

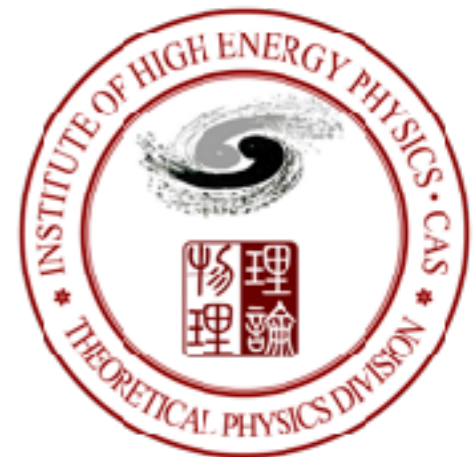
Numerical approach to multi-scale multi-loop integrals

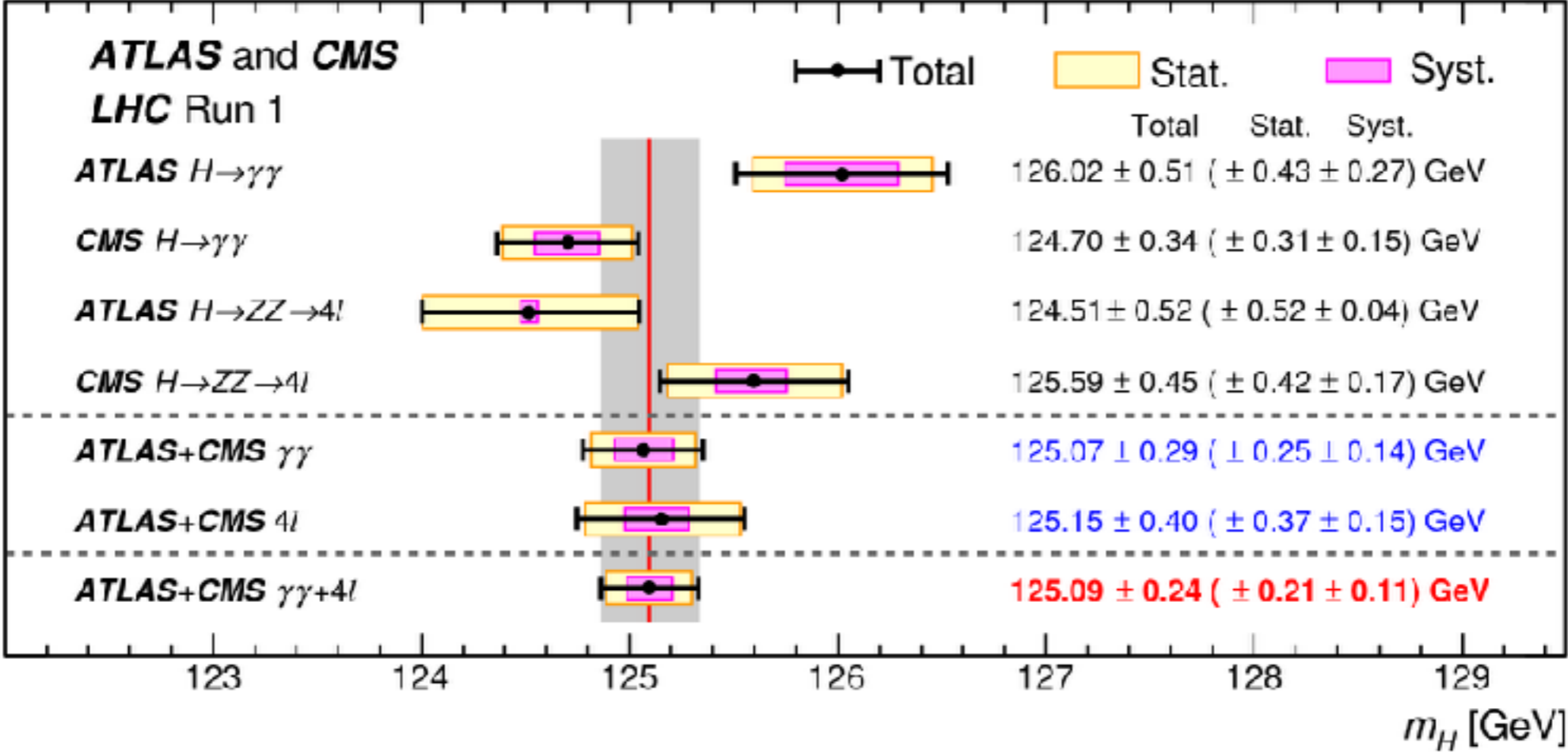
Zhao Li

Institute of High Energy Physics, Chinese Academy of Science

June 1st, 2017 @ LOOPFEST XVI, ANL

i.c.w. Gexing Li, Jian Wang, Yan Wang, Xiaoran Zhao





State-of-Higgs

- Anastasiou et al. JHEP 1605 (2016) 058
N³LO inclusive Higgs Xsection in infinite m_t (finite m_t @NLO)
- Dulat et al. arXiv:1704.08220
NNLO differential Higgs Xsection in EFT
- Banfi et al. JHEP 1604 (2016) 049
N³LO+NNLL Jet vetoed Higgs Xsection in EFT
- Chen et al. JHEP 1610 (2016) 066
NNLO Higgs+Jet in EFT (finite m_t @LO)
- Grigo et al. NPB888 (2014) 17
NNLO Higgs pair in EFT
- Grigo et al. NPB900 (2015) 412
NNLO Higgs pair in $1/m_t$ expansion
- Borowka et al. JHEP 1610 (2016) 107
NLO Higgs pair with finite m_t
- Many other calculations.....

CEPC-SPPC



Perimeter: 50~70km
~1million Higgs bosons @ CEPC

100TeV SppC?

时间	主要工作内容
2015 ~ 2020	CEPC 设计、预研
2021 ~ 2027	CEPC 建造
2027 ~ 2035	CEPC 运行
2015 ~ 2030	SppC 设计、预研
2030 ~ 2040	SppC 建造
2040 ~ 2050	SppC 运行

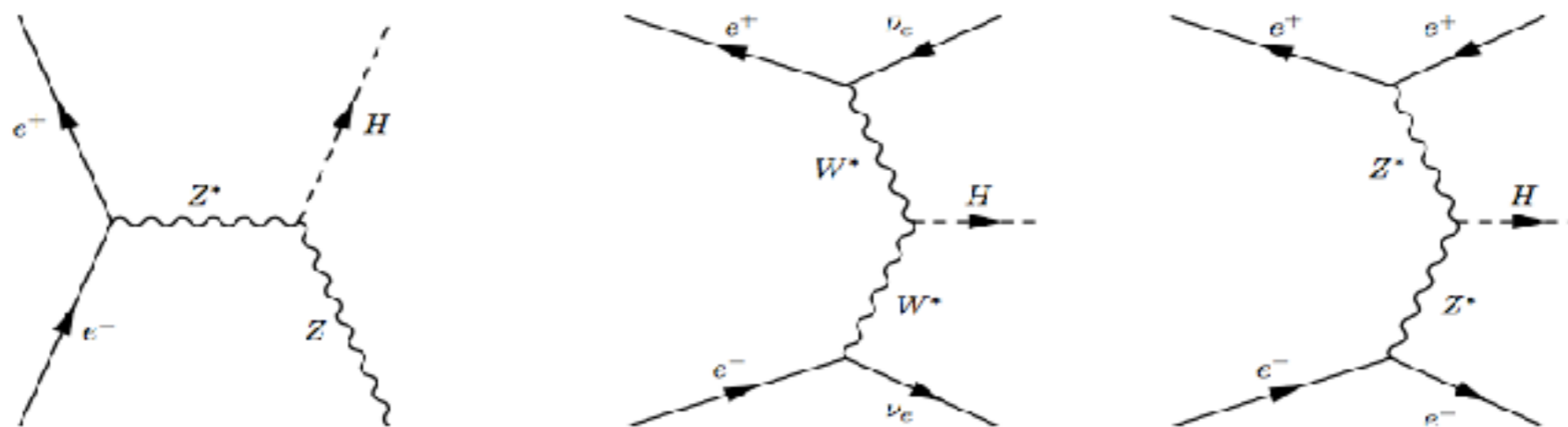
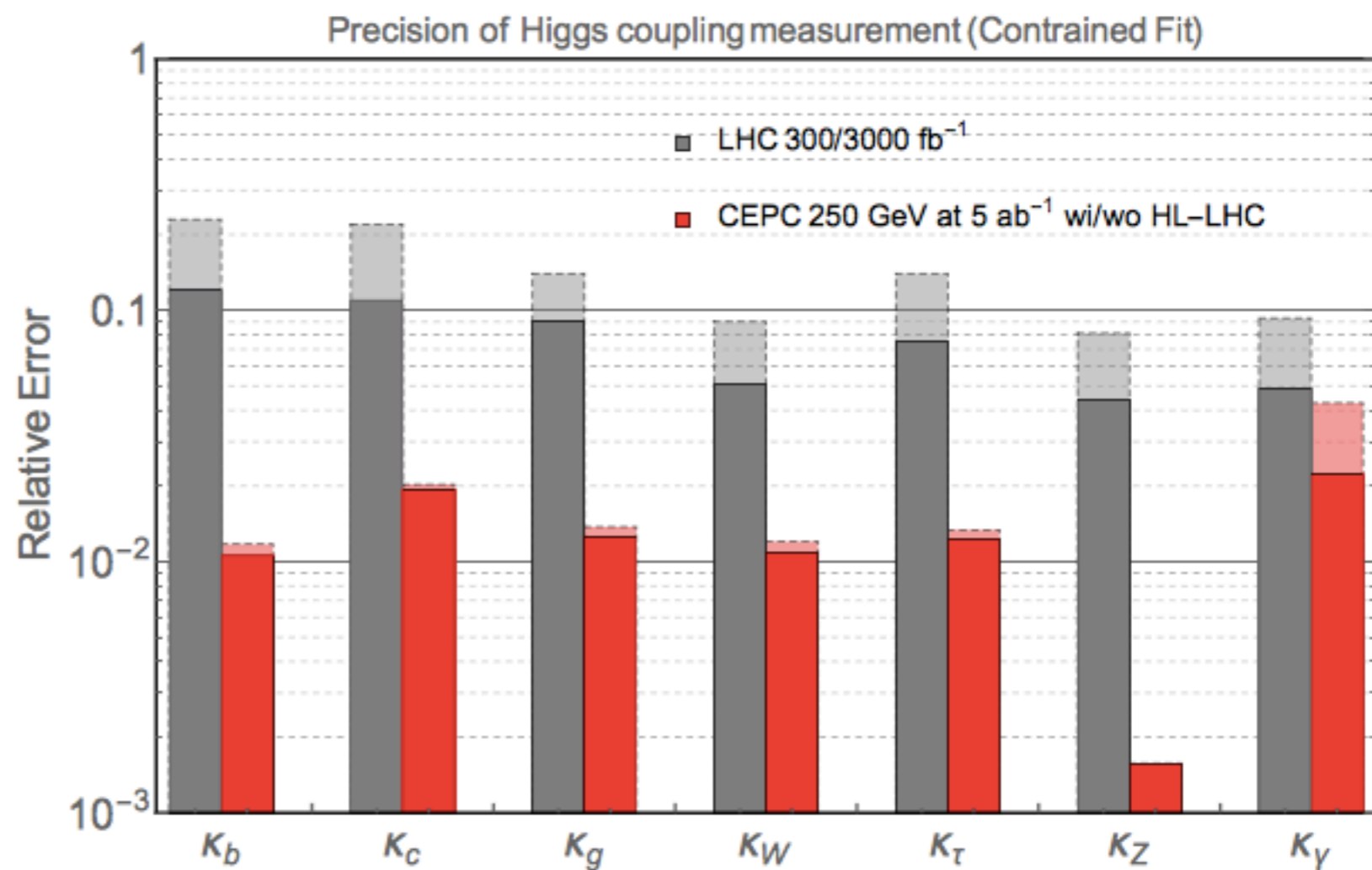
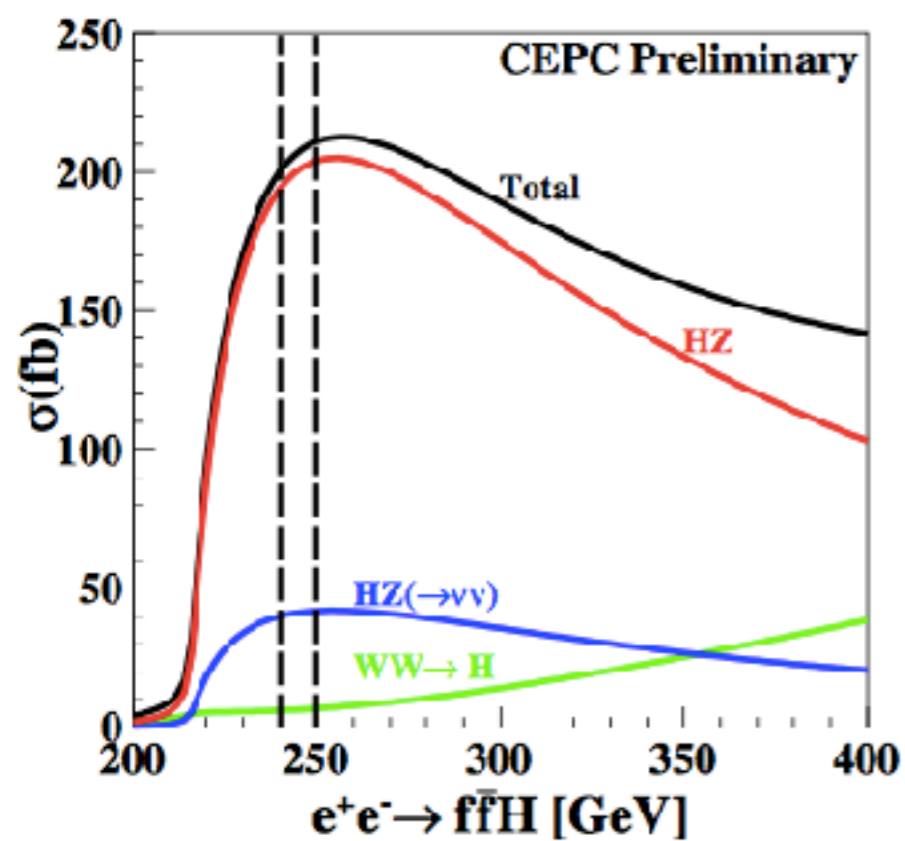


Figure 3.6 Feynman diagrams of the $e^+e^- \rightarrow ZH$, $e^+e^- \rightarrow \nu\bar{\nu}H$ and $e^+e^- \rightarrow e^+e^-H$ processes.



Needs & Obstacles

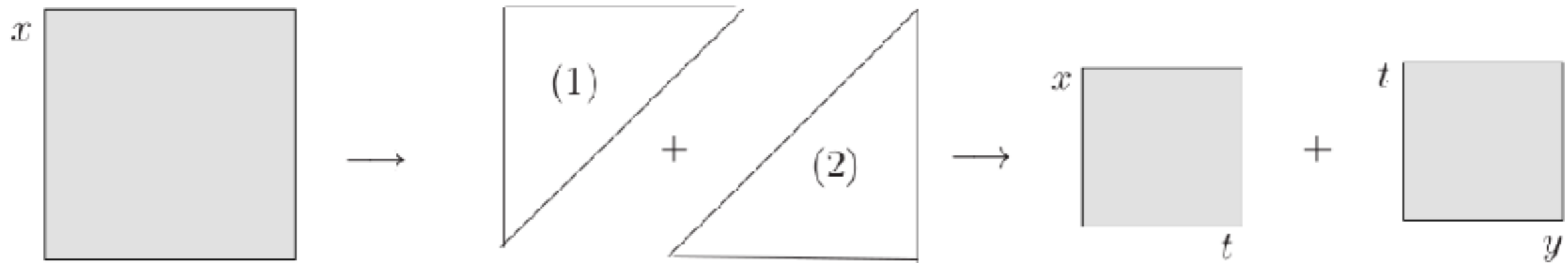
- Higher accuracy of data from LHC, HL-LHC, ILC/CEPC/FCC-ee (high order (SUSY)EW corrections?)
- Multiple scales induce problems on analytical evaluation of higher-loop.
- Analytical expressions can reveal important behaviors, but progress is getting slower. Special math may be behind, but how? when? where?
- By demand from experiments, practically more theoretical predictions can be obtained by numerical approaches.

Numerical approaches for multi-scale multi-loop

- **Mellin-Barnes Representation**
Many tools, faster, difficult on many scales.
- **Sector Decomposition**
More general, slower, okay for many scales.

Sector Decomposition

$$I = \int_0^1 dx \int_0^1 dy x^{-1-\epsilon} y^{-\epsilon} (x + (1-x)y)^{-1}$$



$$I = \int_0^1 dx x^{-1-\epsilon} \int_0^1 dt t^{-\epsilon} (1 + (1-x)t)^{-1}$$

$$+ \int_0^1 dy y^{-1-2\epsilon} \int_0^1 dt t^{-1-\epsilon} (1 + (1-y)t)^{-1}$$

$$G_{l_1 \dots l_R}^{\mu_1 \dots \mu_R} = \int \prod_{l=1}^L d^D \kappa_l \frac{k_{l_1}^{\mu_1} \dots k_{l_R}^{\mu_R}}{\prod_{j=1}^N P_j^{\nu_j}(\{k\}, \{p\}, m_j^2)},$$

$$d^D \kappa_l = \frac{\mu^{4-D}}{i\pi^{\frac{D}{2}}} d^D k_l, \quad P_j(\{k\}, \{p\}, m_j^2) = (q_j^2 - m_j^2 + i\delta),$$

Feynman parameterization

$$\frac{1}{\prod_{j=1}^N P_j^{\nu_j}} = \frac{\Gamma(N_\nu)}{\prod_{j=1}^N \Gamma(\nu_j)} \int_0^\infty \prod_{j=1}^N dx_j x_j^{\nu_j-1} \delta\left(1 - \sum_{i=1}^N x_i\right) \frac{1}{\left[\sum_{j=1}^N x_j P_j\right]^{N_\nu}},$$

where $N_\nu = \sum_{j=1}^N \nu_j$, leads to

$$G_{l_1 \dots l_R}^{\mu_1 \dots \mu_R} = \frac{\Gamma(N_\nu)}{\prod_{j=1}^N \Gamma(\nu_j)} \int_0^\infty \prod_{j=1}^N dx_j x_j^{\nu_j-1} \delta\left(1 - \sum_{i=1}^N x_i\right) \int d^D \kappa_1 \dots d^D \kappa_L$$

$$\times k_{l_1}^{\mu_1} \dots k_{l_R}^{\mu_R} \left[\sum_{i,j=1}^L k_i^T M_{ij} k_j - 2 \sum_{j=1}^L k_j^T \cdot Q_j + J + i\delta \right]^{-N_\nu},$$

Integrate out loop momenta

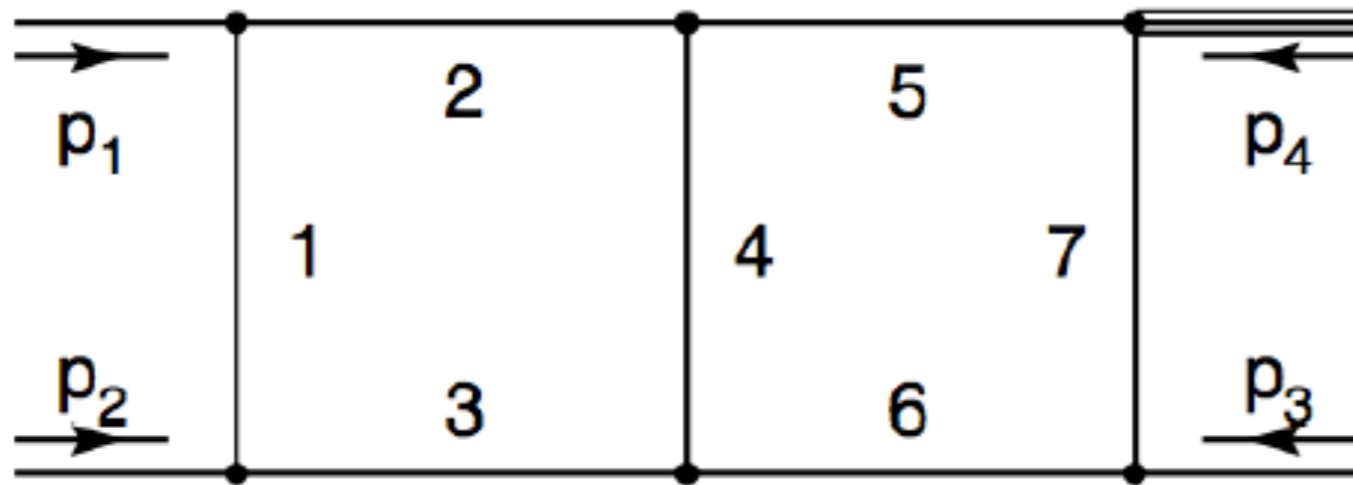
$$\begin{aligned}
G_{l_1 \dots l_R}^{\mu_1 \dots \mu_R} &= (-1)^{N_\nu} \frac{1}{\prod_{j=1}^N \Gamma(\nu_j)} \int_0^\infty \prod_{j=1}^N dx_j x_j^{\nu_j-1} \delta\left(1 - \sum_{l=1}^N x_l\right) \\
&\times \sum_{m=0}^{[R/2]} \left(-\frac{1}{2}\right)^m \Gamma(N_\nu - m - LD/2) [(\tilde{M}^{-1} \otimes g)^{(m)} \tilde{l}^{(R-2m)}]_{\Gamma_1, \dots, \Gamma_R} \\
&\times \frac{\mathcal{U}^{N_\nu - (L+1)D/2 - R}}{\mathcal{F}^{N_\nu - LD/2 - m}}, \tag{7}
\end{aligned}$$

where

$$\mathcal{F}(\mathbf{x}) = \det(M) \left[\sum_{j,l=1}^L Q_j M_{jl}^{-1} Q_l - J - i\delta \right], \tag{8}$$

$$\mathcal{U}(\mathbf{x}) = \det(M), \quad \tilde{M}^{-1} = \mathcal{U} M^{-1}, \quad \tilde{l} = \mathcal{U} v$$

U and F can
be determined
geometrically



$$\mathcal{U}(\mathbf{x}) = \sum_{T \in \mathcal{T}_1} \left[\prod_{j \in \mathcal{C}(T)} x_j \right],$$

$$\mathcal{F}_0(\mathbf{x}) = \sum_{\hat{T} \in \mathcal{T}_2} \left[\prod_{j \in \mathcal{C}(\hat{T})} x_j \right] (-s_{\hat{T}}),$$

$$\mathcal{F}(\mathbf{x}) = \mathcal{F}_0(\mathbf{x}) + \mathcal{U}(\mathbf{x}) \sum_{j=1}^N x_j m_j^2.$$

$$\mathcal{U} = x_{123}x_{567} + x_4x_{123567},$$

$$\begin{aligned} \mathcal{F} = & (-s_{12})(x_2x_3x_{4567} + x_5x_6x_{1234} + x_2x_4x_6 + x_3x_4x_5) \\ & + (-s_{23})x_1x_4x_7 + (-p_4^2)x_7(x_2x_4 + x_5x_{1234}), \end{aligned}$$

where $x_{iik\dots} = x_i + x_i + x_k + \dots$ and $s_{ij} = (p_i + p_j)^2$.

First generate primary sectors to eliminate Delta function

$$\int_0^\infty d^N x = \sum_{l=1}^N \int_0^\infty d^N x \prod_{\substack{j=1 \\ j \neq l}}^N \theta(x_l \geq x_j).$$

$$x_j = \begin{cases} x_l t_j & \text{for } j < l, \\ x_l & \text{for } j = l, \\ x_l t_{j-1} & \text{for } j > l \end{cases}$$

$$G_l = \int_0^1 \prod_{j=1}^{N-1} dt_j \frac{\mathcal{U}_l^{N_\nu - (L+1)D/2}(\mathbf{t})}{\mathcal{F}_l^{N_\nu - LD/2}(\mathbf{t})}, \quad l = 1, \dots, N.$$

Determine a sub-set of parameters t_i

$$\mathcal{S} = \{t_{\alpha_1}, \dots, t_{\alpha_r}\}$$

Then divide into r sub-sectors

$$\prod_{j=1}^r \theta(1 \geq t_{\alpha_j} \geq 0) = \sum_{k=1}^r \prod_{\substack{j=1 \\ j \neq k}}^r \theta(t_{\alpha_k} \geq t_{\alpha_j} \geq 0).$$

$$t_{\alpha_j} \rightarrow \begin{cases} t_{\alpha_k} t_{\alpha_j} & \text{for } j \neq k, \\ t_{\alpha_k} & \text{for } j = k. \end{cases}$$

$$G_{lk} = \int_0^1 \left(\prod_{j=1}^{N-1} dt_j t_j^{a_j - b_j \epsilon} \right) \frac{\mathcal{U}_{lk}^{N_\nu - (L+1)D/2}}{\mathcal{F}_{lk}^{N_\nu - LD/2}}, \quad k = 1, \dots, r.$$

$$\mathcal{U}_{lk_1 k_2 \dots} = 1 + u(\mathbf{t}), \quad \mathcal{F}_{lk_1 k_2 \dots} = -s_0 + \sum_{\beta} (-s_{\beta}) f_{\beta}(\mathbf{t}),$$

All the coefficients of divergences are finite (complicated).

Decomposition strategies

- **Hironaka's polyhedra game**

Bogner and Weinzierl, Comput.Phys.Commun. 178 (2008) 596; A. V. Smirnov and V. A. Smirnov, JHEP 05 (2009) 004;

- **Geometric method**

Kaneko and Ueda, Comput.Phys.Commun. 181 (2010) 1352

Iteration of certain strategy will show explicitly dimensional regulators, where poles can be extracted.

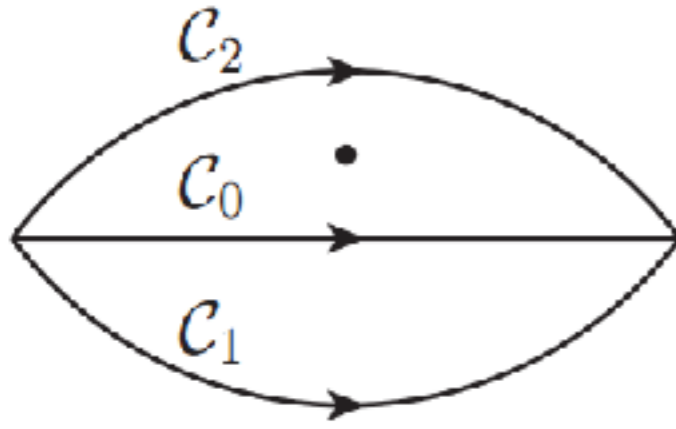
$$I_j = \int_0^1 dt_j t_j^{(a_j - b_j \epsilon)} \mathcal{I}(t_j, \{t_{i \neq j}\}, \epsilon) ,$$

$$I_j = \sum_{p=0}^{|a_j|-1} \frac{1}{a_j + p + 1 - b_j \epsilon} \frac{\mathcal{I}_j^{(p)}(0, \{t_{i \neq j}\}, \epsilon)}{p!} + \int_0^1 dt_j t_j^{a_j - b_j \epsilon} R(\vec{t}, \epsilon) .$$

$$I_j = -\frac{1}{b_j \epsilon} \mathcal{I}_j(0, \{t_{i \neq j}\}, \epsilon) + \int_0^1 dt_j t_j^{-1 - b_j \epsilon} \left(\mathcal{I}(t_j, \{t_{i \neq j}\}, \epsilon) - \mathcal{I}_j(0, \{t_{i \neq j}\}, \epsilon) \right) ,$$

Contour Deformation

$$I_s = C(\epsilon) \lim_{\delta \rightarrow 0} \int_0^1 \frac{\mathcal{D}(\vec{x}, \epsilon) \mathcal{H}_s(\vec{x}, \epsilon)}{[\mathcal{F}_s(\vec{x}, m_i^2, s_{jk}) - i\delta]^{a+b\epsilon}}$$



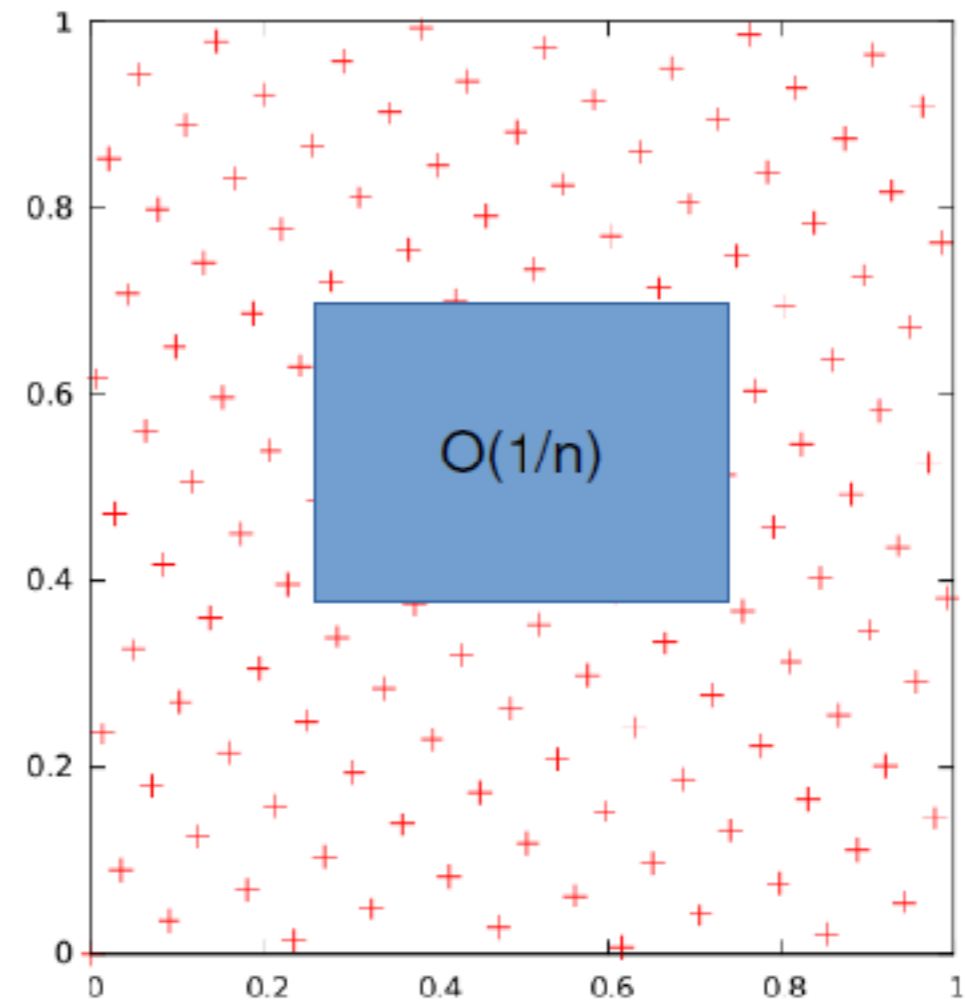
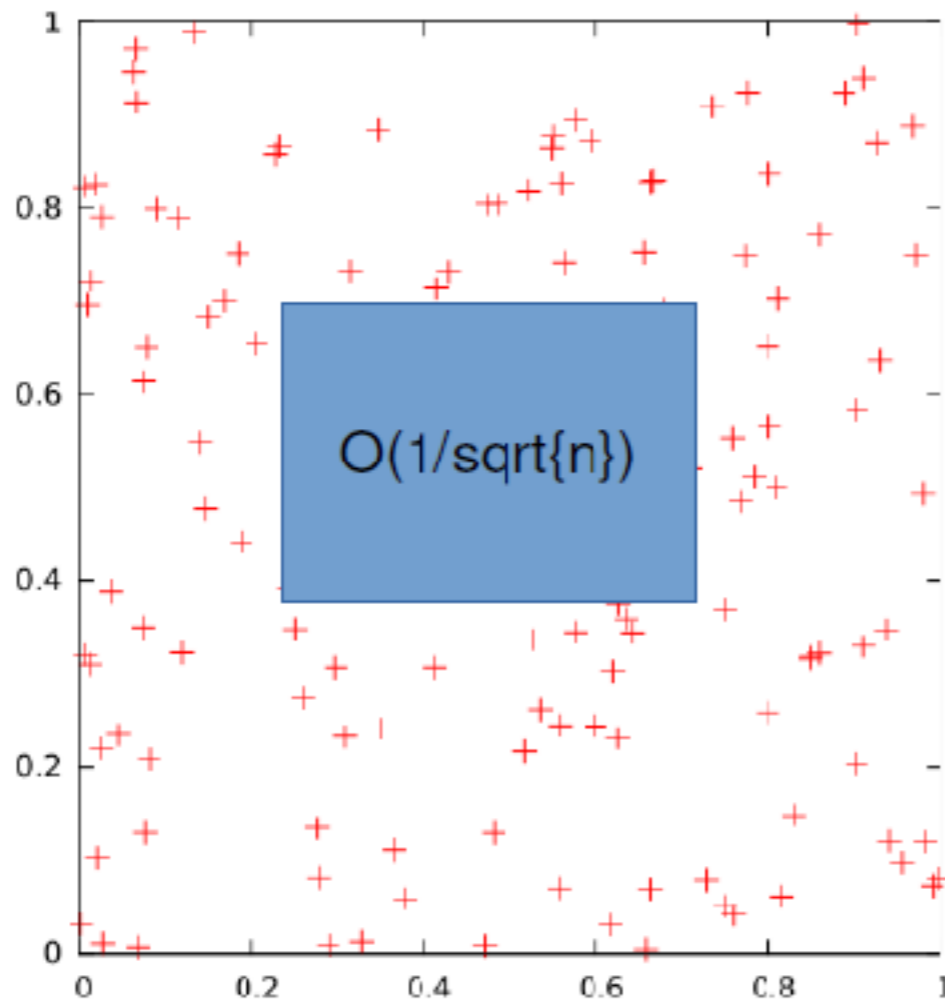
$$z_i = x_i - i\lambda x_i^\alpha (1 - x_i)^\beta \frac{\partial \mathcal{F}_s}{\partial x_i}$$

$$\lim_{\delta \rightarrow 0} \int_0^1 \frac{\mathcal{D}(\vec{x}, \epsilon) \mathcal{H}_s(\vec{x}, \epsilon)}{[\mathcal{F}_s(\vec{x}, m_i^2, s_{jk}) - i\delta]^{a+b\epsilon}} = \int_C \frac{\mathcal{D}(\vec{z}, \epsilon) \mathcal{H}_s(\vec{z}, \epsilon)}{[\mathcal{F}_s(\vec{z}, m_i^2, s_{jk})]^{a+b\epsilon}}$$

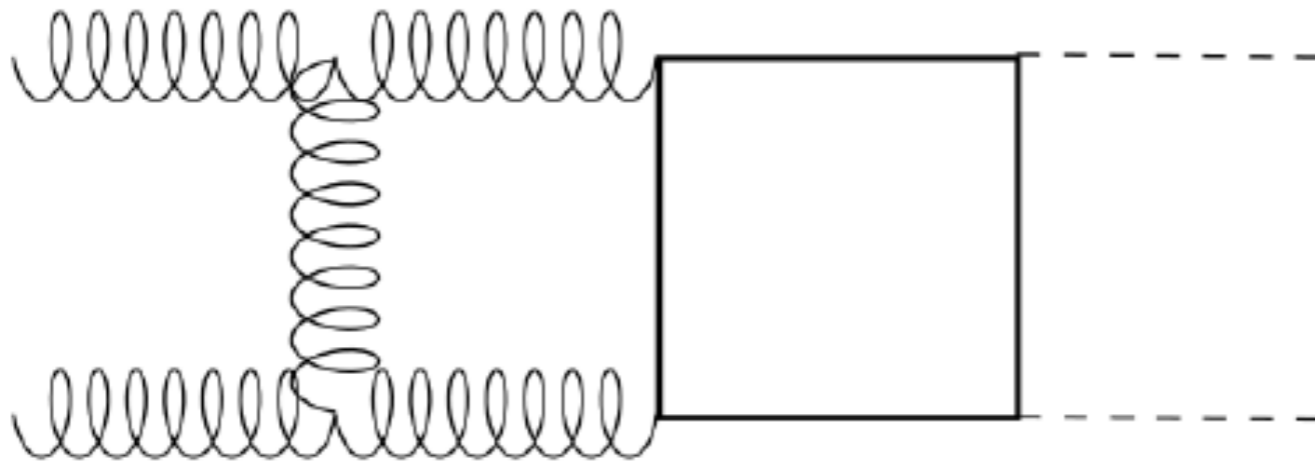
Improve via quasi-Monte-Carlo

$$I(f) = \int_0^1 d^s x f(\vec{x})$$

$$I_{estimate}(f) = \sum_{i=0}^{n-1} f(\vec{x}_i)$$



Implementation on GPU

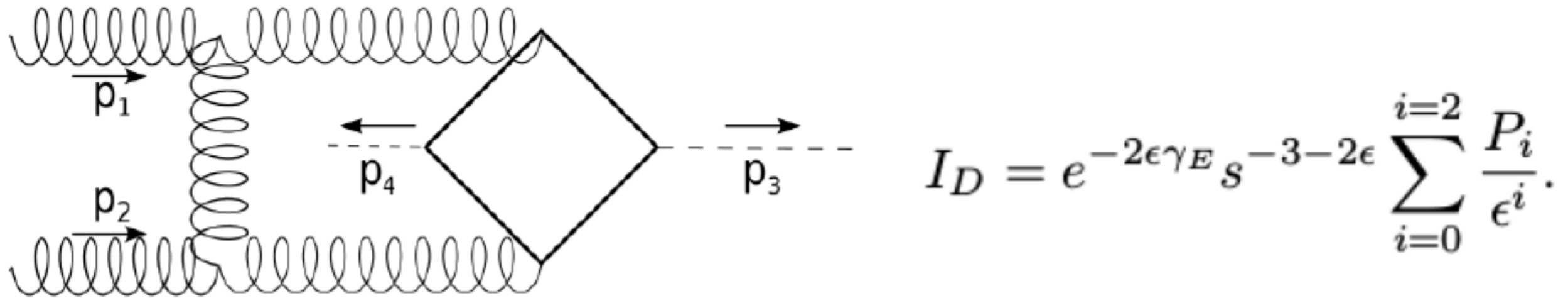


$$I_C = e^{-2\epsilon\gamma_E} s^{-3-2\epsilon} \sum_{i=0}^{i=2} \frac{P_i}{\epsilon^i}.$$

	Vegas/CPU	QMC/GPU
P_2	$-7.959 \pm 0.009 - 10.586i \pm 0.009i$	$-7.949 \pm 0.003 - 10.585i \pm 0.005i$
P_1	$3.9 \pm 0.1 - 28.1i \pm 0.1i$	$3.831 \pm 0.005 - 28.022i \pm 0.005i$
P_0	$-3.9 \pm 0.8 + 92.3i \pm 0.8i$	$-4.63 \pm 0.07 + 92.13i \pm 0.07i$
Integration Time	45540s	19s

Z. Li et al., Chin.Phys. C40 (2016) no.3, 033103

Implementation on GPU



	Vegas/CPU	QMC/GPU
P_2	$-3.848 \pm 0.004 + 0.0005i \pm 0.003i$	$-3.8482 \pm 0.0007 + 0.0004i \pm 0.0003i$
P_1	$3.81 \pm 0.03 - 6.41i \pm 0.03i$	$3.83 \pm 0.02 - 6.40i \pm 0.02i$
P_0	$77.2 \pm 0.2 + 20.1i \pm 0.2i$	$77.2 \pm 0.1 + 19.9i \pm 0.1i$
Integration Time	54290s	20s

Z. Li et al., Chin.Phys. C40 (2016) no.3, 033103

Practical Process calculations

- Successful example: Borowka et al. JHEP 1610 (2016) 107
NLO Higgs pair with finite m_t
- Option #1: IBP reduction?
- Option #2: GPU or MPI?
- Option #3: Best choice of parameters of contour deformation?
- Option #4: Best lattice of QMC?
-

Mixed QCD-EW corrections for Higgs boson production at e^+e^- colliders

\sqrt{s} (GeV)	σ_{LO} (fb)	σ_{NLO} (fb)	σ_{NNLO} (fb)	$\sigma_{\text{NNLO}}^{\text{exp.}}$ (fb)
240	256.3(9)	228.0(1)	230.9(4)	230.9(4)
250	256.3(9)	227.3(1)	230.2(4)	230.2(4)
300	193.4(7)	170.2(1)	172.4(3)	172.4(3)
350	138.2(5)	122.1(1)	123.9(2)	123.6(2)
500	61.38(22)	53.86(2)	54.24(7)	54.64(10)

TABLE I. The NNLO predictions for the total cross sections at various collider energies.

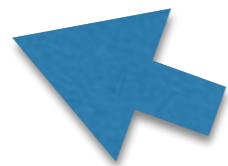
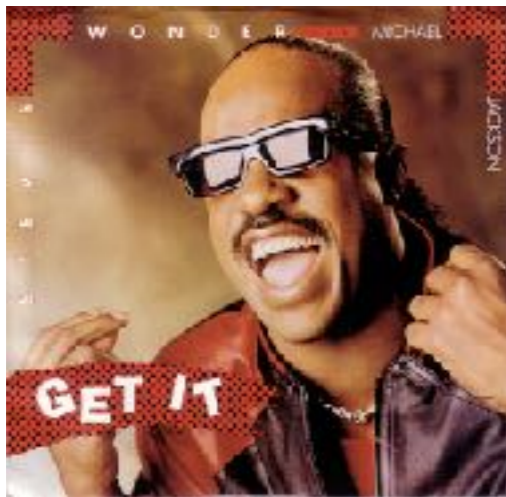
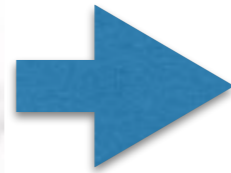
Y.Gong et al., Phys.Rev. D95 (2017) no.9, 093003

Toward automatic NNLO calculations (MIRACLE)

- Feynman diagram generation
- Automatic amplitude manipulation
- Feynman parameterization
- Sector Decomposition
- Pole extraction
- Classification of sectors
- Export expressions to numerical code
- Deploy to cluster -> Money/Funding

FORM, Mathematica, Fermat, GiNaC, C++, CUDA/MPI+AVX,
QGRAF, FIESTA, latBuilder, WPO

Goal of MIRACLE



Prospect

- Numerical approach can give numbers for multi-scale multi-loop processes.
- Hard to identify physics structures, eg. large logs.
- Many precise predictions can be expected in a few years, Higgs/single-top/ttH/dijet/new-physics?.
- Still hope breakthrough on analytical approach.



Thank you!