

Broadband Microstrip-Coupled TiN KIDs: A Scalable Detector Technology for CMB-S4

S. Golwala

Caltech

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Outline

Review of needs for CMB-S4

Why microstrip coupling

Why TiN KIDs

Overview of design

Expected performance

Review of Needs for CMB-S4

Raw Sensitivity:

10^5 - 10^6 detectors required to achieve desired sensitivity
bec. limited by irreducible photon noise

Multiplexability is critical!

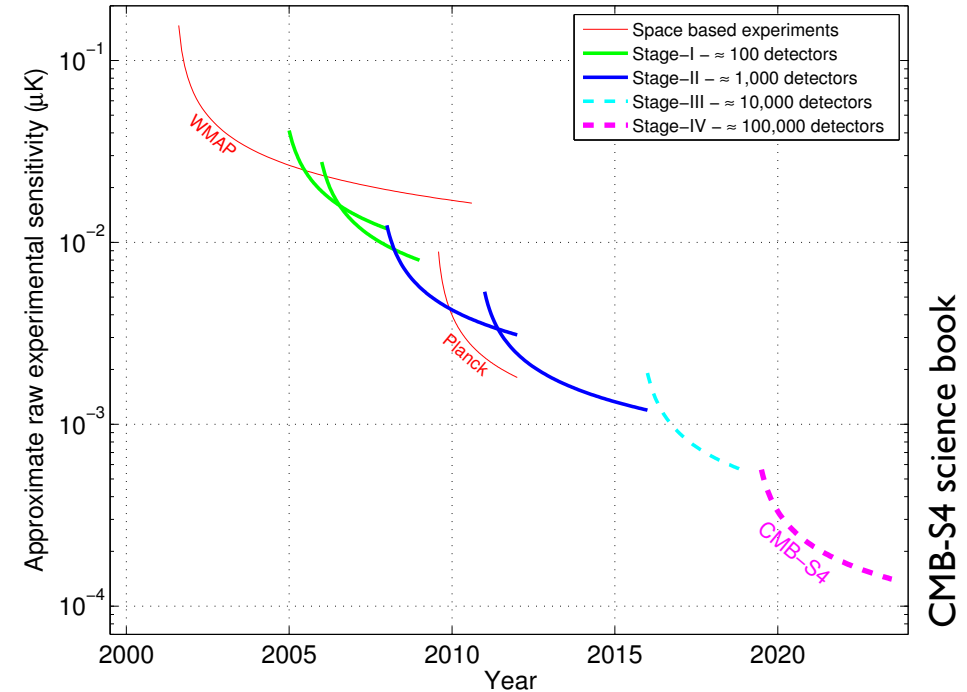
Systematics Control:

Now known that foregrounds exceed inflationary signal even in cleanest sky:
need multiple spectral bands

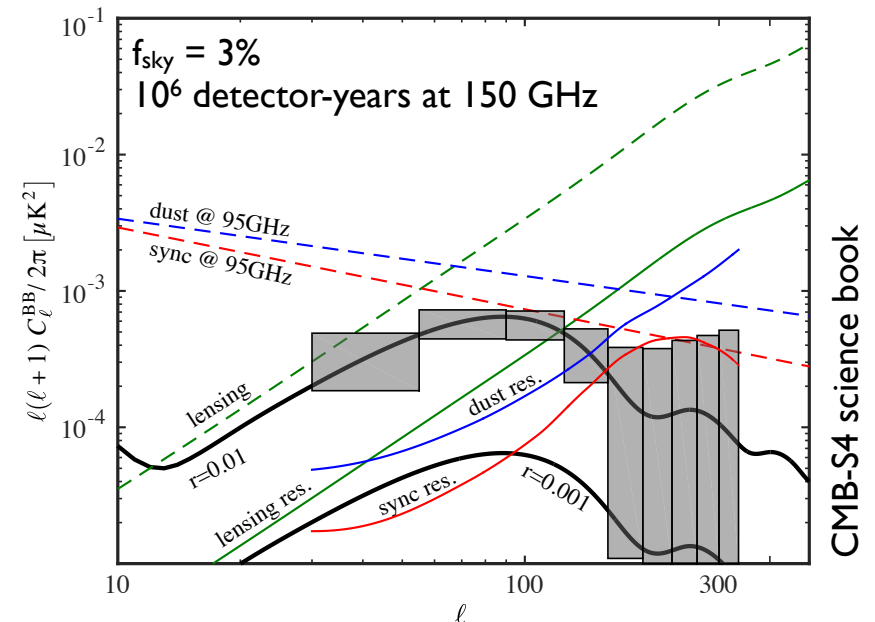
Polarization systematics must be controlled and measured precisely

antenna/bandpass filter systematics requirements from EPIC-IM study
(same sensitivity goal as CMB-S4, $\sigma(r) = 0.0005$)
100 GHz, 8.4' FWHM (some relax for finer beams)

| Systematic Error | Description | Knowledge to Meet Requirement |
|----------------------|--|-------------------------------|
| Δ Beam Size | $\text{FWHM}_E \neq \text{FWHM}_H, \sigma_E - \sigma_H /\langle\sigma\rangle$ | 1.8×10^{-3} |
| Δ Beam Offset | Pointing $E \neq H, \theta_E - \theta_H /\sigma$ | 3×10^{-3} |
| Δ Ellipticity | $e_E \neq e_H, \psi = 0^\circ, e_E - e_H $ | 5×10^{-1} |
| | $e_E \neq e_H, \psi = 45^\circ, e_E - e_H $ | 5×10^{-3} |
| Δ Rotation | $E \& H$ not orthogonal, $ \theta_E - \theta_H - 90^\circ $ | $13.4'$ |
| Pixel Rotation | $E \perp H$ but rotated w.r.t. beam major axis | $2.9'$ |
| Bandpass Mismatch | Variation in filters, $\delta\nu_c/\nu_c$ | 10^{-3} |



CMB-S4 science book



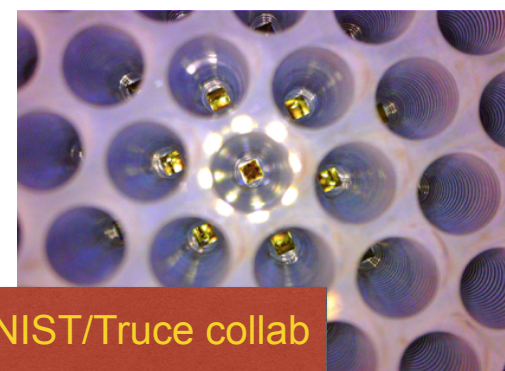
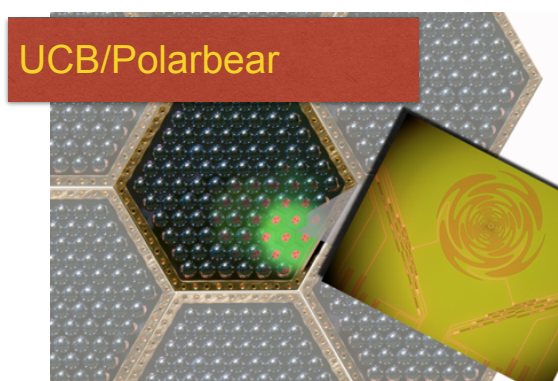
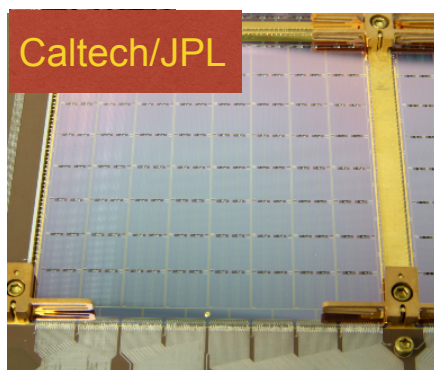
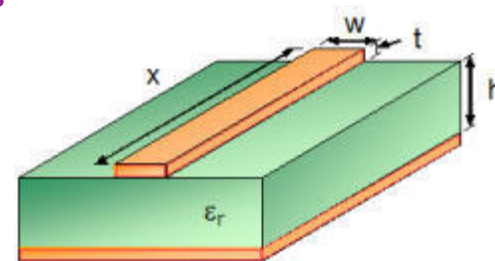
CMB-S4 science book

Why Microstrip Coupling?

Flexible

Many optical reception architectures can output onto superconducting microstripline

Microstrip enables use of bandpass filters, decouples detector size from optical reception element



Excellent systematics control demonstrated in the field for these technologies via BICEP/Keck, Polarbear, SPTpol, ACTpol

Scalable

Fully photolithographic (phased arrays), photolithographic + mass-hybridized (lens-coupled, feedhorn-coupled)

All amenable to scaling to 6" wafer fab modulo yield

Why TiN KIDs?

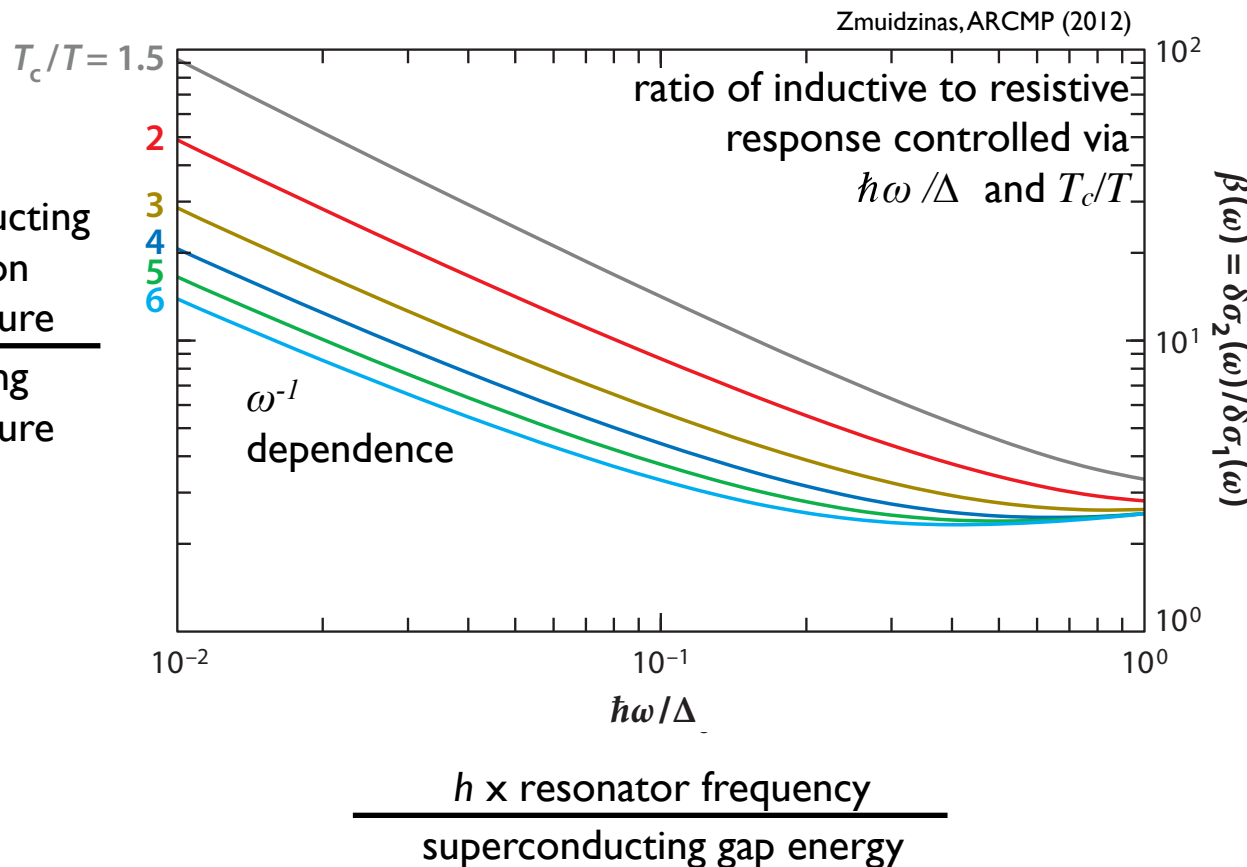
High kinetic inductance fraction due to high normal state resistivity

responsivity in recombination-limited,
mm-wave quasiparticle-limited regime

$$\frac{d(\delta f/f)}{dP_{opt}} = -\frac{\beta(f_r, T)}{4 P_{opt} Q_i}$$

$\beta(f_r, T) > 2$
= ratio of imaginary (freq) to
real (Q) responsivity to qps

$$\beta \propto f_r^{-1} \propto \text{kinetic inductance} \propto \text{superconducting penetration depth} \propto \rho_n^{1/2}$$



TiN: $\rho_n \sim 100 \mu\Omega \text{ cm}$
Al: $\rho_n \sim 1 \mu\Omega\text{-cm}$

imaginary responsivity
real responsivity

Why TiN KIDs?

High kinetic inductance

→ low operating
frequency (10s of MHz)

→ demonstrated

$N_{\text{mux}} \sim 1000$ at low cost

e.g. ROACH1

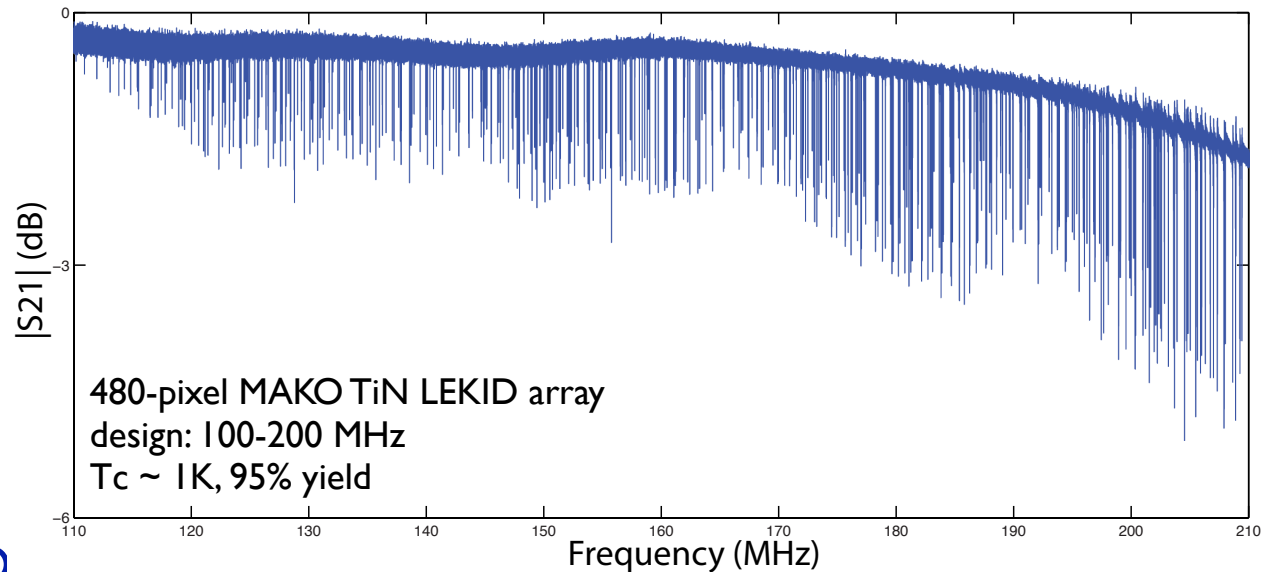
+ MUSIC/ARCHONS

ADC/DAC board + MAKO

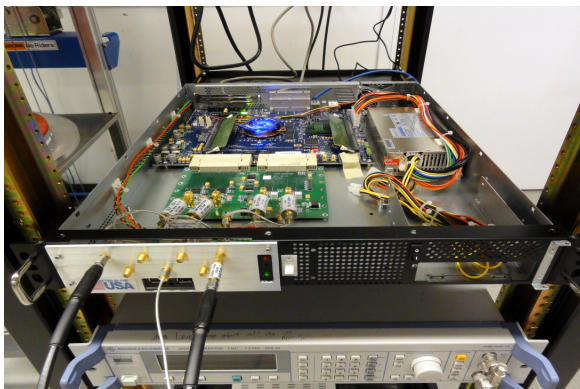
\$4/detector → \$1/detector if optimized

e.g. BLAST-TNG ROACH2 + MUSIC/ARCHONS ADC/DAC: 1K detectors also

e.g. NIKA2, A-MKID readouts



MAKO readout (no IF)



MUSIC thermal-engineered
readout module



MUSIC 8-module readout rack



Why TiN KIDs?

High kinetic inductance

- low operating frequency (10s of MHz) → high β
- amplifier and two-level-system noise subdominant to fundamental noises (photon + recombination)

$$\text{NEP}_{\text{fund}}^2 = \text{NEP}_{\text{shot}}^2 + \text{NEP}_{\text{Bose}}^2 + \text{NEP}_{\text{rec}}^2 = 2 P_{\text{opt}} h\nu \left[1 + \frac{P_{\text{opt}}}{h\nu \Delta\nu} + \frac{2 \Delta}{\eta_{\text{abs}} \eta_{\text{ph}} h\nu} \right]$$

In recombination-dominated, mm-wave quasiparticle-dominated regime:

Amplifier noise:

$$\text{NEP}_{\text{amp}}^2 = 2 P_{\text{opt}} h\nu \frac{32}{[\beta(f_r, T)]^2} \frac{P_{\text{opt}}}{P_{\text{read}}} \frac{k_B T_N}{h\nu}$$

$$\begin{aligned} P_{\text{opt}}/P_{\text{read}} &< 1 && \text{generally holds} \\ k_B T_N/h\nu &\approx 2 && \text{for } T_N = 4 \text{ K} \\ &&& \text{and } \nu = 45 \text{ GHz} \end{aligned}$$

$$\boxed{\beta \approx 55 \quad \text{for our design}}$$

Two-Level-System Noise:

$$\text{NEP}_{\text{TLS}}^2 = \frac{16 P_{\text{opt}}^2 Q_i^2}{[\beta(f_r, T)]^2} S_{\delta f/f}^{\text{TLS}}(f_r^0, f_a^0, T_0, V_{d,0}, E_0) \frac{E_0 V_{d,0}}{\ell_C^2} \sqrt{\frac{2 \pi f_r C}{Q_i P_{\text{read}}}} \sqrt{\frac{f_a^0}{f_a}} \left(\frac{T_0}{T}\right)^{1.7} g\left(\frac{f_r}{f_r^0}\right)$$

TLS noise level
measured for material
in specific geometry
scalings to desired geometry and
operating parameters

high $\beta \approx 55$ for our design makes this subdominant

Why TiN KIDs?

Fundamental noise level comparable to TESs:

$$\text{NEP}_{\text{fund}}^2 = \text{NEP}_{\text{shot}}^2 + \text{NEP}_{\text{Bose}}^2 + \text{NEP}_{\text{rec}}^2 = 2 P_{\text{opt}} h \nu \left[1 + \frac{P_{\text{opt}}}{h \nu \Delta \nu} + \frac{2 \Delta}{\eta_{\text{abs}} \eta_{\text{ph}} h \nu} \right]$$

KIDs suffer recombination noise, TESs suffer phonon noise:

both are noise on coupling to bath needed for optical power to escape

$$\text{NEP}_{\text{rec}}^2 = 2 P_{\text{opt}} \frac{h \nu_{\text{min}}}{\eta_{\text{abs}} \eta_{\text{ph}}} \quad \text{NEP}_{\text{G}}^2 = 4 k_B \gamma (T_c^{\text{T}})^2 G_c \approx 2 P_{\text{opt}} (2 S \gamma k_B T_c^{\text{T}}) \frac{T_c^{\text{T}}}{T_c^{\text{T}} - T}$$

Calculate for

ground-based optical loads

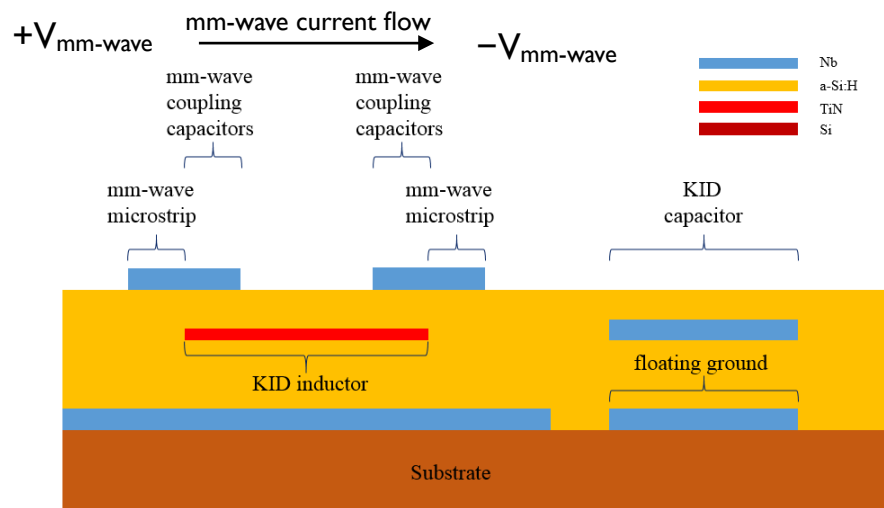
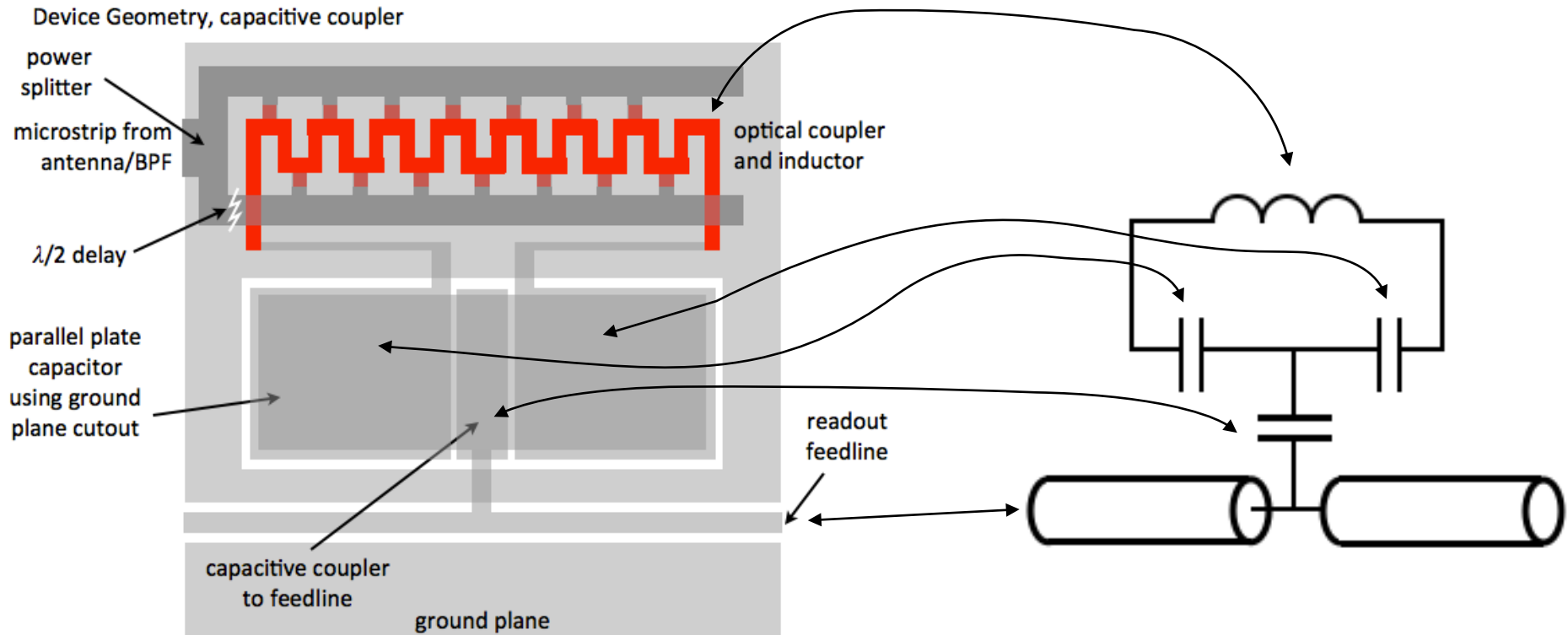
KID: $T_c = 0.52$ K, $T = 0.05$ K, η_{abs} near unity, $\eta_{\text{ph}} = 1.0$ at $\nu_{\text{min}} \rightarrow 0.57$ as $\nu \rightarrow \infty$

TES: $T_c = 0.1$ K, $T = 0.05$ K, $\gamma = 0.5$, safety factor $S = 3$

| ν [GHz] | $\Delta \nu$ [GHz] | P_{opt} [pW] | η_{ph} | NEP_{shot} [aW/ $\sqrt{\text{Hz}}$] | NEP_{Bose} | NEP_{rec} [aW/ $\sqrt{\text{Hz}}$] | $\text{NEP}_{\text{fund}}^{\text{KID}}$ [aW/ $\sqrt{\text{Hz}}$] | NEP_{G} [aW/ $\sqrt{\text{Hz}}$] | $\text{NEP}_{\text{fund}}^{\text{TES}}$ | $\frac{\text{NEP}_{\text{fund}}^{\text{KID}}}{\text{NEP}_{\text{fund}}^{\text{TES}}}$ |
|----------------|-----------------------|--------------------------|--------------------|---|----------------------------|--|--|--|---|---|
| 45 | 13.5 | 1.0 | 0.85 | 8 | 12 | 8 | 16 | 6 | 15 | 1.05 |
| 100 | 30 | 1.7 | 0.57 | 15 | 14 | 12 | 24 | 8 | 22 | 1.10 |
| 150 | 47 | 2.5 | 0.57 | 22 | 17 | 15 | 32 | 9 | 29 | 1.08 |

Negligible difference in theory; real-world effects will dominate

Overview of Design



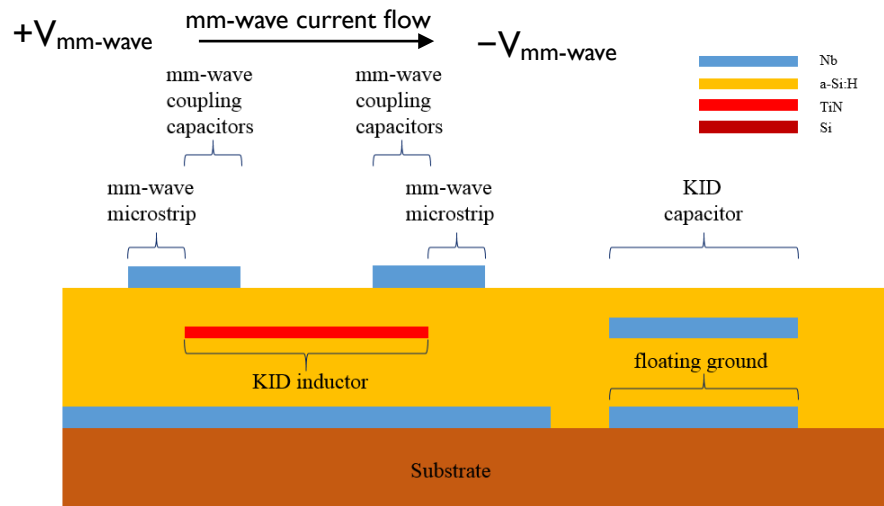
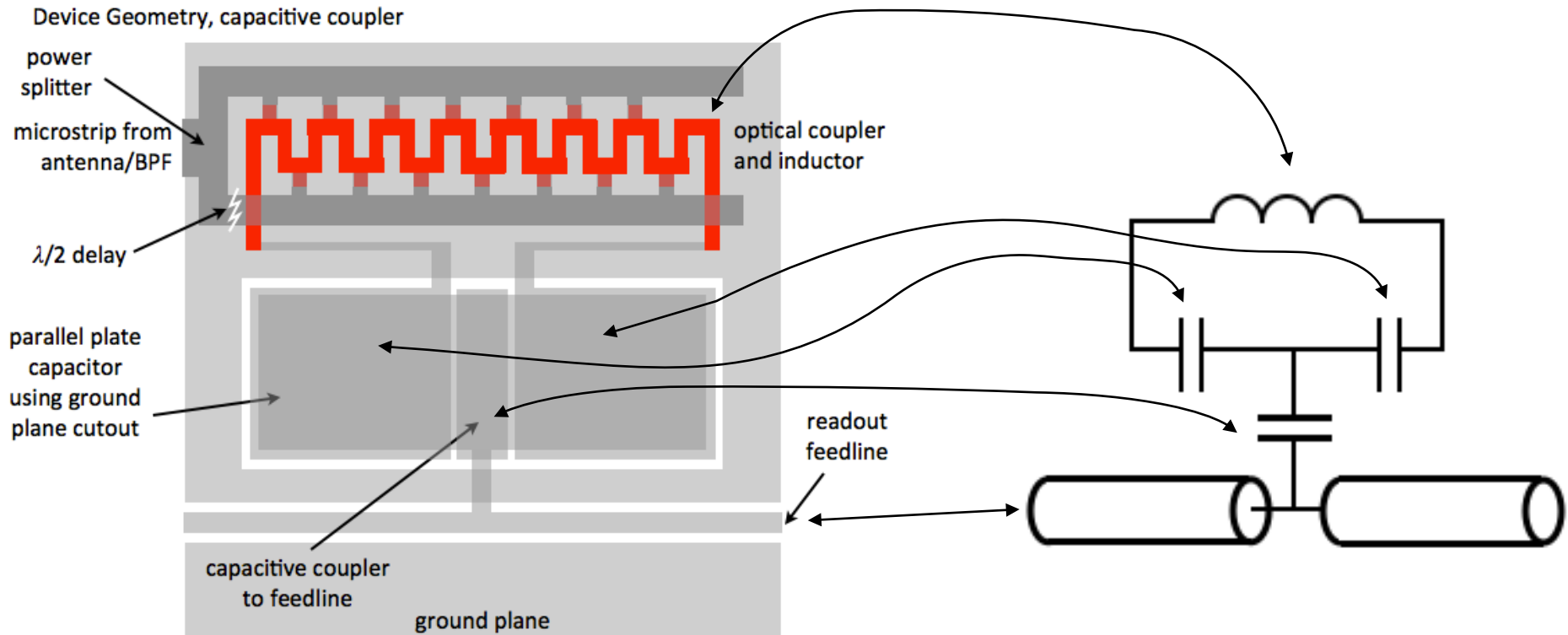
Microstrip coupling

High sheet resistance of TiN makes it good match to free space ($Z_0 = 377 \Omega$)

Poor match to microstrip ($Z_0 = 10 \Omega$)

Solution: adiabatic coupler, achieves η_{abs} near unity

Overview of Design



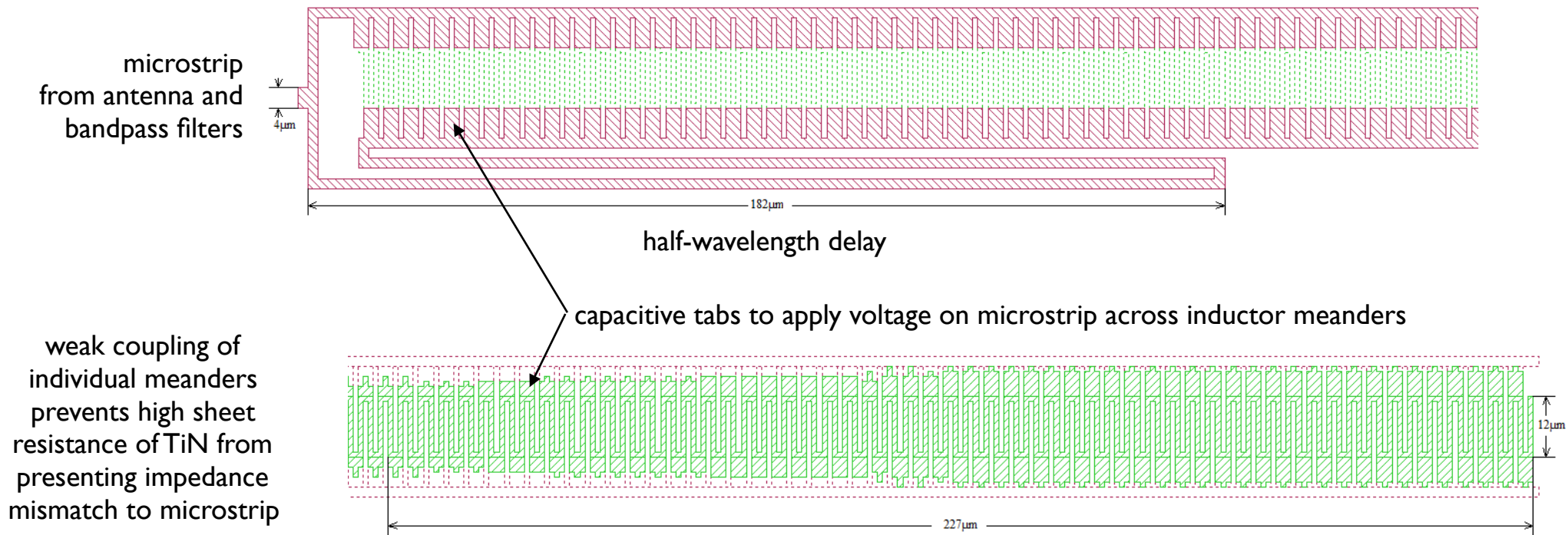
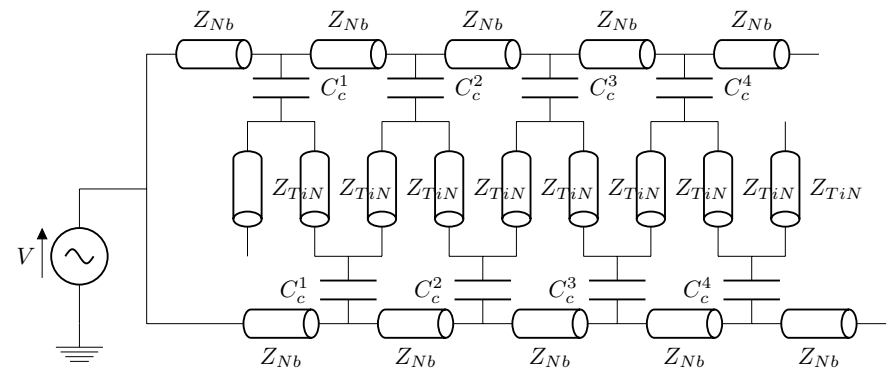
Prevents direct absorption

Use parallel plate rather than interdigitated capacitor

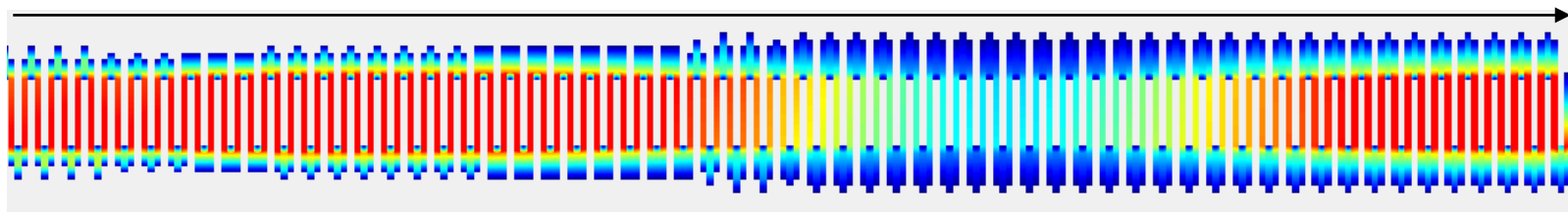
TLS noise of a-Si:H (Mazin et al 2010)
good enough (v. little from inductor)

Ground plane shields inductor, capacitor via short-circuit condition

Millimeter-Wave Coupler

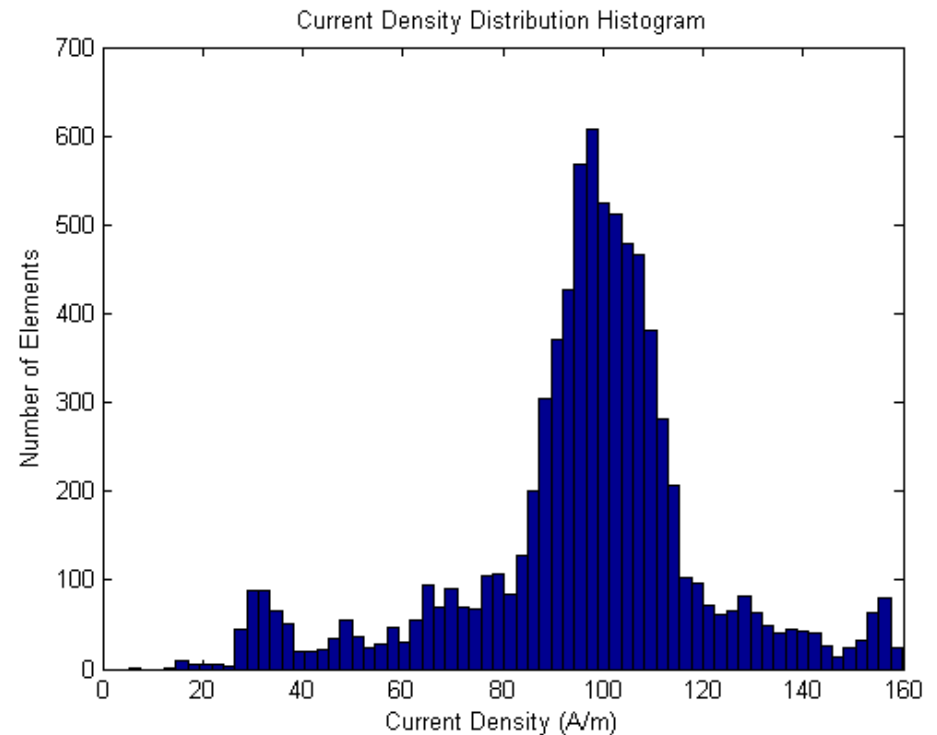
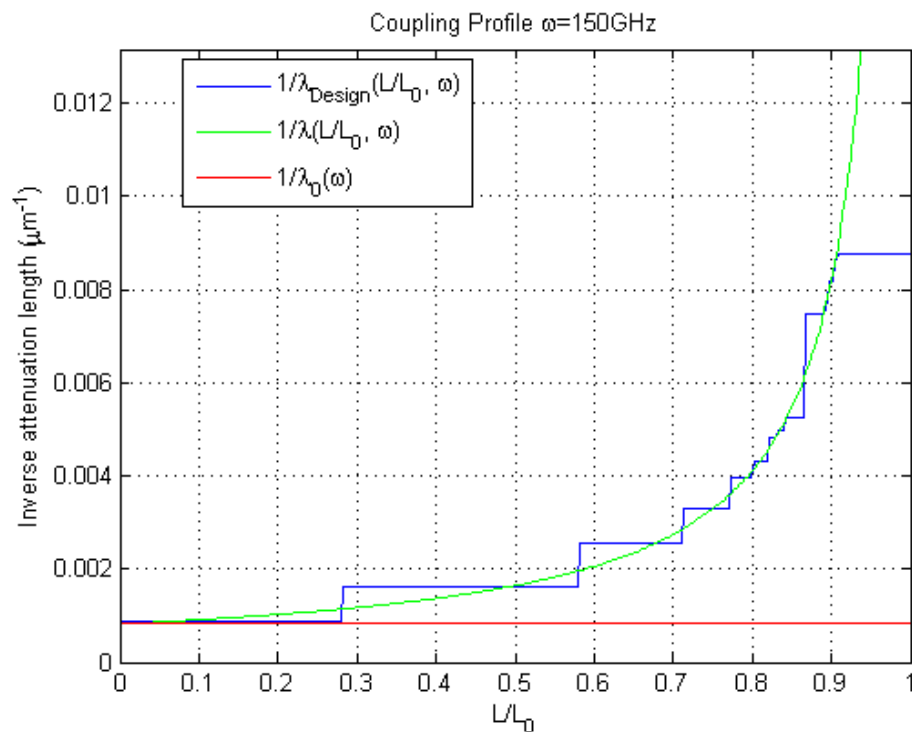
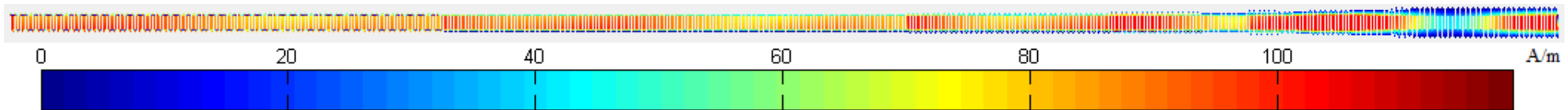
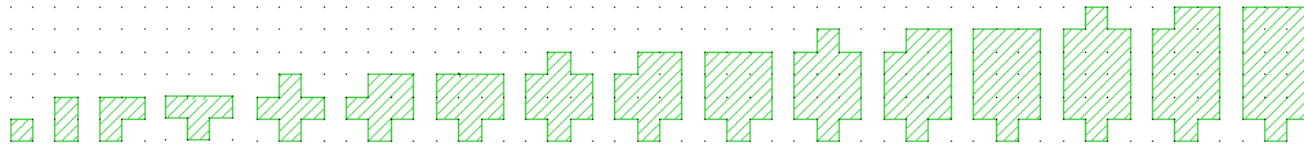


increase coupling left to right by increasing capacitor size: ensure uniform mm-wave power absorption



Millimeter-Wave Coupler

Increasing coupling adiabatically yields uniform mm-wave current density, power absorption



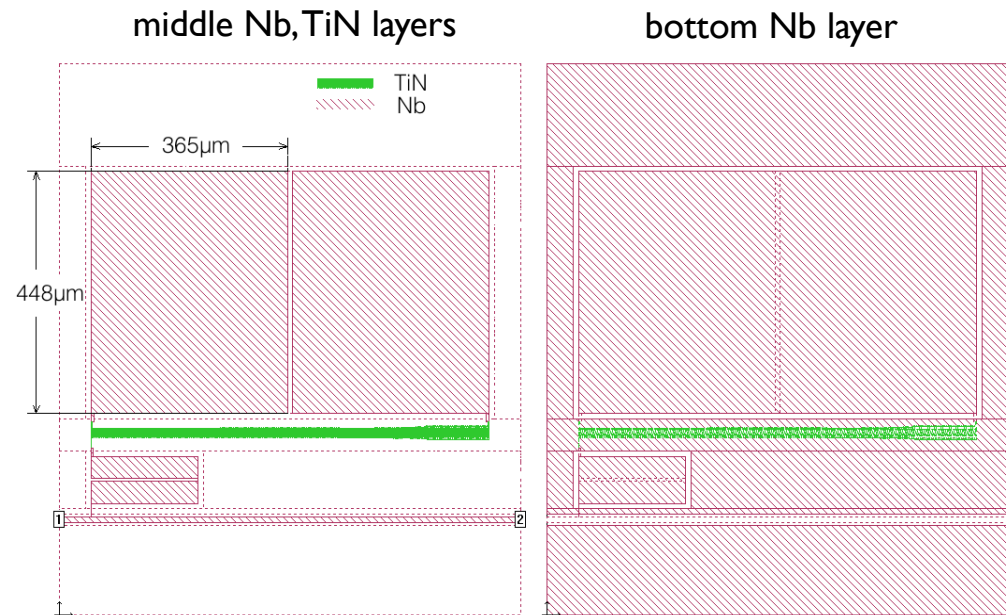
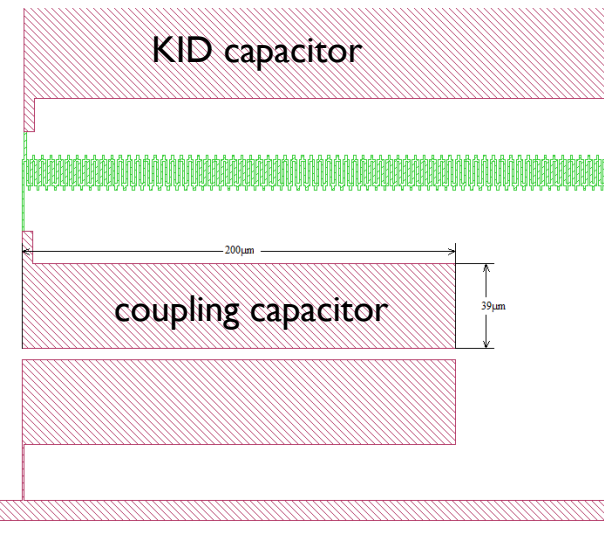
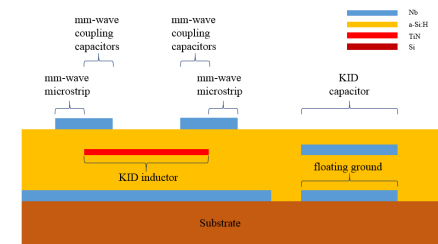
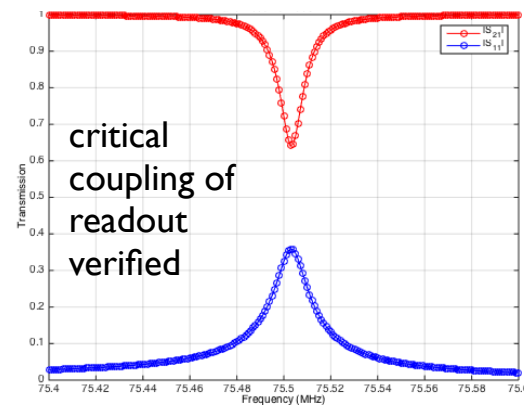
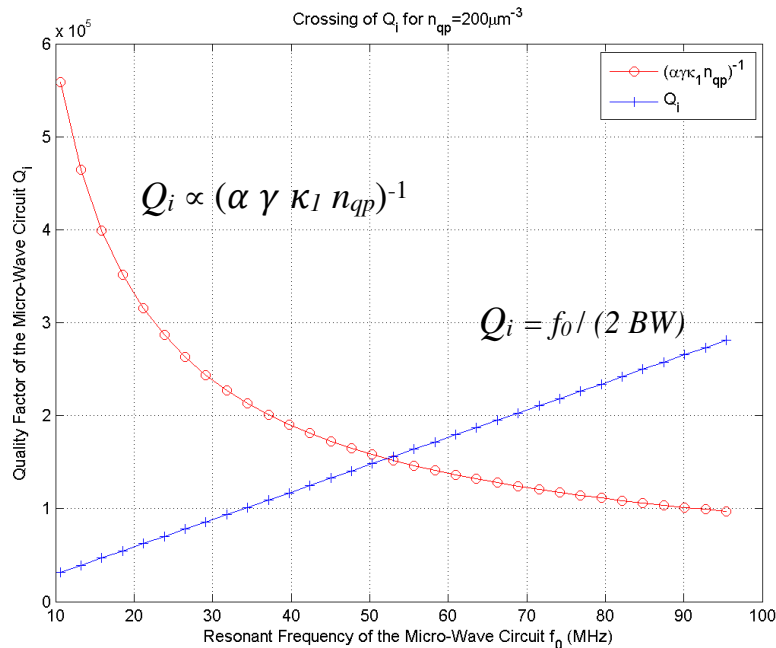
KID Design

Coupling to feedline

Use parallel plate capacitor
like KID capacitor

Readout frequency

Maximize Q_i subject to constraint
that resonator bandwidth
 $f_0/2Q_i > \text{astronomical signal BW}$
(50 Hz for $\ell = 5000$ and $1^\circ/\text{s}$;
 Q_i is mm-wave qp-limited)



Optimization Pressures

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TiN thickness and inductor width saturate at min allowed values, 20 nm and 1 μm

Limits have been chosen ensure robust fabrication (thinner/narrower possible, but concern about yield)

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At fixed capacitor area, increasing thickness reduces the TLS noise ($\delta C/C$ or $\delta f/f$ units) bec. TLS noise $\propto I/\sqrt{N_{\text{fluctuators}}} \propto I/\sqrt{V_{\text{cap}}}$, so choosing the maximum value is sensible

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TLS noise pushes for high P_{read} . Require $J < \text{critical current density}$ to avoid nonlinearity

Expected Sensitivity

| Quantity | 100 GHz | 150 GHz |
|--|---------|---------|
| NEP_{ph} [aW Hz ^{-1/2}] | 21 | 38 |
| NEP_{r} [aW Hz ^{-1/2}] | 17 | 30 |
| NEP_{amp} [aW Hz ^{-1/2}] | 2 | 4 |
| NEP_{TLS} [aW Hz ^{-1/2}] | 9 | 20 |
| NEP_{fund} [aW Hz ^{-1/2}] | 27 | 48 |
| NEP_{tot} [aW Hz ^{-1/2}] | 28 | 52 |
| $\text{NEP}_{\text{tot}}^{\text{eff}}$ [aW Hz ^{-1/2}] | 29 | 53 |
| $\text{NEP}_{\text{tot}}^{\text{eff}}/\text{NEP}_{\text{ph}}$ | 1.40 | 1.41 |
| $\text{NEP}_{\text{tot}}^{\text{eff}}/\text{NEP}_{\text{fund}}$ | 1.07 | 1.10 |
| $\text{NET}_{\text{CMB}}^{\text{eff}}$ [μK_{CMB} s ^{1/2}] | 200 | 250 |

$T_c = 1 \text{ K}, T = 0.1 \text{ K}$
 evaluated at 1 Hz
 $\rightarrow \ell = 100 \text{ at } 1^\circ/\text{s}$
 $\Delta\nu/\nu = 0.3$
 $T_{\text{load}} = 15 \text{ K}$
 $\eta_{\text{opt}} = 0.26, 0.40$
 $P_{\text{opt}} = 1.9, 3.9 \text{ pW}$

Expected Sensitivity

Yields very competitive
expected sensitivities

c.f on-sky sensitivities:

BICEP2/Keck Array: $400 \mu\text{K}_{\text{CMB}} \text{ s}^{1/2}$

SPTpol: $500 \mu\text{K}_{\text{CMB}} \text{ s}^{1/2}$

POLARBEAR: $550 \mu\text{K}_{\text{CMB}} \text{ s}^{1/2}$

ACTPol: $330 \mu\text{K}_{\text{CMB}} \text{ s}^{1/2}$

Substantial headroom for
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Some other aspects of design:

capacitor area $\sim 2 \text{ mm}^2 \gg A_{\text{feed}} = 44 \text{ mm}^2$ assumed

inductor length $\sim 7 \text{ mm}$, inductor volume $\sim 150 \mu\text{m}^3$

qp density $\sim 2000 \mu\text{m}^{-3} \gg$ “irreducible qp density” $\sim 100 \mu\text{m}^{-3}$

$\beta \approx 55$; $Q_i \sim 70000$, dominated by optical load

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Such low loss tangent a-Si:H not demonstrated at JPL

But crystalline silicon tests show comparable or better ($\times 10$) loss tangent, can integrate into design if necessary

Caveats/Modeling To Do

Assumes TiN obeys Mattis-Bardeen; known to be not true

But TiN responsivity seems to decrease less quickly with increasing P_{opt} than expected (or even increase!); our calculations are conservative

Need to add empirical model for TiN responsivity

Assumes reasonable uniformity, behavior of TiN

Try other high-resistivity materials if a problem

Assumes bulk material η_{ph}

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P_{read} generation of qps not included

$P_{\text{read}} = -90 \text{ dBm}$ ($= 1 \text{ pW}$) assumed; heating effects not usually seen at this low power (e.g., deVissers et al 2010, Hubmayr et al 2015)

Need to incorporate in modeling

Conclusions

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CMB-S4 detectors need to be highly multiplexable and amenable to good polarization systematics control

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- Provides microstrip coupling

- Will be multiplexable at $N_{\text{mux}} \sim 1000$ for approach \$1/detector

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Plans

- Mask in hand; will make first pass at fabricating this FY on JPL R&TD (seed) funds

- DOE detector R&D proposal submitted

- Also working on crystalline silicon capacitors (and microstrip) as backup for a-Si:H

Backups

Optical Power Density and Responsivity

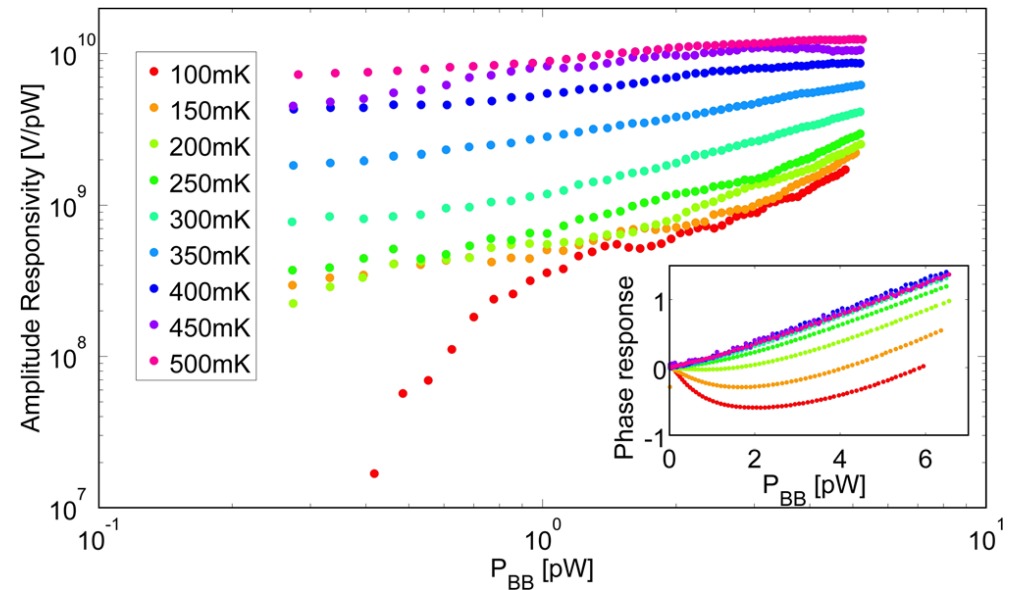
Bueno et al 2014

Ignore TLS-dominated phase response

Amplitude response increases with load rather than decreasing as $I/\sqrt{P_{\text{opt}}}$ as expected from M-B

This would improve margin relative to amplifier and TLS noise!

Strange discrepancy between optical and electrical NEP observed



Hubmayr et al 2015

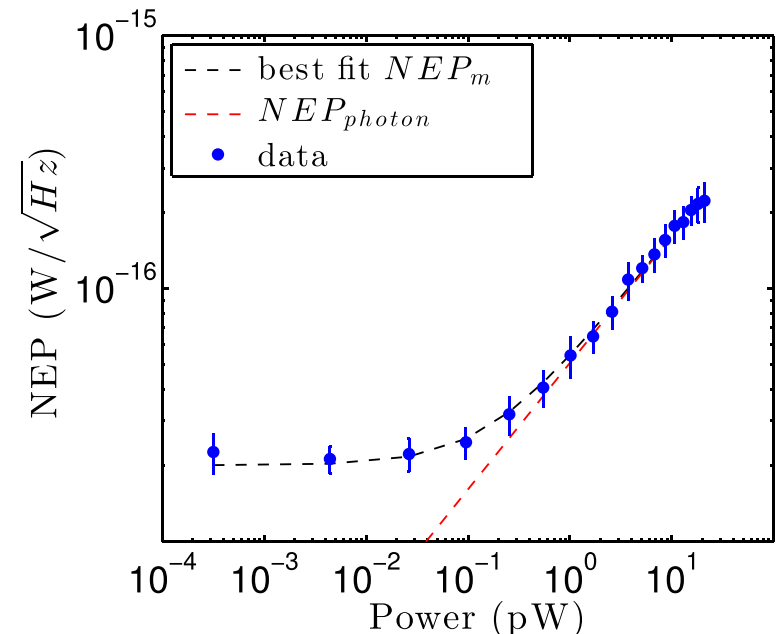
$-97 \text{ dBm} < P_{\text{read}} < -87 \text{ dBm}$, still linear

P_{res} comparable to our design

$V_{\text{inductor}} = 86 \text{ } \mu\text{m}^3$, only 2x smaller than ours

Shows photon noise limited behavior for $P_{\text{opt}} > 1 \text{ pW}$

Measured at 1.2 THz, so NEP_{rec} negligible



Optimization Detail

Table 2: Parameters of KID design.

External (fixed) parameters at left, optimized parameters and parameters derived from them at right. Parameters not defined in the text: A_{feed} is the assumed area of the feed and corresponds to $2.2 f \lambda$ at 150 GHz and $1.3 f \lambda$ at 90 GHz for $f/\# = 1.5$. These numbers are typical of current systems (*e.g.*, Keck Array [24], SPTpol [64, 32]); $A_L = \ell_L w$; $V = A_L t$; $A_C = \ell_C^2$; $A_{KID} = A_L + A_C$; $L = \mu_0 \lambda_{pen} \ell_L/w$; $C = \epsilon_d \epsilon_0 \ell_c^2/t_d$ where $\epsilon_d = 12$ is the dielectric constant of a-Si:H; E = electric field in capacitor; $\ell_{coupler}$ = length of inductive coupler to feedline.

Band-Independent External Parameters

| Quantity | Value |
|------------------------|--|
| T_c | 1.0 K |
| 2Δ | 300 μeV |
| $2\Delta/h$ | 73 GHz |
| N_0 | $3.5 \times 10^{10} \mu\text{m}^{-3} \text{eV}^{-1}$ |
| R | 100 $\mu\text{m}^3 \text{s}^{-1}$ |
| τ_0 | 100 μs |
| T | 0.1 K |
| t_d | 400 nm |
| T_N | 4.0 K |
| P_{read} | 1.0 pW |
| J^* | 6.0 mA μm^{-2} |
| $S_{\delta f/f}^{TLS}$ | see text |
| f_a | 1 Hz |

Band-Dependent External Parameters

| Quantity | 90 GHz | 150 GHz |
|------------------------------|--------|---------|
| $\Delta\nu$ [GHz] | 35 | 47 |
| T_L [K] | 15 | 15 |
| A_{feed} [mm^2] | 44.4 | 44.4 |
| η_{opt} | 0.26 | 0.40 |
| P_{opt} | 1.9 | 3.9 |
| η_{ph} | 0.81 | 0.57 |

Optimized Parameters

| Quantity | 90 GHz | 150 GHz |
|---|--------|---------|
| Varied Parameters | | |
| ℓ_L [mm] | 6.3 | 7.4 |
| w [μm] | 1.01 | 1.04 |
| t [nm] | 20 | |
| Derived Geometrical Parameters | | |
| ℓ_C [mm] | 1.42 | 1.13 |
| A_L [mm^2] | 0.006 | 0.008 |
| V [μm^3] | 128 | 154 |
| A_C [mm^2] | 2.00 | 1.28 |
| $A_{feed}/(A_{feed} + A_{KID})$ | 0.96 | 0.97 |
| Derived RF Parameters | | |
| λ_{pen} [μm] | 30 | |
| L [μH] | 0.24 | 0.27 |
| C [pF] | 133 | 85 |
| f_r [MHz] | 28 | 33 |
| Q_i [10^4] | 8.3 | 6.8 |
| J_{read} [mA μm^{-2}] | 2.2 | 1.7 |
| E [kV/m] | 4.7 | 4.9 |
| $S_{\delta f/f}^{TLS}$ [10^{-19} Hz $^{-1}$] | 2.1 | 3.1 |
| $\ell_{coupler}$ [mm] | 3.4 | 3.7 |
| Derived Quasiparticle Parameters | | |
| n_{qp} [μm^{-3}] | 1850 | 2000 |
| τ_{dyn} [μs] | 2.6 | 2.4 |
| Derived Responsivity | | |
| η_{abs} | 0.81 | 0.78 |
| $d(\delta f/f)/dP_{opt}$ [10^6 W $^{-1}$] | 95 | 50 |
| Derived Sensitivity | | |
| NEP _{ph} [aW Hz $^{-1/2}$] | 21 | 38 |
| NEP _r [aW Hz $^{-1/2}$] | 17 | 30 |
| NEP _{amp} [aW Hz $^{-1/2}$] | 2 | 4 |
| NEP _{TLS} [aW Hz $^{-1/2}$] | 9 | 20 |
| NEP _{fund} [aW Hz $^{-1/2}$] | 27 | 48 |
| NEP _{tot} [aW Hz $^{-1/2}$] | 28 | 52 |
| NEP _{tot} ^{eff} [aW Hz $^{-1/2}$] | 29 | 53 |
| NEP _{tot} ^{eff} /NEP _{ph} | 1.40 | 1.41 |
| NEP _{tot} ^{eff} /NEP _{fund} | 1.07 | 1.10 |
| NET _{CMB} ^{eff} [μK_{CMB} sec $^{1/2}$] | 200 | 250 |

Crystalline Si and a-Si:H

Loss Tangent Comparisons

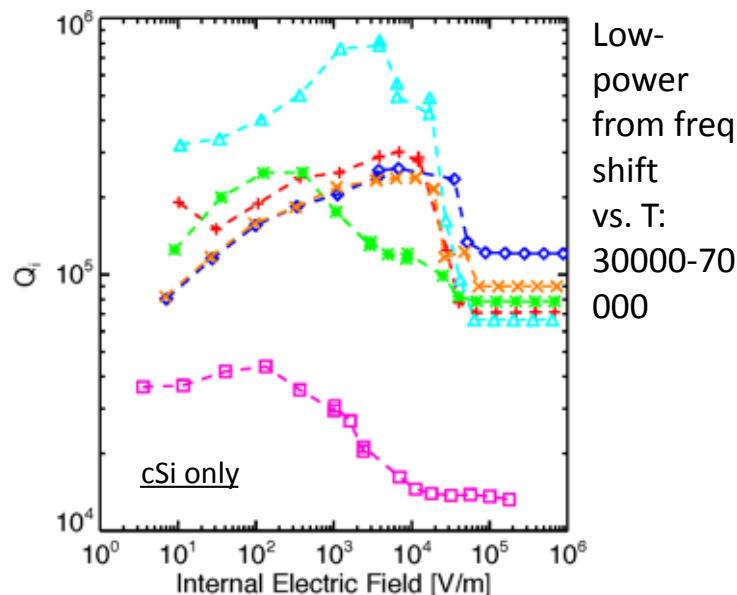
Virgin cSi results comparable w/Siddiqi gp. (though not precisely)

Currently, wafer-bonded cSi shows poorer high-field Q_i than a-Si:H, comparable low-field Q_i

Wafer-bonded cSi comparable to virgin cSi w/native oxide

a-Si:H comparable to virgin cSi at high field, degrades more at low field

MUSIC a-SiN_x had high-field $Q_i \sim 10000$ low-field $Q_i \sim 1000$, all would be vast improvements



Next steps:

- noise measurements to correlate TLS noise w/high-field Q_i
- mm-wave loss measurements to correlate w/low-power Q_i

Low-power from freq shift vs. T: 8000-20000

