Halo scale probes of Cosmology-Splashback and halo profiles in Modified gravity and SIDM

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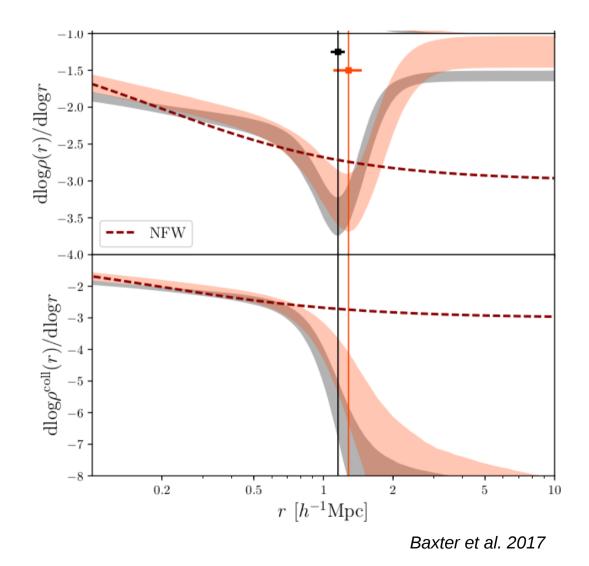
Small scale probes and simulations

- Signatures of modified gravity
- Interactions of dark matter Isotropic and Anisotropic interactions

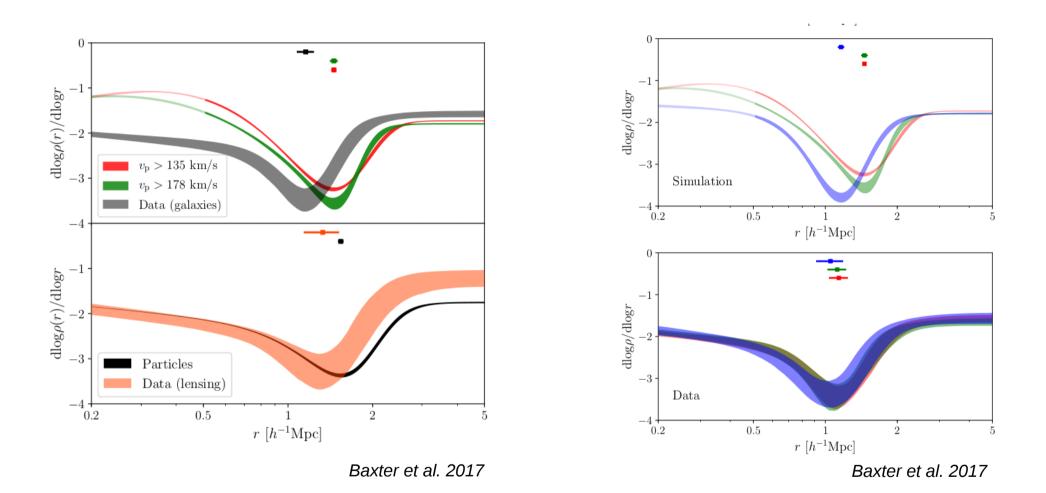
Probes :

- Halo/Void profiles
- Splashback radius
- Shapes
- Velocity profiles
- Galaxy morphology

Splashback radius - current status of measurements (DES)



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- Measured in galaxy profiles around RedMaPPer clusters and in lensing
- Differences with simulations in ΛCDM Smaller than predicted and no movement with galaxy magnitude

Splashback radius in Modified Gravity (with Jeremy Sakstein, Bhuvnesh Jain, Neal Dalal, Baojiu Li)

• Does the location of splashback radius change in modified gravity?

Splashback in nDGP models :

Extra scalar degree of freedom :

$$\nabla^2 \pi + \frac{r_c^2}{3} \left[(\nabla^2 \pi)^2 - \nabla_i \nabla_j \pi \nabla^i \nabla^j \pi \right] = 8\pi \alpha G \rho$$

$$F_{\pi} = -\alpha\pi' = -4\alpha^2 \frac{GM}{r^2} g\left(\frac{r}{r_*}\right)$$

$$g(\zeta) = \zeta^3 \left(\sqrt{1 + \zeta^{-3}} - 1 \right)$$

- Splashback radius is a scale inside the halo that can be predicted from simple Newtonian dynamics
- Orbit of an infalling shell of matter in an accreting halo in an expanding universe
- Acceleration modified in nDGP by extra force:

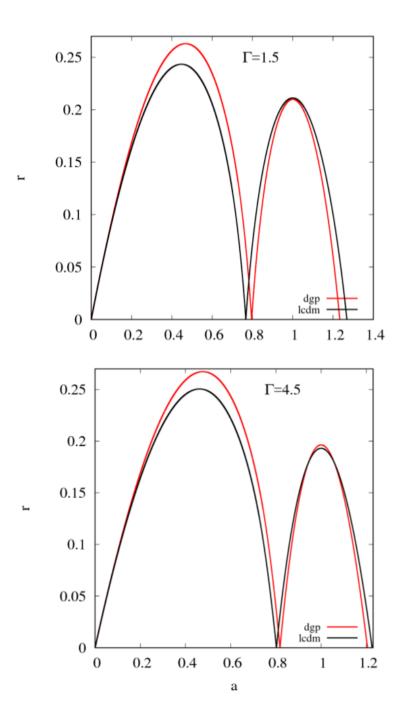
$$\ddot{r} = -\frac{GM}{r^2} + \frac{\Lambda r}{3} + F_{\pi}$$

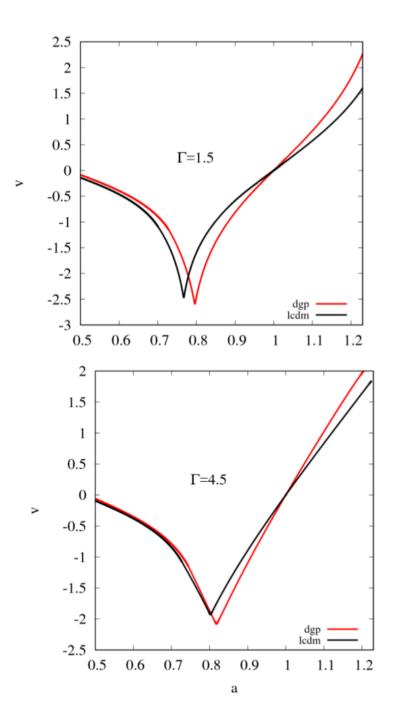
$$M_{tot} = a^s$$
 $M = M_{tot} \frac{f_{\rm NFW}(r/r_s)}{f_{\rm NFW}(R/r_s)}$

$$F_{\pi} = -\alpha\pi' = -4\alpha^2 \frac{GM}{r^2} g\left(\frac{r}{r_*}\right)$$

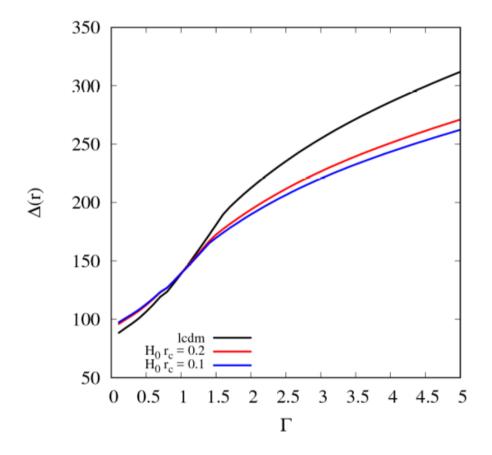
• Shells are initialized from half the turnaround radius from the spherical collapse model.

Orbits of shells in nDGP

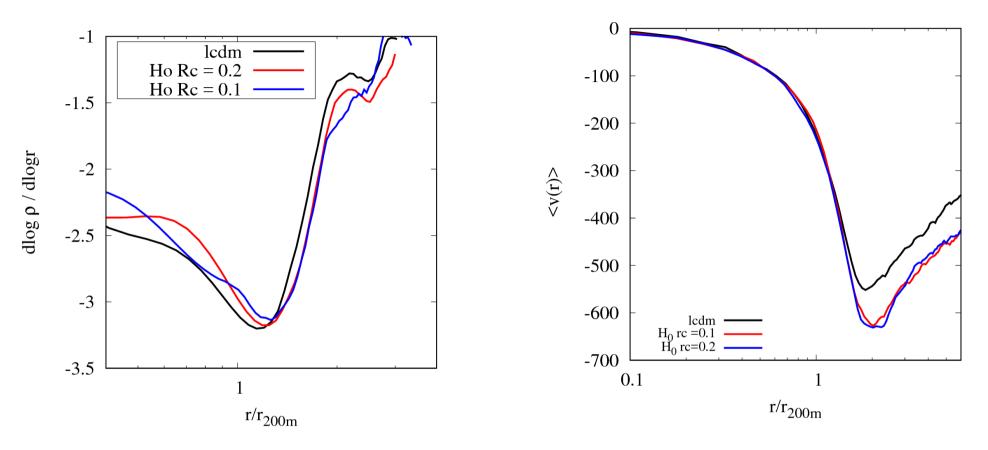




Enclosed density at splashback as a function of accretion rate of the halo

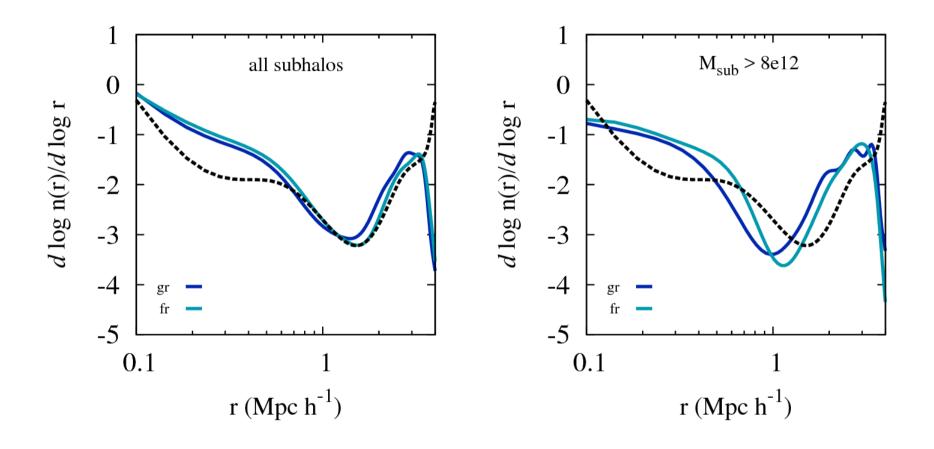


Results from simulations



Splashback radius and velocity profile of halos stacked in mass bins of 1-4e14 M_{sun} h⁻¹ ECOSMOG – Efficient code for simulating modified Gravity (*Baojiu Li et al. 2011*) 400 Mpc h⁻¹ box with 512³ particles

Subhalos in f(R) simulations with $f_R = 10^{-5}$



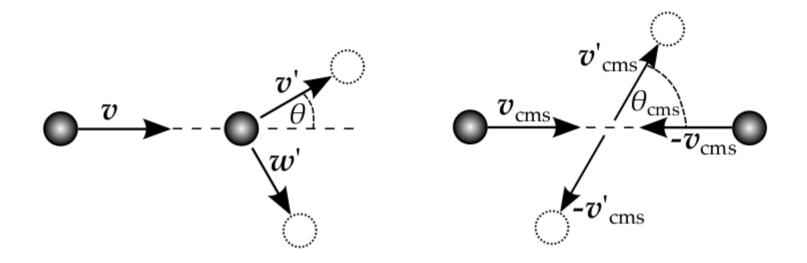
- 1024³ particles, 1Gpc h⁻¹ box
- Massive subhalos splashback at different locations in GR and f(R)
- Differences in strength of dynamical friction?

Further work

- Run larger simulation boxes to get better statistics for massive clusters.
- Follow the dynamics of subhalos to check effects of dynamical friction.
- Extend analytical calculations to other models, eg. beyond Horndeski.
- Cosmological simulations are available for f(R) and DGP mostly.

Dark Matter Self-Interactions

(with Arka Banerjee, Neal Dalal, Suhrud More, Andrey Kravtsov)



Possible signatures:

- i) Evaporation of subhalos
- ii) Displacement of lensing profile and light profile
- iii) Halo profiles (Splashback, velocity structure)
- iv) Warping of disks (NEW)

Anisoropic scattering vs. Isotropic scattering

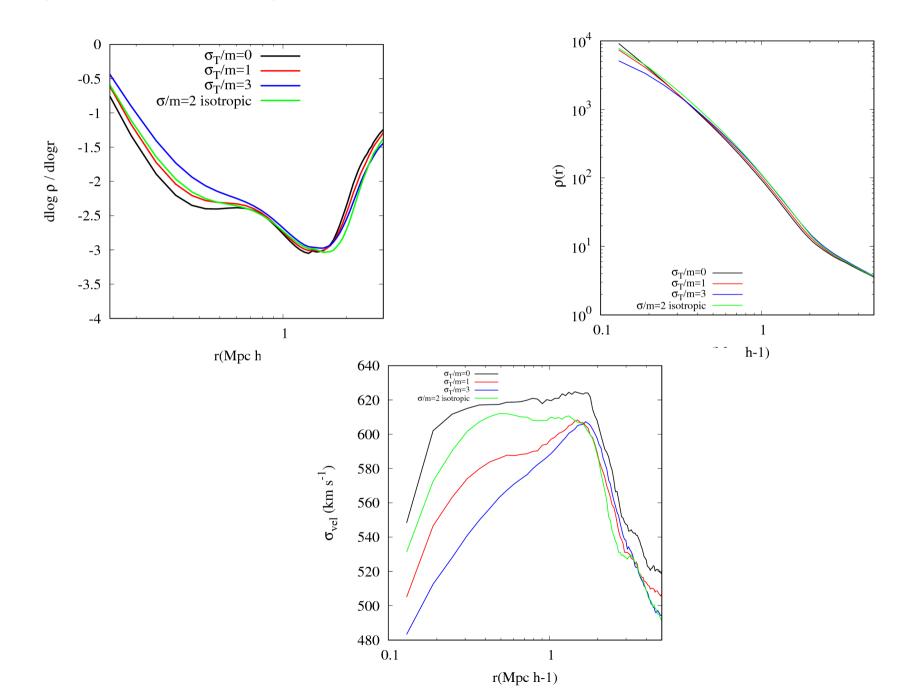
- Scattering probability is high only at small angles
- Small momentum transfer at each interaction

• Net drag force
$$\frac{F_{\rm drag}}{m_{\rm DM}} = \frac{\tilde{\sigma}}{4 \, m_{\rm DM}} \rho \, v_0^{2m}$$

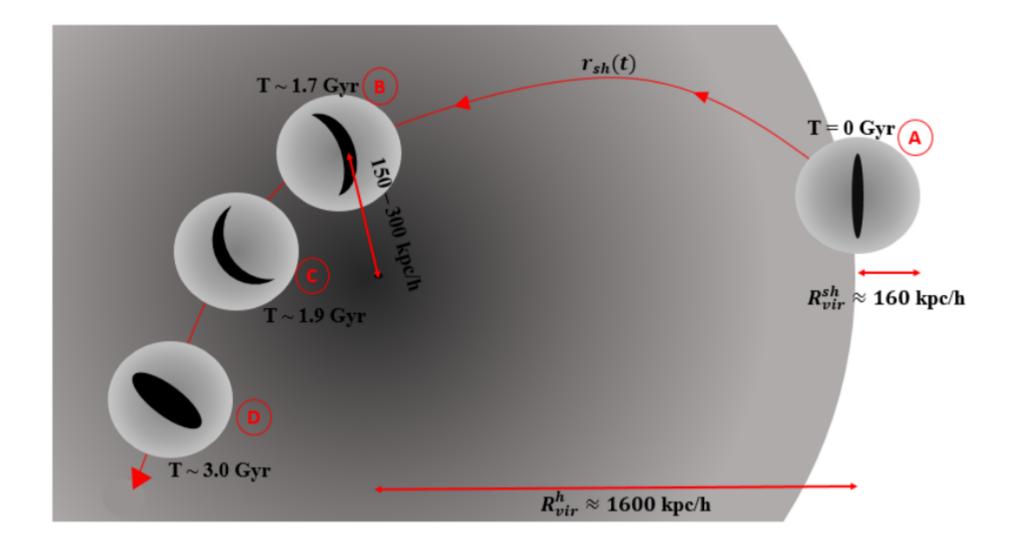
- Cumulative evaporation still important but not as strong as isotropic scattering
- Cosmological simulations are available for Isotropic scattering cross-sections
- Cosmological simulations for anisotropic scattering (work in progress)

Cosmological simulations of anisotropic SIDM

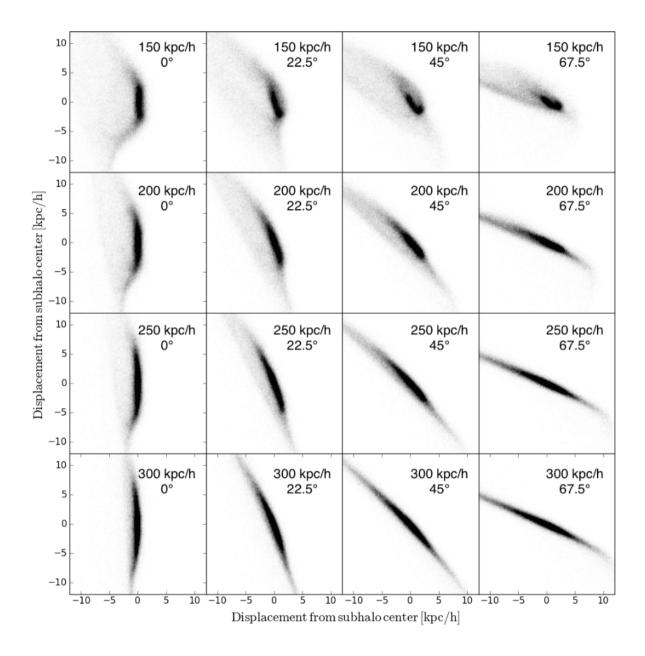
1Gpc h⁻¹ box with 1024³ particles



Warping of Disks in anisotropic SIDM (lead by Lucas Secco)



Edge on galaxies falling in at different impact parameters



Summarizing...

• Simulating Velocity dependent anisotropic cross-sections.

1/v² dependence of drag force, becomes important in lower mass halos than clusters

- Zoom in simulations of subhalo+galaxy infalling into clusters to measure displacements.
- Effect of gas on disk morphologies (Illustris+SIDM)
- Orbital kinematics of subhalos and particles in SIDM.