

S-wave π Exchange And Its Implications for Quarkonium Spectroscopy



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Outline

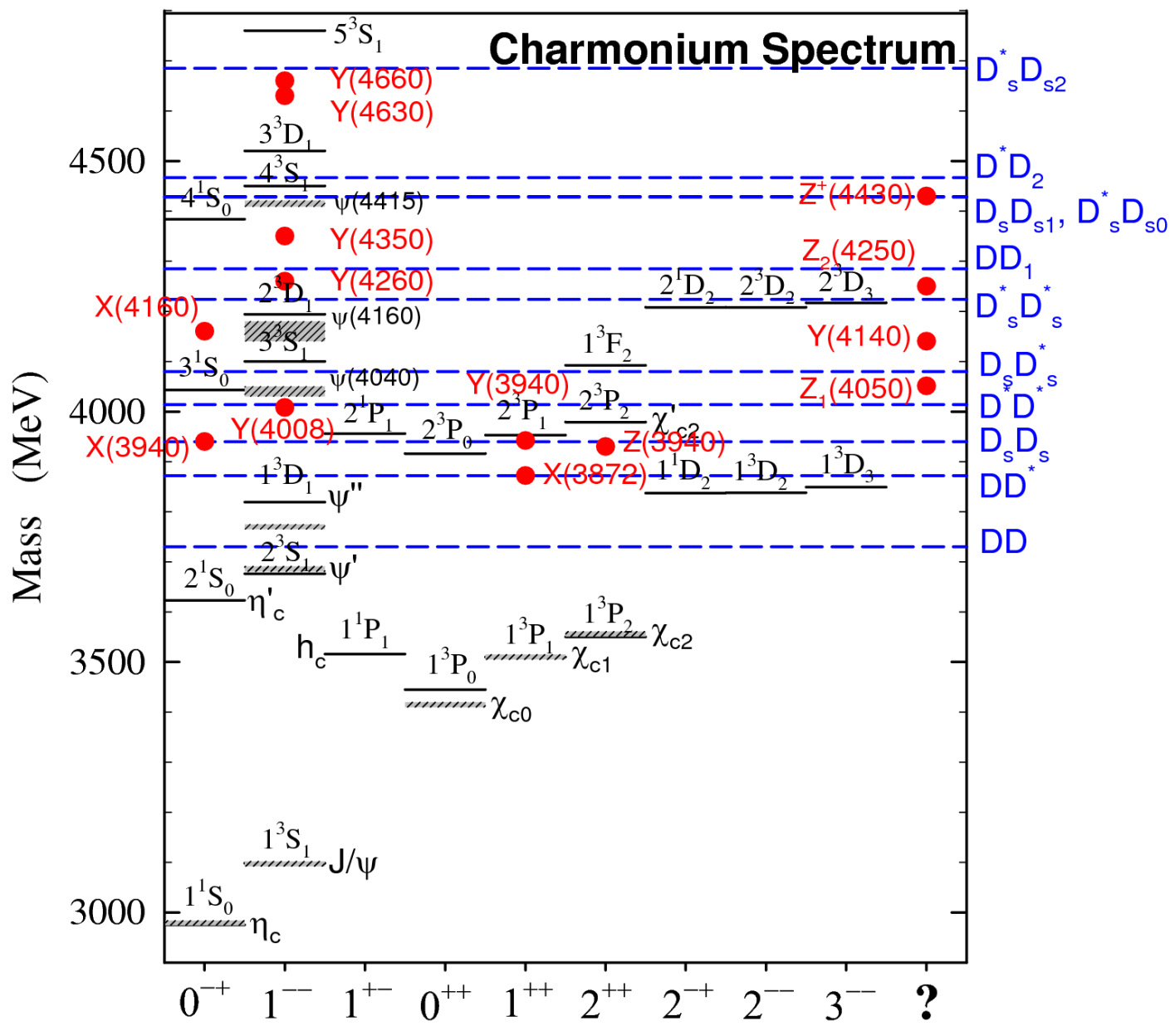
- Review of Charmonium Spectra
- General Phenomenon of S-wave π Exchange
- Application to the $D_1 D^*$ System (1^- , $I=0$)
 - Close, CD, Thomas PRD 81 074033 (2010)
- Generalization to Other Sectors
- Experimental Signature
- Theoretical Uncertainties
 - Filin, *et al.* arXiv:1004.4789
- Summary

States of Interest

- Visha Bhardwaj -- Recent Results on $X(3872)$ at Belle
- Estia Eichten -- Hadronic Transitions in Quarkonium
- Kai Yi -- Quarkonium Spectroscopy Results at CDF
- Kenkichi Miyabayashi -- Other Charmonium and cc -like States at Belle
- Arafat Gabareen Mokhtar -- Babar update on new charmonia in B decays
- Cristoph Hanhart -- Hadron molecules
- Mikhail Voloshin -- Hadroquarkonium
- Chiara Sabelli -- Tetraquarks
- Pierre Artoisenet -- Production of the $X(3872)$ at the Tevatron and the LHC

There are Many Anomalous States

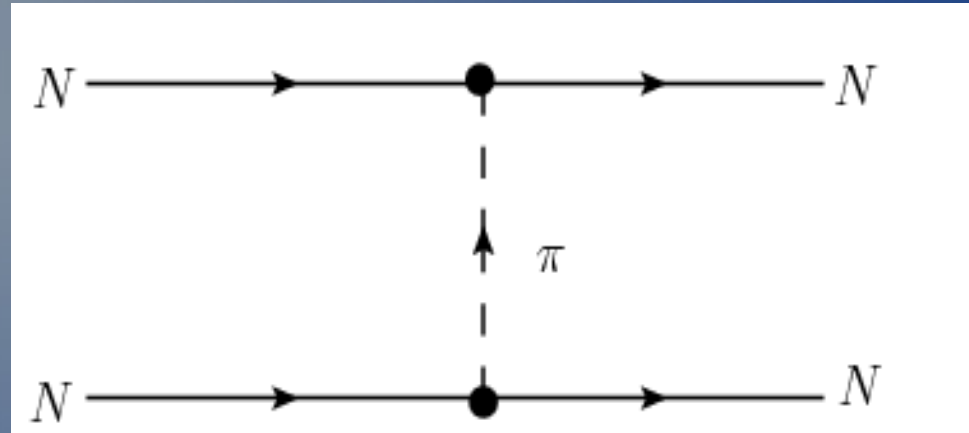
Godfrey Proceedings of DPF-2009



Unique Opportunity to Study QCD Dynamics

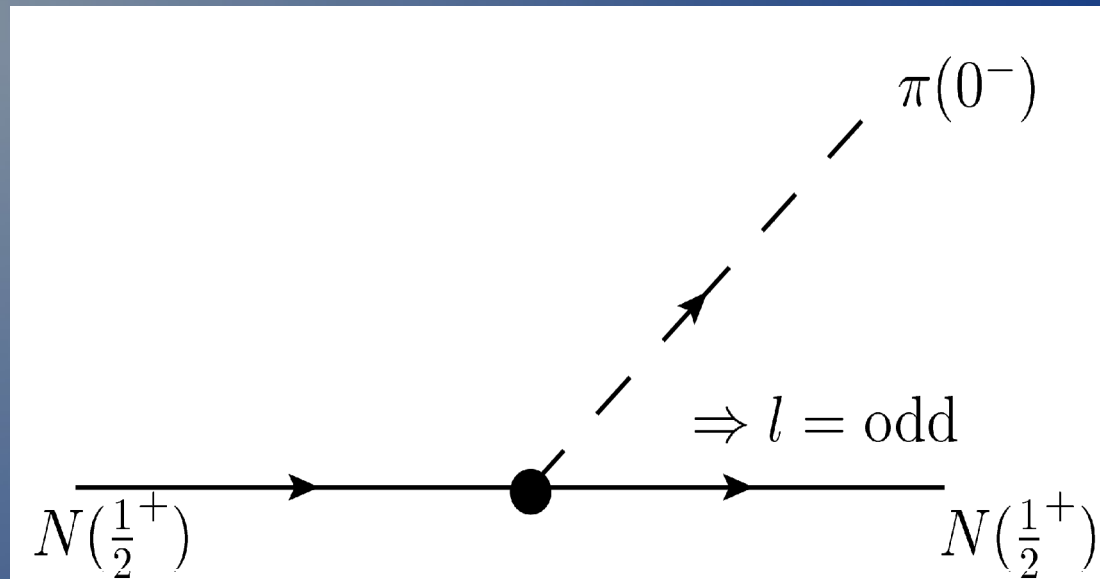
- QCD is very challenging and difficult
- The naïve quark model ($q\bar{q}$) picture of mesons basically worked...
 - Particularly the charm spectrum seemed well understood
- However, many non $q\bar{q}$ states are possible and “predicted”
- Unambiguous identification of these particles will give critical tests of our models.

Historical π Exchange

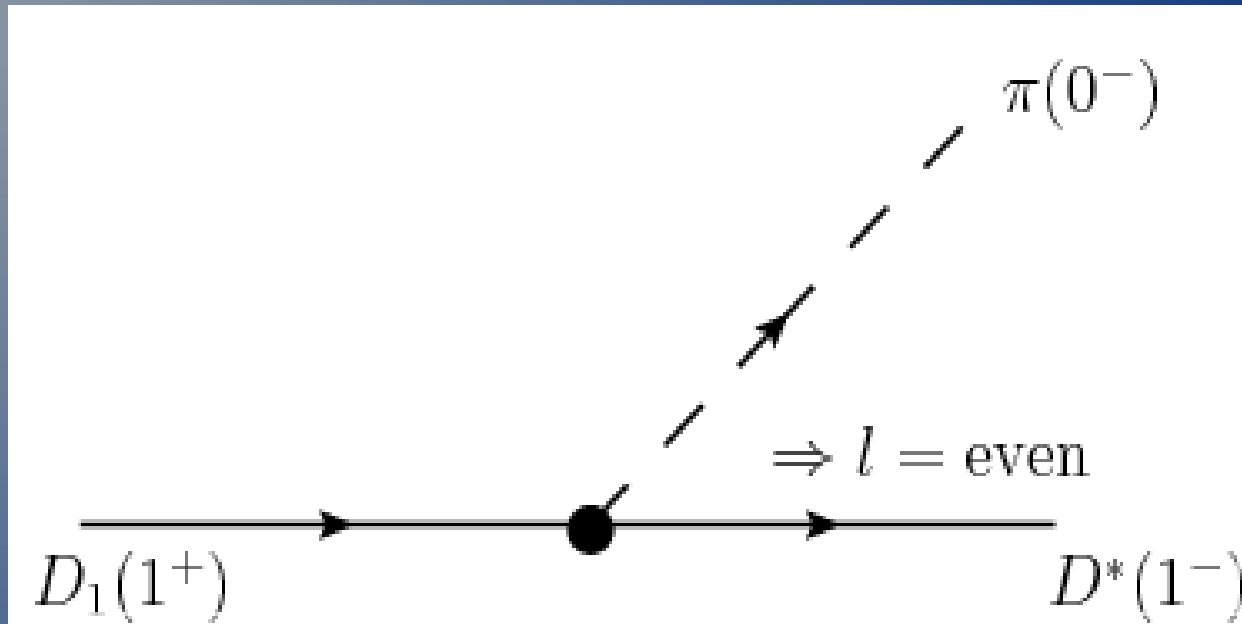


- Familiar from NN Application (Yukawa)
- By Analogy Törnqvist suggested “Deusons”
- Like the NN case Binding Energies were of order 1 MeV
- Discovery of a possible deuson X(3872) led to a great deal of interest in π -bound meson molecules (and other meson exchanges).

P-wave π Emission From Parity Conservation



S-wave π Emission



S-wave π Emission (Cont.)

Chiral Perturbation Theory Model

$$\Gamma(D_1^0 \rightarrow D^{*+} \pi^-) = \frac{h^2}{8\pi f_\pi^2} \frac{|\mathbf{q}| m_{D^*}}{m_{D_1}^3} (m_{D_1}^2 - m_{D^*}^2)^2 \times \frac{1}{3} \left(2 + \frac{(m_{D_1} + m_{D^*})^2}{4m_{D_1}^2 m_{D^*}^2} \right)$$

Quark- π Coupling

$$\bar{\psi} \gamma_5 \psi \mapsto \boldsymbol{\sigma} \cdot \left(\mathbf{q} + \frac{\omega}{m} \mathbf{p} \right)$$

3P_0 Model

$$H_{3P_0} = \gamma \boldsymbol{\sigma} \cdot \vec{p} |q\bar{q}\rangle \langle 0|$$

Empirical Support

PDG Live 2009

$P = +$ Meson	Γ [MeV]
$D_1(2430)$	384^{+130}_{-110}
$D_0^*(2400)^0$	261 ± 50
$D_0^*(2400)^\pm$	283 ± 40
$K_1(1400)^\pm$	174 ± 13
$K_0(1430)^\pm$	270 ± 80
$a_1(1260)$	250 to 600
$b_1(1235)$	142 ± 9

$D_1 D^*$ Study

Liu *et. al.* PRD77,034003(2008)

– $Z(4430)^+$ as a $D_1 D^*$ bound state → Isovector

$$V(r) = \frac{h^2 q_0^2}{8\pi f_\pi^2} \frac{\cos(\mu r)}{r} \approx \frac{h^2 (m_{D_1} - m_{D^*})^2}{8\pi f_\pi^2} \frac{\cos(\mu r)}{r}$$

$$\mu^2 = q_0^2 - m_\pi^2 \approx (m_{D_1} - m_{D^*})^2 - m_\pi^2$$

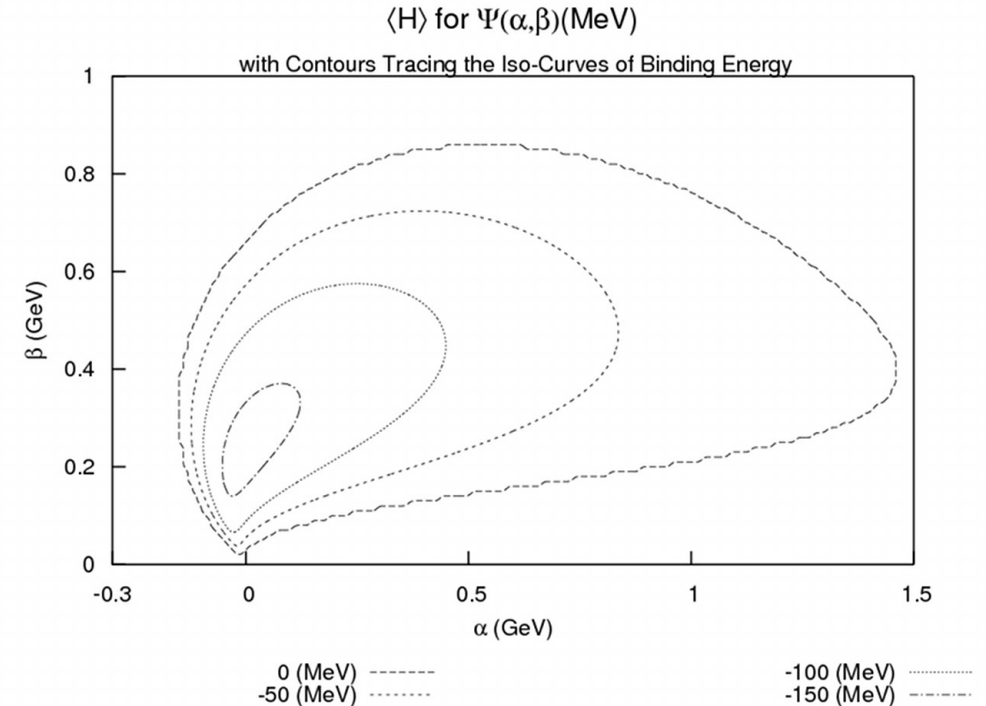
$$\psi(r) = (1 - \alpha r^2) e^{-\beta r^2}$$

Conclusion: No Binding

$D_1 D^*$ Study (Cont.)

$$\Rightarrow V(r) = -3 \frac{h^2 (m_{D_1} - m_{D^*})^2 \cos(\mu r)}{8\pi f_\pi^2 r}$$

$$\psi(r) = (1 - \alpha r^2) e^{-\beta r^2}$$



Conclusion: No Reason to Expect Isovector Binding
Robust Isoscalar Binding

$D_1 D^*$ Study (Cont.) – χ Coupling

$$\Gamma(D_1^0 \rightarrow D^{*+} \pi^-) = \frac{h^2}{8\pi f_\pi^2} \frac{|\mathbf{q}| m_{D^*}}{m_{D_1}^3} (m_{D_1}^2 - m_{D^*}^2)^2 \times \frac{1}{3} \left(2 + \frac{(m_{D_1} + m_{D^*})^2}{4m_{D_1}^2 m_{D^*}^2} \right)$$

$$\Gamma(D_1^0 \rightarrow D^* \pi^-) = \frac{2}{3} \Gamma(D_1 \rightarrow D^* \pi) \approx \frac{2}{3} \Gamma(D_1) = \frac{2}{3} 384_{-110}^{+130} \text{ MeV}$$

$$h = 0.80_{-0.17}^{+0.20}$$

$$h \mapsto h \left(1 - \frac{2}{9} \frac{q^2}{\beta^2} \right) \exp \left\{ -\frac{p^2}{12\beta^2} \right\}$$

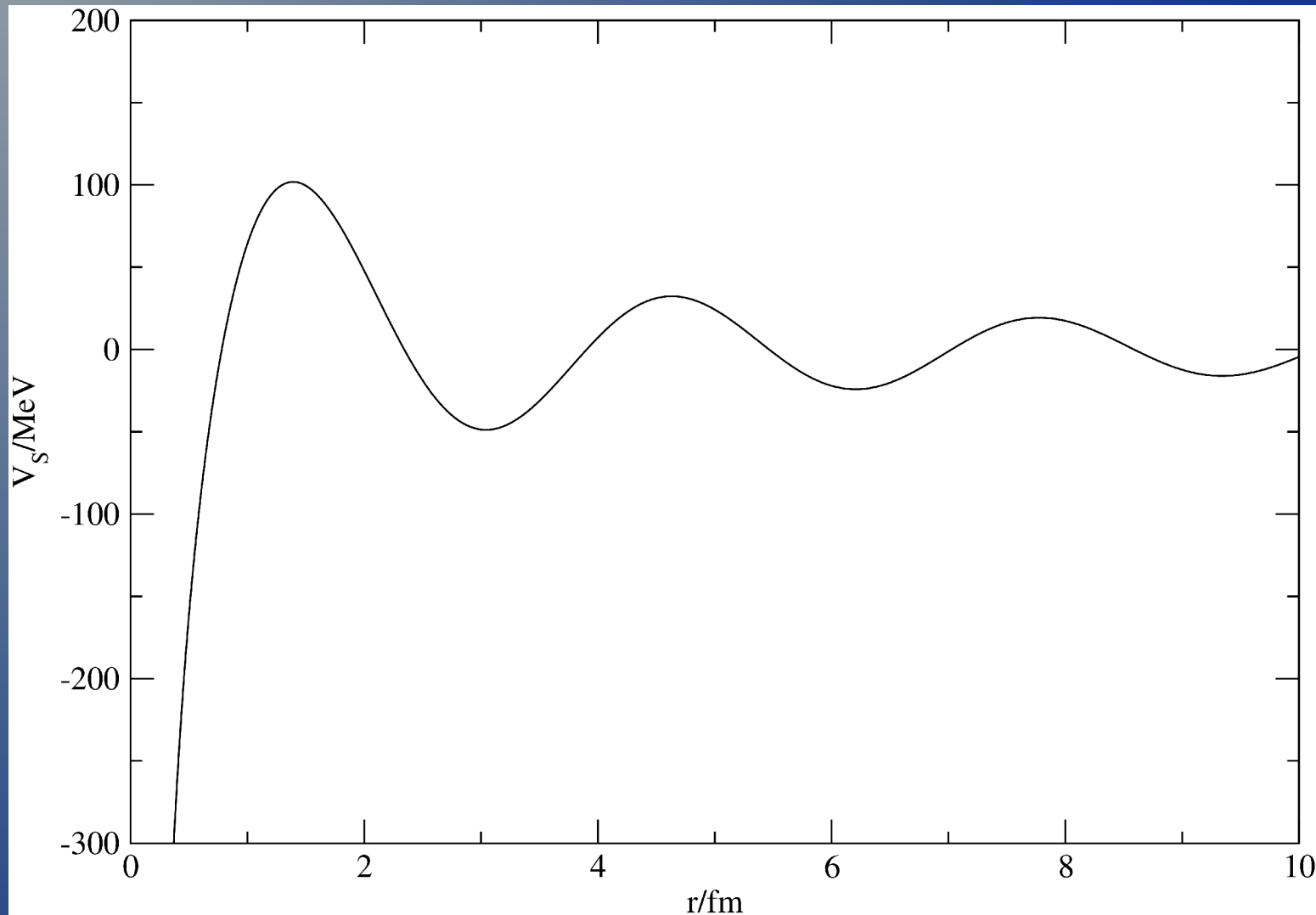
$$h = 1.0_{-0.2}^{+0.3}$$

$D_1 D^*$ Study (Cont.) – Schrödinger Eq.

$$\Rightarrow V(r) = -3 \frac{h^2 (m_{D_1} - m_{D^*})^2 \cos(\mu r)}{8\pi f_\pi^2 r}$$

State	Binding Energy / MeV					
	$h = 0.8$	$h = 0.9$	$h = 1.0$	$h = 1.1$	$h = 1.2$	$h = 1.3$
1S (0, 1, 2) ⁻⁻	230	415	680	1000	1500	2100
2S	12	20	29	39	76	210
3S	1.5	3.6	6.7	11	51	65
1P (0, 1, 2, 3) ^{+ -}	8.2	16	25	35	120	260
2P	1	2.8	5.8	22	48	61
3P	–	1	1.5	9.7	15	20
1D (0, 1, 2, 3, 4) ⁻⁻	1.4	8.3	17	28	40	54
2D	–	1.3	4.1	8.0	13	18
3D	–	–	1	2.6	4.9	7.8

$D_1 D^*$ Study (Cont.) – The Potential

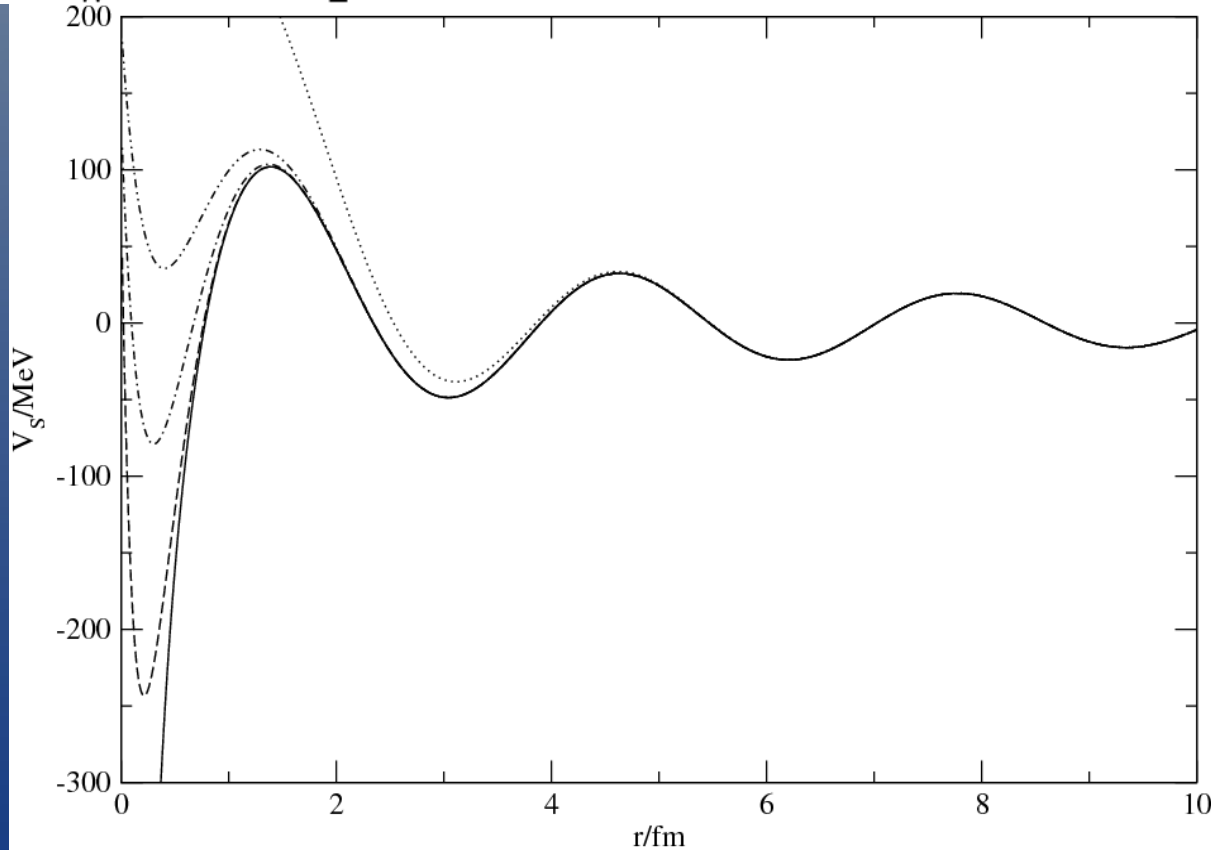


$$r_{ms} \text{ [fm]} = 0.6(1S) ; 3(2S) ; 6.5(3S)$$

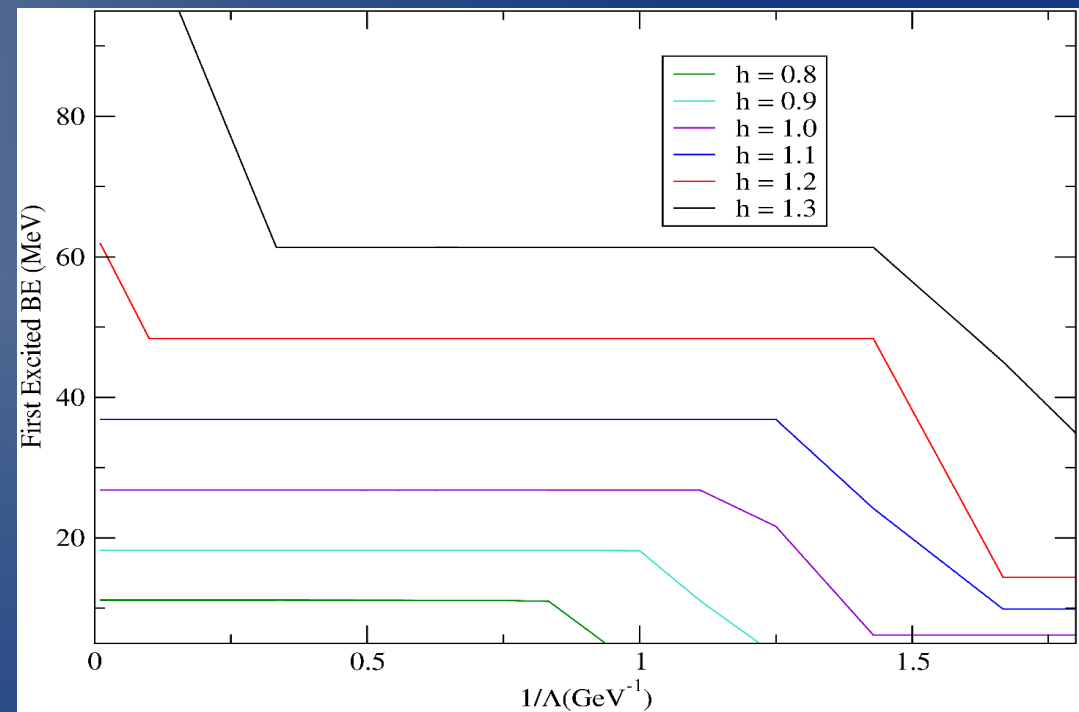
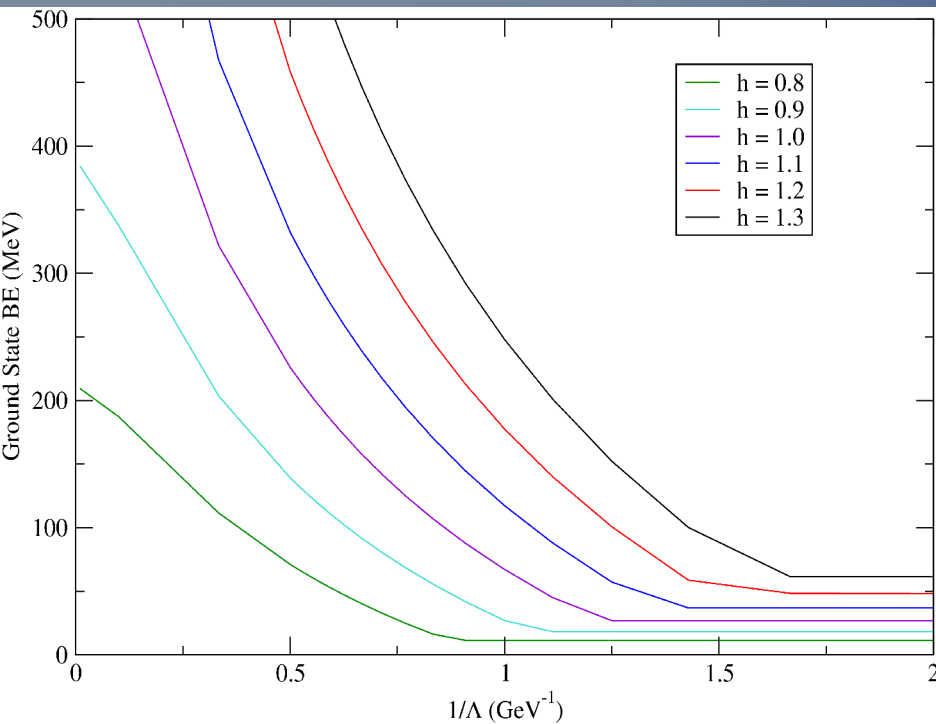
$D_1 D^*$ Study (Cont.) – Form Factor

$$\mathcal{F} = \left(\frac{\Lambda^2 - m_\pi^2}{\Lambda^2 - q^2} \right) \approx \left(\frac{\Lambda^2 - m_\pi^2}{\Lambda^2 + \mu^2 - m_\pi^2 + \mathbf{q}^2} \right)$$

$$V_S(r) = \frac{h^2(m_{D_1} - m_{D^*})^2}{8\pi f_\pi^2} \left[\frac{\cos(\mu r)}{r} - \frac{e^{-Xr}}{r} - \frac{(\Lambda^2 - m_\pi^2)}{2X} e^{-Xr} \right] (\tau_i \cdot \tau_j)$$



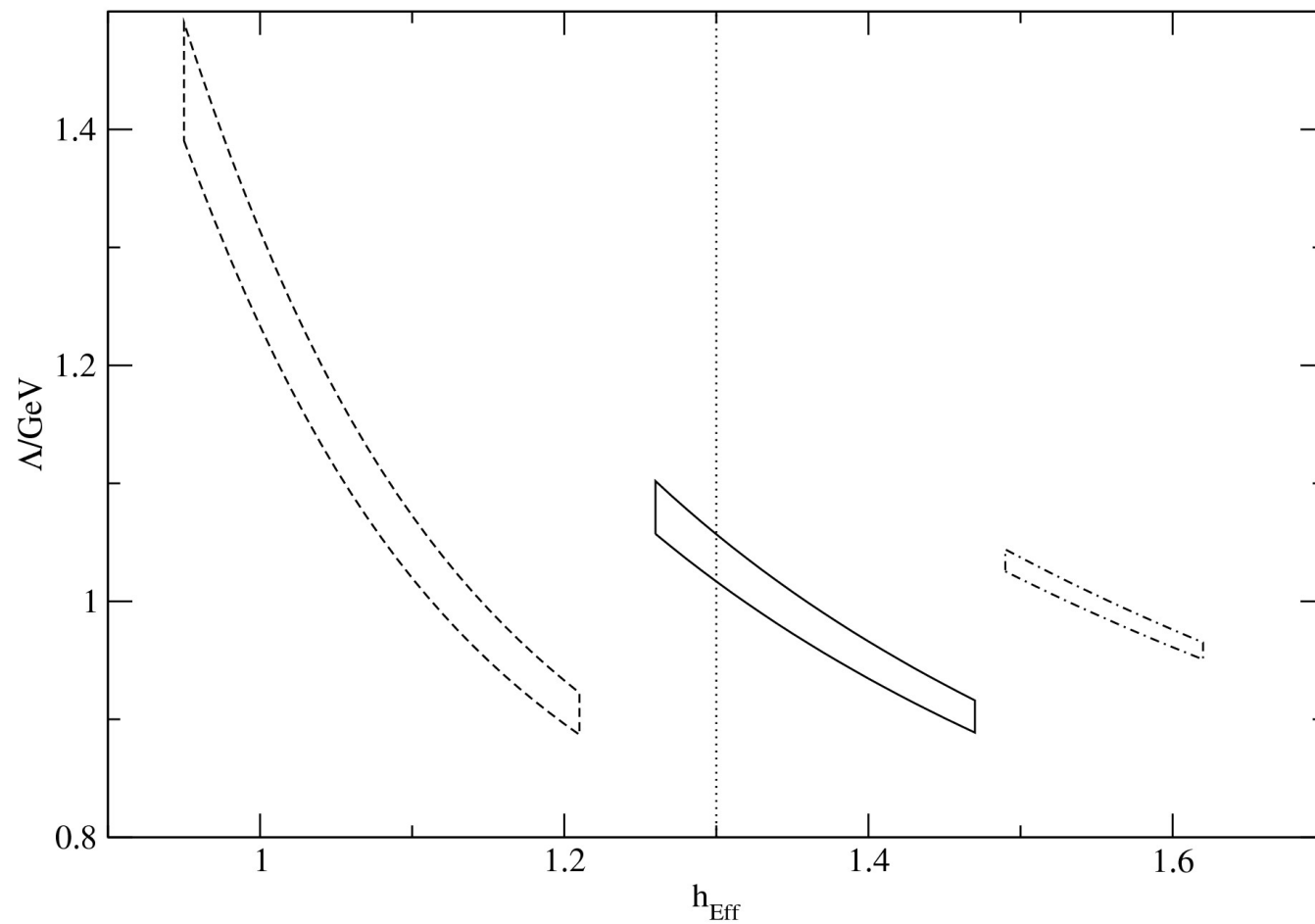
$D_1 D^*$ Study (Cont.) – Parameter Flexibility



D₁D* Study (Cont.) – More Sensible Results ($\Lambda = 1$ GeV)

State	Binding Energy / MeV					
	$h = 0.8$	$h = 0.9$	$h = 1.0$	$h = 1.1$	$h = 1.2$	$h = 1.3$
1S (0, 1, 2) ⁻⁻	12	20	29	60	106	160
2S	1.5	3.6	23	39	51	65
3S	---	---	6.7	11	15	21
1P (0, 1, 2, 3) ⁺⁻⁻	8.2	16	25	35	48	61
2P	1	2.8	5.8	9.7	15	20
3P	–	---	1.5	3.3	6	8.6
1D (0, 1, 2, 3, 4) ⁻⁻	1.4	8.3	17	28	40	54
2D	–	1.3	4.1	8.0	13	18
3D	–	–	1	2.6	4.9	7.8

$D_1 D^*$ Study (Cont.) – Tune to Fit Y(4260) and Y(4360)



$D_1 D^*$ Study (Cont.) – Generalize (Other C/I Channels; $\Lambda = 1$ GeV)

State	Isospin	Binding Energy / MeV					
		$h = 0.8$	$h = 0.9$	$h = 1.0$	$h = 1.1$	$h = 1.2$	$h = 1.3$
1S $(0, 1, 2)^{--}$	0	12	20	29	60	110	160
2S		1.6	3.6	23	39	57	65
3S		–	0.6	6.7	11	15	21
1S $(0, 1, 2)^{--}$	1	4.2	8.8	15	21	29	38
2S		0.0	0.0	0.2	0.7	1.5	2.8
3S		–	–	–	–	0.0	0.2
1S $(0, 1, 2)^{-+}$	0	47	67	90	120	150	180
2S		4.2	8.1	13	19	27	35
3S		0.5	1.7	3.5	6.1	10	14
1S $(0, 1, 2)^{-+}$	1	0.1	0.5	1.6	3.4	5.9	8.9
2S		–	–	–	0.1	0.4	0.9

Theoretical Uncertainty

Pretend S-wave π is the only thing...

Make a static approximation to q_0

$$\mu^2 = q_0^2 - m_\pi^2 \approx (m_{D_1} - m_{D^*})^2 - m_\pi^2$$

Normal Yukawa Process...

$$V(q) \propto \frac{1}{q^2 + m^2 + i\epsilon} \rightarrow V(r) \propto \frac{e^{-mr}}{r}$$

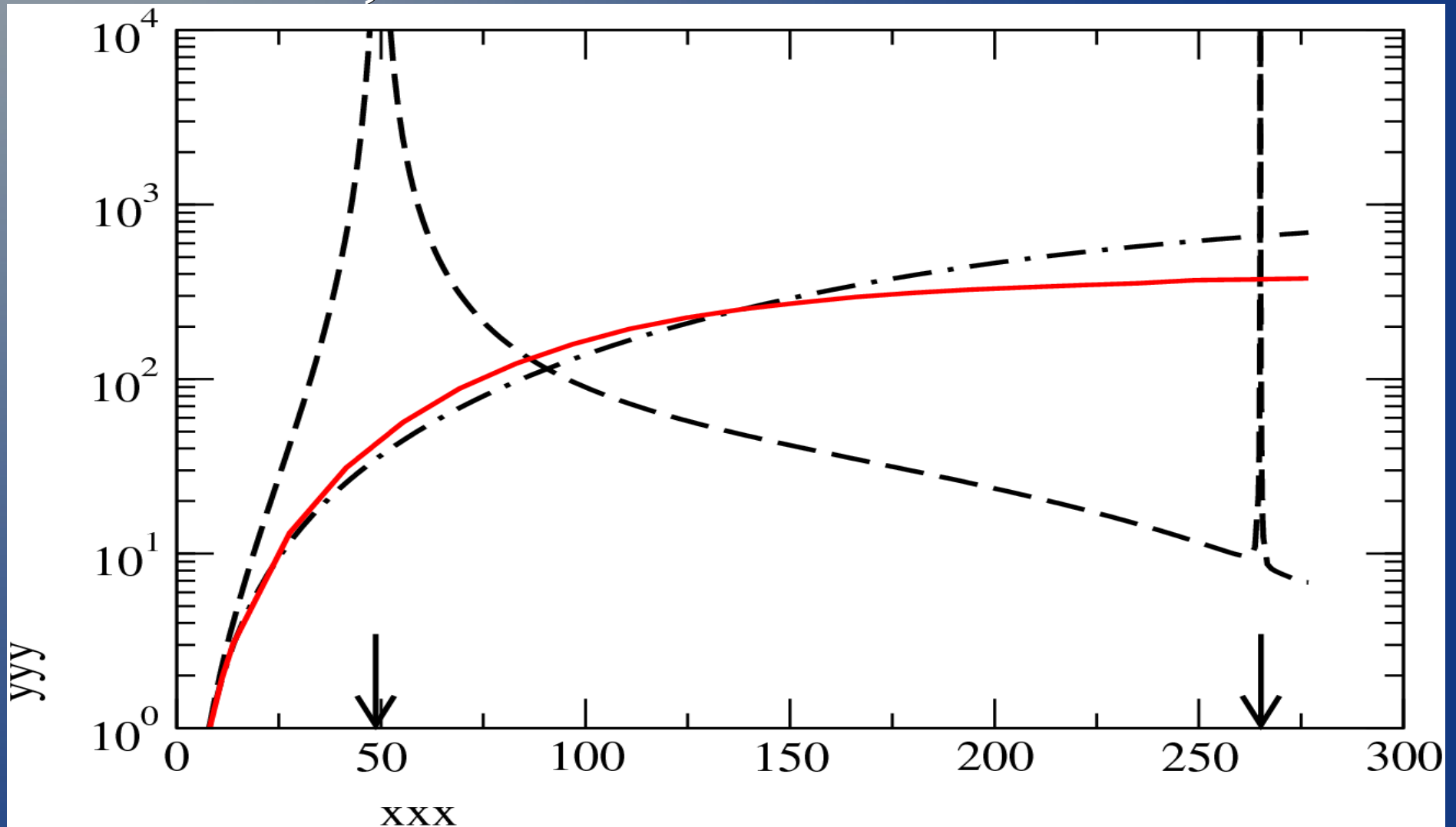
Parity Changing π Exchange Prop...

$$V(q) \propto \frac{1}{q^2 - \mu^2 + i\epsilon} \rightarrow V(r) \propto \frac{e^{-i|\mu|r}}{r} \propto \frac{\cos(|\mu|r) - i \sin(|\mu|r)}{r}$$

$$E_n \mapsto \tilde{E}_n + i \frac{\Gamma}{2}$$

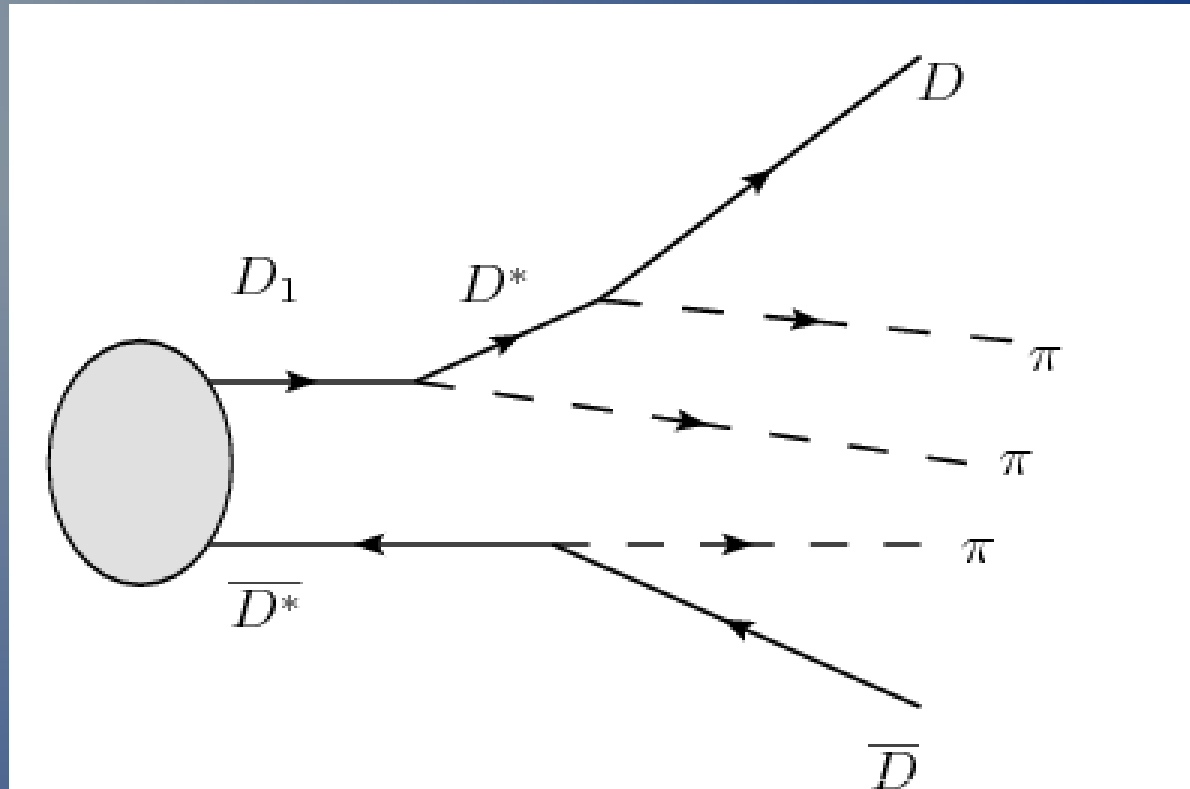
Recent Developments

Filin, *et al.* arXiv:1004.4789



Conclusion: Our States Do Not Exist

Experimental Searches



- A deeply bound D_1D^* system requires strength in the $DD3\pi$ channel.
- We recommended looking in $Y(4260)$ and $Y(4360)$ decays

Summary

- S-wave π Exchange as an interaction has general applicability ($D_1 D^*$, DD_0 , $B_1 B^*$, $K_1 K^*$, $K_0 K$, $a_1 \rho$, $b_1 \omega$, etc...)
- Uncertainties of OBE Molecular Models are large
→ Study patterns and unique decay channels.
 - $Y(4260) \rightarrow DD3\pi$
- Filin *et al.* find that the states are, at best, too broad to be observed. We should confirm.
- Is there a pair of mesons with suitable masses for S-wave π Exchange to work well?