

Zeroing in on Supersymmetric Radiation Amplitude Zeros

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SUSY11

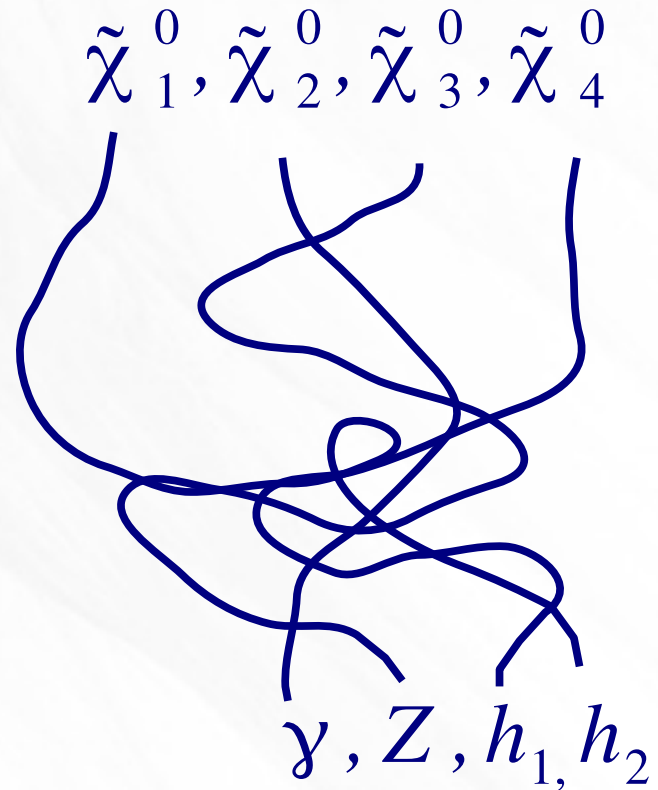
based on arXiv:1109.XXXX
(JoAnne Hewett, AI, Tom Rizzo)

Outline

- Motivation: neutralinos in the MSSM
- Radiation amplitude zeros (RAZ)
- Observing supersymmetric RAZ at the LHC

Motivation

- The **neutralino** sector of the Minimal Supersymmetric Standard Model has a rich phenomenology that is largely dependent on the mixing of the gaugino and Higgsino eigenstates



Motivation

- As the LHC makes more neutralinos, it will become important to know their properties
- While specific breaking scenarios predict various mixing arrangements for the neutralinos, **bottom-up** approaches to determining their content are few and far between (Kane et al., 1105.3742; Allanach et al., 1010.4261; Tata, talk on Sunday)
- How to untangle neutralino mixing in a model-independent fashion?

Radiation Amplitude Zeros

- Vanishing amplitudes in specific regions of phase space for processes with external gauge bosons
- RAZ were first seen, unexpectedly, in the calculation of $d\bar{u} \rightarrow W\gamma$ production

Magnetic Moment of Weak Bosons Produced in pp and $p\bar{p}$ Collisions

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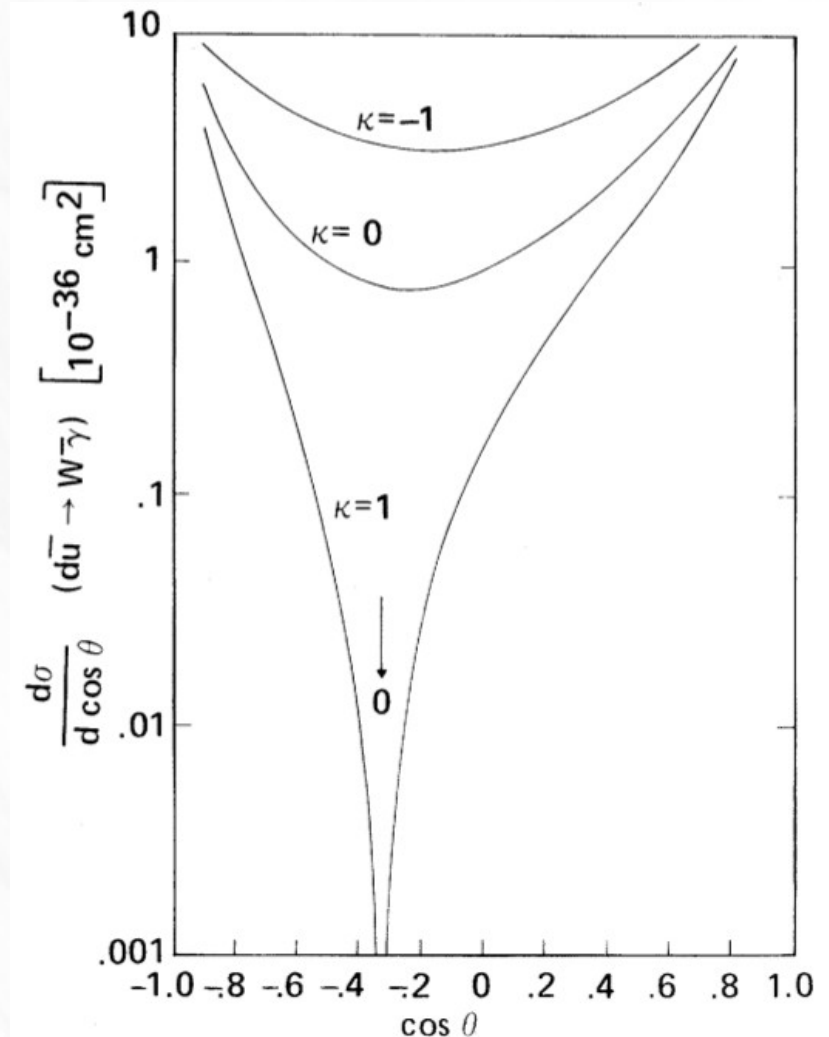
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(Received 5 June 1979)

We suggest that the reactions $pp \rightarrow W^\pm \gamma X$ and $p\bar{p} \rightarrow W^\pm \gamma X$ are good candidates for measuring the magnetic moment parameter κ in $\mu_W = (e/2M_W)(1 + \kappa)$. The angular distribution of the W bosons in $p\bar{p} \rightarrow W^\pm \gamma X$ is particularly sensitive to this parameter. For the gauge-theory value of $\kappa = 1$, we have found a peculiar zero in $d\sigma(d\bar{u} \rightarrow W^- \gamma)/d \cos\theta$ at $\cos\theta = -\frac{1}{3}$, the location of this zero depending on the quark charge through $\cos\theta = -(1 + 2Q_d)$. A similar zero occurs in $d\sigma(u\bar{d} \rightarrow W^+ \gamma)/d \cos\theta$. We can offer no explanation for this behavior.

Radiation Amplitude Zeros

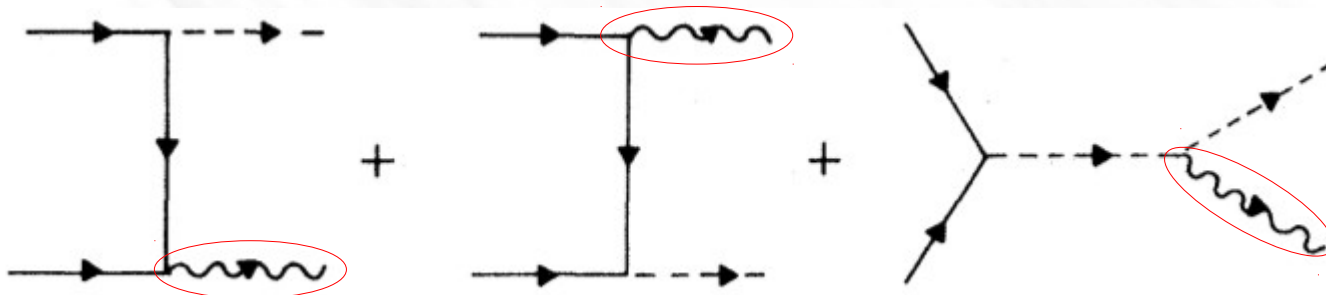
- This amplitude zero has been studied at the **Tevatron**
- Presence of the RAZ at $\cos \theta = -\frac{1}{3}$ allows us to constrain **anomalous couplings** of the W



Why do RAZ happen?

- Diagrams for processes with external photons can be thought of as arising from attaching photons to a simpler diagram in all possible places
- Adding a photon introduces a factor to a diagram which only depends on the spin structure, charges, and momenta

$$M(d\bar{u} \rightarrow W \gamma) \sim f(Q_i, p_i^\mu) \cdot M(d\bar{u} \rightarrow W)$$



Why do RAZ happen?

- In specific regions of phase space, this extra factor **vanishes** since it corresponds to a Lorentz transformation (Brown, Kowalski, Brodsky)
- For 2-2 scattering with a photon in the final state, the kinematic region with an amplitude zero is given by $\cos \theta = \frac{Q_1 - Q_2}{Q_1 + Q_2}$
- e.g. for **W** γ production, we have $Q_d = -\frac{1}{3}$, $Q_{\bar{u}} = -\frac{2}{3}$

Why do RAZ happen?

- Amplitude zeros also happen for processes with external gauge bosons other than photons

$$\cos \theta = \frac{g_1 - g_2}{g_1 + g_2}$$

- For massive gauge bosons, RAZ are only approximate, because of production of longitudinally polarized bosons
- e.g. $q_1 \bar{q}_2 \rightarrow W Z$ production (Baur, Han, Ohnemus)

RAZ in Supersymmetry

- Given the zero in $W\gamma$ production, it's natural to expect a zero in **chargino-neutralino production**
- But the existence of a RAZ here depends on the **neutralino composition!** (need RAZ at $\cos\theta < 1$)

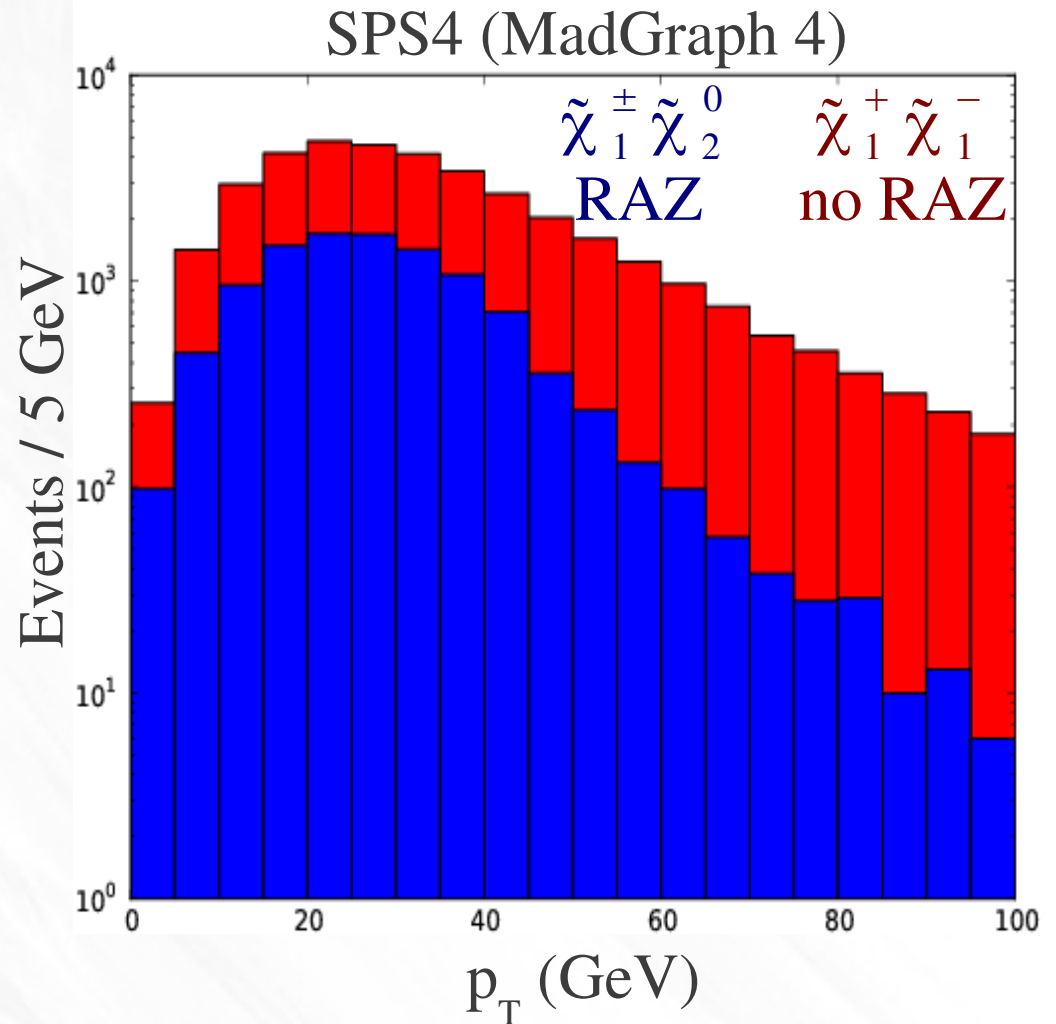
	\tilde{W}^\pm	\tilde{H}^\pm
\tilde{B}	No physical RAZ	No tree level diagrams
\tilde{W}^3	RAZ at $\cos\theta^* = 0$	No tree level diagrams
$\tilde{H}_{1,2}$	No tree level diagrams	No physical RAZ

RAZ in Supersymmetry

- There is no amplitude zero in chargino-neutralino production except when the neutralino is a wino, so we can use the **presence of a RAZ to probe the wino content of the neutralino**
- We now turn to methods of observing the zero experimentally in $\tilde{\chi}_1^\pm \tilde{\chi}_2^0 \rightarrow 3l + MET$
- Clean signal (WZ background), can't do with $\tilde{\chi}_1^\pm \tilde{\chi}_1^0$
- We assume 1 ab^{-1} of LHC data at 14 TeV; not for discovery but exploration after SUSY is found

Turning to Trileptons

- Consider chargino-neutralino production with **trilepton** decay
- When the neutralino is wino-like, we expect fewer events at low $\cos\theta$ because of the amplitude zero, and hence a **lower number of high- p_T leptons**



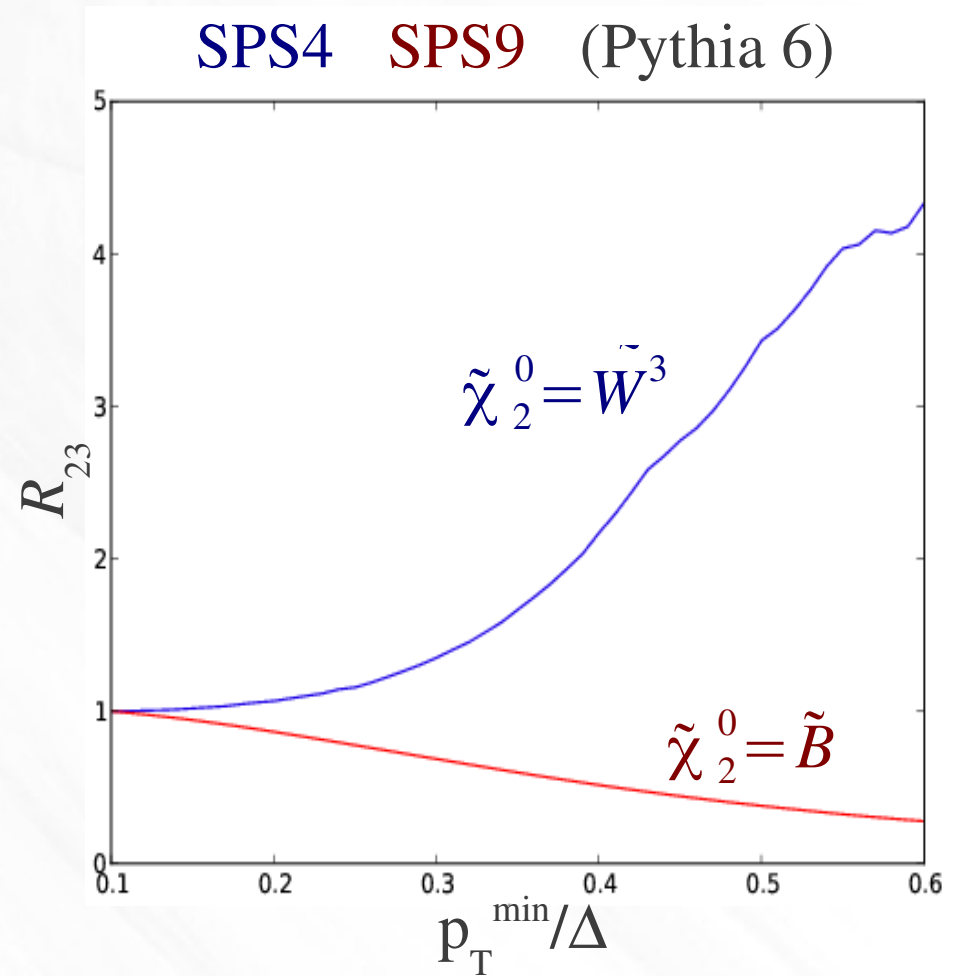
Turning to Trileptons

- We choose to look at the **softest lepton p_T** , seeing how many trilepton events from chargino-neutralino production remained as the minimum lepton p_T was increased
- To calibrate how fast the number of events drops with increasing p_T cut, we compare to a process that never exhibits an amplitude zero, chargino pair production

$$R_{23}(p_T) = \frac{N(\tilde{\chi}_1^+ + \tilde{\chi}_1^- \rightarrow 2l + MET)}{N(\tilde{\chi}_1^+ + \tilde{\chi}_2^0 \rightarrow 3l + MET)}$$

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- Events without ISR, FSR, or fragmentation (negligible effect)
- Require that leptons have rapidity < 2.5
- Δ is chargino-LSP mass splitting
- Ratio **rises** with increasing cut when neutralino is wino-like



A Potential Problem?

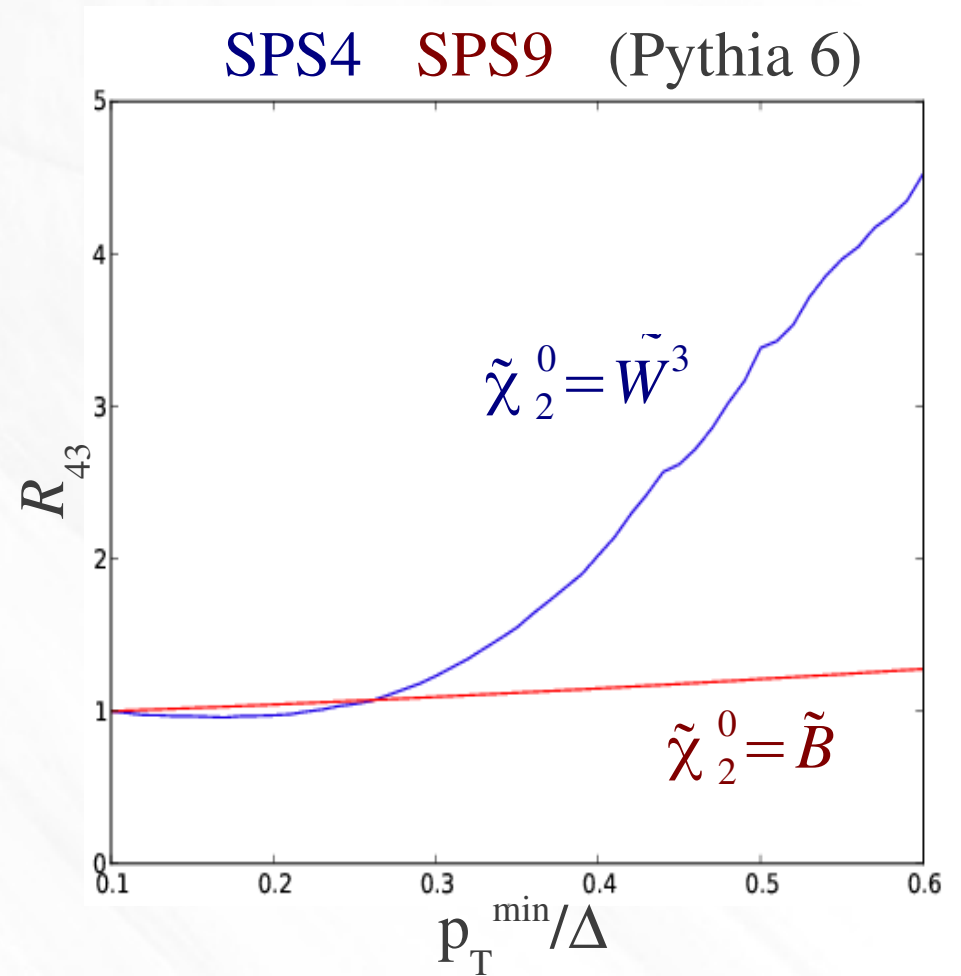
- When the chargino-LSP mass splitting is much greater than the neutralino-LSP mass splitting, R_{23} compares the spectra of leptons that are generally coming from different decays
- To account for different possible SUSY spectra, we also construct an analogous ratio comparing chargino-neutralino associated production to

neutralino pair production

$$R_{43}(p_T) = \frac{N(\tilde{\chi}_2^0 + \tilde{\chi}_2^0 \rightarrow 4l + MET)}{N(\tilde{\chi}_1^+ + \tilde{\chi}_2^0 \rightarrow 3l + MET)}$$

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- R_{23} and R_{43} complement each other
- For a wino-like neutralino, both ratios tend to rise sharply with increasing leptonic transverse momentum cut



Investigating Neutralino Mixing

- The more leptons there are in the final state, the more events a transverse momentum cut removes

$$R_{23}(p_T) = \frac{N(\tilde{\chi}_1^+ + \tilde{\chi}_1^- \rightarrow 2l + MET)}{N(\tilde{\chi}_1^+ + \tilde{\chi}_2^0 \rightarrow 3l + MET)}$$

$$R_{43}(p_T) = \frac{N(\tilde{\chi}_2^0 + \tilde{\chi}_2^0 \rightarrow 4l + MET)}{N(\tilde{\chi}_1^+ + \tilde{\chi}_2^0 \rightarrow 3l + MET)}$$

- **Low $R_{23} \rightarrow \text{low } |N_{22}|$**
- **High $R_{43} \rightarrow \text{high } |N_{22}|$**

$$N_{22} = \text{wino content of } \tilde{\chi}_2^0$$

Using the pMSSM

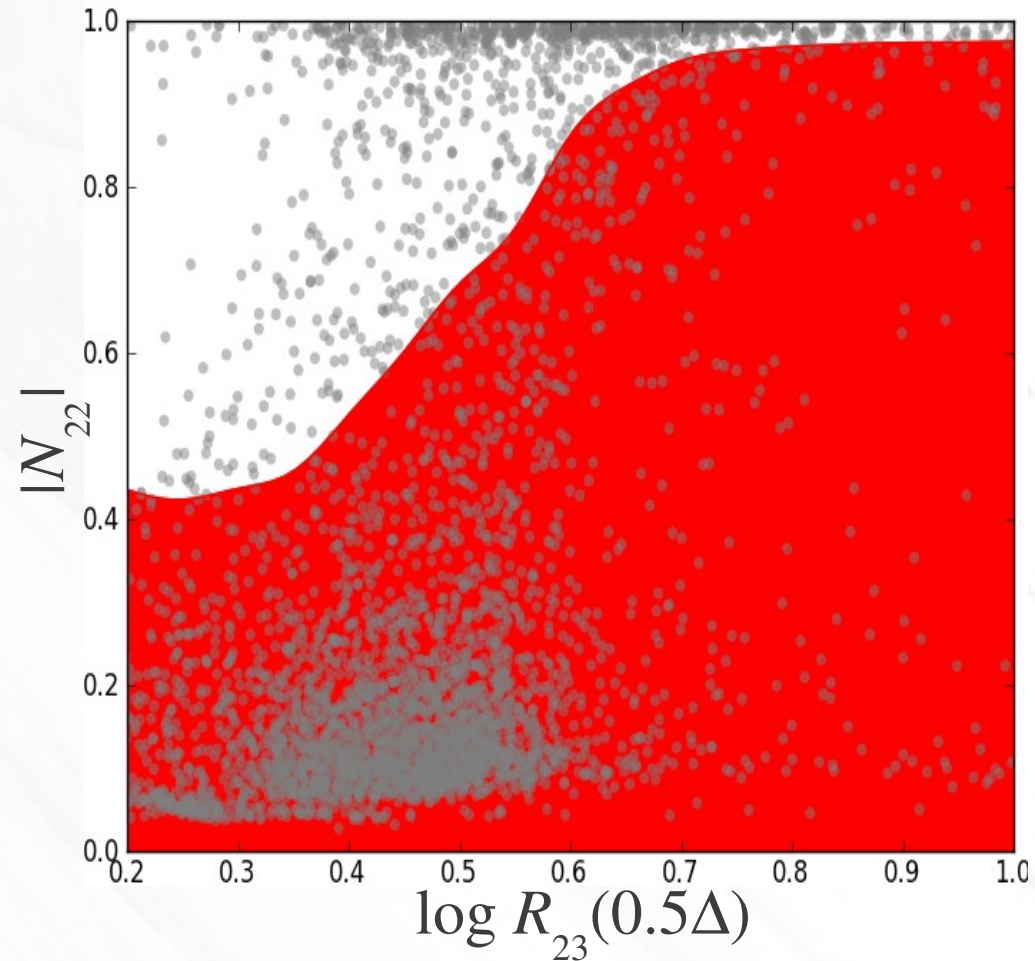
- To test our techniques, we applied them to the results of a previous scan of the **phenomenological MSSM**, which imposes only those constraints on the full MSSM parameter space that are experimentally motivated, e.g. minimal flavor violation
- This large set of pMSSM models contains a **wide variety of neutralino mixing schemes**, providing an ideal testing ground for our method
- For more details, see Berger et al., arXiv:0812.0980

Using the pMSSM

- $50 \text{ GeV} \leq |M_{1,2}, \mu| \leq 1 \text{ TeV}$
- $100 \text{ GeV} \leq M_3 \leq 1 \text{ TeV}$
- $1 \leq \tan \beta \leq 50$
- $43.5 \text{ GeV} \leq m_A \leq 1 \text{ TeV}$
- $100 \text{ GeV} \leq m_f \leq 1 \text{ TeV}$
- $|A_{t,b,\tau}| \leq 1 \text{ TeV}$
- After applying many theoretical and experimental constraints, ~ 70000 models that were still viable before the LHC
- Here, we only consider models with chargino-LSP mass splittings **above 50 GeV** ($\sim 16\%$)

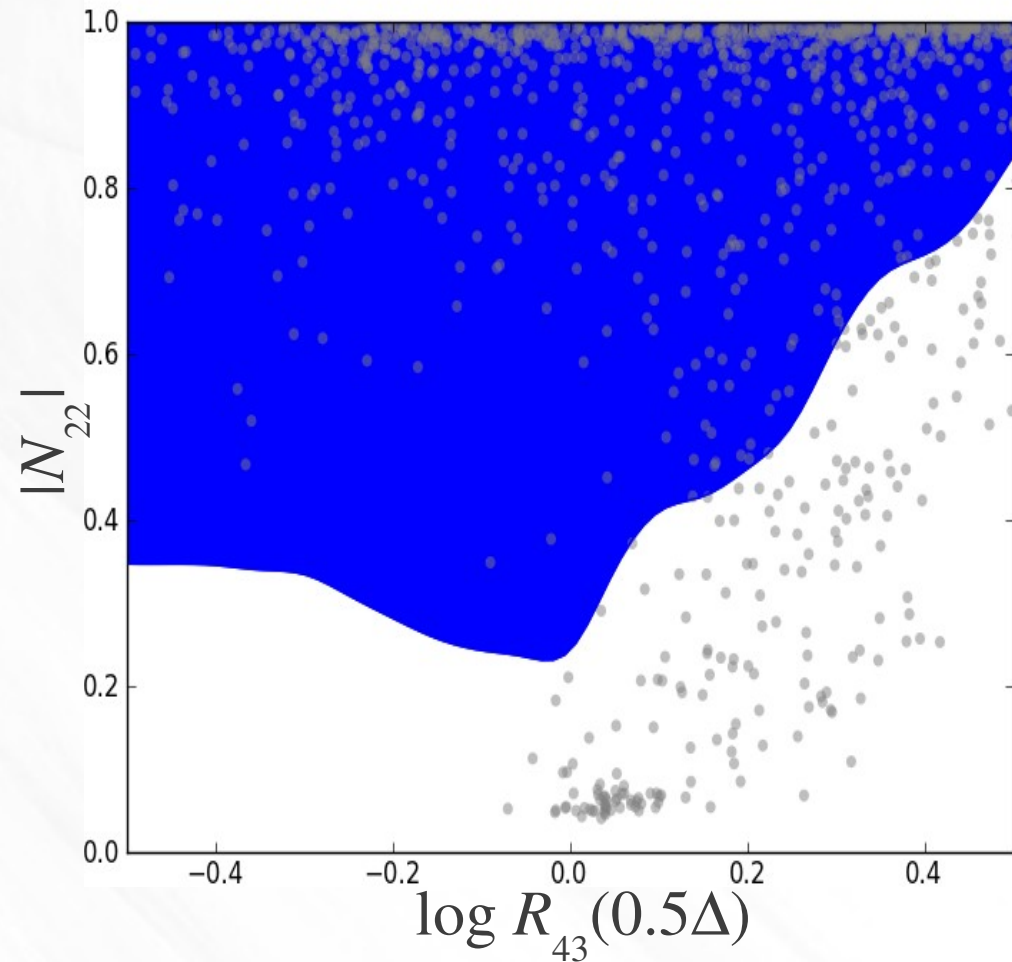
Results from the pMSSM

- Only use models that should have sufficient statistics with 1 ab^{-1} (~ 4400) at LHC-14
- Within our model set, the 90% upper limits on the wino content of the second neutralino are indicated by the shaded region



Results from the pMSSM

- Second neutralino tends to be heavier than lightest chargino in the model set, so only ~ 1200 models for looking at R_{43}
- R_{23} and R_{43} are useful for setting upper and lower bounds on the wino-ness of the second neutralino, respectively



Summary

- **Radiation amplitude zeros**, having previously been used to test the Standard Model, show promise in probing the coupling structure of supersymmetry
- Looking at the potential effects of a RAZ in associated **chargino-neutralino production**, we have demonstrated the power of a technique to investigate the neutralino mixing matrix
- As the LHC ramps up, RAZ should prove useful once again in testing SUSY gauge theory