



Decaying Dark Matter from Dark Instantons

Reinard Primulando

The College of William and Mary


Work with Chris Carone and Josh Erlich

Phys.Rev. D **82**, 055028 (2010)

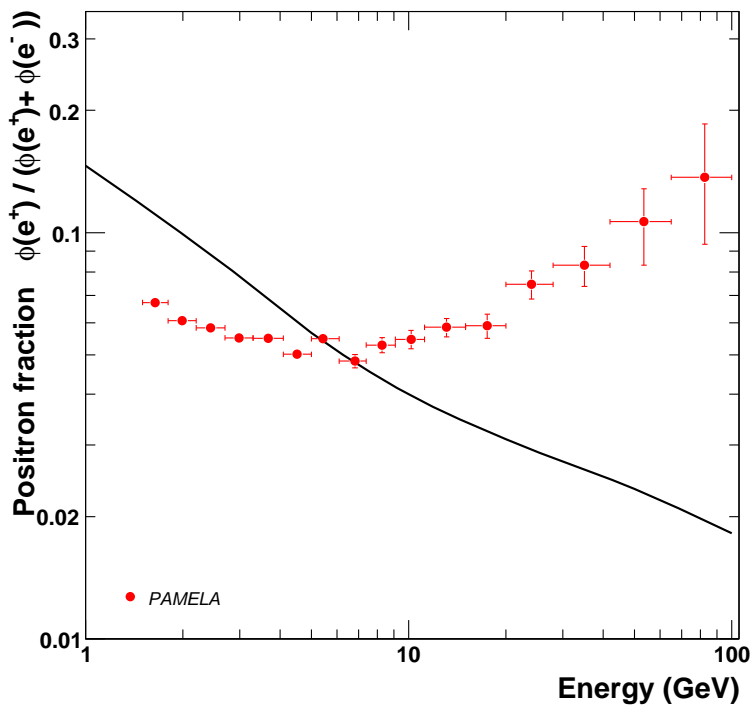




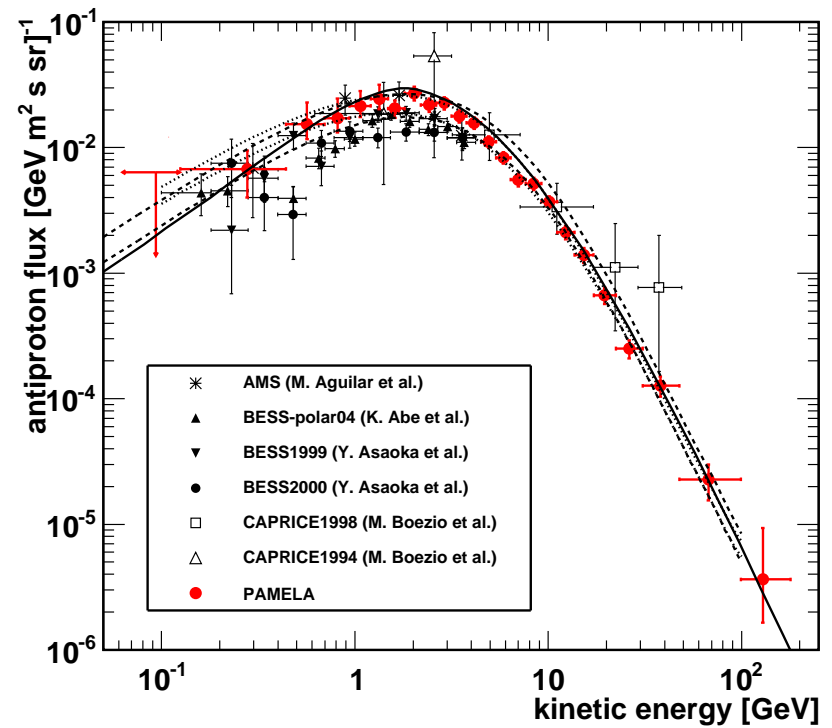
Outline of the Talk

- 
- ⑥ Cosmic ray positron excess from PAMELA and FERMI-LAT data
 - ⑥ Decaying dark matter
 - ⑥ Dark instantons and decaying dark matter
 - ⑥ Conclusions

Cosmic Rays Anomalies

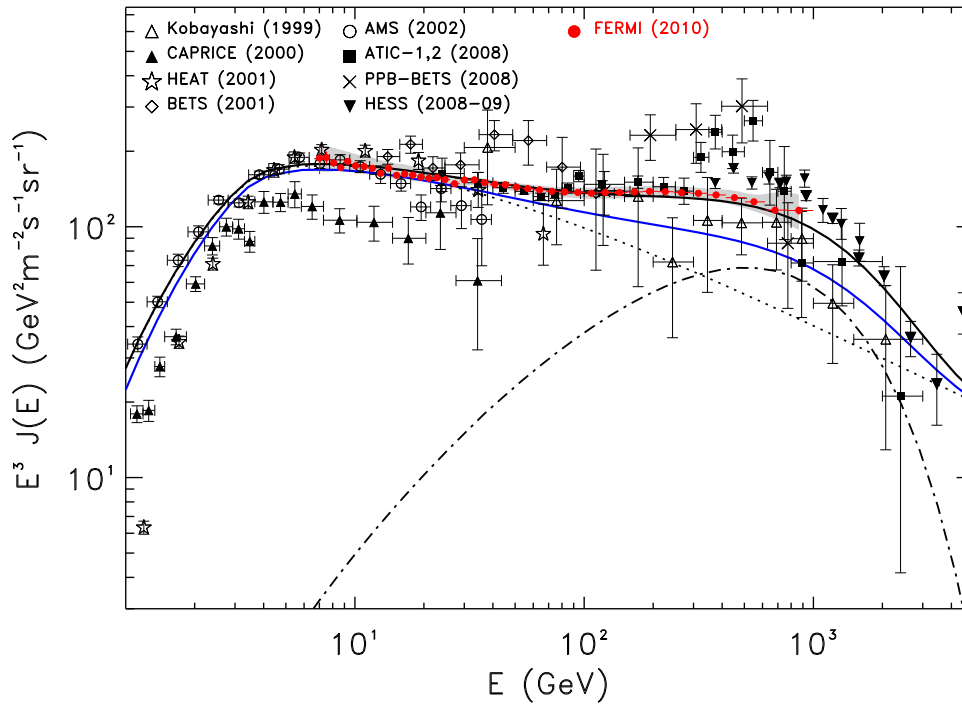


PAMELA positron fraction



PAMELA antiproton flux

Cosmic Rays Anomalies



FERMI-LAT positron + electron flux

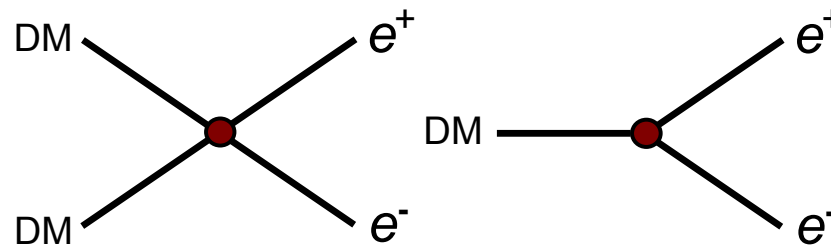
Possible explanations for the anomalies

- ⑥ Astrophysical: Pulsar.

Hooper, et.al. *JCAP*0901,025(2009); Yuksel, et.al. *Phys.Rev.Lett.*103,051101(2009).

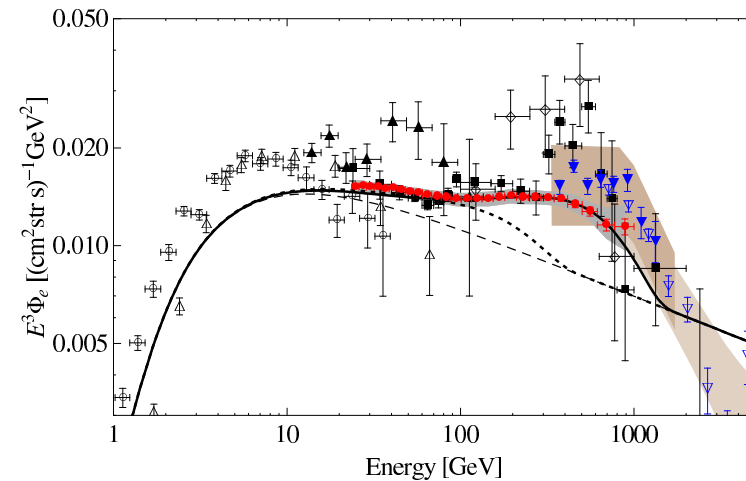
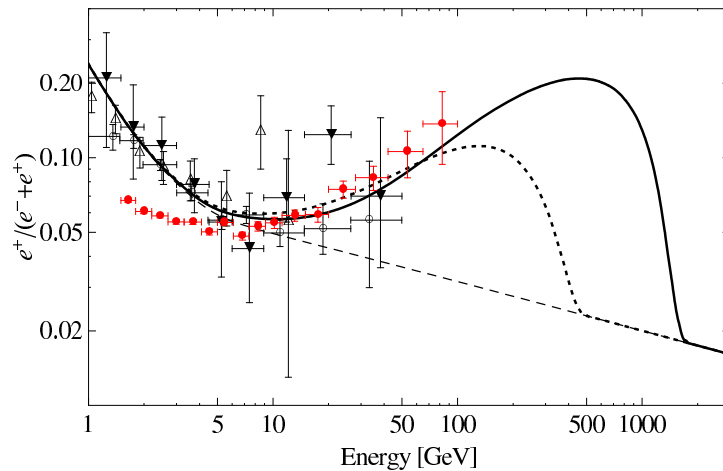
- ⑥ Dark matter annihilating/decaying into positrons.

For a review: Fan, et.al. *Int.J.Mod.Phys.D*19,2011(2010)



Decaying Dark Matter

- Decaying dark matter can explain PAMELA positron excess.
[Ibarra, et.al. JCAP1001,009\(2010\) and references therein.](#)
- $DM \rightarrow l^+ l^-, l^+ l^- \nu; DM \not\rightarrow 2q, Wl$ etc...
- $\mathcal{O}(\text{TeV})$ dark matter.
- The lifetime of DM is $\mathcal{O}(10^{26})$ s.

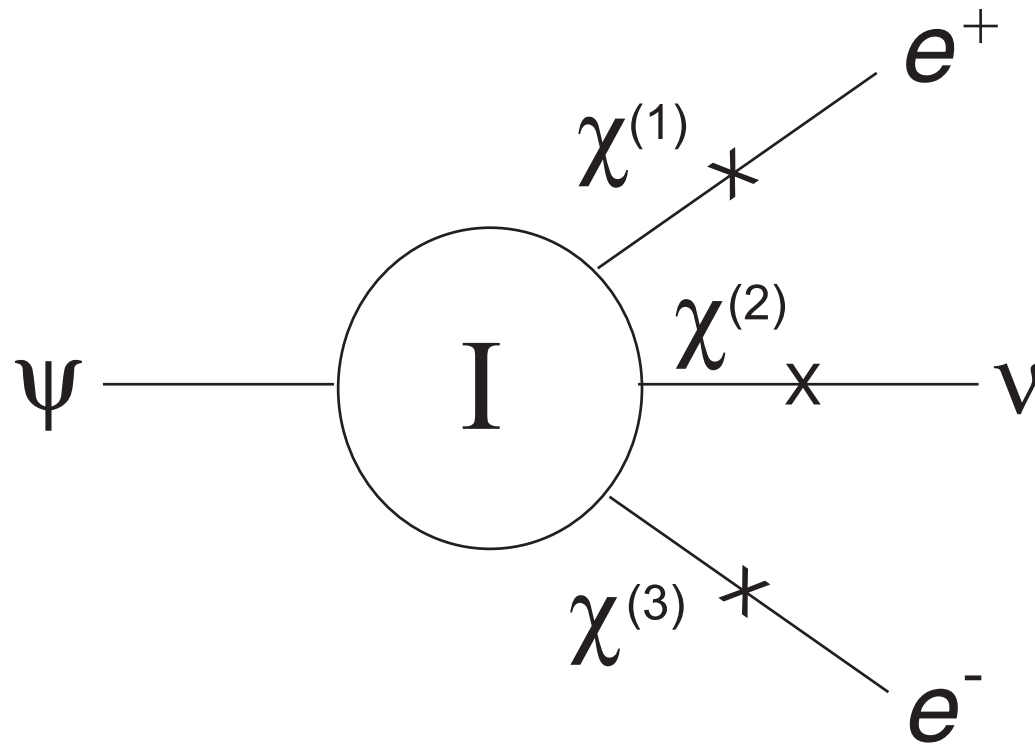


Instanton

- ⑥ Consider a set of N_f massless fermions, ψ_s^α , transform in fundamental representation of $SU(N)$.
- ⑥ The global symmetry of the model is $SU(N_f)_L \times SU(N_f)_R \times U(1)_V \times U(1)_A$.
- ⑥ 't Hooft pointed out that an instanton-induced operator violates the anomalous $U(1)_A$ global symmetry.
['t Hooft, Phys.Rev.Lett.37,8\(1976\)](#)
- ⑥ The effective operator is given by

$$\mathcal{L}^{eff} = C g^{-8} e^{-\frac{8\pi^2}{g^2}} (2\delta_{\alpha_1\beta_1}\delta_{\alpha_2\beta_2} - \delta_{\alpha_1\beta_2}\delta_{\alpha_2\beta_1}) \epsilon^{st} \times \bar{\psi}_1^{\alpha_1} (1 + \gamma^5) \psi_s^{\beta_1} \bar{\psi}_2^{\alpha_2} (1 + \gamma^5) \psi_t^{\beta_2} + \text{h.c.}$$

The Model



The Model

- The gauge group is $G_{SM} \times SU(2)_D \times U(1)_D$.
- New particles charge assignments:

ψ_L	$(\mathbf{2}, -1/2)_0$	ψ_{uR}, ψ_{dR}	$(\mathbf{1}, -1/2)_0$
$\chi_L^{(1)}$	$(\mathbf{2}, +1/6)_+$	$\chi_{uR}^{(1)}, \chi_{dR}^{(1)}$	$(\mathbf{1}, +1/6)_+$
$\chi_L^{(2)}$	$(\mathbf{2}, +1/6)_0$	$\chi_{uR}^{(2)}, \chi_{dR}^{(2)}$	$(\mathbf{1}, +1/6)_0$
$\chi_L^{(3)}$	$(\mathbf{2}, +1/6)_-$	$\chi_{uR}^{(3)}, \chi_{dR}^{(3)}$	$(\mathbf{1}, +1/6)_-$
E_L	$(\mathbf{1}, 0)_-$	E_R	$(\mathbf{1}, 0)_-$
H_D	$(\mathbf{2}, 0)_0$	η	$(\mathbf{1}, 1/6)_0$

$U(1)_\psi$ **Accidental Symmetry**

- With the charge assignments, there is an accidental global $U(1)_\psi$ symmetry.

Combination	$U(1)_\psi$	$U(1)_D$
$\bar{\chi}\psi$	+1	-1/3
$\bar{\chi}\psi^c$	-1	+2/3

ψ_L	$(\mathbf{2}, -1/2)_0$	ψ_{uR}, ψ_{dR}	$(\mathbf{1}, -1/2)_0$
$\chi_L^{(1)}$	$(\mathbf{2}, +1/6)_+$	$\chi_{uR}^{(1)}, \chi_{dR}^{(1)}$	$(\mathbf{1}, +1/6)_+$
$\chi_L^{(2)}$	$(\mathbf{2}, +1/6)_0$	$\chi_{uR}^{(2)}, \chi_{dR}^{(2)}$	$(\mathbf{1}, +1/6)_0$
$\chi_L^{(3)}$	$(\mathbf{2}, +1/6)_-$	$\chi_{uR}^{(3)}, \chi_{dR}^{(3)}$	$(\mathbf{1}, +1/6)_-$
E_L	$(\mathbf{1}, 0)_-$	E_R	$(\mathbf{1}, 0)_-$
H_D	$(\mathbf{2}, 0)_0$	η	$(\mathbf{1}, 1/6)_0$

$U(1)_\psi$ **Accidental Symmetry**

- ⑥ With the charge assignments, there is an accidental global $U(1)_\psi$ symmetry.

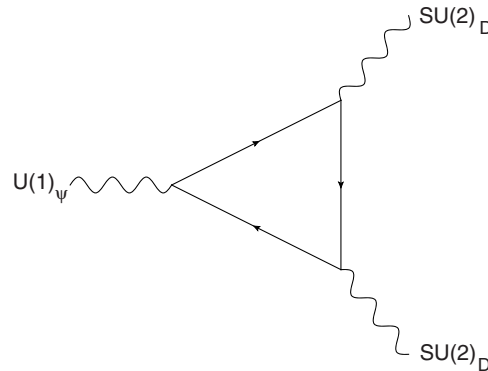
Combination	$U(1)_\psi$	$U(1)_D$
$\bar{\chi}\psi$	+1	-1/3
$\bar{\chi}\psi^c$	-1	+2/3

ψ_L	$(\mathbf{2}, -1/2)_0$	ψ_{uR}, ψ_{dR}	$(\mathbf{1}, -1/2)_0$
$\chi_L^{(1)}$	$(\mathbf{2}, +1/6)_+$	$\chi_{uR}^{(1)}, \chi_{dR}^{(1)}$	$(\mathbf{1}, +1/6)_+$
$\chi_L^{(2)}$	$(\mathbf{2}, +1/6)_0$	$\chi_{uR}^{(2)}, \chi_{dR}^{(2)}$	$(\mathbf{1}, +1/6)_0$
$\chi_L^{(3)}$	$(\mathbf{2}, +1/6)_-$	$\chi_{uR}^{(3)}, \chi_{dR}^{(3)}$	$(\mathbf{1}, +1/6)_-$
E_L	$(\mathbf{1}, 0)_-$	E_R	$(\mathbf{1}, 0)_-$
H_D	$(\mathbf{2}, 0)_0$	η	$(\mathbf{1}, 1/6)_0$

- ⑥ Decay of the lightest ψ is prevented in all order of perturbation theory.

Instanton Induced Operator

Triangle anomaly for $SU(2)_D^2 U(1)_\psi$.



$$\text{tr}[t^a t^b N_\psi] = \frac{1}{2} \delta^{ab} \sum N_\psi$$

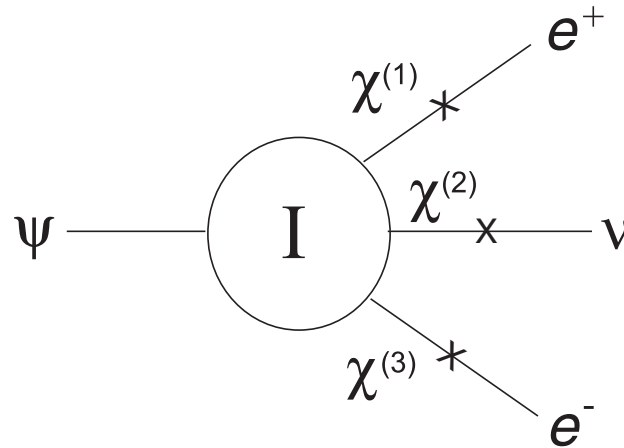
Instanton Induced Operator

Following 't Hooft

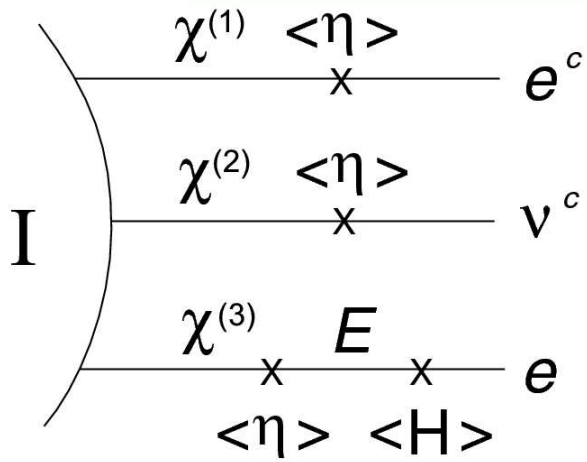
't Hooft, Phys.Rev.D14,3432(1976)

$$\mathcal{L}_I = \frac{C}{6 g_D^8} \exp\left(-\frac{8\pi^2}{g_D^2}\right) \left(\frac{m_\psi}{v_D}\right)^{35/6} \frac{1}{v_D^2} (2 \delta_{\alpha\beta} \delta_{\gamma\sigma} - \delta_{\alpha\sigma} \delta_{\beta\gamma})$$

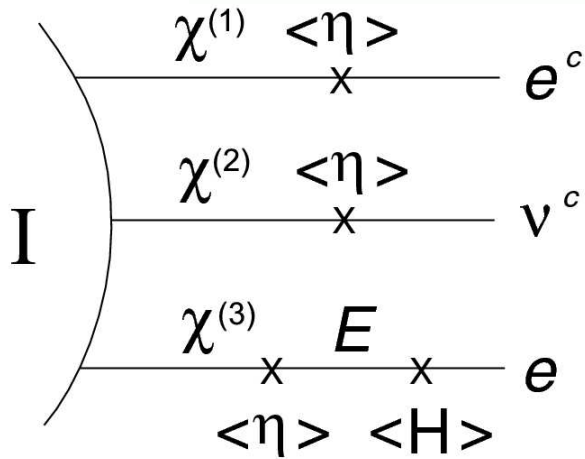
$$\cdot \left[\overline{(\chi_{L\beta}^{(2)c} \psi_L^\alpha)} \overline{(\chi_{L\sigma}^{(1)c} \chi_L^{(3)\gamma})} - \overline{(\chi_{L\beta}^{(1)c} \psi_L^\alpha)} \overline{(\chi_{L\sigma}^{(2)c} \chi_L^{(3)\gamma})} \right] + \text{h.c.}$$



Leptons- χ Mixings



Leptons- χ Mixings



Example: $\mathcal{L} \supset \eta \bar{\chi}_{dR}^{(1)} e_R^c$.

ψ_L	$(\mathbf{2}, -1/2)_0$	ψ_{uR}, ψ_{dR}	$(\mathbf{1}, -1/2)_0$
$\chi_L^{(1)}$	$(\mathbf{2}, +1/6)_+$	$\chi_{uR}^{(1)}, \chi_{dR}^{(1)}$	$(\mathbf{1}, +1/6)_+$
$\chi_L^{(2)}$	$(\mathbf{2}, +1/6)_0$	$\chi_{uR}^{(2)}, \chi_{dR}^{(2)}$	$(\mathbf{1}, +1/6)_0$
$\chi_L^{(3)}$	$(\mathbf{2}, +1/6)_-$	$\chi_{uR}^{(3)}, \chi_{dR}^{(3)}$	$(\mathbf{1}, +1/6)_-$
E_L	$(\mathbf{1}, 0)_-$	E_R	$(\mathbf{1}, 0)_-$
H_D	$(\mathbf{2}, 0)_0$	η	$(\mathbf{1}, 1/6)_0$

Dark Matter Lifetime

- ⑥ The decay width of the dark matter is

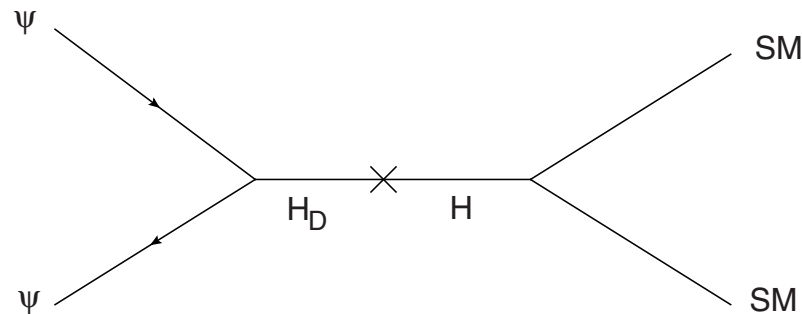
$$\Gamma(\psi \rightarrow \ell^+ \ell^- \nu) \approx \frac{1}{g_D^{16}} \exp(-16\pi^2/g_D^2) \left(\frac{m_\psi}{v_D}\right)^{47/3} m_\psi .$$

- ⑥ For $m_\psi = 3.5$ TeV and $v_D = 4$ TeV, we need $g_D \approx 1.15$.

Relic Density

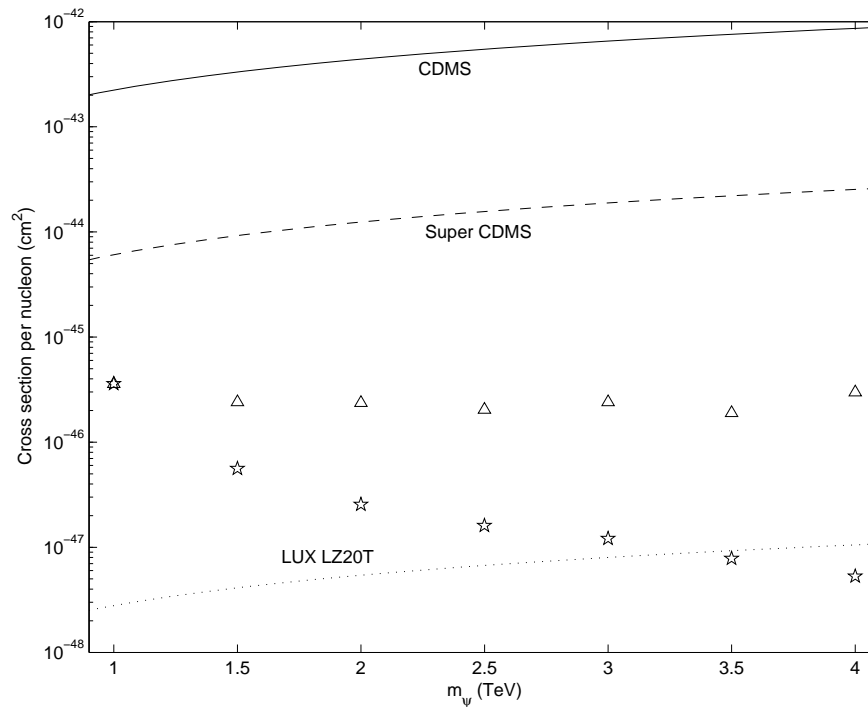
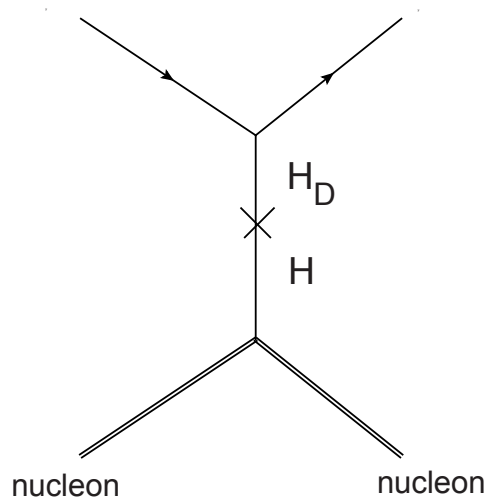
- Dark matter thermal equilibrium in the early universe is maintained by Higgs portal.

$$V = -\mu^2 H^\dagger H + \lambda (H^\dagger H)^2 - \mu_D^2 H_D^\dagger H_D + \lambda_D (H_D^\dagger H_D)^2 + \lambda_{mix} (H^\dagger H) (H_D^\dagger H_D).$$



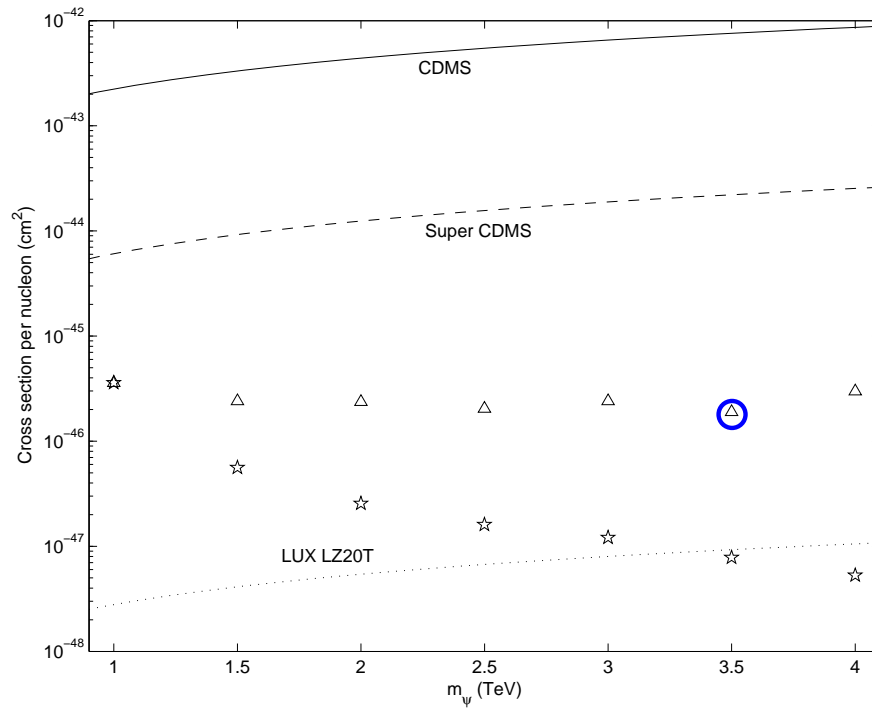
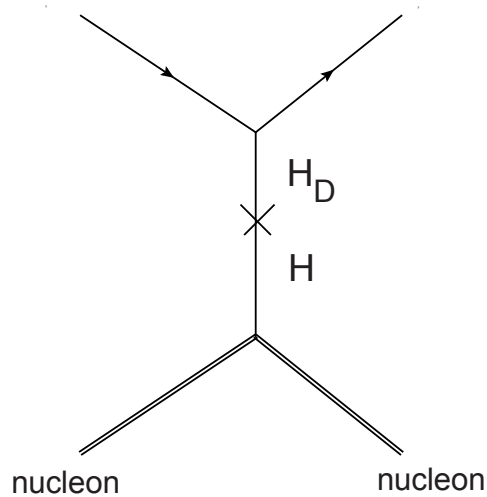
Direct Detection

- ⑥ We found that DM-nucleon cross section is $\mathcal{O}(10^{-46})\text{cm}^2$.
- ⑥ Three orders of magnitude lower than the current bound.



Direct Detection

- ⑥ We found that DM-nucleon cross section is $\mathcal{O}(10^{-46})\text{cm}^2$.
- ⑥ Three orders of magnitude lower than the current bound.



Conclusions

- ⑥ PAMELA positron excess can be explained by decaying dark matter.
- ⑥ The lifetime of the dark matter is $\mathcal{O}(10^{26})$ s.
- ⑥ Instanton-induced operator can explain the long lifetime of the dark matter.
- ⑥ We constructed a model of decaying dark matter inspired by instanton-induced operator.
- ⑥ The model gives the correct dark matter relic density and satisfies direct detection bounds.



Thank You

