

Double charmonium production in exclusive bottomonia decays

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Outline:

- Introduction
 - Leading twist bottomonia decays
 - $n_b \rightarrow J/\psi J/\psi$ decay
-

Introduction

Hard exclusive processes

Exclusive processes :

1. Decays: $Y \rightarrow \rho\pi, \eta_c \rightarrow \rho\rho, \dots$
2. Annihilation: $e^+e^- \rightarrow J/\Psi, J/\Psi, \dots$
3. Different formfactors: $F_\pi(Q^2), \dots$

General property :

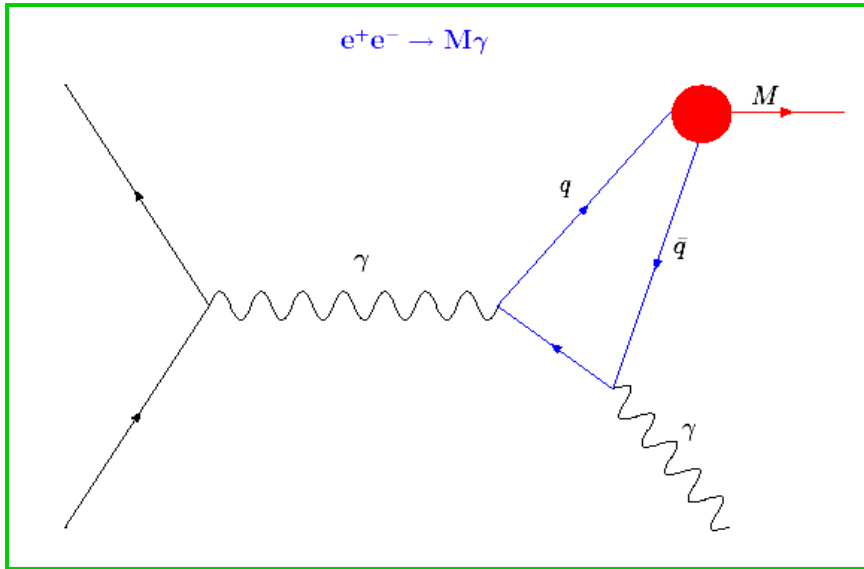
$$E_h \gg \Lambda_{QCD}, M_h$$

$$\text{Expansion parameter} : \sim \frac{M_{cc}^2}{M_{bb}^2} \sim \frac{1}{10}$$

$$\sigma = \underbrace{\frac{a_n(E_h = \infty)}{E_h^n}}_{\text{Leading twist}} + \underbrace{\frac{a_{n+1}(E_h = \infty)}{E_h^{n+1}}}_{\text{Next-to-leading twist}} + \dots$$

Light cone expansion formalism

Factorization



Factorization formula :

$$T = \sum_n \overbrace{C_n}^{\text{Short Dist.}} \times \underbrace{\langle M | O_n(0) | 0 \rangle}_{\text{Large Dist.}}$$

Different Contributions:

1. Short distance contribution : C_n (perturbative QCD)
2. Large distance contribution : $\langle M | O_n | 0 \rangle$ (nonperturbative effects)

The leading twist distribution amplitude

Operators that contribute at the leading order approximation:

$$\langle M(p) | \bar{q} \gamma_\mu \gamma_5 D_{\mu_1} \dots D_{\mu_n} q | 0 \rangle z^\mu z^{\mu_1} \dots z^{\mu_n} \sim (pz)^{n+1} \int_{-1}^1 d\xi \xi^n \varphi(\xi), \quad z^2 = 0, \quad \xi = x_1 - x_2$$

Distribution amplitude $\varphi(\xi)$ can be considered as a meson wave function.

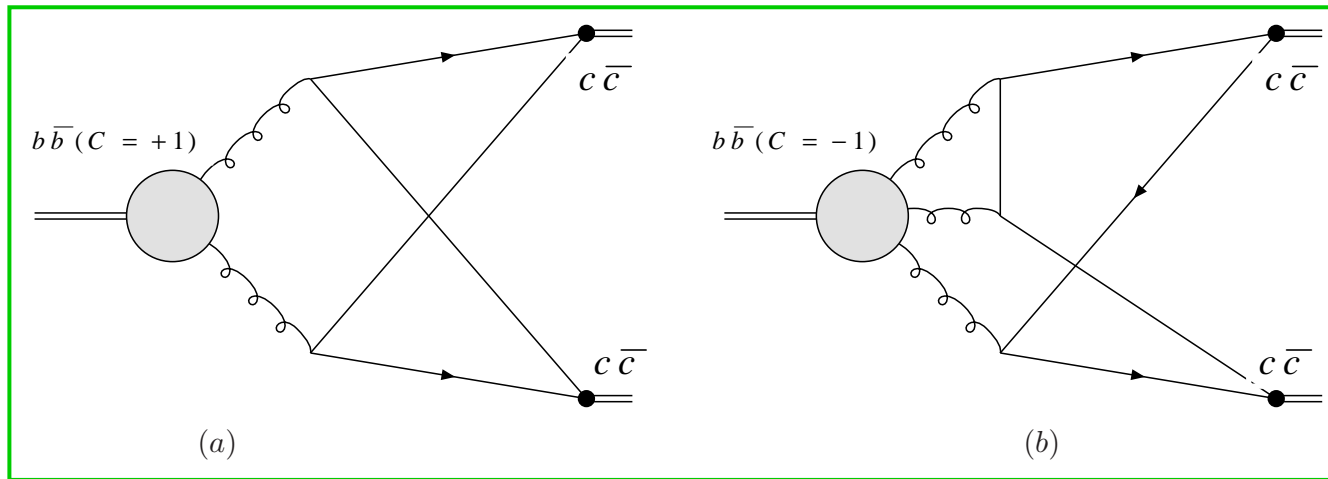
Distribution amplitudes

- Resume infinite series of operators
- Resume leading logarithmic radiative corrections $\left(\alpha_s(E_h) \times \text{Log} \left(\frac{E_h^2}{m_c^2} \right) \right)^n$

$$T = \int_{-1}^1 d\xi H(\xi, \mu) \varphi(\xi, \mu), \quad \mu \sim E_{\text{hard}}$$

Leading twist bottomonia decays

Leading twist decays



- $C = -1$ bottomonia are suppressed
- Leading twist decays $J_{bb} = 0, 2$

The amplitudes of the decays

$$T = \int d\xi_1 d\xi_2 H(\xi_1, \xi_2) \times \varphi_1(\xi_1) \varphi_2(\xi_2)$$

Charmonia distribution amplitudes

- Distribution amplitudes of the S-wave charmonia
3 DAs, Phys.Lett.B646,80,(2007), Phys.Rev.D75:094016,(2007), Phys.Rev.D77:034026,(2008)
- Distribution amplitudes of the P-wave charmonia
7 DAs, Phys.Rev.D79:074004,(2009)

$$\langle v^2 \rangle_{1S} = 0.21 \pm 0.04 \quad \langle v^2 \rangle_{2S} = 0.54^{+0.33}_{-0.36} \quad \langle v^2 \rangle_{2P} = 0.30 \pm 0.10$$

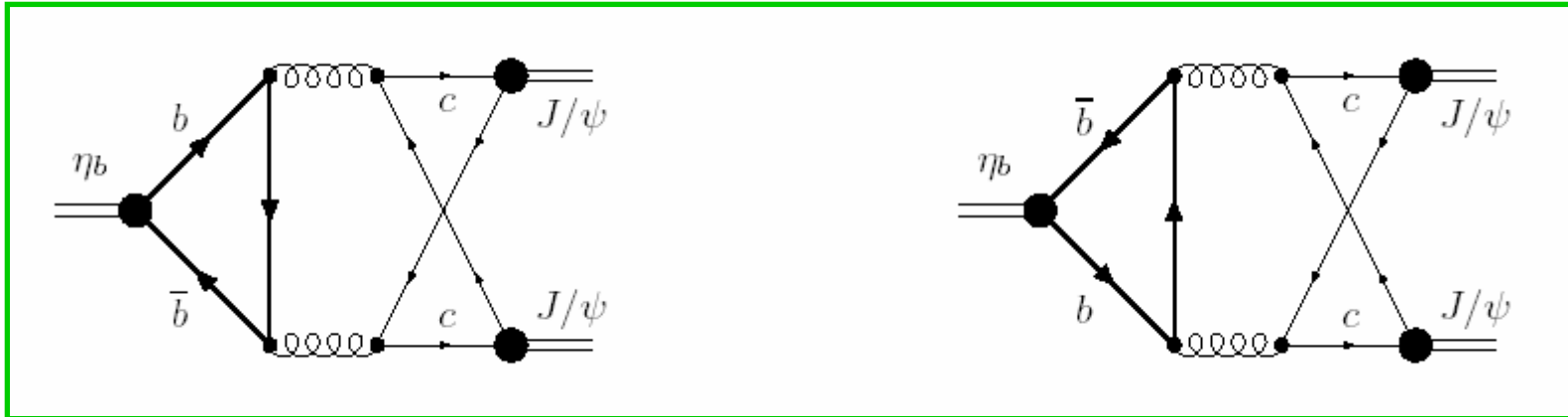
Charmonia are unique mesons since
one can build realistic models of
charmonia distribution amplitudes

$M_1 M_2$	$\Gamma_{\text{NRQCD}}, \text{eV}$	$\Gamma_{\text{LC}}, \text{eV}$	$\text{Br}_{\text{LC}}, \%$
$\eta_b \rightarrow h_c J/\psi$	$73_{-11}^{+15} \pm 38 \pm 37$	$520_{-43}^{+44} \pm 59 \pm 140$	0.0047
$\eta_b \rightarrow h_c \psi(2S)$	$39_{-6}^{+7} \pm 26 \pm 19$	$300_{-26}^{+28} \pm 33 \pm 78$	0.0027
$\eta_b \rightarrow \eta_c \chi_{c0}$	$56_{-14}^{+18} \pm 29 \pm 28$	$49_{-4.1}^{+4.1} \pm 25 \pm 71$	4.5×10^{-4}
$\eta_b \rightarrow \eta_c(2S) \chi_{c0}$	$30_{-7.4}^{+9.3} \pm 20 \pm 15$	$26_{-2.4}^{+2.4} \pm 18 \pm 37$	2.3×10^{-4}
$\eta_b \rightarrow \eta_c \chi_{c2}$	$17_{-5.2}^{+5.2} \pm 8.9 \pm 8.6$	$97_{-8.1}^{+8.2} \pm 48 \pm 71$	8.9×10^{-4}
$\eta_b \rightarrow \eta_c(2S) \chi_{c2}$	$9.1_{-2.7}^{+2.6} \pm 6.1 \pm 4.6$	$51_{-4.6}^{+4.6} \pm 35 \pm 37$	4.6×10^{-4}
$\eta_b \rightarrow \chi_{c0} \chi_{c1}$	$11_{-2}^{+2} \pm 8.4 \pm 5.5$	$26_{-2.2}^{+2.2} \pm 13 \pm 38$	2.4×10^{-4}
$\eta_b \rightarrow \chi_{c1} \chi_{c2}$	$4.7_{-1.1}^{+1.2} \pm 3.6 \pm 2.4$	$51_{-4.3}^{+4.4} \pm 26 \pm 37$	4.7×10^{-4}
$\chi_{b0} \rightarrow \eta_c \chi_{c1}$	$12_{-2.4}^{+2.4} \pm 6.1 \pm 5.8$	$53_{-2.3}^{+2.4} \pm 26 \pm 17$	0.0064
$\chi_{b0} \rightarrow \eta_c(2S) \chi_{c1}$	$6.2_{-1.3}^{+1.3} \pm 4.2 \pm 3.1$	$39_{-6.1}^{+5.2} \pm 27 \pm 13$	0.0048
$\chi_{b0} \rightarrow \chi_{c0} \chi_{c2}$	$(5.9_{-4.1}^{+4.2} \pm 4.5 \pm 2.9) \times 10^{-4}$	$0.81_{-0.18}^{+0.19} \pm 0.41 \pm 1.8$	10^{-4}
$\chi_{b0} \rightarrow \eta_c \eta_c$	$49_{-8.6}^{+11} \pm 35 \pm 25$	$51_{-1.9}^{+2.1} \pm 25 \pm 16$	0.0062
$\chi_{b0} \rightarrow \eta_c \eta_c(2S)$	$26_{-4.6}^{+5.7} \pm 17 \pm 13$	$76_{-11}^{+9.7} \pm 52 \pm 24$	0.0093
$\chi_{b0} \rightarrow \eta_c(2S) \eta_c(2S)$	$14_{-2.6}^{+2.5} \pm 15 \pm 6.9$	$30_{-5.4}^{+7.2} \pm 25 \pm 9.6$	0.0036
$\chi_{b0} \rightarrow J/\psi J/\psi$	$27_{-4.2}^{+5} \pm 19 \pm 13$	$79_{-2.9}^{+3.2} \pm 2.7 \pm 25$	0.0096
$\chi_{b0} \rightarrow J/\psi \psi(2S)$	$14_{-2.2}^{+2.6} \pm 9.3 \pm 7.1$	$130_{-19}^{+16} \pm 4 \pm 41$	0.016
$\chi_{b0} \rightarrow \psi(2S) \psi(2S)$	$7.5_{-1.2}^{+1.4} \pm 8.3 \pm 3.8$	$54_{-15}^{+13} \pm 1.5 \pm 17$	0.0066
$\chi_{b0} \rightarrow h_c h_c$	$0.094_{-0.029}^{+0.023} \pm 0.072 \pm 0.047$	$15_{-3.3}^{+3.5} \pm 2.3 \pm 3.9$	0.0018
$\chi_{b0} \rightarrow \chi_{c0} \chi_{c0}$	$0.048_{-0.034}^{+0.061} \pm 0.037 \pm 0.024$	$0.21_{-0.046}^{+0.048} \pm 0.11 \pm 1.1$	2.5×10^{-3}
$\chi_{b0} \rightarrow \chi_{c1} \chi_{c1}$	$0.56_{-0.16}^{+0.26} \pm 0.43 \pm 0.28$	$14_{-0.67}^{+0.66} \pm 6.8 \pm 15$	0.0017
$\chi_{b0} \rightarrow \chi_{c2} \chi_{c2}$	$0.026_{-0.013}^{+0.034} \pm 0.02 \pm 0.013$	$0.81_{-0.18}^{+0.19} \pm 0.4 \pm 1.8$	9.9×10^{-5}
$\chi_{b1} \rightarrow h_c J/\psi$	$1.2_{-0.19}^{+0.22} \pm 0.64 \pm 0.61$	$14_{-1.7}^{+1.7} \pm 1.6 \pm 4.6$	8.9×10^{-4}
$\chi_{b1} \rightarrow h_c \psi(2S)$	$0.65_{-0.1}^{+0.12} \pm 0.44 \pm 0.32$	$15_{-4.4}^{+4.4} \pm 1.7 \pm 4.7$	9.2×10^{-4}
$\chi_{b1} \rightarrow \eta_c \chi_{c0}$	$0.24_{-0.05}^{+0.061} \pm 0.13 \pm 0.12$	$1.3_{-0.18}^{+0.16} \pm 0.67 \pm 0.43$	8.3×10^{-5}
$\chi_{b1} \rightarrow \eta_c(2S) \chi_{c0}$	$0.13_{-0.027}^{+0.027} \pm 0.087 \pm 0.064$	$1.3_{-0.36}^{+0.35} \pm 0.89 \pm 0.41$	8×10^{-5}
$\chi_{b1} \rightarrow \eta_c \chi_{c2}$	$0.74_{-0.12}^{+0.13} \pm 0.38 \pm 0.37$	$2.7_{-0.31}^{+0.33} \pm 1.3 \pm 0.42$	1.7×10^{-4}
$\chi_{b1} \rightarrow \eta_c(2S) \chi_{c2}$	$0.39_{-0.063}^{+0.069} \pm 0.26 \pm 0.19$	$2.5_{-0.75}^{+0.69} \pm 1.8 \pm 0.41$	1.6×10^{-4}
$\chi_{b1} \rightarrow \chi_{c0} \chi_{c1}$	$0.6_{-0.18}^{+0.21} \pm 0.45 \pm 0.3$	$0.68_{-0.086}^{+0.088} \pm 0.35 \pm 3.4$	4.3×10^{-5}
$\chi_{b1} \rightarrow \chi_{c1} \chi_{c2}$	$0.13_{-0.019}^{+0.024} \pm 0.098 \pm 0.064$	$1.4_{-0.17}^{+0.17} \pm 0.68 \pm 0.43$	8.4×10^{-5}
$\chi_{b2} \rightarrow \eta_c \chi_{c1}$	$1.8_{-0.29}^{+0.31} \pm 0.92 \pm 0.88$	$17_{-0.52}^{+0.53} \pm 8.2 \pm 2.7$	0.0069
$\chi_{b2} \rightarrow \eta_c(2S) \chi_{c1}$	$0.93_{-0.15}^{+0.17} \pm 0.63 \pm 0.47$	$11_{-1.3}^{+1.1} \pm 7.6 \pm 1.8$	0.0046
$\chi_{b2} \rightarrow \chi_{c0} \chi_{c2}$	$0.58_{-0.17}^{+0.21} \pm 0.44 \pm 0.29$	$1.4_{-0.2}^{+0.21} \pm 0.72 \pm 6.1$	5.8×10^{-4}
$\chi_{b2} \rightarrow \eta_c \eta_c$	$1.6_{-0.49}^{+0.5} \pm 1.1 \pm 0.78$	$16_{-0.43}^{+0.47} \pm 7.8 \pm 10$	0.0066
$\chi_{b2} \rightarrow \eta_c(2S) \eta_c$	$0.83_{-0.26}^{+0.26} \pm 0.54 \pm 0.41$	$21_{-2.5}^{+2.1} \pm 15 \pm 14$	0.0088
$\chi_{b2} \rightarrow \eta_c(2S) \eta_c(2S)$	$0.44_{-0.14}^{+0.14} \pm 0.49 \pm 0.22$	$8_{-2}^{+1.7} \pm 6.7 \pm 5.1$	0.0033
$\chi_{b2} \rightarrow J/\psi J/\psi$	$65_{-12}^{+14} \pm 46 \pm 32$	$270_{-9.1}^{+10} \pm 41 \pm 93$	0.11
$\chi_{b2} \rightarrow \psi(2S) J/\psi$	$34_{-6.2}^{+7.6} \pm 23 \pm 17$	$380_{-53}^{+45} \pm 120 \pm 130$	0.16
$\chi_{b2} \rightarrow \psi(2S) \psi(2S)$	$18_{-3.3}^{+4} \pm 20 \pm 9.1$	$140_{-53}^{+34} \pm 59 \pm 50$	0.059
$\chi_{b2} \rightarrow h_c h_c$	$0.42_{-0.1}^{+0.11} \pm 0.32 \pm 0.21$	$37_{-4.2}^{+4.4} \pm 8 \pm 24$	0.015
$\chi_{b2} \rightarrow \chi_{c0} \chi_{c0}$	$0.014_{-0.0038}^{+0.0035} \pm 0.011 \pm 0.0072$	$0.36_{-0.051}^{+0.053} \pm 0.18 \pm 0.11$	1.5×10^{-4}
$\chi_{b2} \rightarrow \chi_{c1} \chi_{c1}$	$0.18_{-0.059}^{+0.057} \pm 0.14 \pm 0.089$	$8_{-0.83}^{+0.86} \pm 4.7 \pm 4$	0.0033
$\chi_{b2} \rightarrow \chi_{c2} \chi_{c2}$	$0.2_{-0.049}^{+0.043} \pm 0.15 \pm 0.098$	$9.8_{-1.6}^{+1.9} \pm 7.7 \pm 0.8$	0.004
$\chi_{b2} \rightarrow h_c J/\psi$	$7.6_{-1.5}^{+1.5} \pm 4 \pm 3.8$	$100_{-4.2}^{+4.4} \pm 42 \pm 24$	0.042
$\chi_{b2} \rightarrow h_c \psi(2S)$	$4_{-0.76}^{+0.79} \pm 2.7 \pm 2$	$71_{-10}^{+8.6} \pm 36 \pm 17$	0.029
$\chi_{b2} \rightarrow \chi_{c1} \chi_{c2}$	$0.32_{-0.039}^{+0.037} \pm 0.24 \pm 0.16$	$11_{-2.1}^{+2.2} \pm 9.5 \pm 0.047$	0.0045

Braguta, Likhoded, Luchinsky, Phys.Rev.D80:094008,2009

$n_b \rightarrow J/\psi J/\psi$ decay

$\eta_b \rightarrow J/\psi J/\psi$ decay



The amplitude of the process

$$T = F e_{\mu\nu\rho\sigma} p_1^\mu \varepsilon_1^\nu p_2^\mu \varepsilon_2^\nu$$

F is the only formfactor of the process

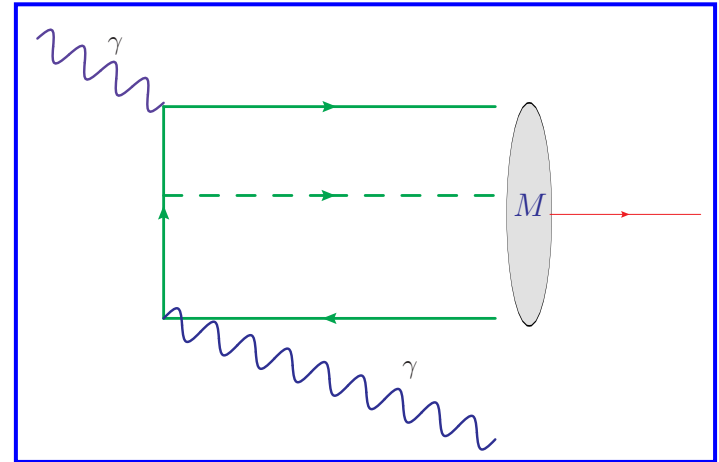
- Helicity conservation + angular momentum conservation $\Rightarrow \Lambda_{J/\psi}=0$
- Helicity flip in gluon-quark-quark vertex leads to the suppression of the amplitude

$\eta_b \rightarrow J/\psi J/\psi$ is a next-to-next-to-leading twist process

Higher twist distribution amplitudes

Amplitude :

$$T = \underbrace{\frac{a_n}{E_h^n}}_{\text{twist } -2} + \underbrace{\frac{a_{n+1}}{E_h^{n+1}}}_{\substack{\text{twist } -2, \\ \text{twist } -3}} + \underbrace{\frac{a_{n+2}}{E_h^{n+2}}}_{\substack{\text{twist } -2, \\ \text{twist } -3, \\ \text{twist } -4}} + \dots$$



Expansion in Fock states

$$|Meson\rangle = \underbrace{|QQ\rangle}_{\substack{\text{twist } -2, \\ \text{twist } -3, \\ \dots}} + \underbrace{|QQg\rangle}_{\substack{\text{twist } -3, \\ \text{twist } -4, \\ \dots}} + \underbrace{|QQgg\rangle}_{\substack{\text{twist } -4, \\ \text{twist } -5, \\ \dots}} + \underbrace{|QQq\bar{q}\rangle}_{\substack{\text{twist } -4, \\ \text{twist } -5, \\ \dots}} + \dots$$

Higher twist decays as a probe of charmonia structures

Distribution amplitudes of J/Ψ up to twist-4

$$\begin{aligned} \langle J/\psi(p, \epsilon) | \bar{c}(x) \gamma_\rho [x, -x] c(-x) | 0 \rangle &= f_V M_V \left[\frac{(\epsilon x)}{(px)} p_\rho \int_{-1}^1 d\xi e^{i\xi(px)} (\varphi_1(\xi, \mu) + \frac{M_V^2 x^2}{4} \varphi_2(\xi, \mu)) \right. \\ &+ (\epsilon_\rho - p_\rho \frac{(\epsilon x)}{(px)}) \int_{-1}^1 d\xi e^{i\xi(px)} \varphi_3(\xi, \mu) \\ &\left. - \frac{1}{2} x_\rho \frac{(\epsilon x)}{(px)^2} M_V^2 \int_{-1}^1 d\xi e^{i\xi(px)} \varphi_4(\xi, \mu) \right], \end{aligned}$$

$$\begin{aligned} \langle J/\psi(p, \epsilon) | \bar{c}(x) \sigma_{\rho\lambda} [x, -x] c(-x) | 0 \rangle &= f_T(\mu) \left[(\epsilon_\rho p_\lambda - \epsilon_\lambda p_\rho) \int_{-1}^1 d\xi e^{i\xi(px)} (\chi_1(\xi, \mu) + \frac{M_V^2 x^2}{4} \chi_2(\xi, \mu)) \right. \\ &+ (p_\rho x_\lambda - p_\lambda x_\rho) \frac{(\epsilon x)}{(px)^2} M_V^2 \int_{-1}^1 d\xi e^{i\xi(px)} \chi_3(\xi, \mu) \\ &\left. + \frac{1}{2} (\epsilon_\rho x_\lambda - \epsilon_\lambda x_\rho) \frac{M_V^2}{(px)} \int_{-1}^1 d\xi e^{i\xi(px)} \chi_4(\xi, \mu) \right], \end{aligned}$$

$$\langle J/\psi(p, \epsilon) | \bar{c}(x) \gamma_\rho \gamma_5 [x, -x] c(-x) | 0 \rangle = f_A(\mu) e_{\rho\lambda\alpha\beta} \epsilon^\lambda p^\alpha x^\beta \int_{-1}^1 d\xi e^{i\xi(px)} \Phi_1(\xi, \mu),$$

$$\langle J/\psi(p, \epsilon) | \bar{c}(x) [x, -x] c(-x) | 0 \rangle = -i f_S(\mu) (\epsilon x) \int_{-1}^1 d\xi e^{i\xi(px)} \Phi_2(\xi, \mu),$$

There are 10 distribution amplitudes needed in the calculation

The result of the calculation

$$F = \int d\xi_1 d\xi_2 H(\xi_1, \xi_2, \mu) \left(f_V f_A(\mu) M_{J/\psi} \varphi_1(\xi_1, \mu) \Phi_1(\xi_2, \mu) + f_V f_A(\mu) M_{J/\psi} \varphi_1(\xi_2, \mu) \Phi_1(\xi_1, \mu) \right. \\ \left. + f_S(\mu) f_T(\mu) \chi_1(\xi_2, \mu) \Phi_2(\xi_1, \mu) + f_S(\mu) f_T(\mu) \chi_1(\xi_1, \mu) \Phi_2(\xi_2, \mu) \right). \\ H(\xi_1, \xi_2, \mu) = \frac{1024\pi^2 \alpha_s^2(\mu)}{27} f_{\eta_b} \frac{1}{M_{\eta_b}^6} \frac{1}{(1-\xi_1^2)(1-\xi_2^2)(1+\xi_1\xi_2)},$$

$$\langle J/\psi(p, \epsilon) | \bar{c}(0) \gamma_\mu c(0) | 0 \rangle = f_V M_{J/\psi} \epsilon_\mu, \quad \langle J/\psi(p, \epsilon) | \bar{c}(0) \sigma_{\mu\nu} c(0) | 0 \rangle = f_T(\mu) (\epsilon_\mu p_\nu - \epsilon_\nu p_\mu) \\ f_A(\mu) = \frac{1}{2} \left(f_V - f_T(\mu) \frac{2m_c(\mu)}{M_{J/\psi}} \right) M_{J/\psi}, \quad f_S(\mu) = \left(f_T(\mu) - f_V \frac{2m_c(\mu)}{M_{J/\psi}} \right) M_{J/\psi}^2$$

Fine tuning between parameters at the leading order approximation in NRQCD

$$\frac{f_T}{f_V} = 1 - \frac{\langle v^2 \rangle_{J/\psi}}{3}, \quad \frac{M_{J/\psi}}{2m_c} = 1 + \frac{\langle v^2 \rangle_{J/\psi}}{2}, \quad \langle v^2 \rangle_{J/\psi} \sim 0.2 \\ F = \frac{256\pi^2 \alpha_s^2}{81} \frac{1}{m_b^6} f_{\eta_b} f_V^2 m_c^2 \langle v^2 \rangle$$

Fine tuning is broken due to relativistic and radiative corrections what leads to the dramatic enhancement of the branching ratio

Numerical results

$$\begin{aligned} Br(\eta_b \rightarrow J/\psi J/\psi) &= (6.2 \pm 3.5) \times 10^{-7}, \\ Br(\eta_b \rightarrow J/\psi \psi') &= (10 \pm 6) \times 10^{-7}, \\ Br(\eta_b \rightarrow \psi' \psi') &= (3.7 \pm 2.8) \times 10^{-7}. \end{aligned}$$

Braguta, Kartvelishvili,
Phys.Rev.D81:014012,2010

$$Br(\eta_b \rightarrow J/\Psi J/\Psi) = (0.5 \times 10^{-8} - 1.2 \times 10^{-5}) \quad \text{Santorelli, Phys.Rev.D77:074012,2008}$$

$$Br(\eta_b \rightarrow J/\Psi J/\Psi) = 2.4_{-1.9}^{+4.2} \times 10^{-8} \quad \text{Jia, Phys.Rev.D78:054003,2008}$$

$$Br(\eta_b \rightarrow J/\Psi J/\Psi) = (2.1 - 18.6) \times 10^{-8} \quad \text{Gong et al., Phys.LettB670:350,2009}$$

Approximately 100 events of the $\eta_b \rightarrow J/\Psi J/\Psi$ decay at LHC per year

Thank you
