

Effect of materials defects on the physics of SRF

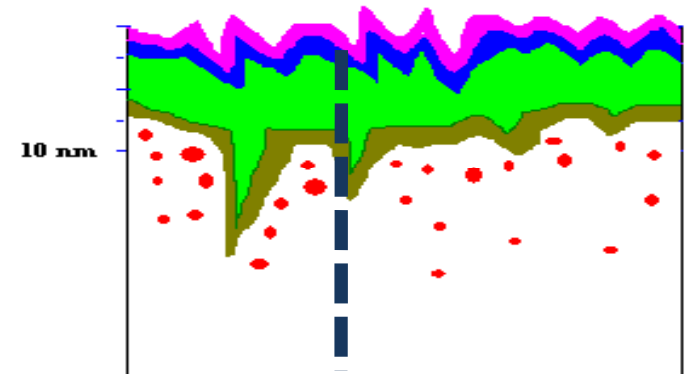
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Tallahassee, FL 32310



What defects may be important

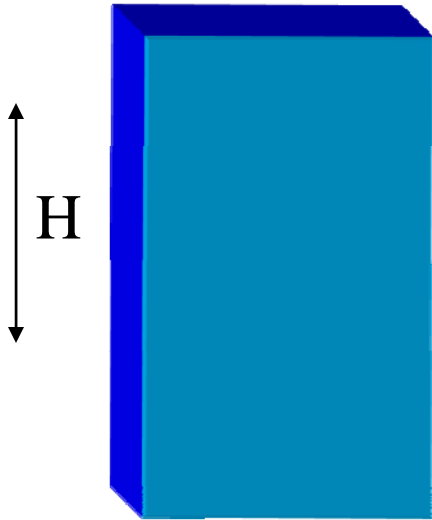
- Only defects in the 100-400 nm surface layer are relevant for SRF
- **Nanoscale oxide layers**
 - Normal surface layers: suppression of superconductivity by proximity effect
 - Supply of impurities which suppress superconductivity at the surface
- **Impurities**
 - nonmagnetic
 - magnetic
- **Grain boundaries**
 - current blocking effects
 - impurity segregation
- **Is pinning of vortices important for SRF?**
 - penetration of vortices
 - trapped vortices
 - vortex hotspots



Hydrocarbons & impurities
Nb hydroxides
 Nb_2O_5 , dielectric
 NbO_x ($0.2 < x < 2$), metallic
 NbO_x precipitates
($0.02 < x < 0.2$)

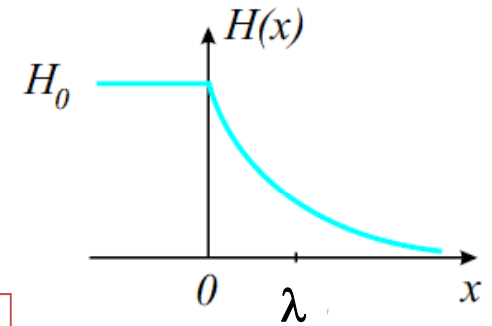
It is always “yes/no, but...”

RF field and current density scales



- Meissner current density $J_s(y)$:

$$H(y) = H_0 e^{-y/\lambda}, \quad J_s(y) = \frac{H_0}{\lambda} e^{-y/\lambda}$$



- For $H = 100$ mT the Meissner current density in Nb:

$$J_s(0) \approx 2 \times 10^8 \text{ A/cm}^2$$

- Maximum (pairbreaking) current density:

$$J_d \cong \frac{H_c(T)}{\lambda(T)}$$

$$J_d(0) \approx 4 \text{ MA/mm}^2 \text{ for pure Nb}$$

Can pinning of vortices be important in SRF?

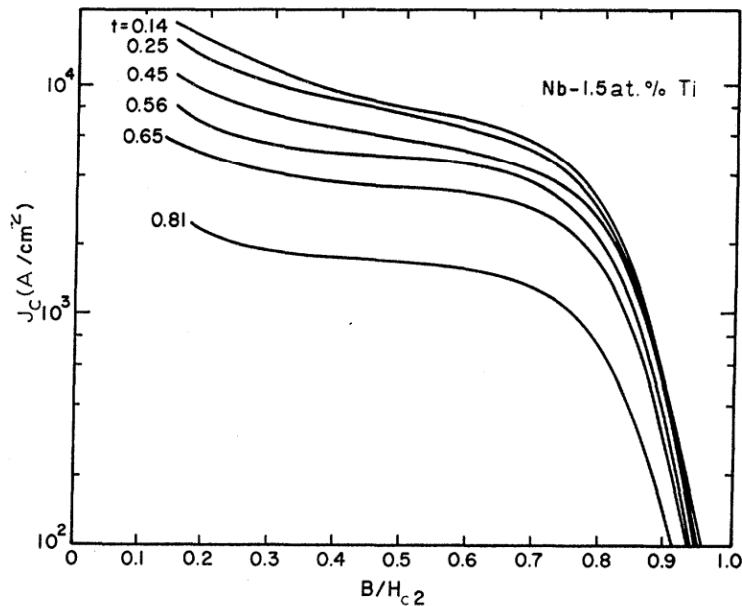
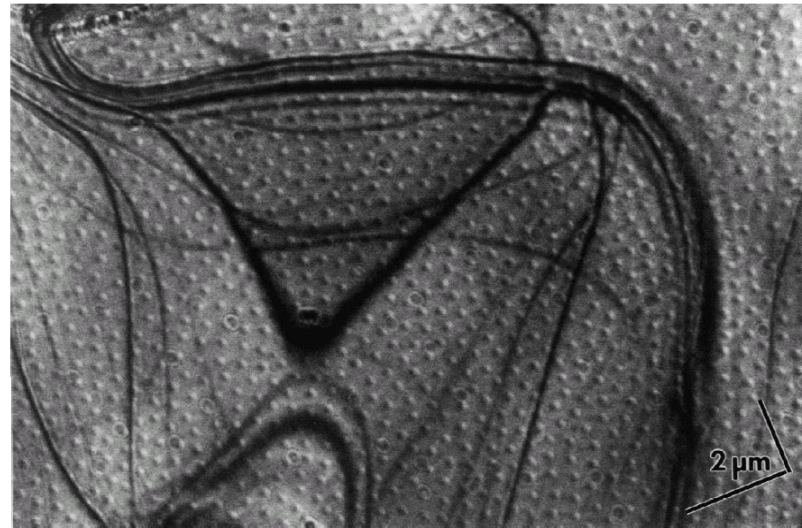


FIG. 4. Logarithmic plot of the critical current density versus reduced field for Nb-1.5 at.% Ti.

Pinning by dislocations in Nb:

W.A. Fietz and W.W. Webb, Phys. Rev. 178, 657 (1967)
A.M. Campbell and J.E. Evetts, Adv. Phys. 20, 199 (1972)

Lorentz microscopy of vortex interaction with dislocations. A. Tonomura et al, 2002



Pinning by single dislocations in clean Nb is very weak

Dislocation networks caused by strong plastic deformation can produce

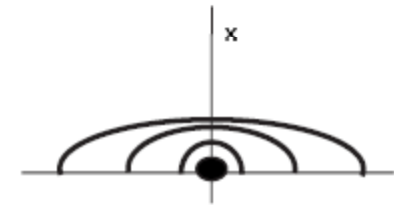
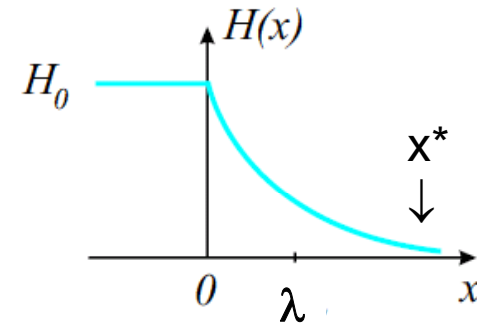
$$J_c \sim 10^4 - 10^5 \text{ A}/\text{cm}^2$$

Depinning J_c is **3-4 orders** of magnitude smaller than the Meissner RF currents at 100mT

Trapped vortices

Vortex can be trapped by pinning at the distances above x^* from the surface where $J_s(x^*) \sim J_c$

$$x^* \approx \lambda \ln \frac{J_d}{J_c} \cong (8 - 10)\lambda = 300 - 400 \text{ nm}$$

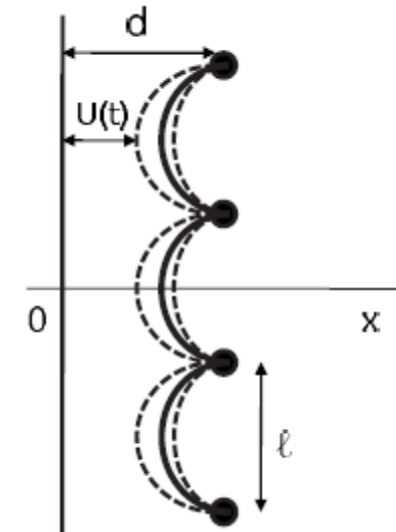
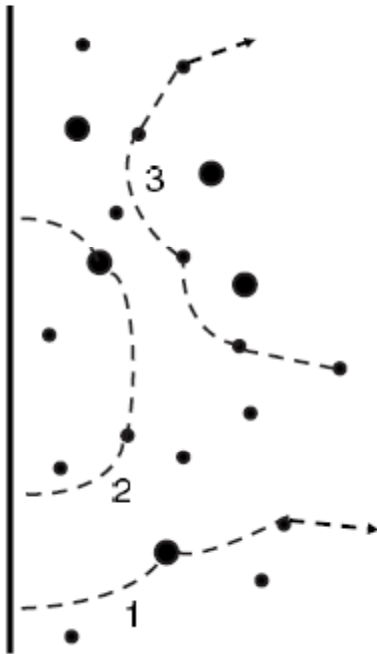


Penetration of single vortices through reduced surface barrier

Oscillating trapped vortices produce hotspots and Q slope

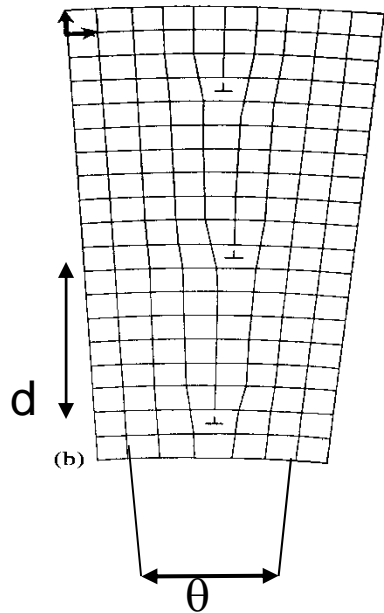
Vortex hotspots can be swiped away by thermal gradients (Ciovati, 2008)

Gurevich and Ciovati, Phys. Rev. B 77, 104501 (2008)
Ciovati and Gurevich Phys. Rev. STAB 11, 122001 (2008)



Types of grain boundaries

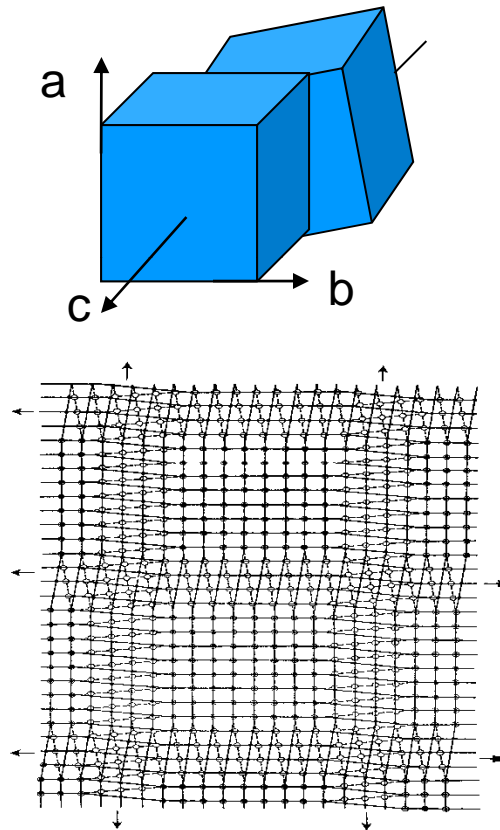
[001] tilt GB



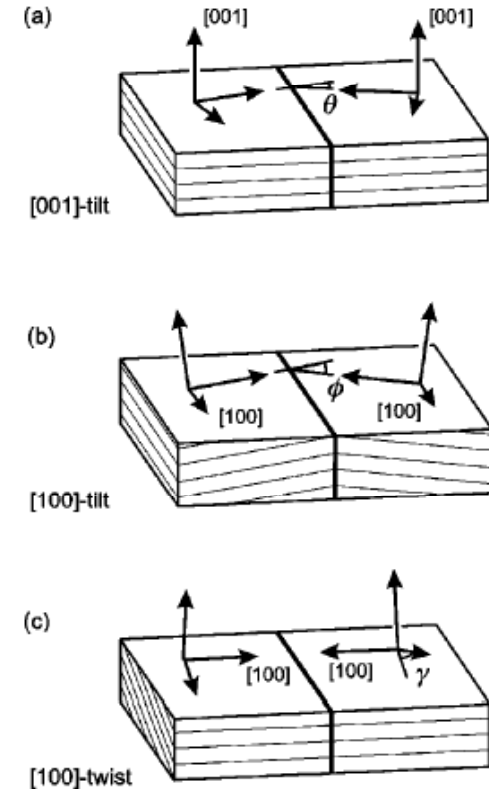
Chain of edge dislocations spaced by

$$d = b/2\sin(\theta/2)$$

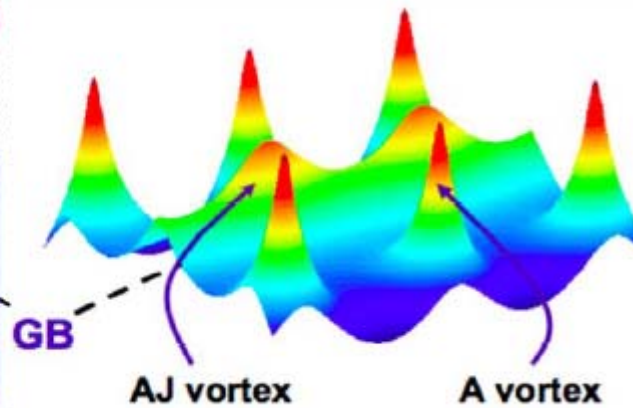
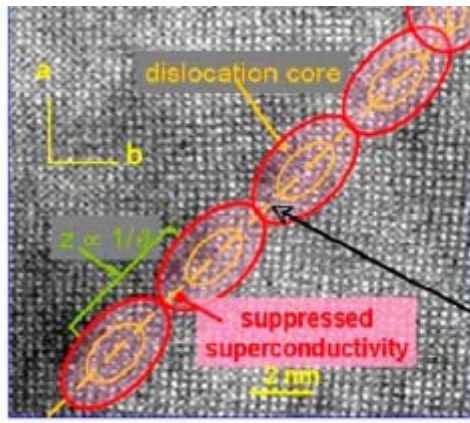
Twist GB



Cellular structure of twist dislocations in the ab plane



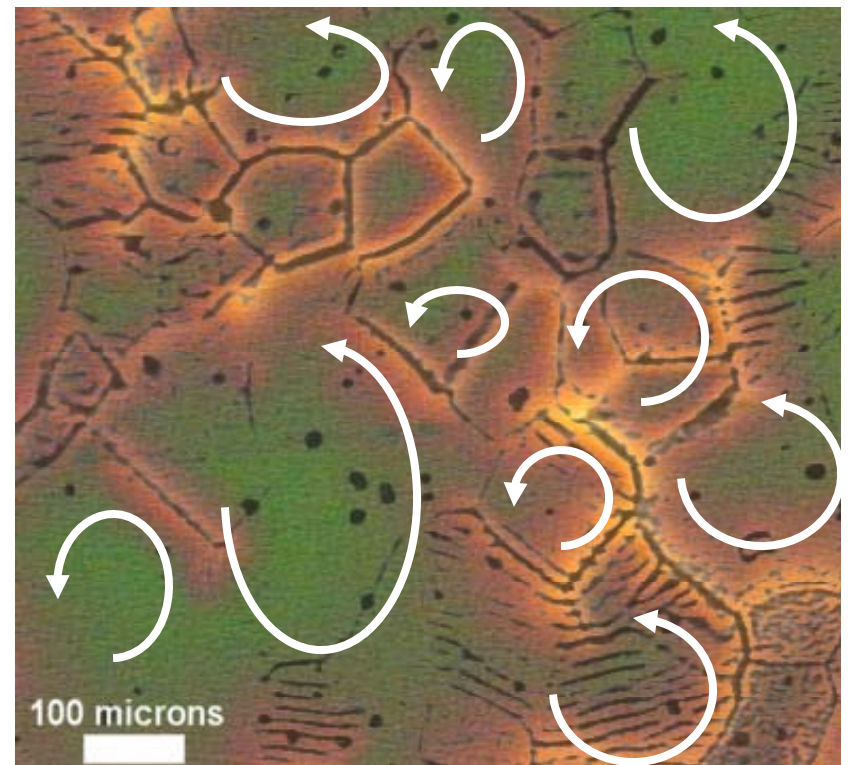
Magnetic granularity in polycrystals



Magneto-optical imaging

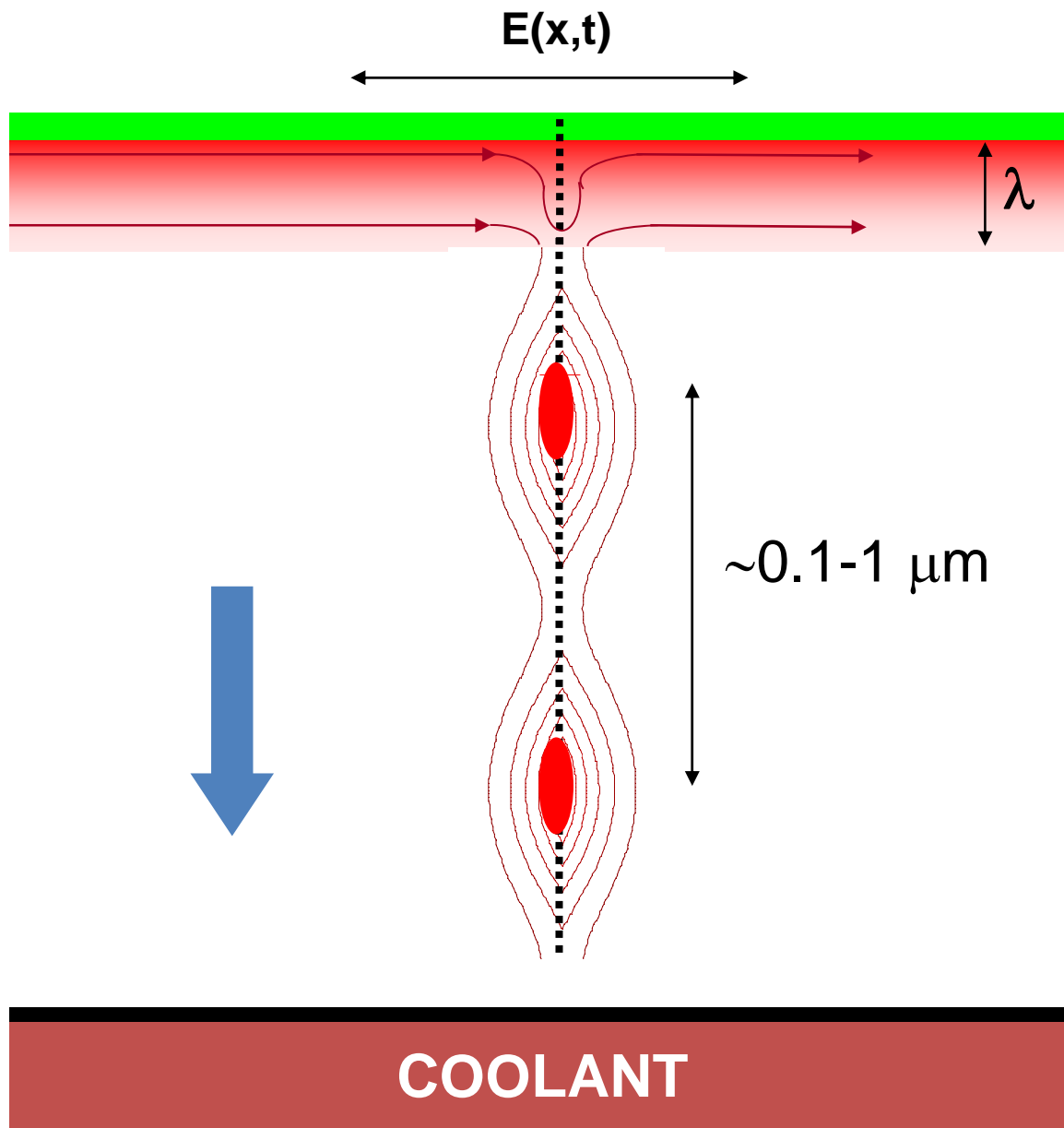
Phys. Rev B, 46, R3187 (1992); Phys. Rev. B 48, 12857 (1993);
Phys. Rev. B 50, 13563 (1994); Phys Rev B 65, 214531 (2002).
Phys Rev Lett 88, 097001 (2002).

- Only small currents can pass through GBs
- Current blocking by GBs is weak if:
 - $\xi = 40\text{nm} \gg$ lattice spacing ✓
 - screening length $\ll \xi$ ✓



Polyanskii, Feldmann, 2001

Penetration of vortices along grain boundaries



Are grain boundaries always weak links?

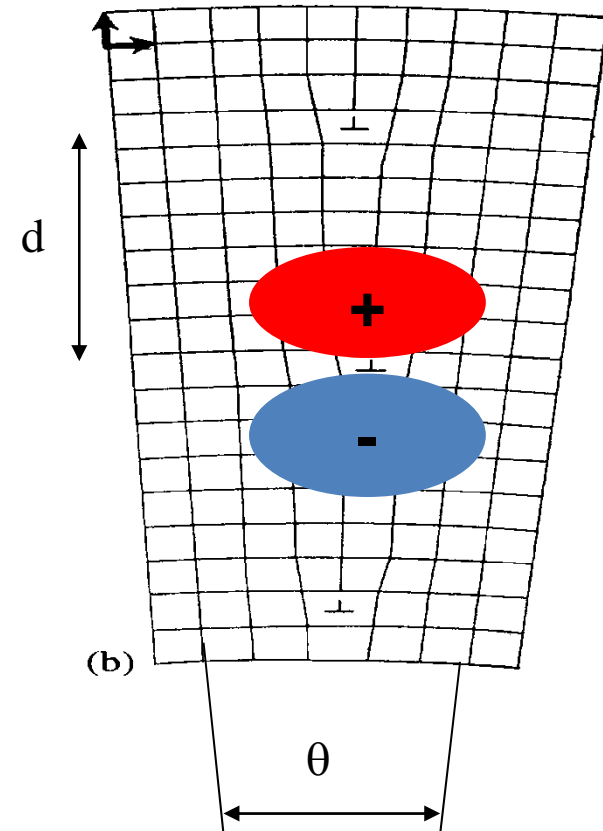
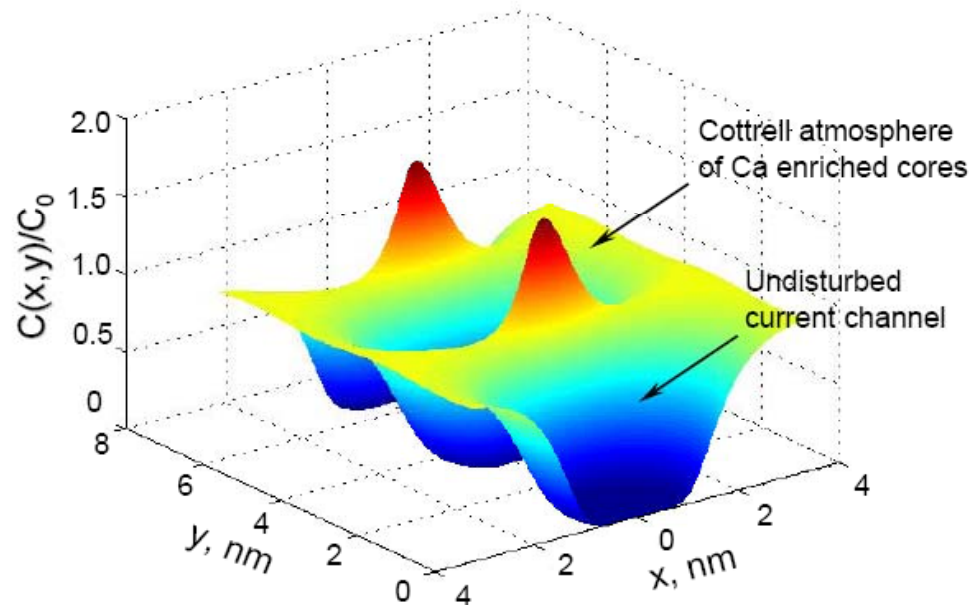
$$H > H_J \cong H_c \sqrt{\frac{\xi J_d}{\lambda J_c}}$$

For $J_c < 0.1 J_d \sim 10 \text{ MA/cm}^2$, the onset of vortex penetration along GBs drops below **60 mT**

Breakdown fields of 160-200 mT of the best polycrystalline Nb cavities seem to rule out the weak link behavior of GBs

Are grain boundaries in Nb always benign ?

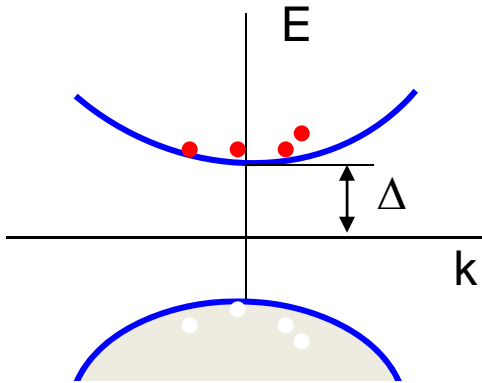
- Impurity poisoning due to segregation and diffusion along GBs
- Cottrell atmospheres of impurities in strain and electric fields of dislocation cores



Segregation of big impurities
in tensile (-) regions

depletion in compressed (+) regions

RF dissipation

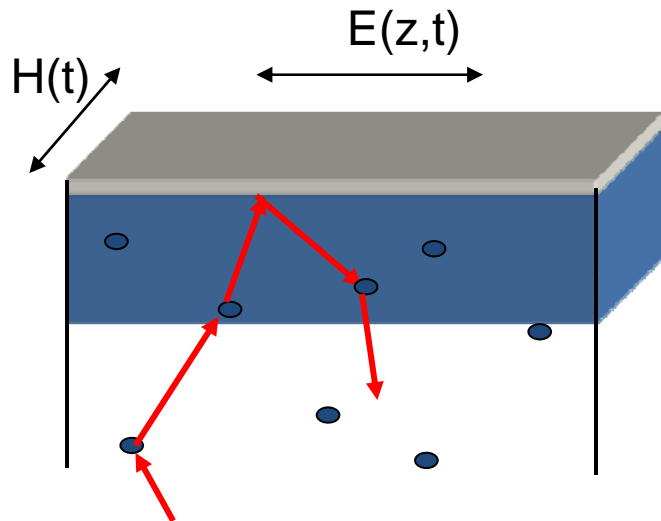


- Thermal activation of normal electrons
- Scattering mechanisms and normal state conductivity: $\sigma_n = e^2 n_0 \ell / \rho_F$, $\rho_F = \hbar (3\pi^2 n_0)^{1/3}$

- Normal skin effect ($\ell \ll \lambda$): multiple impurity scattering in the λ - belt:

$$R_s \sim (\mu_0^2 \omega^2 \lambda^3 \sigma_n \Delta / T) \exp(-\Delta / T)$$

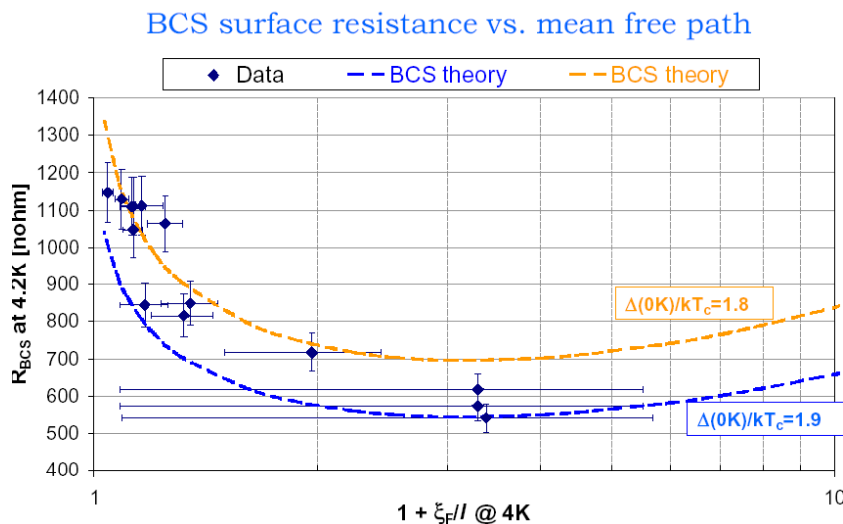
- Anomalous skin effect in the clean limit ($\ell \gg \lambda$): scattering by the gradient of the rf field:



$$R_s \propto \frac{\mu_0^2 \omega^2 \lambda^4 \Delta n_0}{k_B T \rho_F} \left[\ln \left(\frac{\Delta}{\hbar \omega} \right) + C_0 \right] \exp \left(- \frac{\Delta}{k_B T} \right)$$

Nonmagnetic impurities

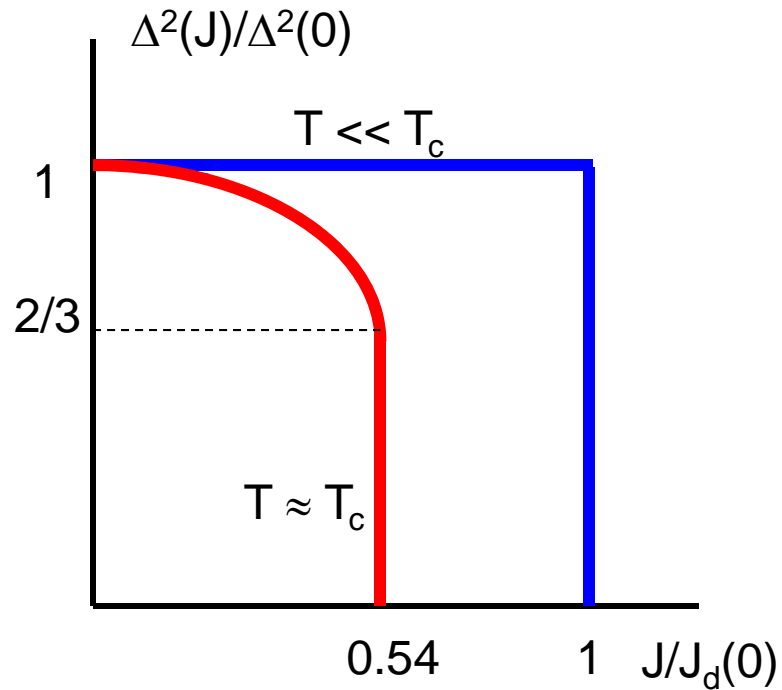
- Effect on the **linear** surface resistance
 - No suppression of the superconducting gap (Anderson theorem)
 - Increase of the London penetration depth
 - Increase of the BCS surface resistance
 - Decrease the lower critical field (the onset of vortex penetration)



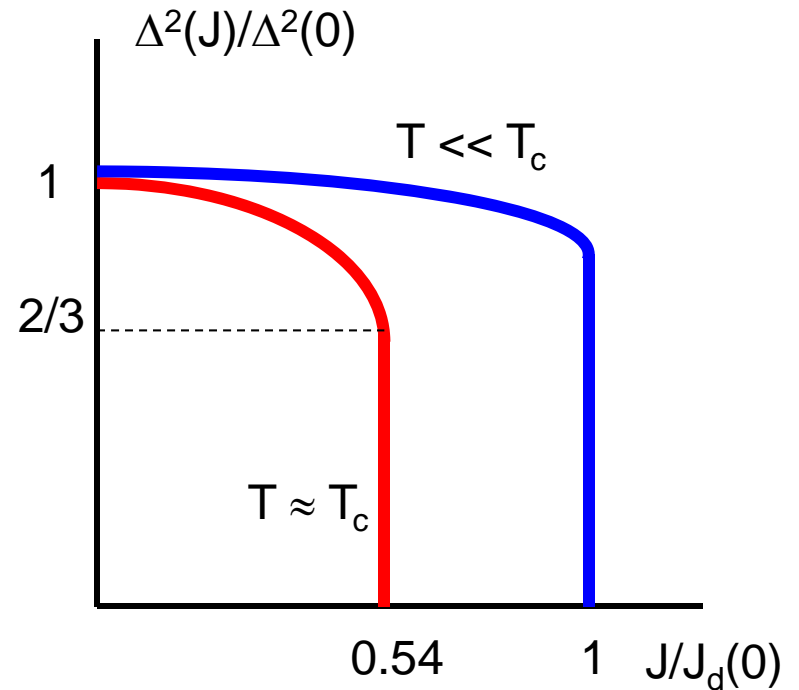
Nonmagnetic impurities appear to be not too bad for R_{BCS} , but are they benign at high rf fields?

Current dependence of the SC gap

- **Clean limit ($\ell \gg \xi_0$)**



- **Dirty limit ($\ell \ll \xi_0$)**



In the clean limit the gap $\Delta(J)$ is independent of RF field at low T

J. Bardeen, Rev. Mod. Phys. 34, 667 (1962).

Nonlinear Meissner effect

- Linear Meissner effect in the clean limit up to $H = H_c$
- Nonlinear Meissner effect in the dirty limit: Field-dependent penetration depth calculated from the Usadel equations:

$$\lambda^2(H) \cong \lambda^2 \left[1 + a \left(\frac{H}{H_c} \right)^2 \right], \quad a = \frac{3\pi}{64} + \frac{1}{4\pi} \approx 0.23$$

- Increase $R_s \star \lambda^3(H)$ and the kinetic inductance L_k of the Cooper pairs in a film of length ℓ , width w and thickness $d < \lambda$:

$$L_k = \mu_0 \lambda^2(H) \ell / wd$$

- Shift of the resonance frequency $f = (C/L)^{1/2}$ of a thin film resonator

$$\delta f / f = -\mu_0 \ell [\lambda^2(H) - \lambda^2] / 2Ldw$$

enabled us to measure the nonlinear Meissner effect in a dirty 65 nm Nb film

Groll, Gurevich and Chiorescu, Phys. Rev. B81, 020504(R) (2010).

Nonlinear Meissner effect in a thin Nb film

$f \approx 20$ GHz, $T = 80$ mK

GROLL, GUREVICH, AND CHIORESCU

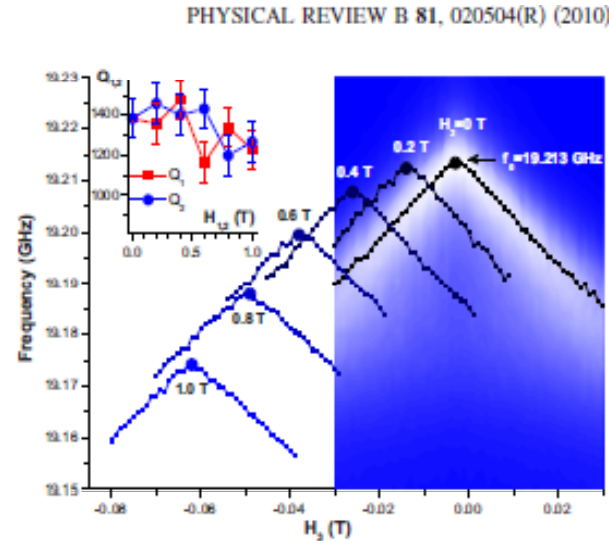
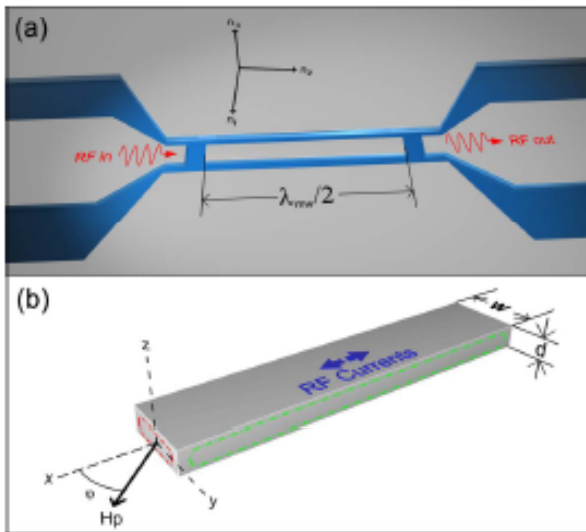
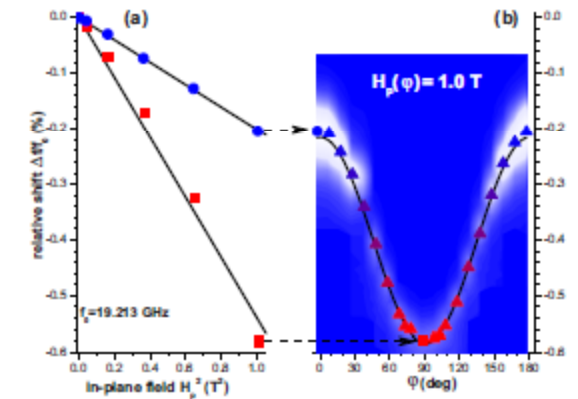


FIG. 2. (Color online) Contour plot of amplified cavity transmission for variable H_s field while $H_z=0$ T (blue/orav: 0 μ W

MEASUREMENT OF THE NONLINEAR MEISSNER EFFECT...



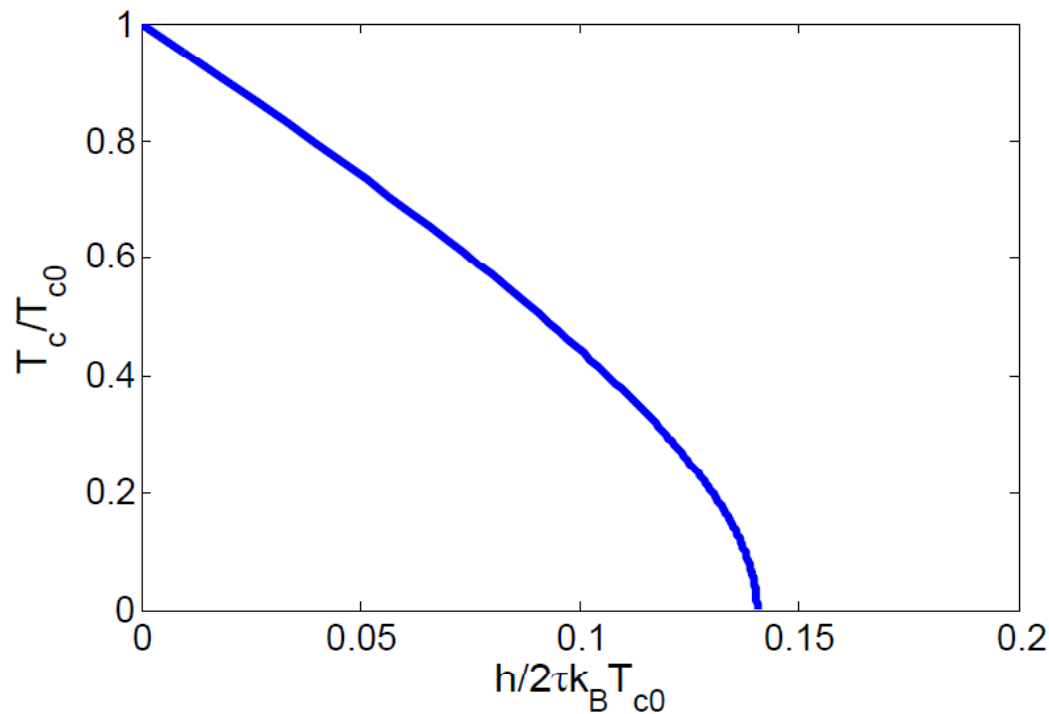
Dependence of δf on the orientation of the in-plane dc magnetic field:

$$\delta f \propto 2 - \cos 2\varphi$$

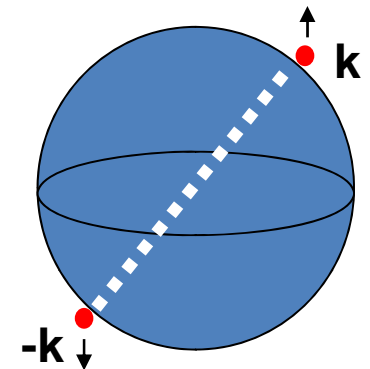
- Excellent agreement between theory and experiment
- Seven-fold increase of the onset of the vortex field penetration in a dirty Nb thin film with $d < \lambda$ up to 1T: $H < H_{c1} = (2\phi_0/\pi d^2)\ln(d/\xi)$
- Proof of principle of the multilayer coating idea

Magnetic impurities

- Suppression of T_c and superconducting gap by spin-flip scattering by magnetic impurities
- Gapless superconductivity



Abrikosov and Gorkov, 1960

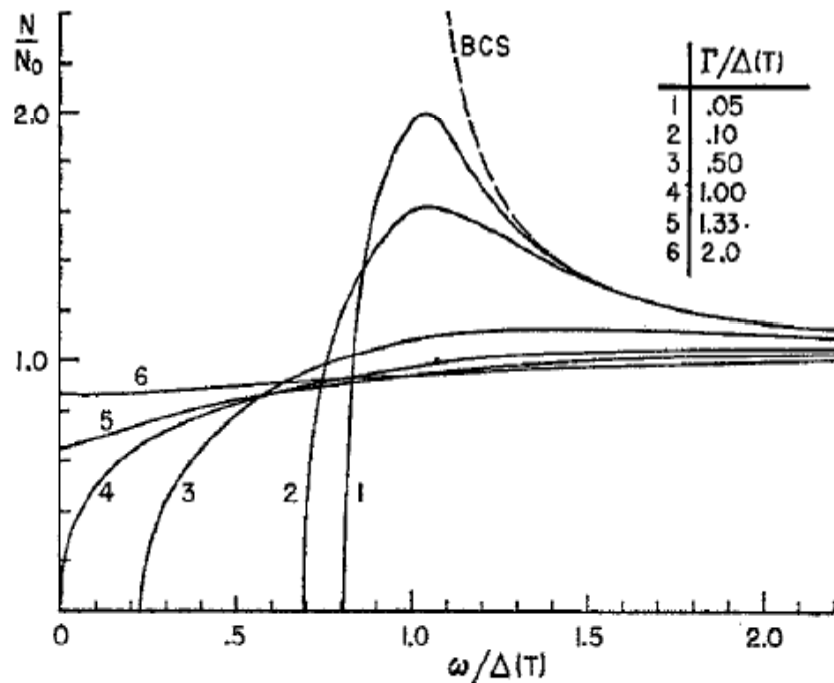


Spin flip scattering time:

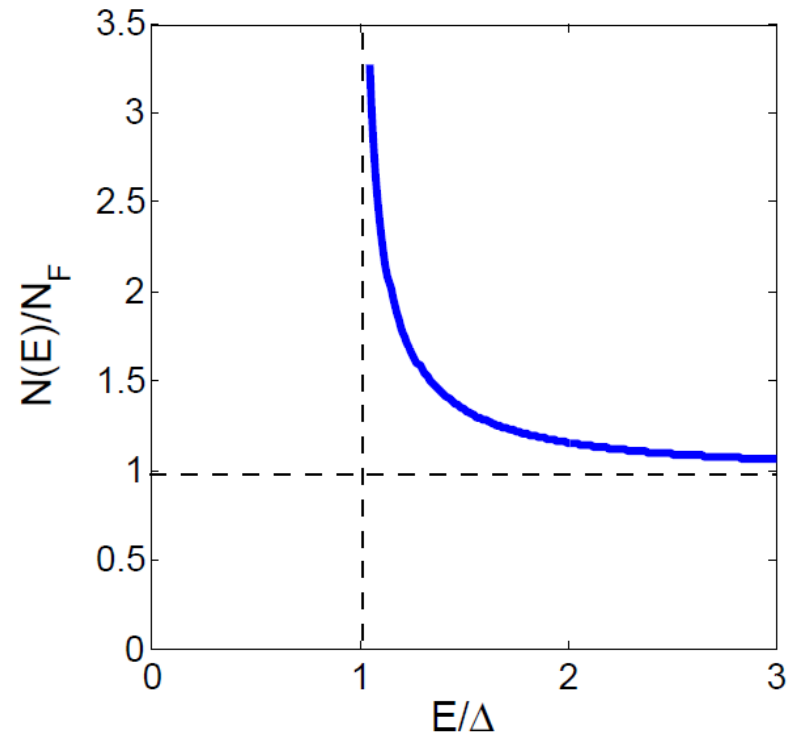
$$\frac{1}{\tau_s} = n_s N_F S(S+1) \int |J(k)|^2 d\Omega$$

Magnetic impurities in the 100 nm surface layer strongly increase R_s

Filling up the superconducting gap



Skalski et al (1964); Ambegaokar and Griffin (1965)



Ideal BCS density of states which provides the BCS surface resistance

Effect of weak magnetic impurities:

- smearing up $N(E)$ and reduction of the gap
- gapless superconductivity at large $\Gamma=1/\tau_s$
- **giant increase of R_s**

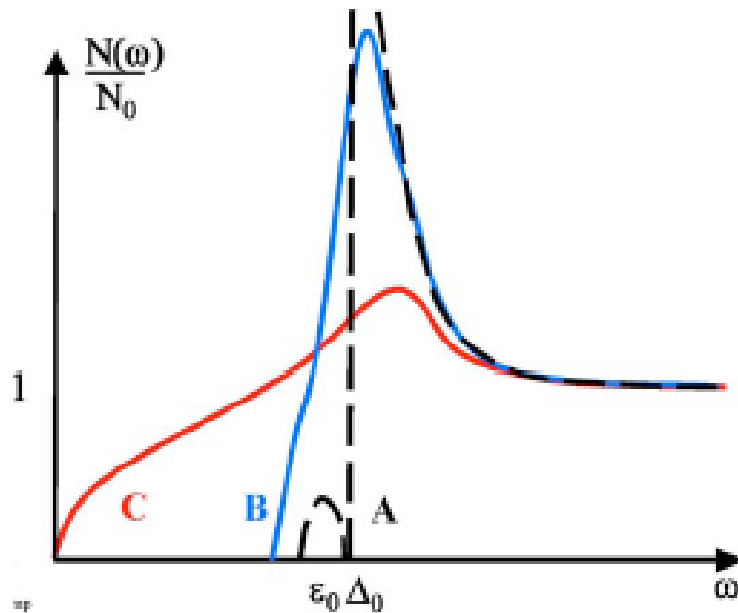
Small concentration of magnetic impurities

Localizes states on magnetic impurities at energies below the SC gap

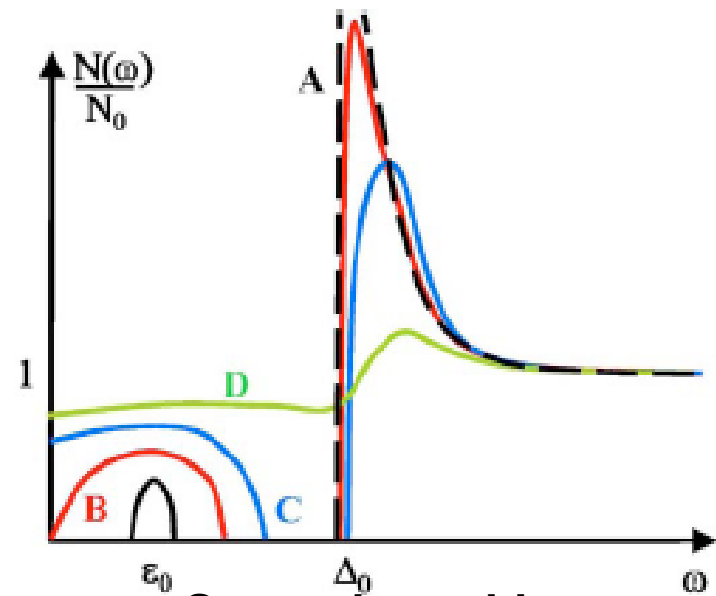
(Shiba (1968); Rusinov (1968))

Can be revealed by point contact tunneling spectroscopy

(Prossler, Zasadzinski, Cooley, Pellin et al, Appl. Phys. Lett. 92, 212505 (2008),



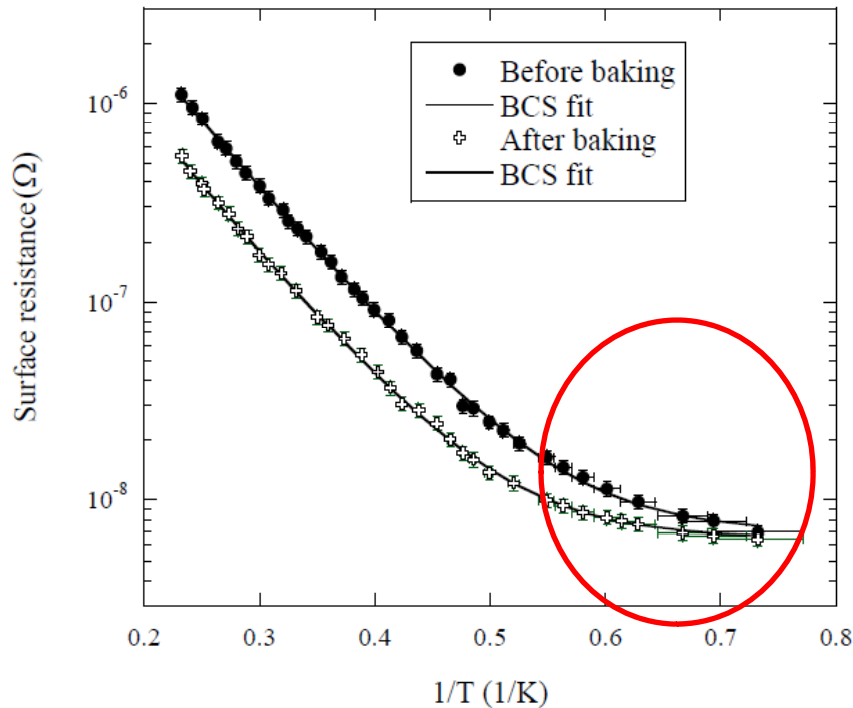
Weak impurities



Strong impurities

- Shiba magnetic impurity band reduces the activation energy of $R_s(T)$
- Finite density of states at $E=0$ causes non-exponential $R_s(T)$
- A mechanism of residual resistance

Defects and residual resistance



Ciovati, 2005

Requires a finite density of states at the Fermi surface

Can result from hotspots caused by:

- Trapped vortices
- Thin normal islands of sub-oxides
- Magnetic impurities at the surface

Neither of these factors can be definitive, to be compatible with $Q = 10^{10}$

Conclusions

- **Point and extended defects can seriously limit the SRF performance but universal answers good for all circumstances can be obtained**
- **Some defects (i.e. impurities or GBs) can be both bad and benign depending on the rf field and temperature**
- **Magnetic impurities and vortices are always bad for SRF**
- **New SRF physics produced by defects (nonlinear Meissner effect, vortex hotspots, residual resistance, etc.)**
- **Measurements of the nonlinear Meissner effect revealed the 5-fold increase of the vortex penetration field in a thin Nb film**