



Search for the resonant production of two Higgs bosons in the $\gamma\gamma b\bar{b}$ final state

CMS-PAS-HIG-13-032

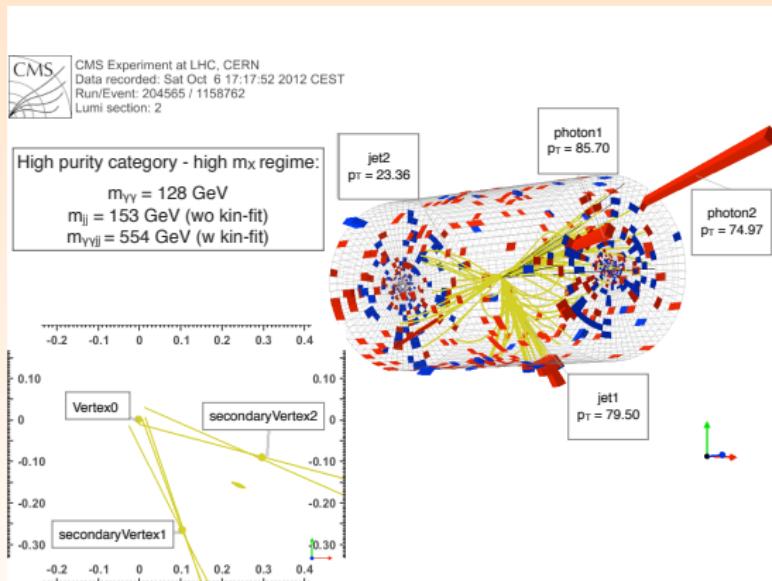
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14 November 2014

Outline

- Overview of the process:
 - ➊ Motivation from theory
 - ➋ Production and decay
- Analysis
 - ➊ Signal extraction
 - ➋ Results



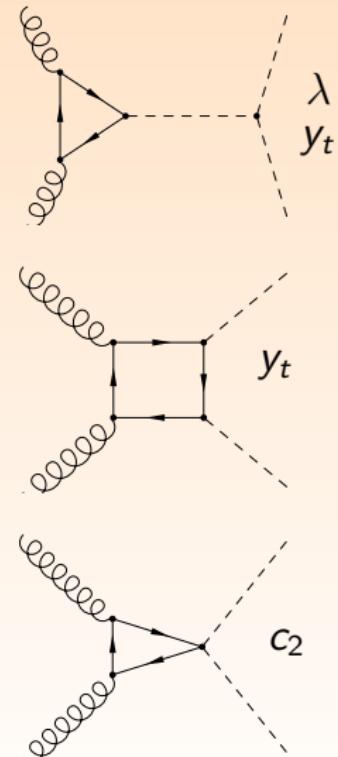
Theoretical motivation

HH production in SM

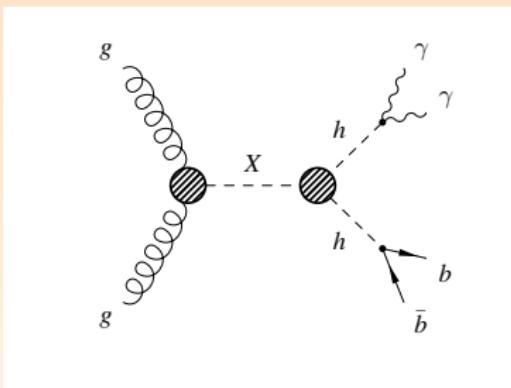
- $\sigma(pp \rightarrow HH) \approx 10 \text{ fb}$ at $\sqrt{s} = 8 \text{ TeV}$.
- Lots of lumi is needed (cf. HL-LHC).

HH production in beyond SM scenarios

- Nonresonant HH can be enhanced by altering the couplings of the Higgs.
 - ▶ This models the presence of new particles.
 - ▶ There are interference effects.
- λ , y_t , and c_2 can vary in $gg \rightarrow HH$.



Theoretical motivation



HH production in beyond SM scenarios

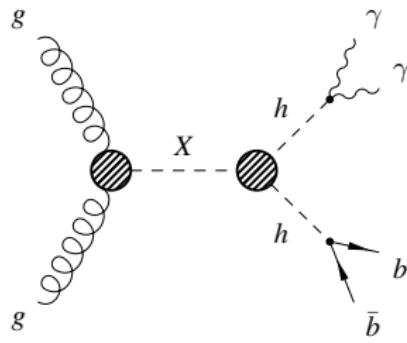
- Resonant HH can be motivated by WED, 2HDM, and (N)MSSM.
 - ▶ Radion (spin 0) or Graviton (spin 2) from perturbations of the ED
 - ▶ Higgs-like scalar from extended Higgs sector in SUSY
- Resonance can decay to HH if it is kinematically allowed.

Production and decay

Enhanced HH resonance

- We search for a heavy resonance (X) produced through gluon fusion.
- The object decays to a pair of SM Higgs, which decay to a pair of photons (high resolution) and a pair of b quarks (high BR).
- This search concerns the non-boosted regime, m_X from 260 to 1100 GeV.

Channel	Freq. (%)
$H(b\bar{b}, c\bar{c}, gg)H(b\bar{b}, c\bar{c}, gg)$	47.86
$H(b\bar{b})H(b\bar{b})$	33.30
$H(b\bar{b}, c\bar{c}, gg)H(VV^*)$	33.40
$H(b\bar{b}, c\bar{c}, gg)H(\tau^+\tau^-)$	8.77
$H(b\bar{b})H(\tau^+\tau^-)$	7.29
$H(VV^*)H(VV^*)$	5.83
$H(I^+I^-)H(VV^*)$	3.06
$H(\tau^+\tau^-)H(\tau^+\tau^-)$	0.40
$H(b\bar{b}, c\bar{c}, gg)H(\gamma\gamma)$	0.32
$H(b\bar{b})H(\gamma\gamma)$	0.26
$H(b\bar{b}, c\bar{c}, gg)H(\mu^+\mu^-)$	0.03
$H(I^+I^-)H(\gamma\gamma)$	0.03



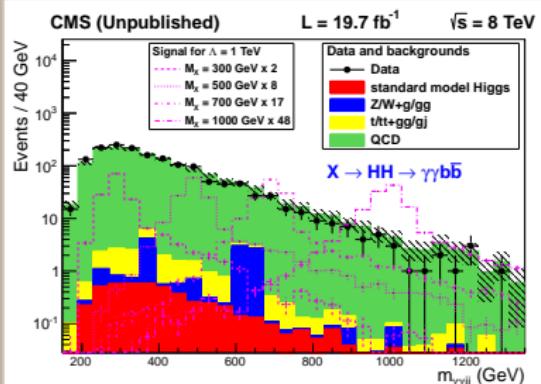
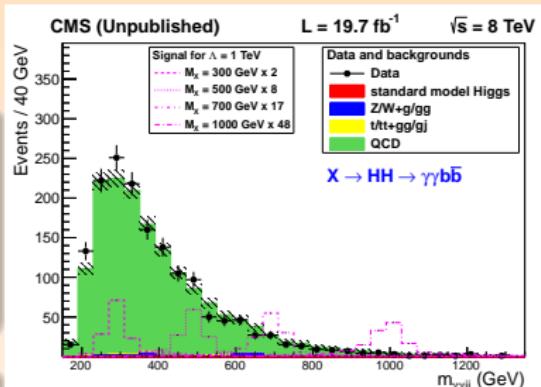
Analysis strategy

Three mass spectra

- A signal will appear as an excess in $M_{\gamma\gamma}$, M_{jj} , and $M_{\gamma\gamma jj}$ spectra.
- Analysis was blinded during design.

High-mass regime

- $M_{\gamma\gamma jj}$ gives a direct handle on m_X . Due to the background peaking around 300 GeV, this spectrum is only used for high mass hypotheses.
- The **high-mass regime** examines the shape of $M_{\gamma\gamma jj}$ to test $m_X \geq 400$ GeV.
- Mass windows are applied to $M_{\gamma\gamma}$ and M_{jj} .



Data/MC comparison after preselection,
event-normalized

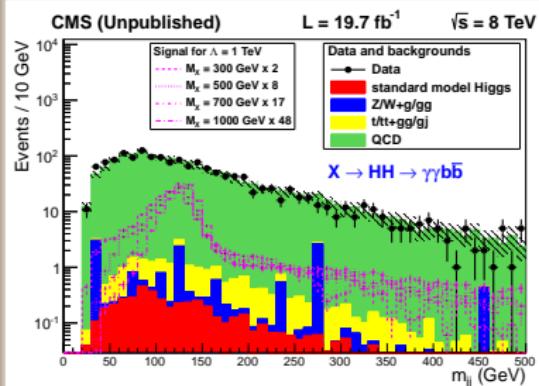
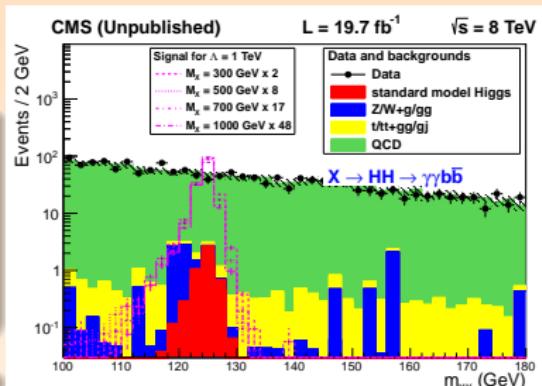
Analysis strategy

Three mass spectra

- A signal will appear as an excess in $M_{\gamma\gamma}$, M_{jj} , and $M_{\gamma\gamma jj}$ spectra.
- Analysis was blinded during design.

Low-mass regime

- To avoid fitting a signal peak on the background peak, $M_{\gamma\gamma jj}$ is not used for fitting m_X hypotheses close to 300 GeV.
- The **low-mass regime** examines the shape of $M_{\gamma\gamma}$ to test $m_X \leq 400$ GeV.
- Mass windows are applied to $M_{\gamma\gamma jj}$ and M_{jj} .

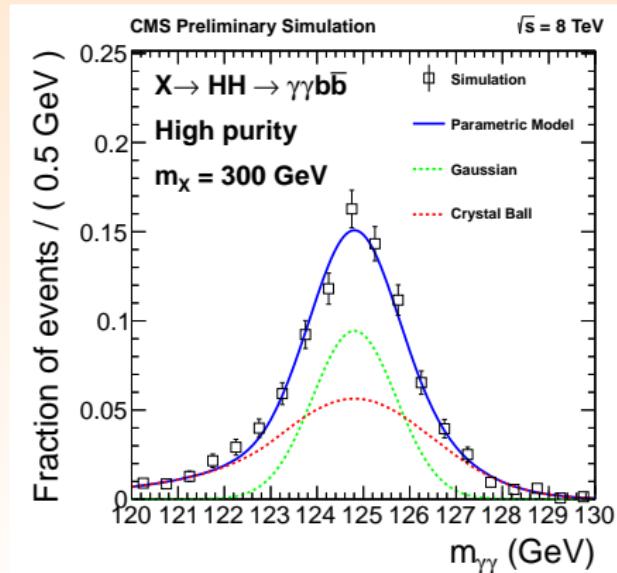
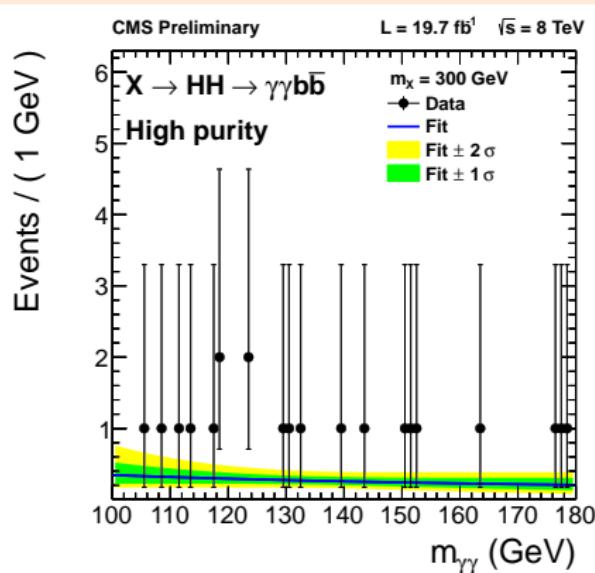


Data/MC comparison after preselection,
event-normalized

Signal extraction

Low mass ($m_X \leq 400$ GeV) modeling

Fit to $M_{\gamma\gamma}$, cut on $M_{\gamma\gamma jj}$ and M_{jj}

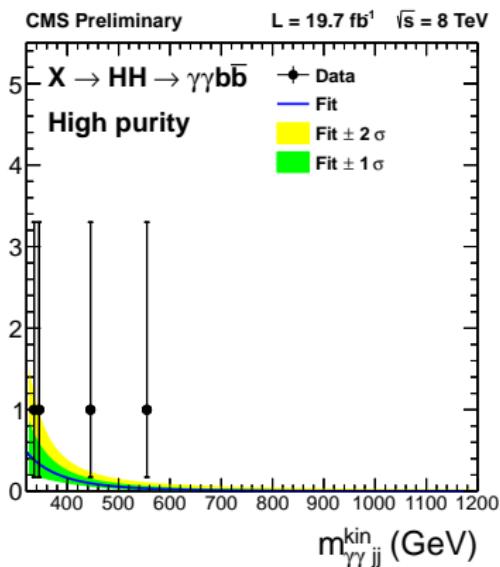


Signal extraction

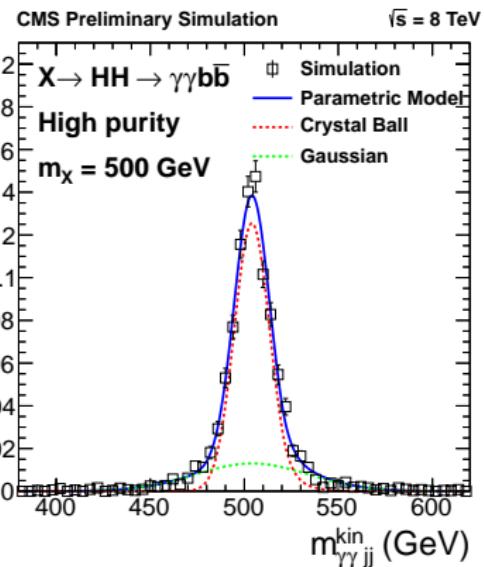
High mass ($m_X \geq 400$ GeV) modeling

Fit to $M_{\gamma\gamma jj}$, cut on $M_{\gamma\gamma}$ and M_{jj}

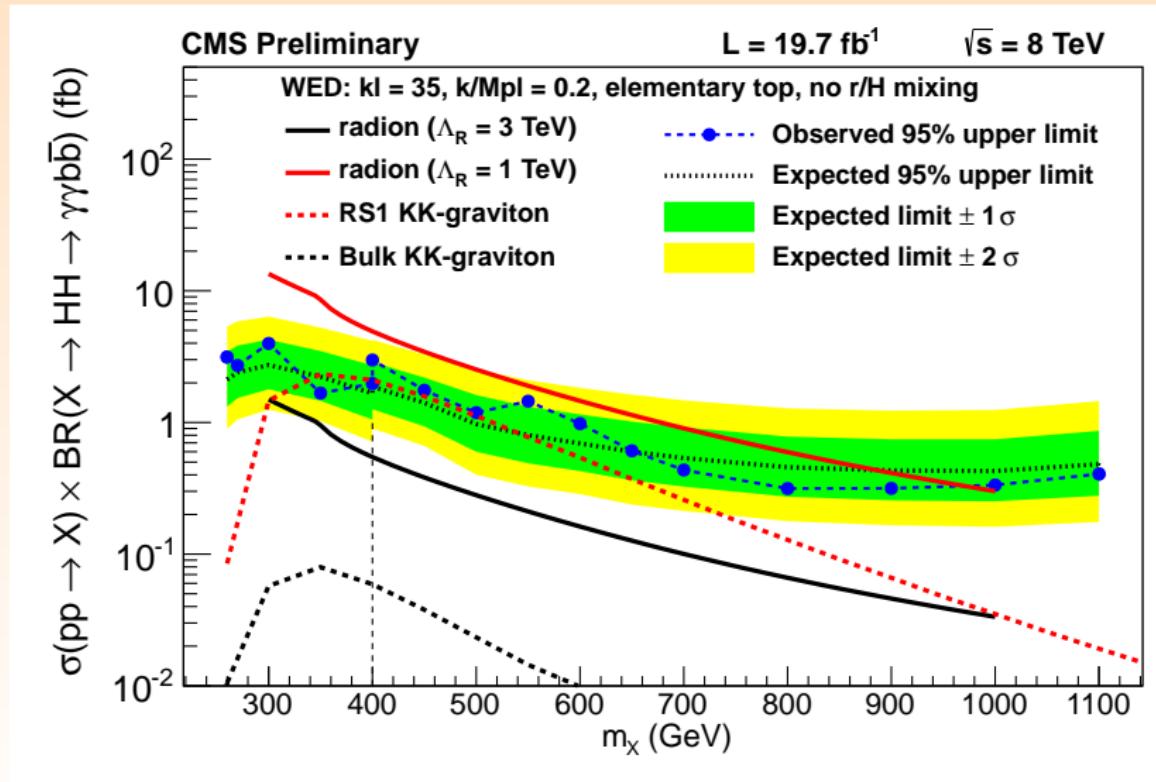
Events / (10 GeV)



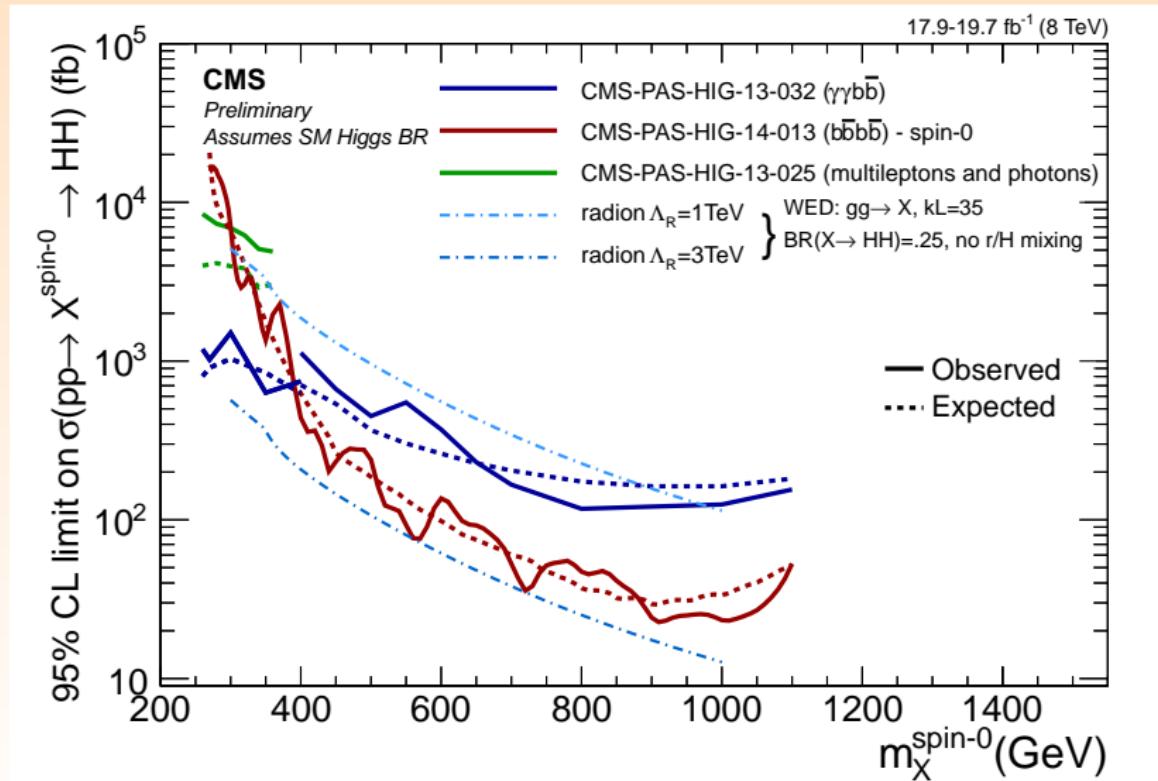
Fraction of events / (4 GeV)



Exclusions



Comparison with other CMS HH results



Conclusions

Results

- $X \rightarrow HH \rightarrow \gamma\gamma b\bar{b}$ for $260 \text{ GeV} < m_X < 1100 \text{ GeV}$
- Analysis is statistically-limited.
- SM Higgs contribution is negligible.
- Analysis is spin-independent.

Outlook

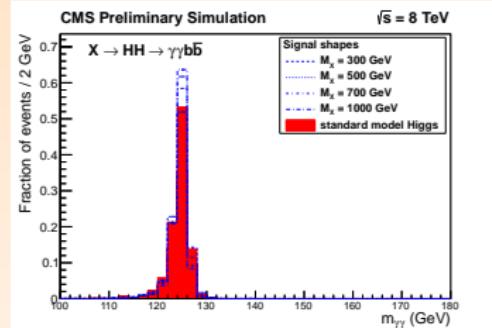
- Look for SM HH analysis to come out soon!
- This will include interpretation for nonresonant models with anomalous couplings (λ, y_t, c_2).

Backup

Objects

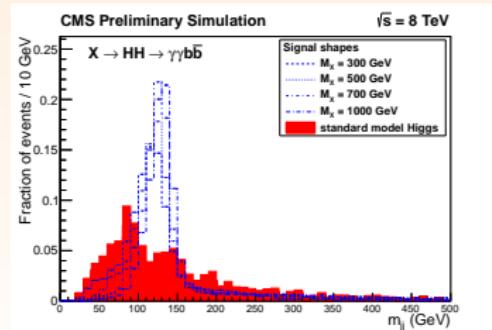
Starting from SM $H \rightarrow \gamma\gamma$ analysis

- Double photon triggers
- Photon preselection cuts
 - ▶ $p_T^{\text{lead}} > \frac{M_{\gamma\gamma}}{3}$
 - ▶ $p_T^{\text{sublead}} > \frac{M_{\gamma\gamma}}{4}$
 - ▶ $|\eta| < 2.5$
 - ▶ $100 \text{ GeV} < M_{\gamma\gamma} < 180 \text{ GeV}$
- Cut-based photon ID based on shower shape and isolation

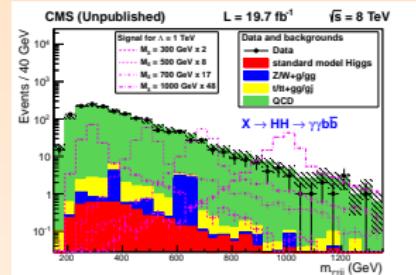
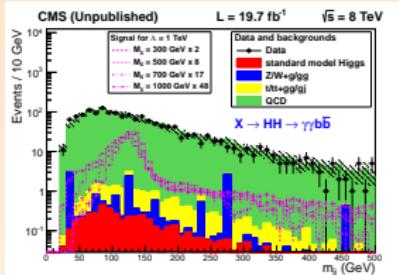
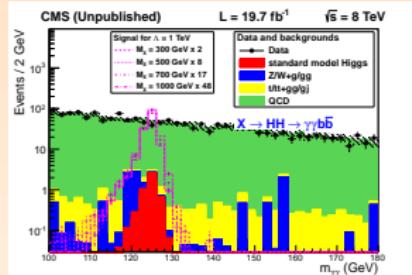


And adding jets

- Anti- k_T jets with $R=0.5$
- Jet preselection cuts
 - ▶ $p_T > 25 \text{ GeV}$
 - ▶ $|\eta| < 2.5$
 - ▶ Loose/pileup jet ID



To the signal region



b-tagging and mass windows

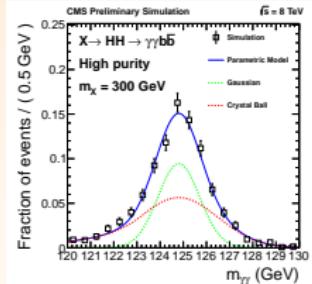
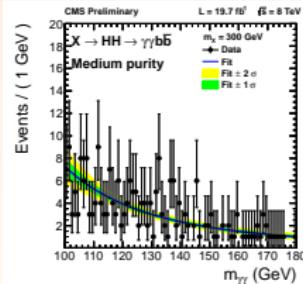
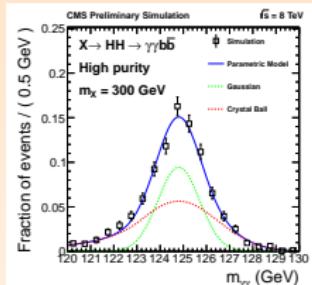
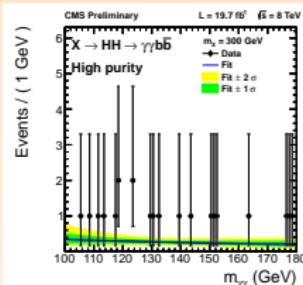
- Displaced tracks allow for identification of jets coming from b quarks.
- The tagger used has 70% signal efficiency and 10% background acceptance.
 - ▶ ≥ 2 b-tagged jets is **high purity**
 - ▶ $= 1$ b-tagged jet is **medium purity** and recovers signal efficiency
- Finally, mass windows are applied to the two spectra not being fit.

Signal extraction

Low mass ($m_X \leq 400$ GeV) modeling

Modeling choice

- Background is modeled by a power law in both categories.
 - ▶ Mass windows are placed on M_{jj} and $M_{\gamma\gamma jj}$, changing background shape for different mass hypotheses.
- Signal is modeled by a sum of Crystal Ball and Gaussian.

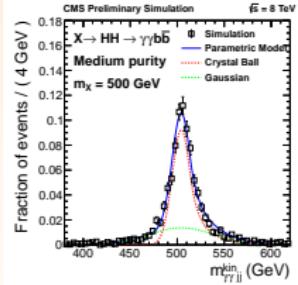
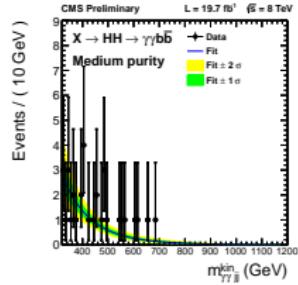
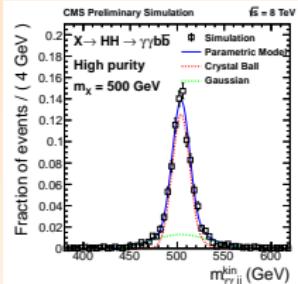
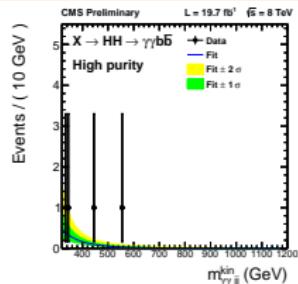


Signal extraction

High mass ($m_X \geq 400$ GeV) modeling

Modeling choice

- Background model depends on category.
 - ▶ For high-purity category:
$$\frac{1}{(M_{\gamma\gamma jj})^{2\alpha}}$$
 - ▶ For medium-purity category:
$$\frac{1}{(M_{\gamma\gamma jj}^2 + b)^{\alpha}}$$
- Signal is modeled by a sum of Crystal Ball and Gaussian.



Signal efficiency

