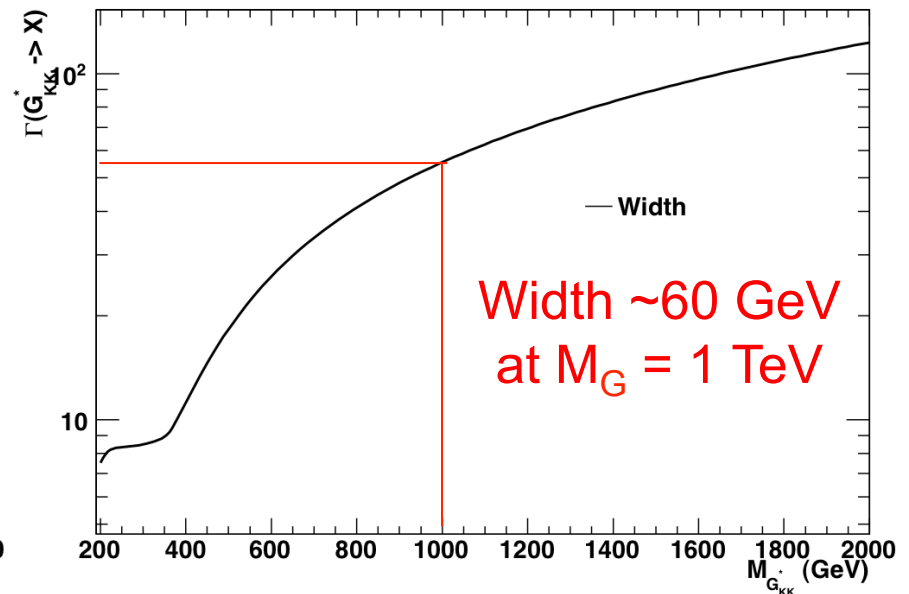
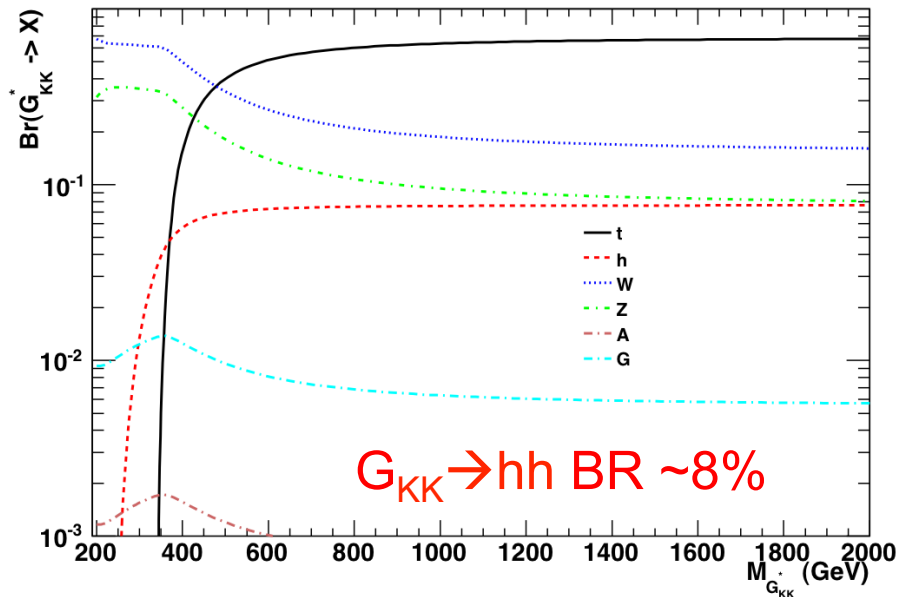
QCD 4-jet production has a massive cross-section

Signal efficiency reduced by ϵ_b^4

We have performed a particle-level study to ascertain whether advantages outweigh the disadvantages

Benchmark Signal Model

Taken from [3] with kind permission



Randall-Sundrum Kaluza-Klein graviton (G_{KK}) in Agashe-Davoudiasl-Perez-Soni (ADPS) model with $k/M_{Pl} = 1.0$ [1,2]

G_{KK} production/decay to light fermions/photons highly suppressed.

Significant $G_{KK} \rightarrow hh$ branching ratio.

Generated using Madgraph + Pythia 8.17 with CTEQ6L1, using the CP³-Origins Madgraph implementation [3] of the ADPS model.

Only the $G_{KK} \rightarrow hh \rightarrow bbbb$ decay mode with $m_h = 125$ GeV.

Particle-Level Study

Cluster final state particles into anti- K_T $R=0.4$ jets.
No detector smearing applied.

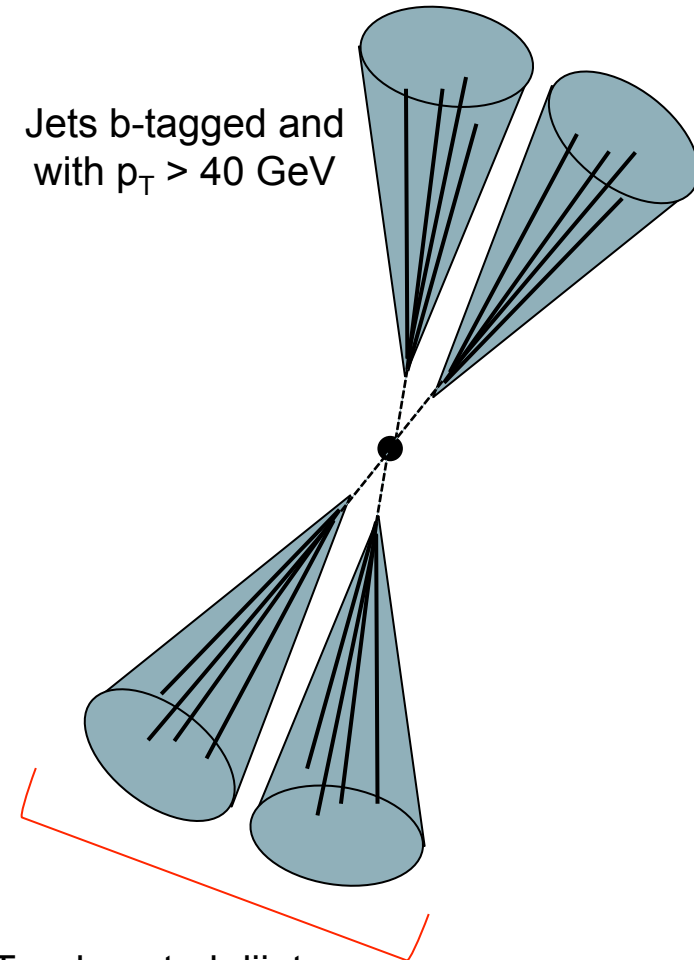
Require at least 4 b-tagged jets with $p_T > 40$ GeV and $|\eta| < 2.5$

Require 2 dijets with $p_T^{\text{dijet}} > 200$ GeV and $\Delta R_{\text{dijet}} < 1.2$

Require m_{dijet} consistent with m_h
 $100 < M_{\text{dijet}} < 130$ GeV

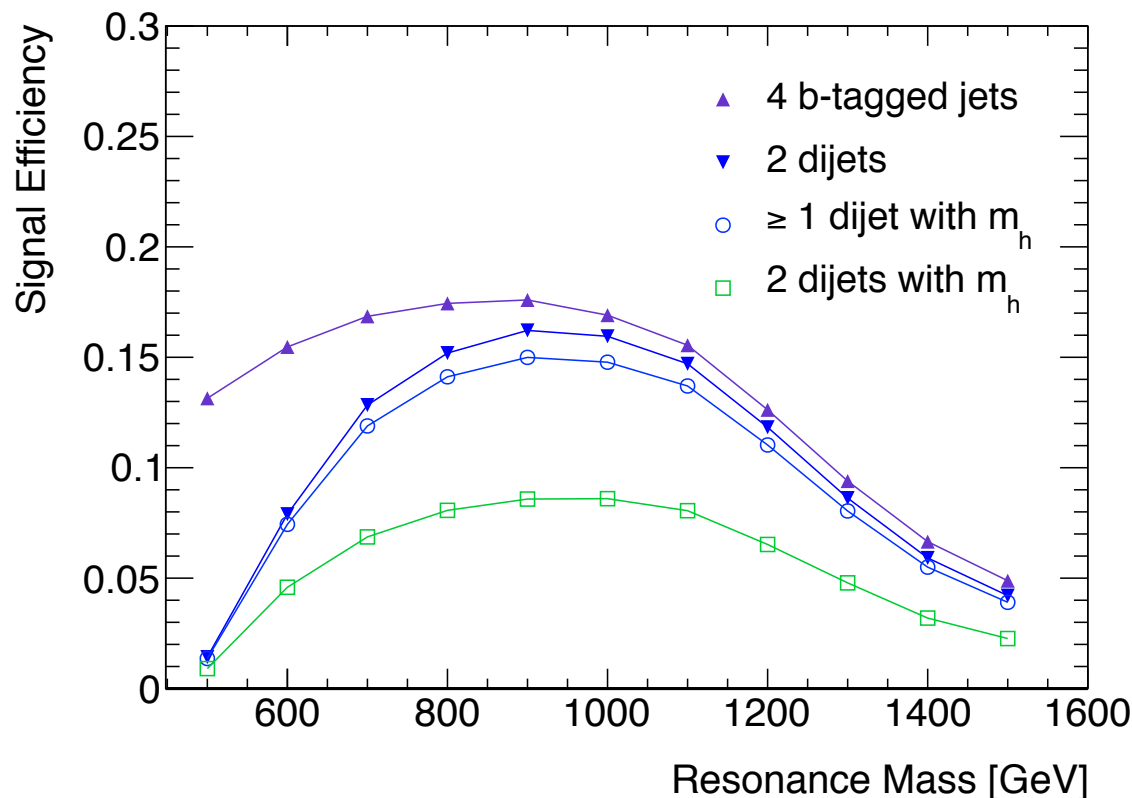
“Truth” b-tagging using simple parameterisation of ATLAS/CMS b-tagging performance:
B-jet 70%
C-jet 20%
Light 1%

Jets b-tagged and with $p_T > 40$ GeV



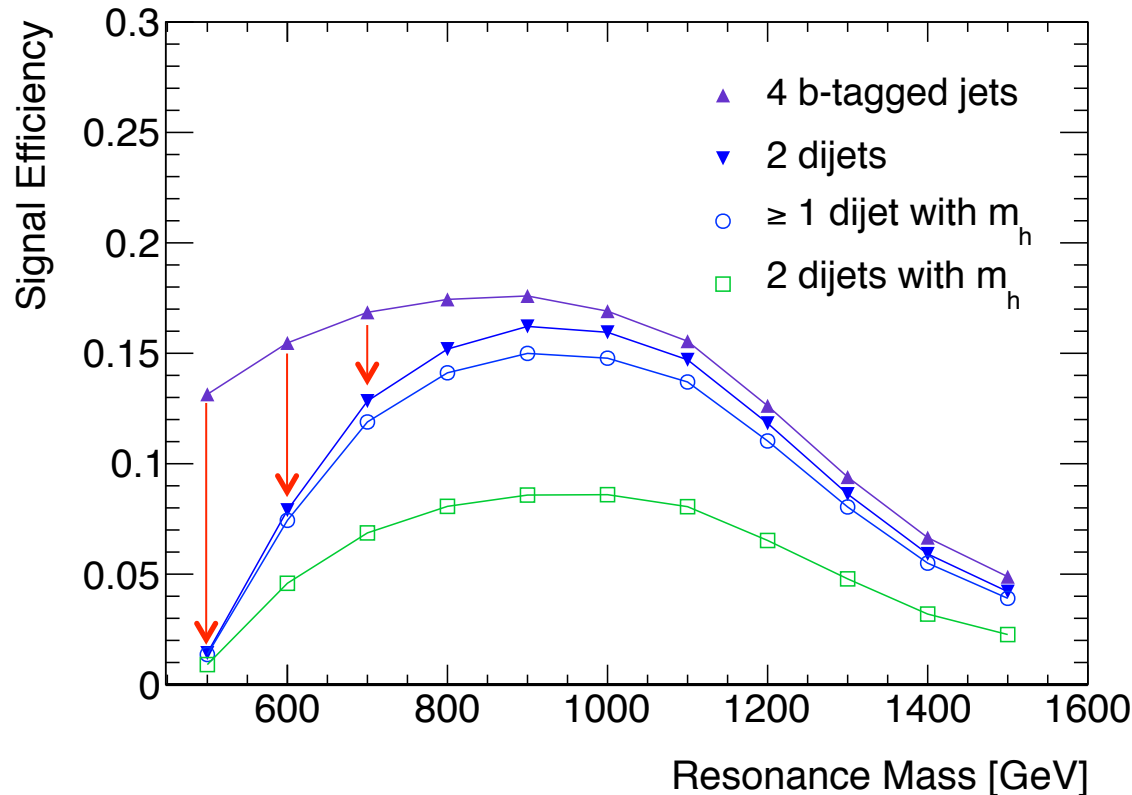
Two boosted dijet systems $p_T > 200$ GeV and $m_{\text{dijet}} \sim m_h$

Signal Efficiency



Efficiency peaks for masses between 0.7 and 1.2 TeV
 b-tagging efficiency probably better than our assumption
 Most jets are central and have p_T , where ATLAS and CMS b-tagging perform best

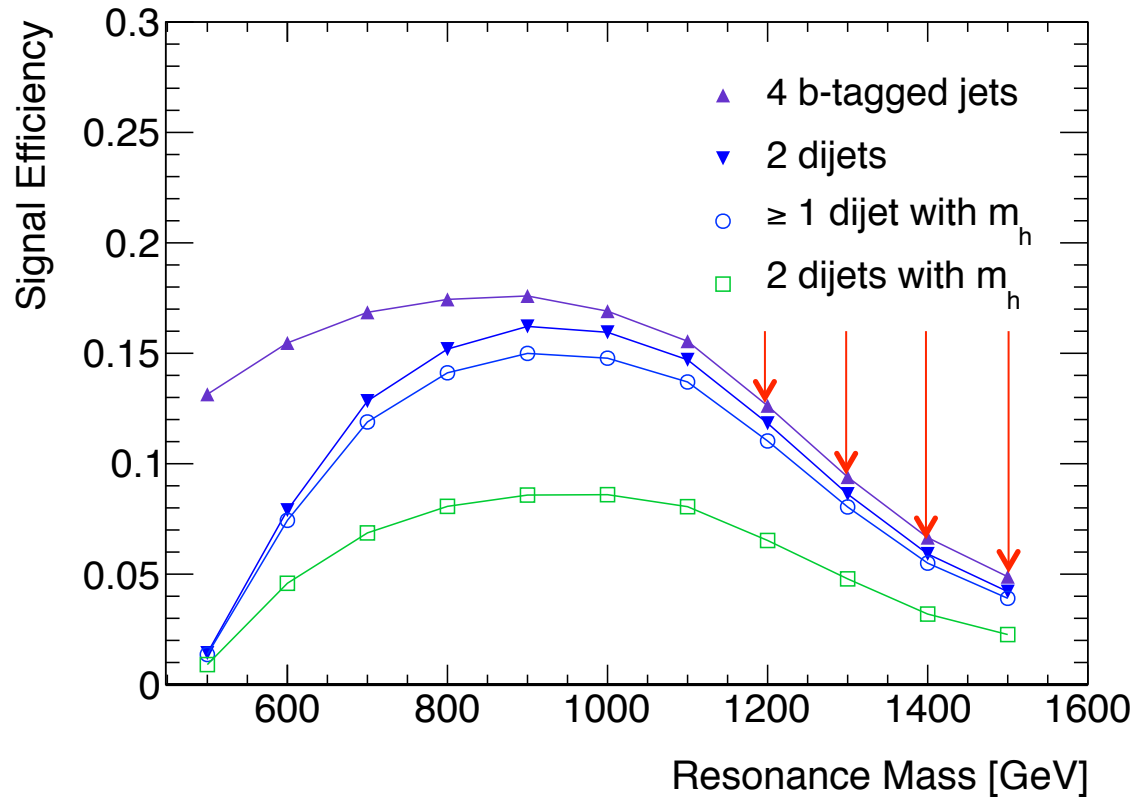
Signal Efficiency



Efficiency loss at low mass because of dijet p_T and ΔR requirements (and jet p_T)

requirements could be optimised to increase efficiency

Signal Efficiency



Efficiency loss at high mass because of jet merging

substructure techniques could regain efficiency

Background for 20 fb⁻¹ at $\sqrt{s} = 8$ TeV

Requirement	G _{KK} (M = 800 GeV)	QCD	$t\bar{t}$
4 b -tagged jets	126	19700	3590
2 dijets	109	414	151
≥ 1 dijet with m_h	102	183	89
2 dijets with m_h	58	28^{+20}_{-11}	21 ± 3

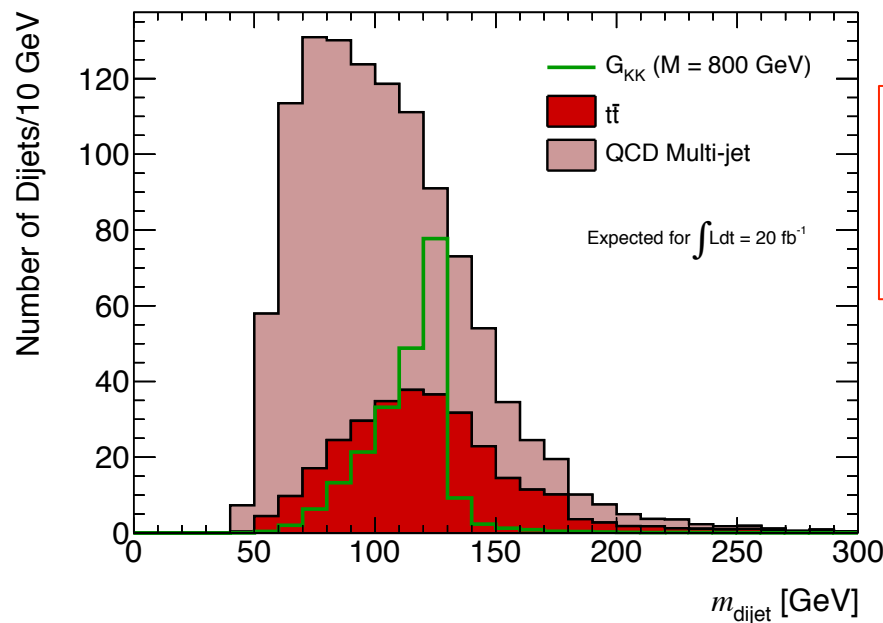
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Backgrounds are dramatically suppressed when we require the *b*-tagged jets form two boosted dijets
 ~50× for QCD multijet and ~25× for $t\bar{t}$

Background for 20 fb⁻¹ at $\sqrt{s} = 8$ TeV

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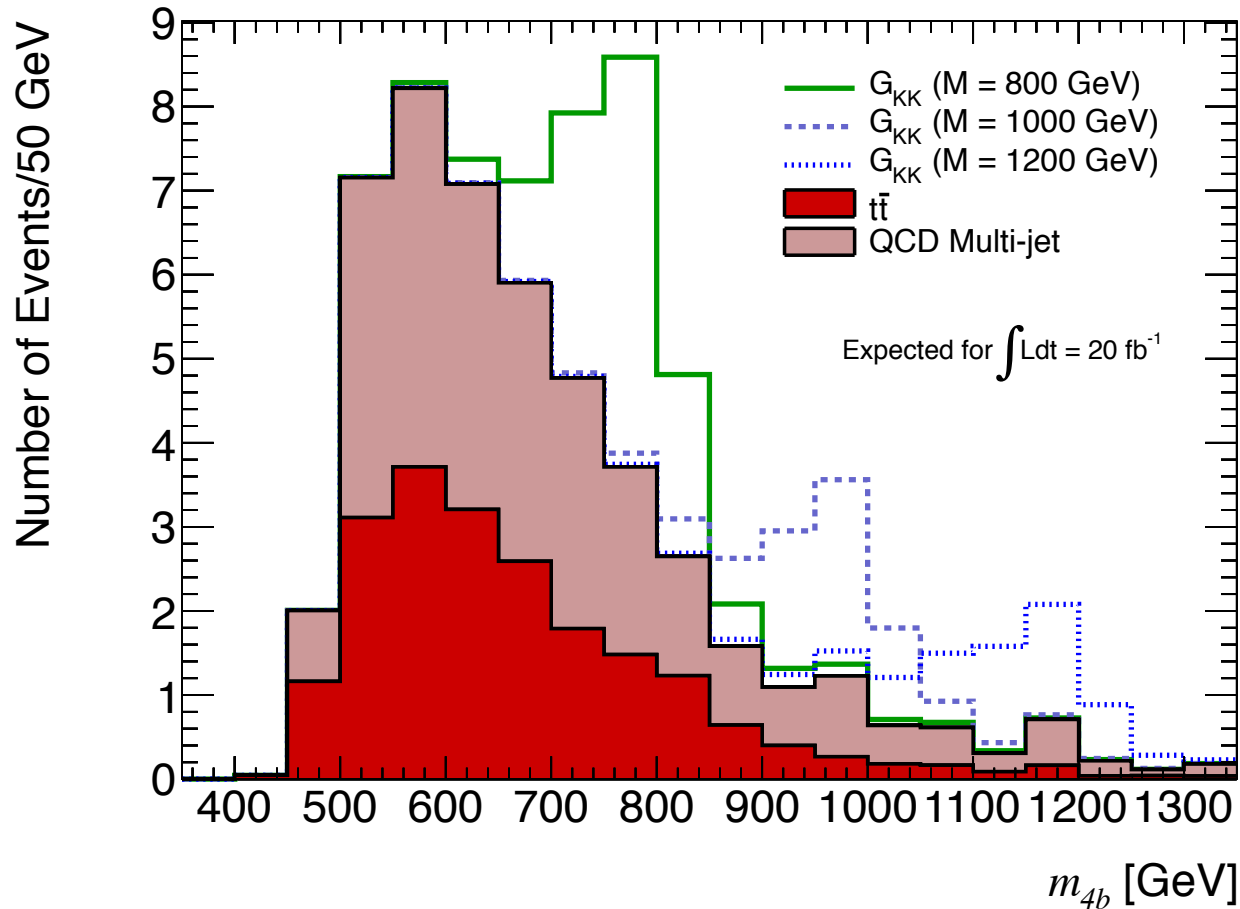
Dijet mass distribution after requirement of two dijets. S/B already 20%!

Background for 20 fb⁻¹ at $\sqrt{s} = 8$ TeV

Requirement	G _{KK} (M = 800 GeV)	QCD	$t\bar{t}$
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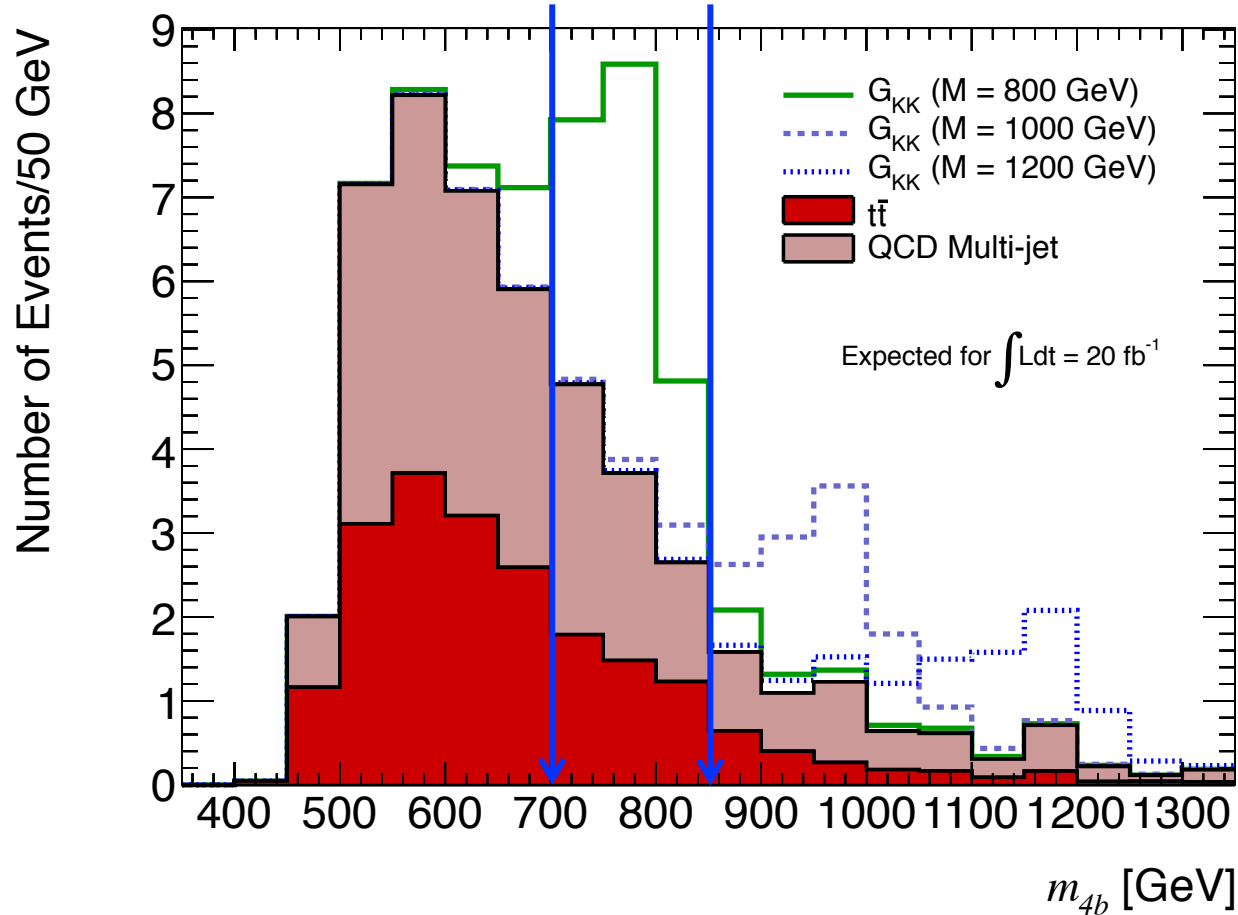
After requiring m_h , very little background remains
 $S/B \approx 1$, despite low signal cross-section $\sigma = 36$ fb
 QCD and $t\bar{t}$ backgrounds of similar size

4b Invariant Mass After Selection



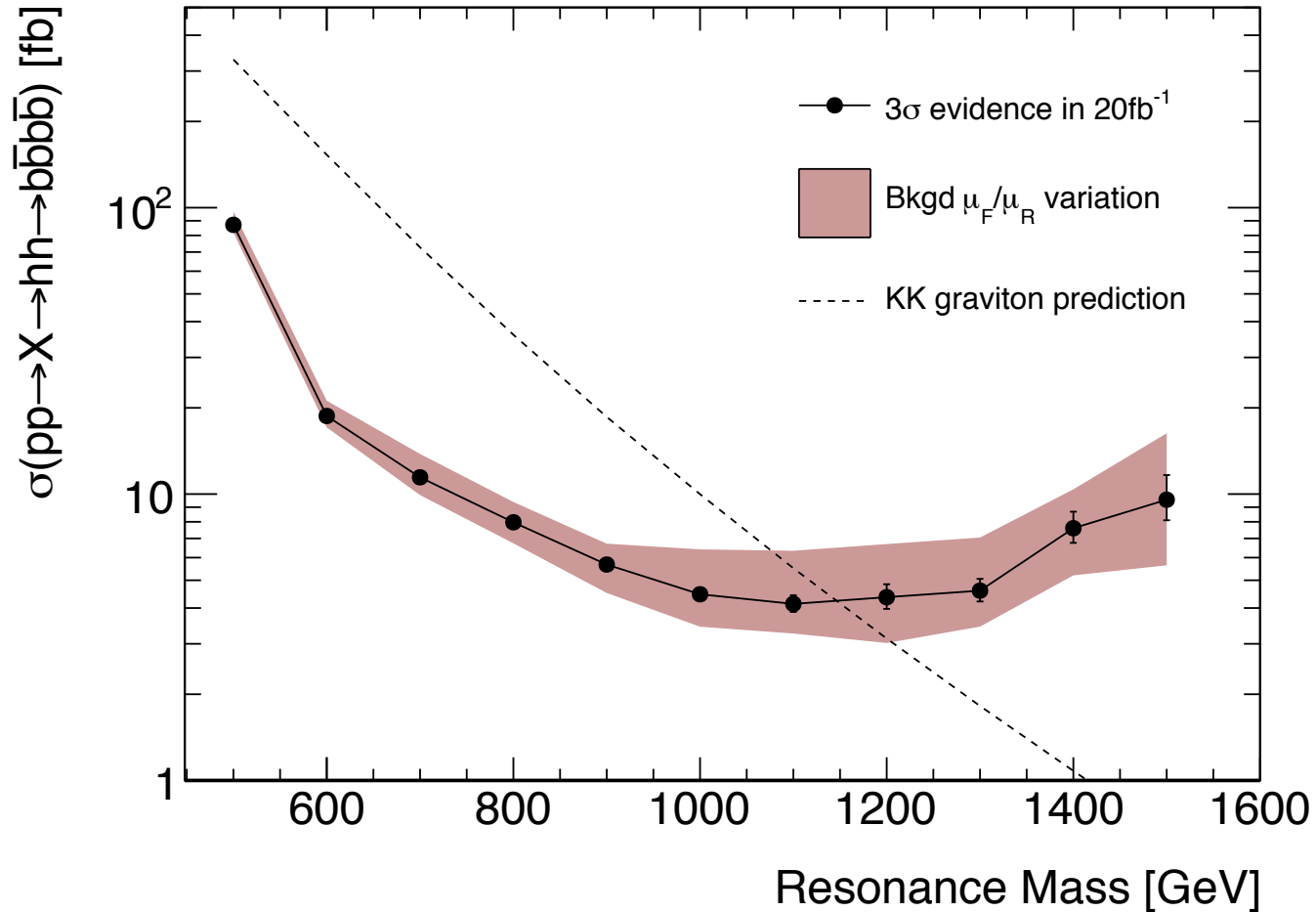
Signals are clearly visible over background

Sensitivity Estimation

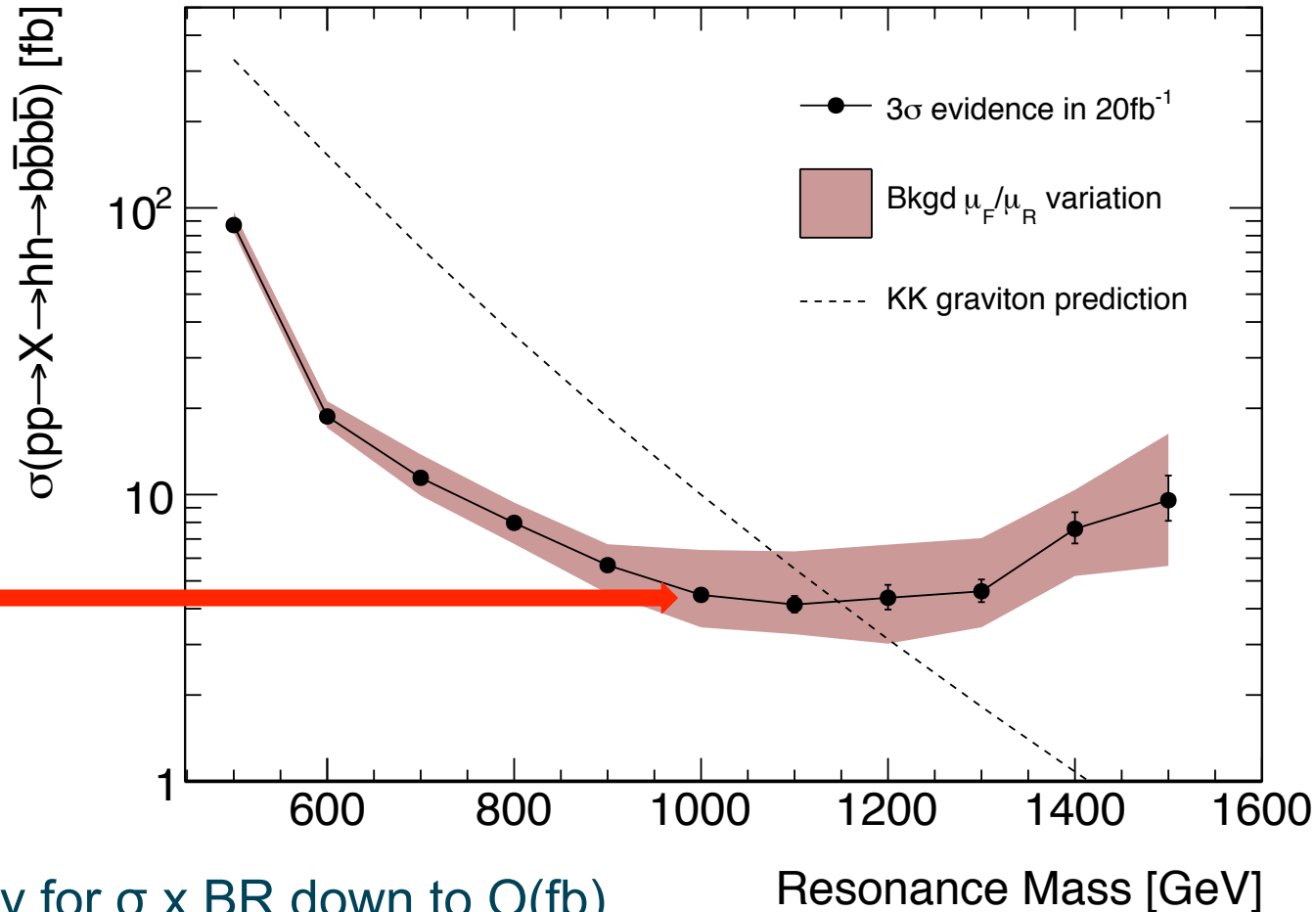


Count background in $[-100, +50]$ GeV window around $m_{G_{kk}}$
 Calculate signal cross-section needed for $s/\sqrt{b} = 3$ in 20fb^{-1}

Cross-section for 3σ Sensitivity

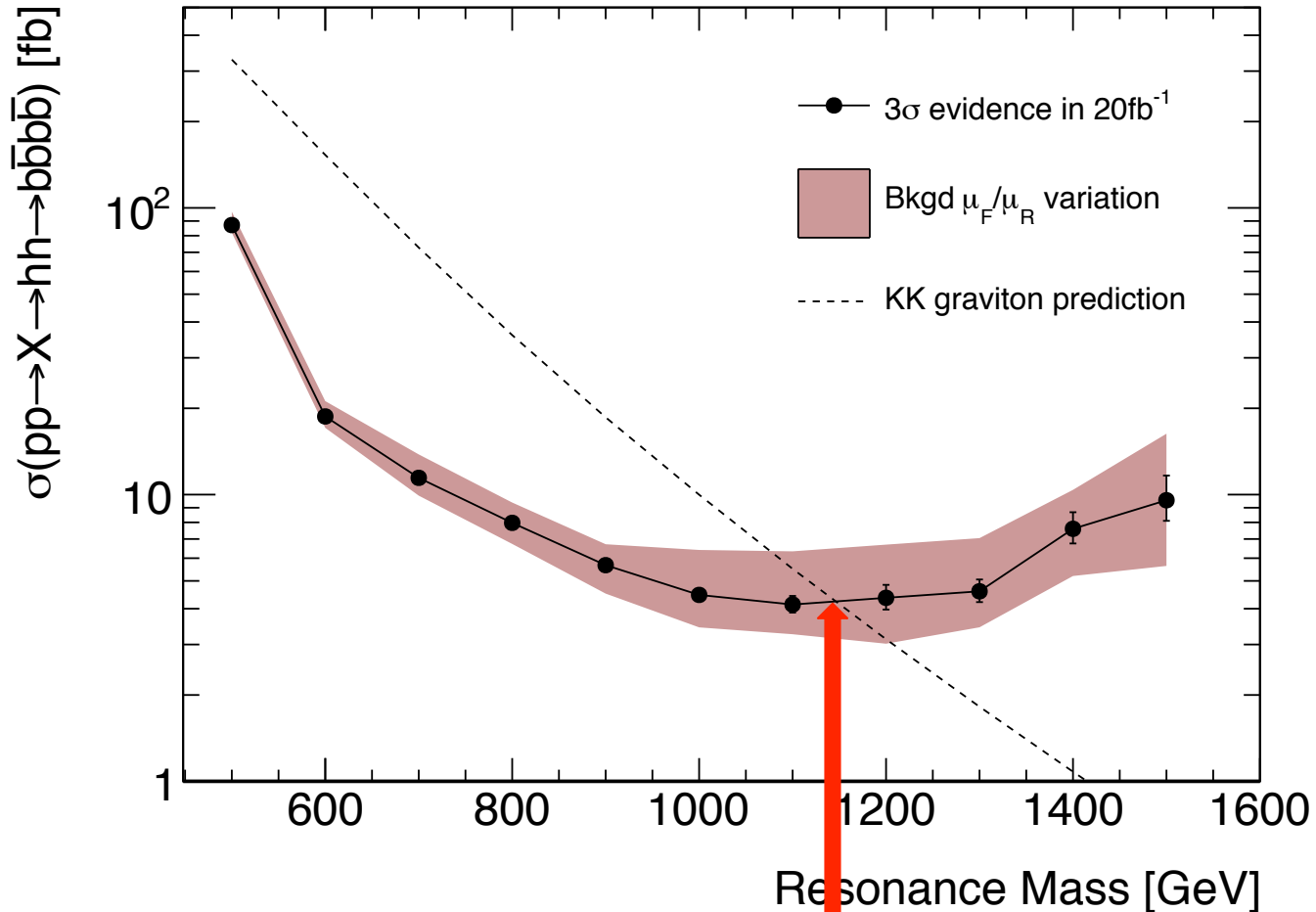


Cross-section for 3σ Sensitivity



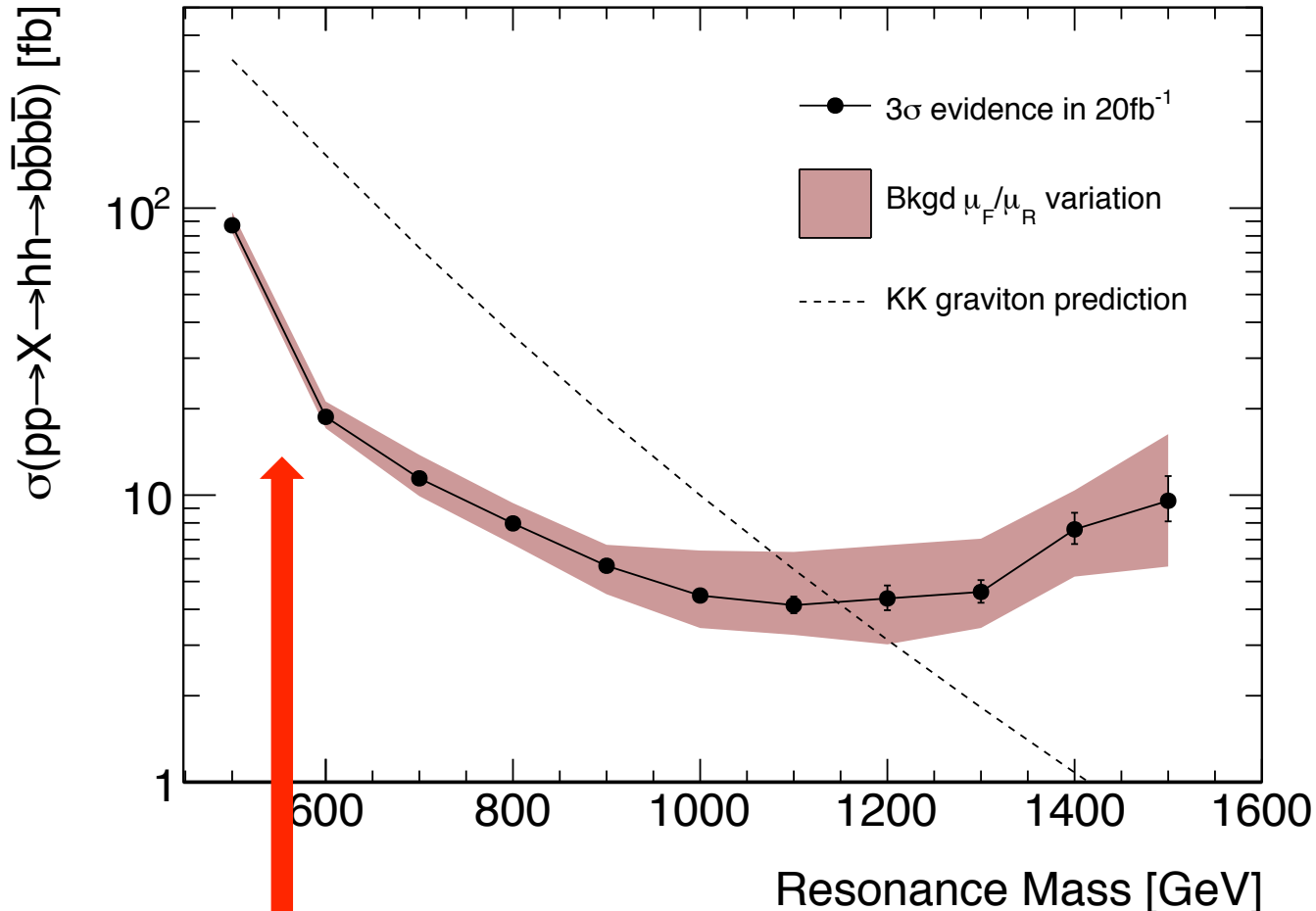
3σ sensitivity for $\sigma \times \text{BR}$ down to $O(\text{fb})$
for TeV resonances with existing LHC
dataset of $\int L dt = 20\text{fb}^{-1}$!

Cross-section for 3σ Sensitivity



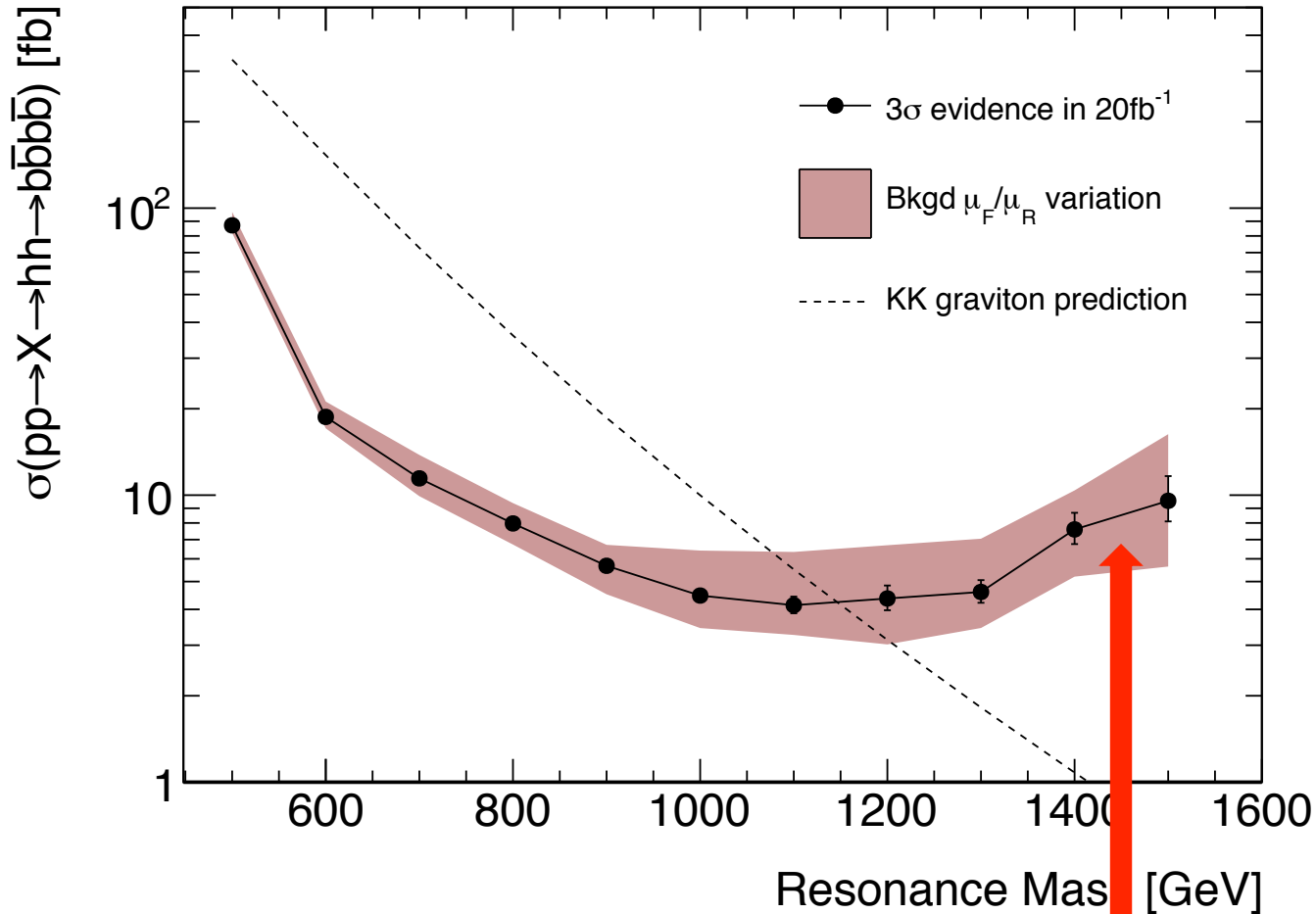
Have discovery sensitivity to KK gravitons up to 1.15 TeV

Cross-section for 3σ Sensitivity



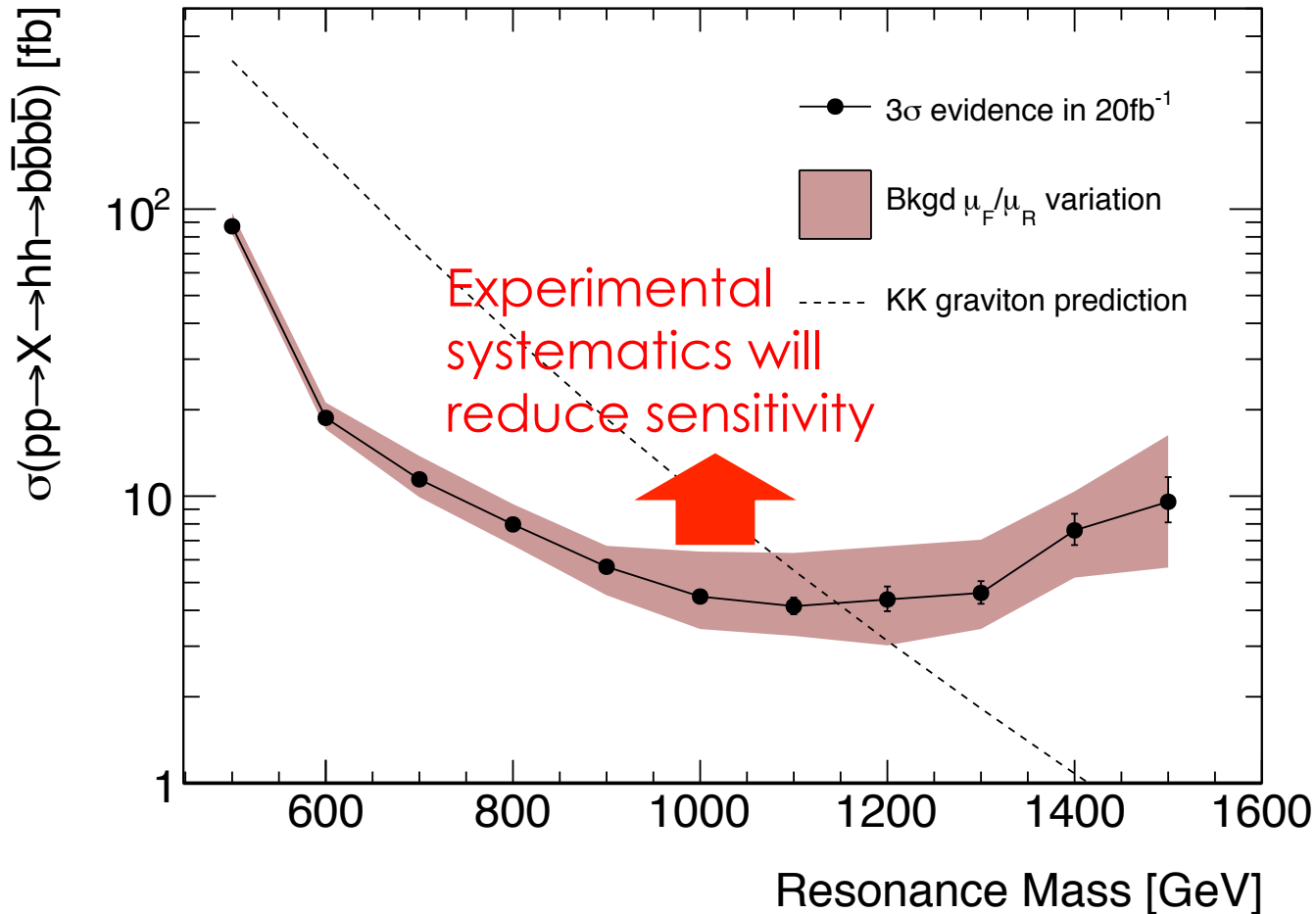
Sensitivity worse at lower mass due to lower signal efficiency and higher background

Cross-section for 3σ Sensitivity

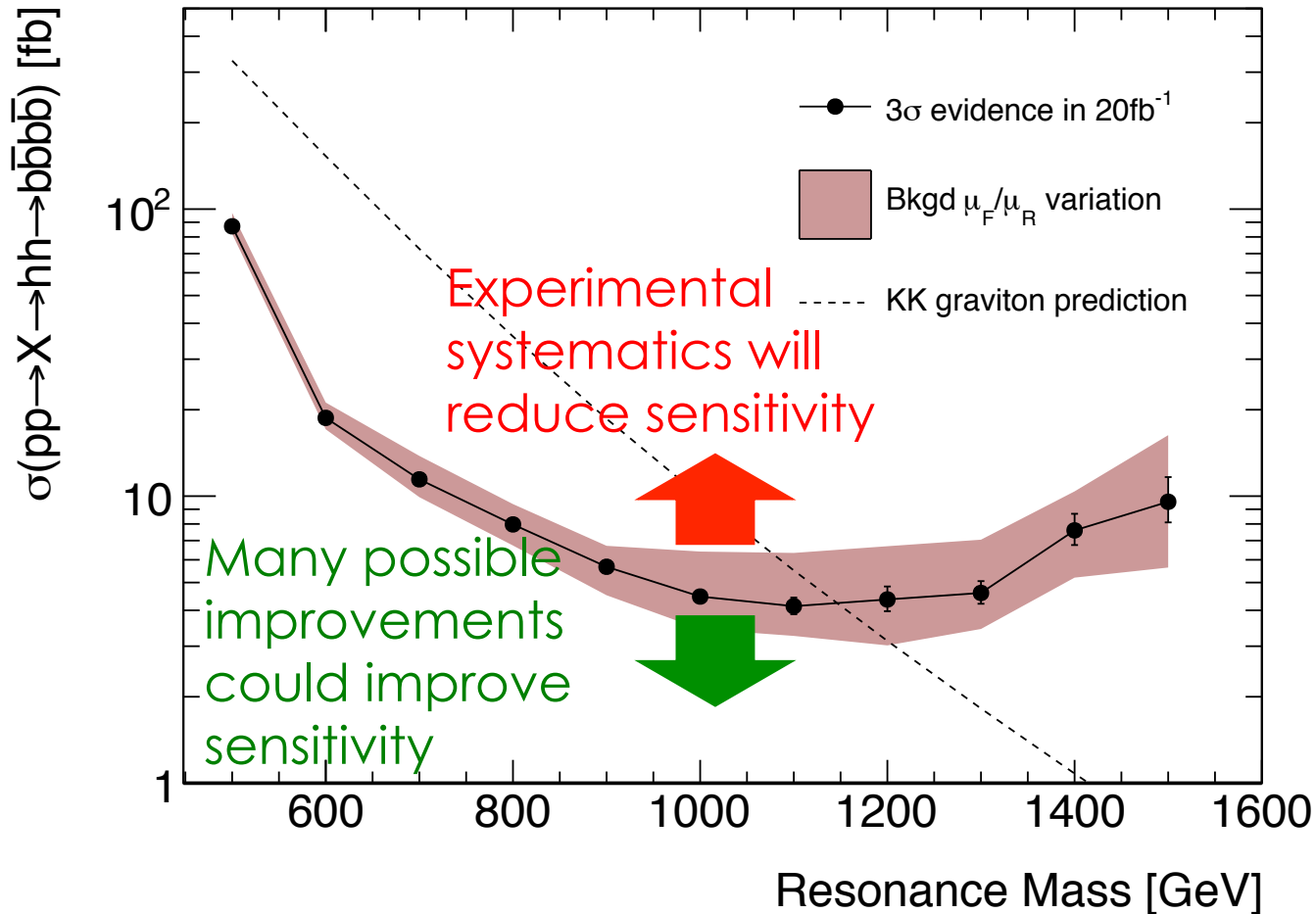


Sensitivity decreases at high mass due to lower signal efficiency
 Counteracted by very low background at these resonance masses

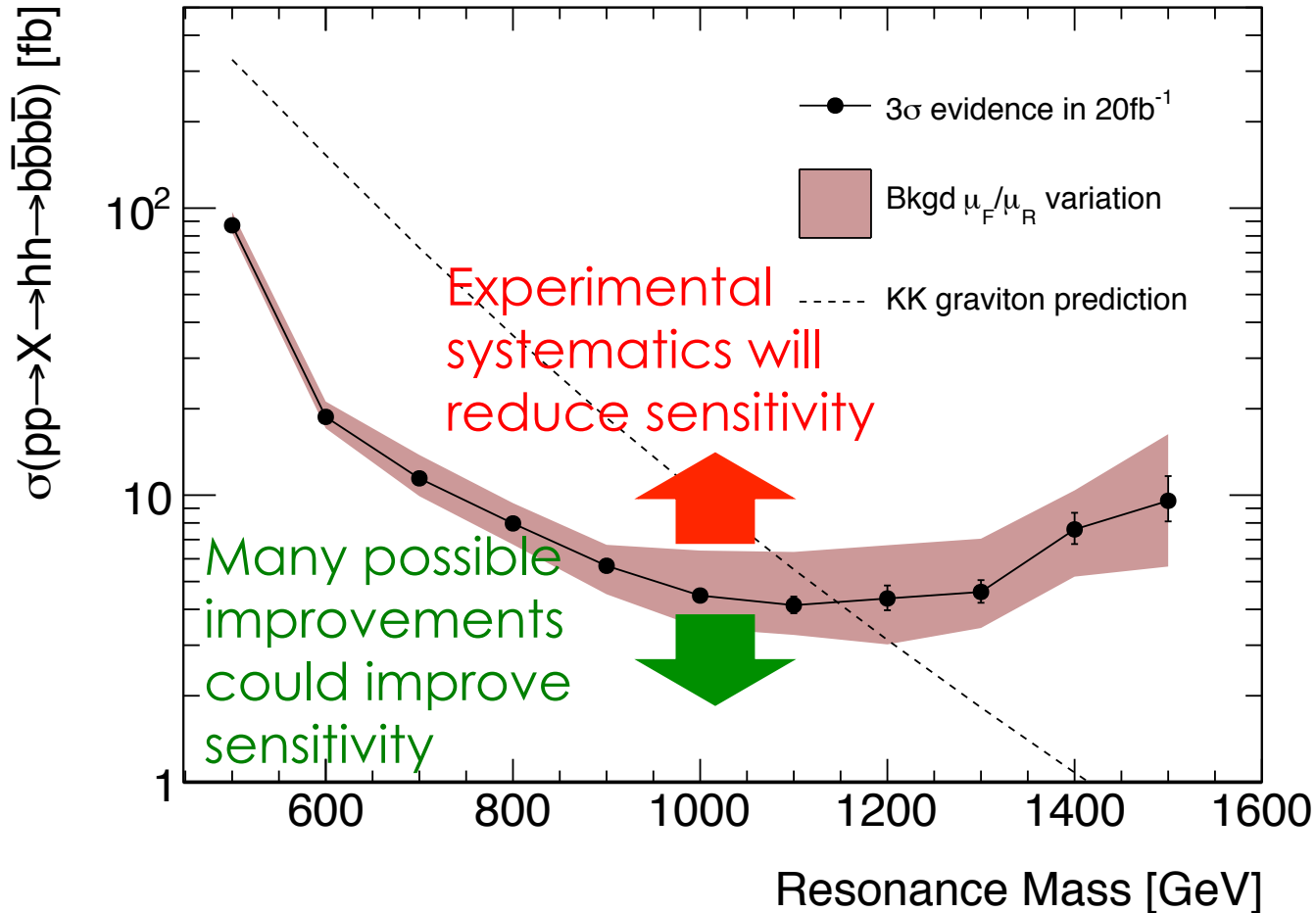
Cross-section for 3σ Sensitivity



Cross-section for 3σ Sensitivity



Cross-section for 3σ Sensitivity



We conclude there is great potential for searches in the boosted $4b$ final state

Potential for Optimisation

Purpose of this simple study is to flag $X \rightarrow hh \rightarrow bbbb$ as a very promising final state for new physics searches

We leave the optimisation for the experiments!

A few suggestions for extracting best possible sensitivity:

Tuning of basic cuts versus mass in the resolved analysis

p_T^{jet} , p_T^{dijet} , ΔR_{dijet} , m_H window

Incorporate large- R jets and substructure into analysis for resonance masses > 1.2 TeV

Additional cuts to reduce $t\bar{t}b\bar{b}$ background

e.g. optimised b -tagging, n_{jets} requirements

Improve m_{4b} resolution using a kinematic fit

take advantage of known m_h

Summary

The boosted bb-bb final state is extremely promising
 Powerful background rejection of the boosted dijet topology enables sensitive searches for $X \rightarrow hh$
 despite final state being fully-hadronic

Should work for $X \rightarrow hZ$ and to complement $X \rightarrow ZZ$

And for ZZ VBS in bb-bb final state

$ZZ \rightarrow llbb$

$hZ \rightarrow \tau \tau bb$

could even extend to VLQ $BB \rightarrow bbbbbb$

Many more we haven't thought of

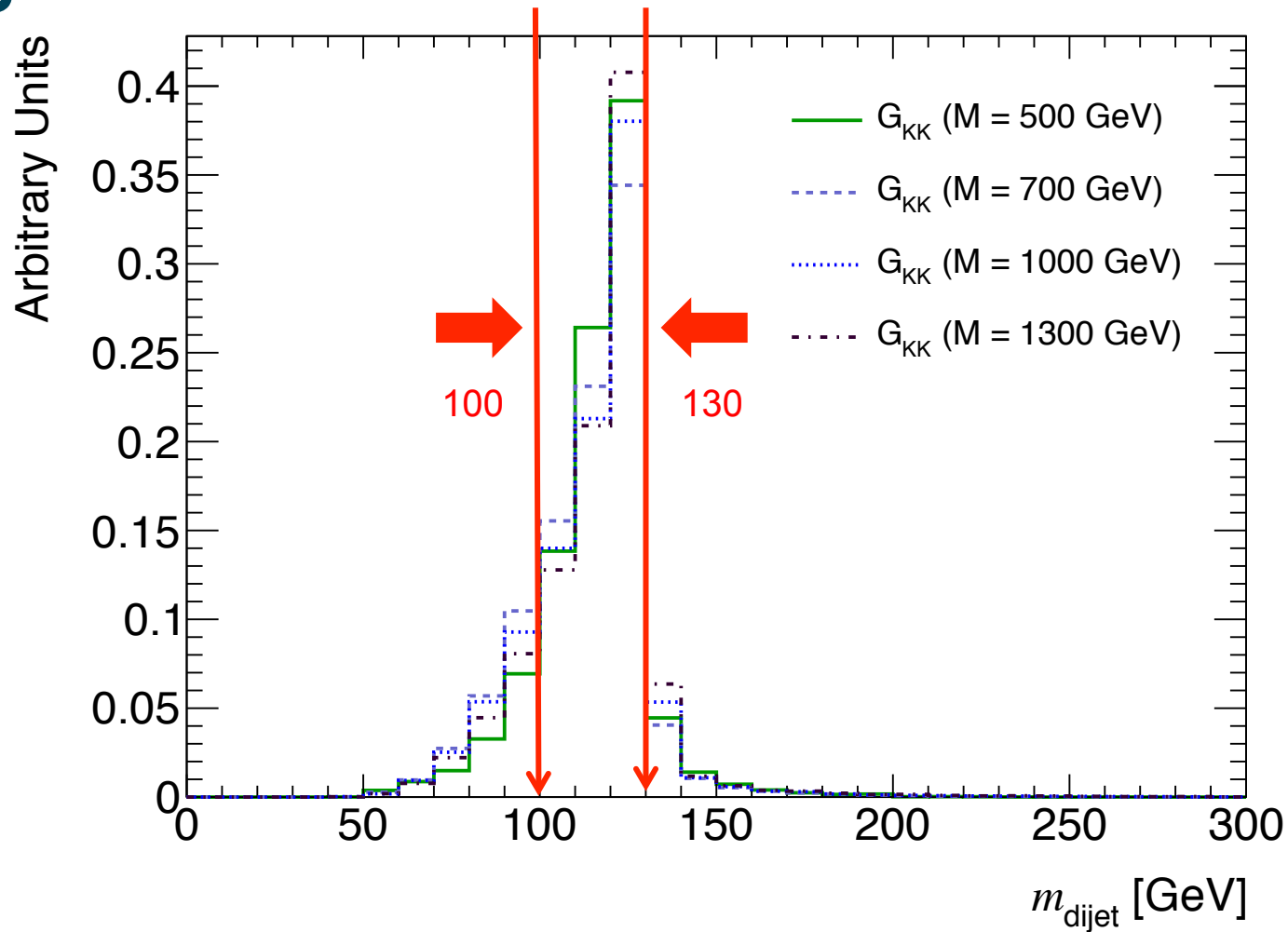
This is uncharted territory!

ADDITIONAL SLIDES

References

- [1] K. Agashe et al Phys.Rev. D76 (2007) 036006.
- [2] L. Fitzpatrick et al, JHEP 2007 (2007) 013.
- [3] <http://cp3-origins.dk/content/uploads/2011/10/kkgrav.pdf>
- [4] N. Greiner et al, Phys. Rev. Lett. 107 (Sep, 2011) 102002.
- [5] G. Bevilacqua et al arXiv:1304.6860 [hep-ph].
- [6] ATLAS Collaboration, ATLAS-CONF-2012-149.
- [7] CMS Collaboration, CMS-PAS-TOP-12-027.
- [8] M. Gouzevitch et al, JHEP 1307 (2013) 148.

Higgs Mass Window

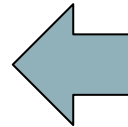
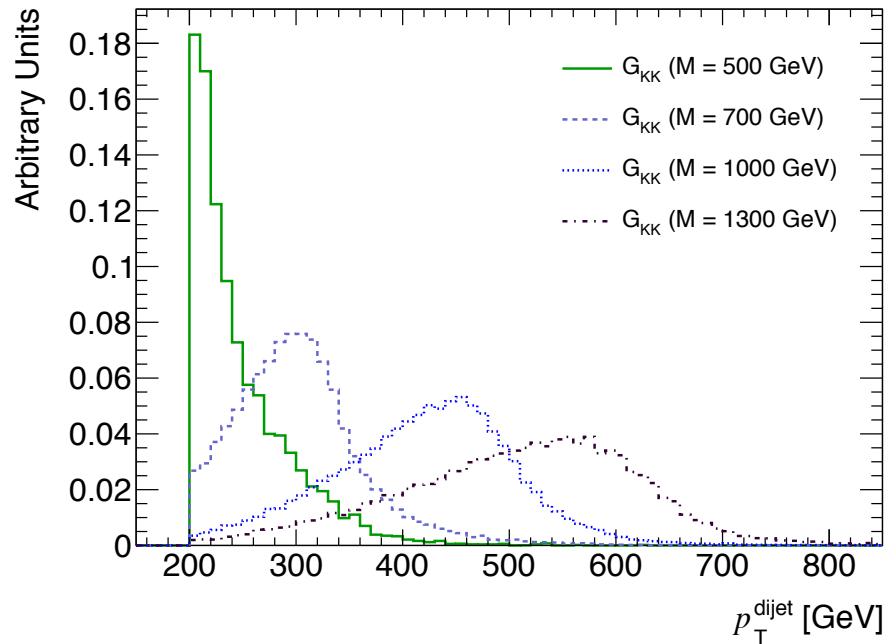


Jets do not include muons or neutrinos, and not corrected for out-of-cone. Asymmetric cut around $m_h = 125$ GeV is appropriate.

Benchmark Signal Model

Graviton Mass [GeV]	$\sigma(pp \rightarrow G_{KK} \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ [fb]	Γ [GeV]
500	329	18.6
700	72.7	33.9
900	18.6	48.6
1100	5.51	62.7
1300	1.82	76.5
1500	0.65	90.0

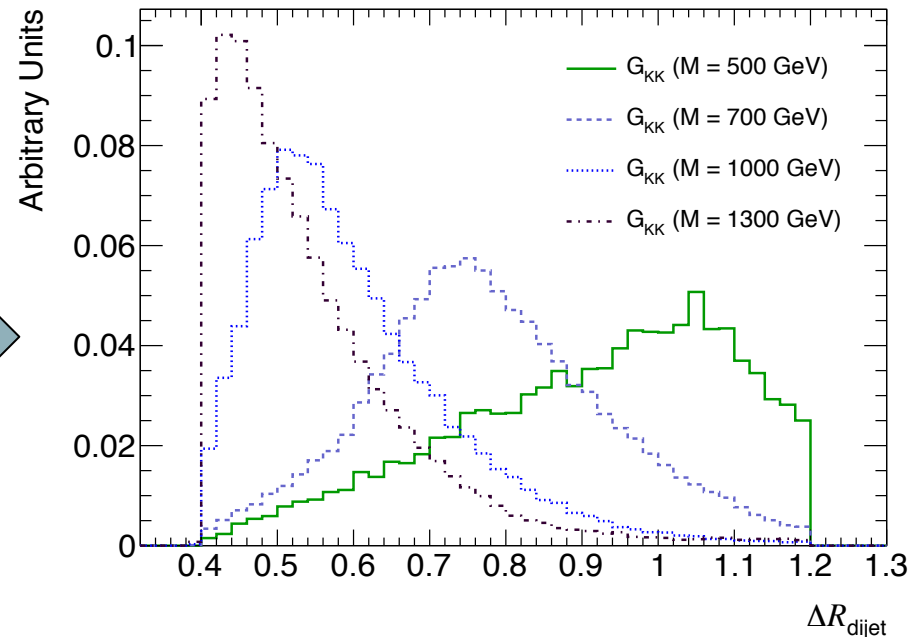
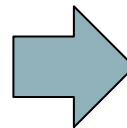
Signal Kinematics



Clear efficiency loss at low G_{KK} masses from dijet p_T requirement
 Optimal dijet p_T cut likely to be higher for higher masses

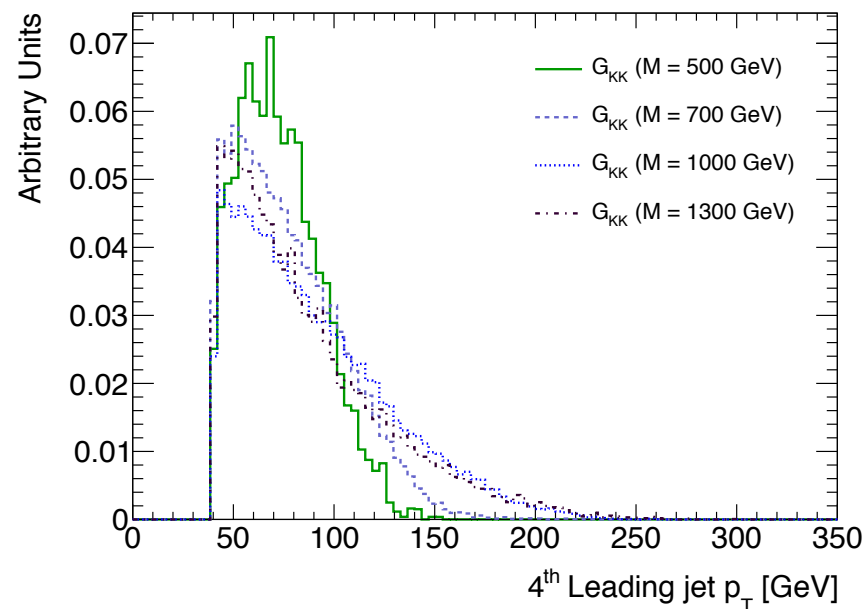
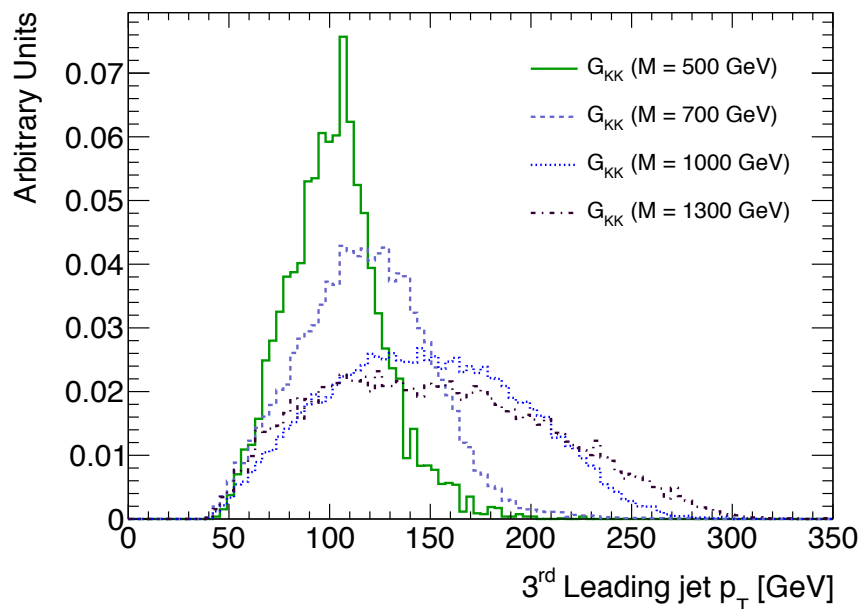
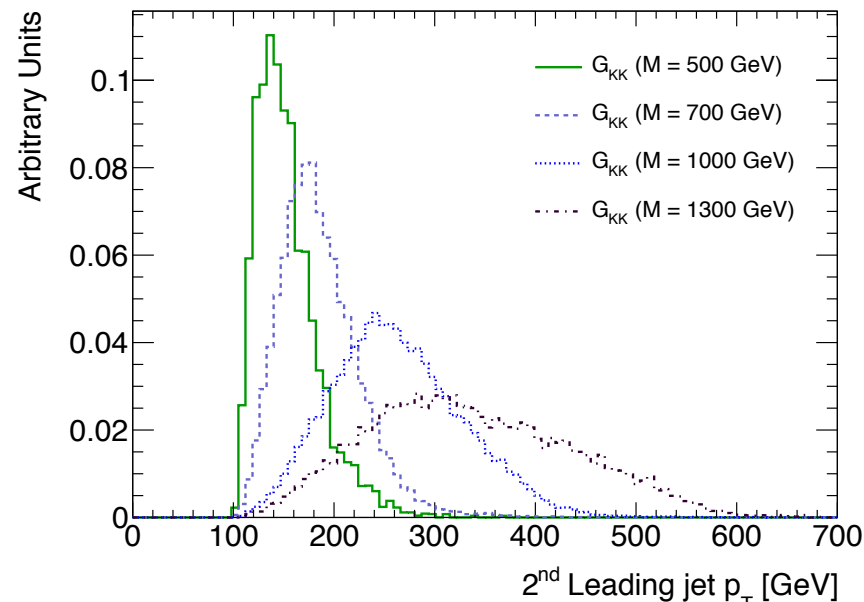
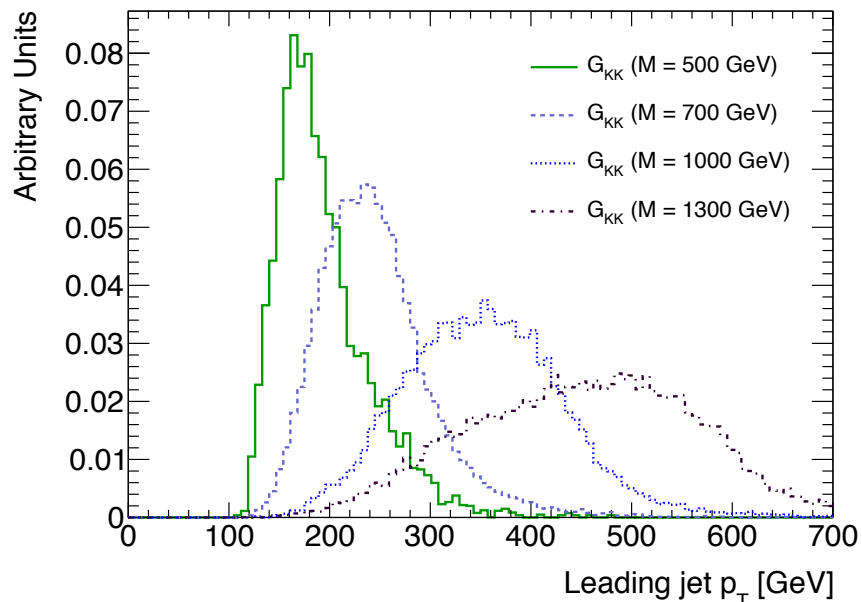
Loss of efficiency from jet merging at high G_{KK} masses in this resolved analysis

For lowest masses the ΔR_{dijet} cut could be optimised



ΔR_{dijet}

Signal Jet Kinematics



QCD Backgrounds

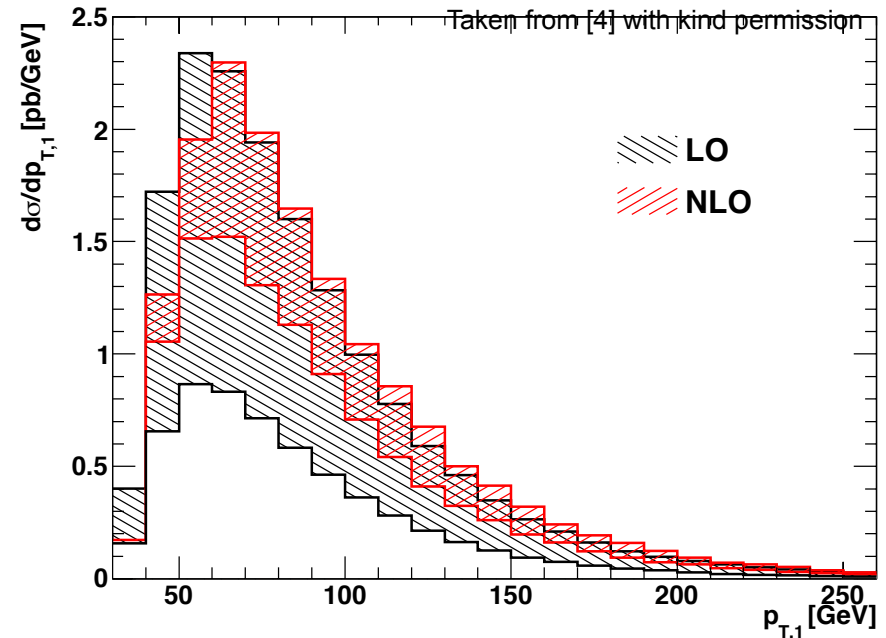
Define uncertainty on our Sherpa background prediction as variation in renormalisation/factorisation scale choice μ_0 by factor $\frac{1}{2}$ and 2:

$$\mu_0 = \mu_F = \mu_R = \frac{1}{4} \sqrt{\sum_i p_{T,i}^2}$$

NLO corrections to LO $pp \rightarrow bbbb$ at $\sqrt{s}=14$ TeV recently calculated in [4] and [5]

NLO/LO corrections are large $\sim 50\%$

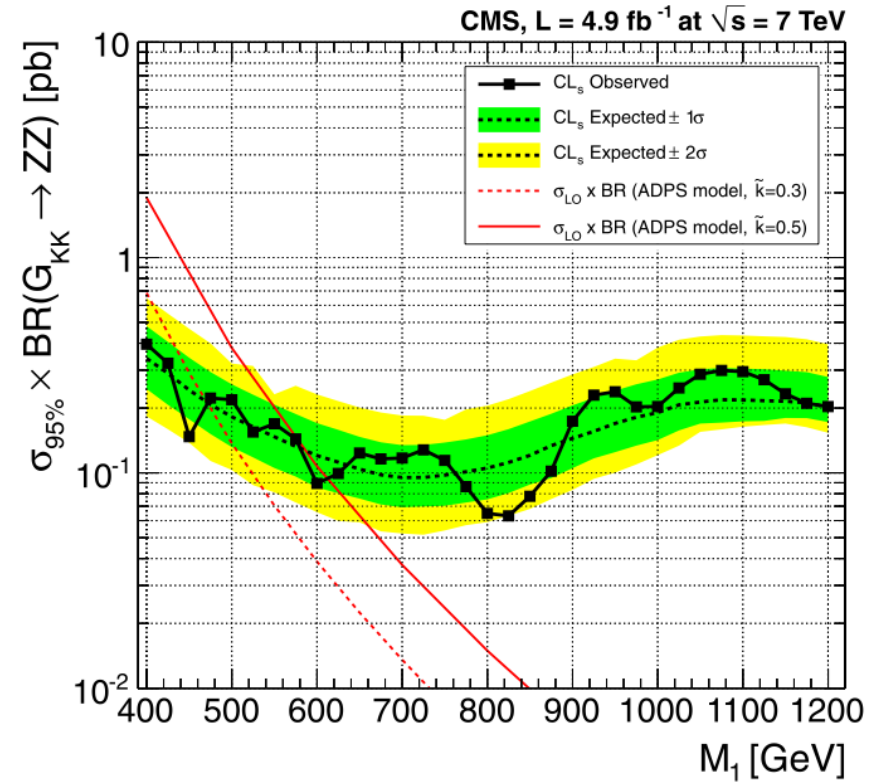
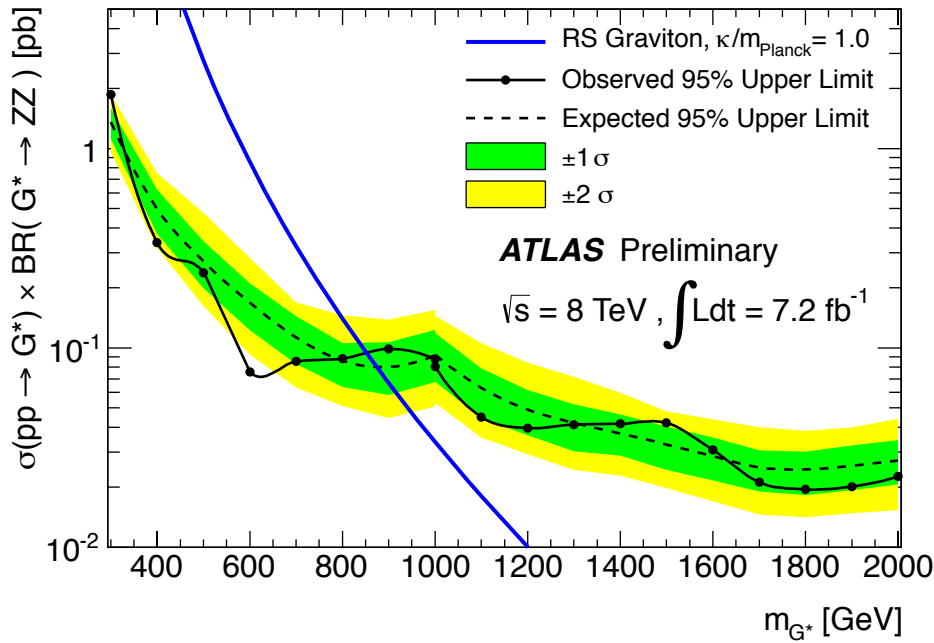
But renormalisation/factorisation scale variations of LO cover the variation at NLO



We successfully reproduced the LO prediction of [4] using Sherpa $bbbb$ at $\sqrt{s}=14$ TeV with the same scale choice μ_0

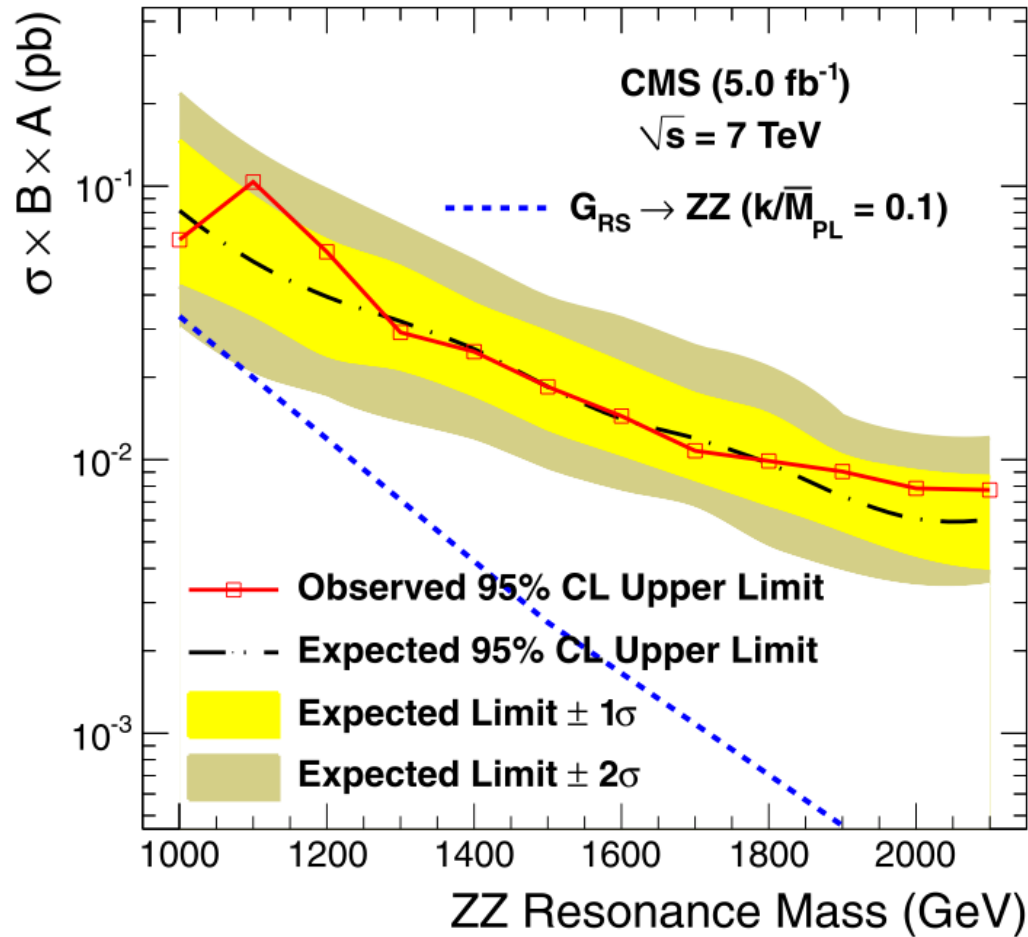
Hence we have some confidence that our scale variations of Sherpa cover NLO corrections

77 Δ Limits on G



95% C.L. upper limits of $\sim 100 \text{ fb}$ at 1 TeV.
 Exclusion up to $m_{G_{KK}} \sim 900 \text{ GeV}$ for $k/M_{\text{Pl}} = 1.0$.

ZZ → jjjj Limits



Don't use ADPS model explicitly.
 95% C.L. upper limits of $\sim 90 \text{ fb}$ at 1 TeV.
 Uses dijet mass of fat-jets with pruning and MDT.