



A MultiTeV Muon Collider: The Physics Case

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Physics Landscape

- Matter fields interact with massless spin one gauge bosons. (The Gauge Principle)

- Global Symmetry \Rightarrow Gauged Local Symmetry

$$\psi \rightarrow \exp\{iQ\} \psi$$

$$\psi(x) \rightarrow \exp\{i e \Lambda(x)\} \psi(x) = G \psi(x); A^\mu(x) \rightarrow G A^\mu(x) G^{-1} - (1/ie)[\partial^\mu, G]G^{-1}$$

$$L = \bar{\psi}(-i \gamma \cdot \partial + M) \psi$$

$$L = \bar{\psi}(-i \gamma \cdot D + M) \psi - (1/4e^2)(F^{\mu\nu})(F_{\mu\nu})$$

$$\text{with } D^\mu = [\partial^\mu + ie A^\mu(x)] \text{ and } (F^{\mu\nu}) = i[D^\mu, D^\nu]$$

- QED: Charged particles ψ (matter fields) interact with the photon A^μ (gauge particle)

- The Standard Model (SM) is based on this principle.

- QCD - SU(3) gauge interactions:

- color octet gluons (g) and color triplet quarks (u, d, s, c, b, t)_(L,R)
- Confinement \rightarrow physical states color singlets.

- Electroweak - $SU(2)_L \times U(1)$ gauge interactions:

- $SU(2)_L$ triplet gauge bosons: (W^\pm, W^0) and a $U(1)_Y$ gauge boson B
- quarks: $SU(2)_L$ doublets: $(u_L, d_L), (c_L, s_L), (t_L, b_L)$; and singlets: q_R
- leptons $SU(2)_L$ doublets: $(\nu_e, e^-_L), (\nu_\mu, \mu^-_L), (\nu_\tau, \tau^-_L)$; and singlets l_R
- No fermion or W, Z masses unless gauge symmetry spontaneously broken.

- Gravity is also a gauge theory. Hope to eventually unify of all four forces: Strong, Electromagnetic, Weak and Gravity. [String Theory]



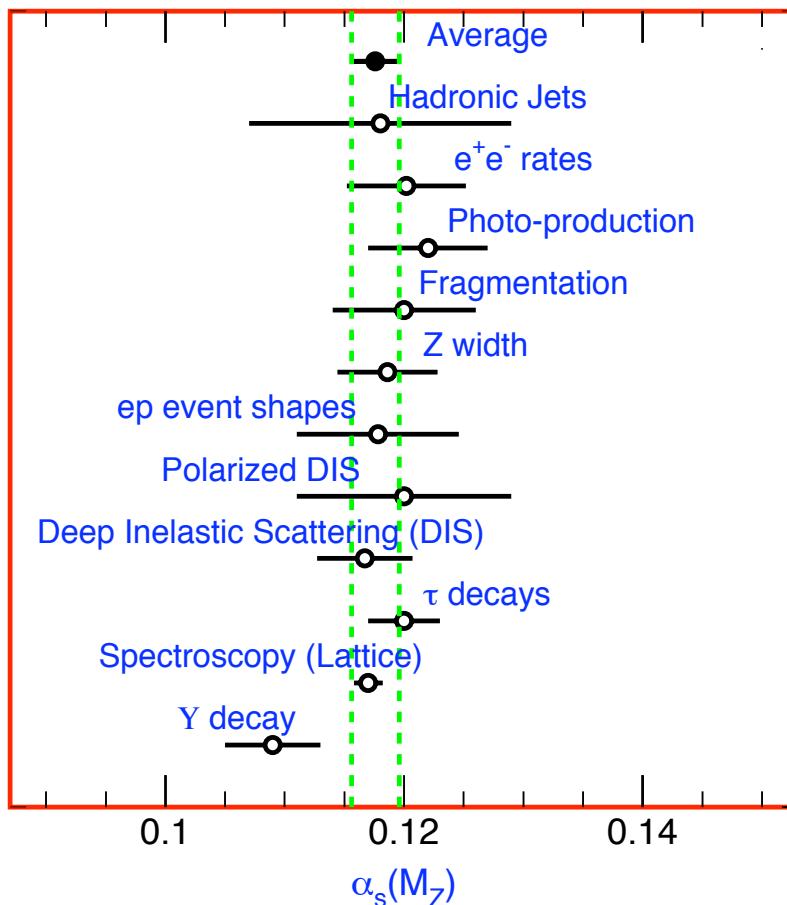
Physics Landscape

- Electroweak Symmetry Breaking (EWSB)
 - Introduce a $SU(2)_L$ complex doublet scalar field Φ , with self interactions
 - $\mu^2 (\Phi^\dagger \Phi) + \lambda (\Phi^\dagger \Phi)^2$ with EWSB $\rightarrow \langle \Phi^\dagger \Phi \rangle = v^2 = -\mu^2 / \lambda$;
one physical Higgs boson (mass $m_H^2 = 2\lambda v^2$)
 - Gauge interactions
 - $D^\mu \Phi^\dagger D_\mu \Phi$ with EWSB \rightarrow massive W^\pm, Z^0 and massless photon γ
 - Yukawa couplings to fermions
 - $\Gamma_{ij} \Psi_{iL}^\dagger \Psi_{jR} \Phi + h.c.$ with EWSB \rightarrow fermion masses and mixing of flavor eigenstates into mass eigenstates. CKM matrix for quarks.
- The Standard Model (SM) has been a spectacular success. For more than 30 years all new observations have fit naturally into this framework.
 - See figure and table
- Basic questions remain
 - There is as of yet no direct evidence for the Higgs boson or its interactions. Is this the correct mechanism for electroweak symmetry breaking?
 - How do the fermion masses and flavor mixings arise?



Physics Landscape

QCD



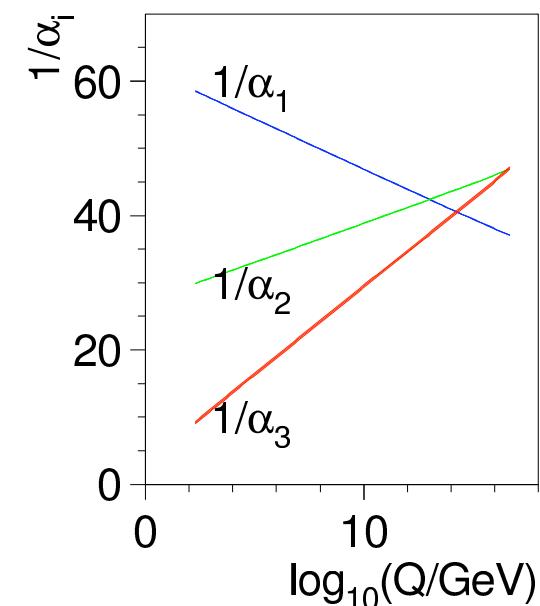
EW

Quantity	Value	Standard Model	Pull
m_t [GeV]	$172.7 \pm 2.9 \pm 0.6$	172.7 ± 2.8	0.0
M_W [GeV]	80.450 ± 0.058	80.376 ± 0.017	1.3
	80.392 ± 0.039		0.4
M_Z [GeV]	91.1876 ± 0.0021	91.1874 ± 0.0021	0.1
Γ_Z [GeV]	2.4952 ± 0.0023	2.4968 ± 0.0011	-0.7
$\Gamma(\text{had})$ [GeV]	1.7444 ± 0.0020	1.7434 ± 0.0010	—
$\Gamma(\text{inv})$ [MeV]	499.0 ± 1.5	501.65 ± 0.11	—
$\Gamma(\ell^+\ell^-)$ [MeV]	83.984 ± 0.086	83.996 ± 0.021	—
σ_{had} [nb]	41.541 ± 0.037	41.467 ± 0.009	2.0
R_e	20.804 ± 0.050	20.756 ± 0.011	1.0
R_μ	20.785 ± 0.033	20.756 ± 0.011	0.9
R_τ	20.764 ± 0.045	20.801 ± 0.011	-0.8
R_b	0.21629 ± 0.00066	0.21578 ± 0.00010	0.8
R_c	0.1721 ± 0.0030	0.17230 ± 0.00004	-0.1
$A_{FB}^{(0,e)}$	0.0145 ± 0.0025	0.01622 ± 0.00025	-0.7
$A_{FB}^{(0,\mu)}$	0.0169 ± 0.0013		0.5
$A_{FB}^{(0,\tau)}$	0.0188 ± 0.0017		1.5
$A_{FB}^{(0,b)}$	0.0992 ± 0.0016	0.1031 ± 0.0008	-2.4
$A_{FB}^{(0,c)}$	0.0707 ± 0.0035	0.0737 ± 0.0006	-0.8
$A_{FB}^{(0,s)}$	0.0976 ± 0.0114	0.1032 ± 0.0008	-0.5
$s_\ell^2(A_{FB}^{(0,g)})$	0.2324 ± 0.0012	0.23152 ± 0.00014	0.7
	0.2238 ± 0.0050		-1.5
A_e	0.15138 ± 0.00216	0.1471 ± 0.0011	2.0
	0.1544 ± 0.0060		1.2
	0.1498 ± 0.0049		0.6
A_μ	0.142 ± 0.015		-0.3
A_τ	0.136 ± 0.015		-0.7
	0.1439 ± 0.0043		-0.7
A_b	0.923 ± 0.020	0.9347 ± 0.0001	-0.6
A_c	0.670 ± 0.027	0.6678 ± 0.0005	0.1
A_s	0.895 ± 0.091	0.9356 ± 0.0001	-0.4
g_L^2	0.30005 ± 0.00137	0.30378 ± 0.00021	-2.7
g_R^2	0.03076 ± 0.00110	0.03006 ± 0.00003	0.6
$g_V^{\nu e}$	-0.040 ± 0.015	-0.0396 ± 0.0003	0.0
$g_A^{\nu e}$	-0.507 ± 0.014	-0.5064 ± 0.0001	0.0
A_{PV}	-1.31 ± 0.17	-1.53 ± 0.02	1.3
$Q_W(\text{Cs})$	-72.62 ± 0.46	-73.17 ± 0.03	1.2
$Q_W(\text{Tl})$	-116.6 ± 3.7	-116.78 ± 0.05	0.1
$\frac{\Gamma(b \rightarrow s\gamma)}{\Gamma(b \rightarrow Xe\nu)}$	$3.35^{+0.50}_{-0.44} \times 10^{-3}$	$(3.22 \pm 0.09) \times 10^{-3}$	0.3
$\frac{1}{2}(g_\mu - 2 - \frac{\alpha}{\pi})$	4511.07 ± 0.82	4509.82 ± 0.10	1.5
τ_τ [fs]	$290.89^{+12.2907}_{-0.58}$	291.87 ± 1.76	-0.4



Physics Landscape

- All data consistent with Standard Model - but it's incomplete
 - dark matter; neutrino masses and mixing \rightarrow new fields or interactions;
 - baryon asymmetry in the universe \rightarrow more CP violation
- Theoretical questions
 - The issue of naturalness and the origin of mass:
 - $\mu^2 (\Phi^\dagger \Phi) + \lambda (\Phi^\dagger \Phi)^2 + \Gamma_{ij} \Psi_{iL}^\dagger \Psi_{jR} \Phi + h.c.$
 - $m_H^2/M_{\text{planck}}^2 \approx 10^{-34}$ Hierarchy problem
 - vacuum stability
 - large range of fermion masses
 - Many theoretical speculations on how to solve the naturalness problems:
SUSY, New Dynamics, Extra Dimensions, ...
 - gauge unification \rightarrow new interactions;
 - gravity: strings and extra dimensions





Theorists' Grand View

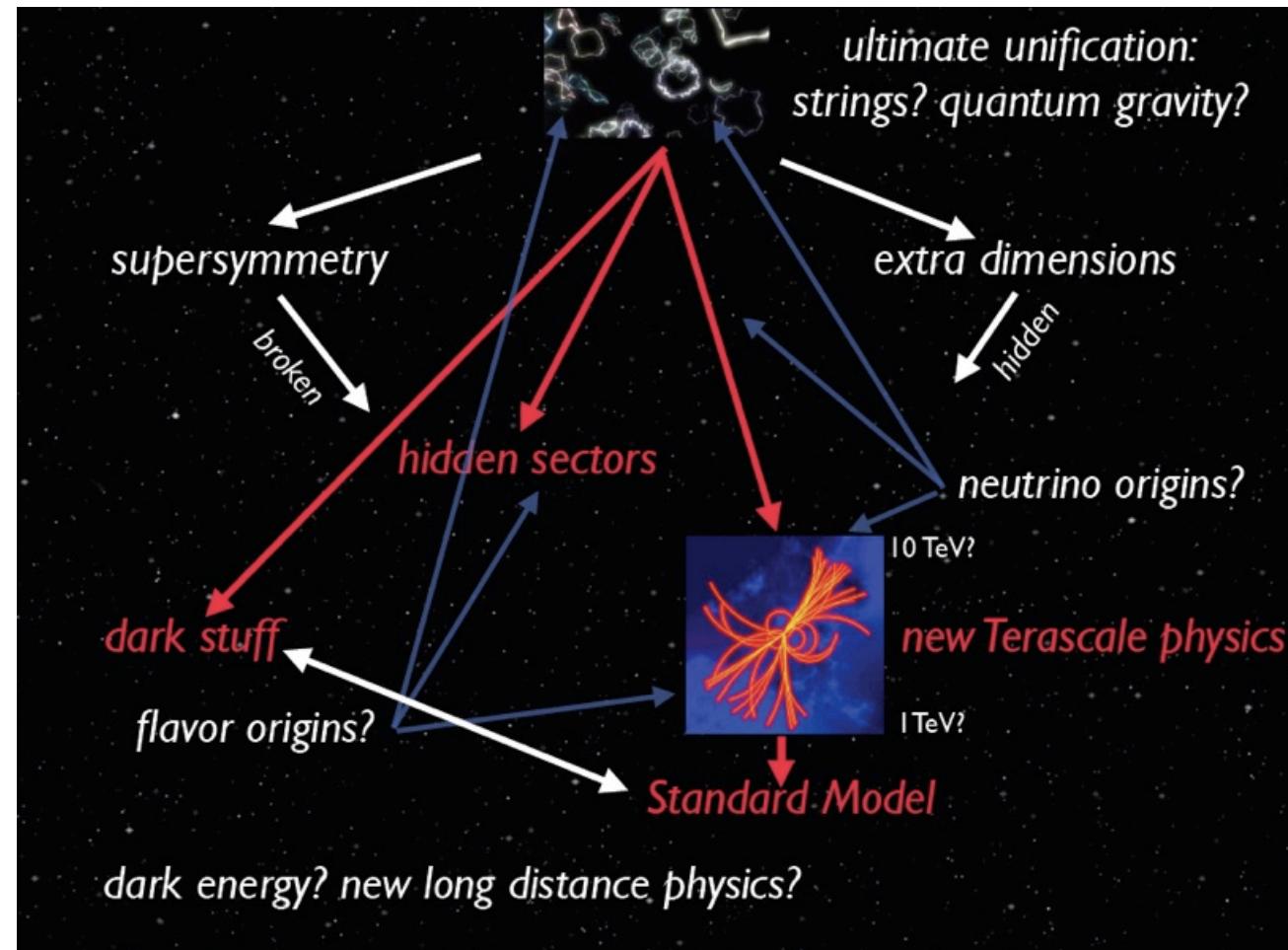
- New Scales and Symmetry Breaking

What is the origin and scale of fermion masses?

Physics Symmetry Scale

QCD confinement m_{glueball}
 χ SB m_{proton}

EW $SU(2)_L \times U_Y(1)$ M_w/g
--> $U_{EM}(1)$



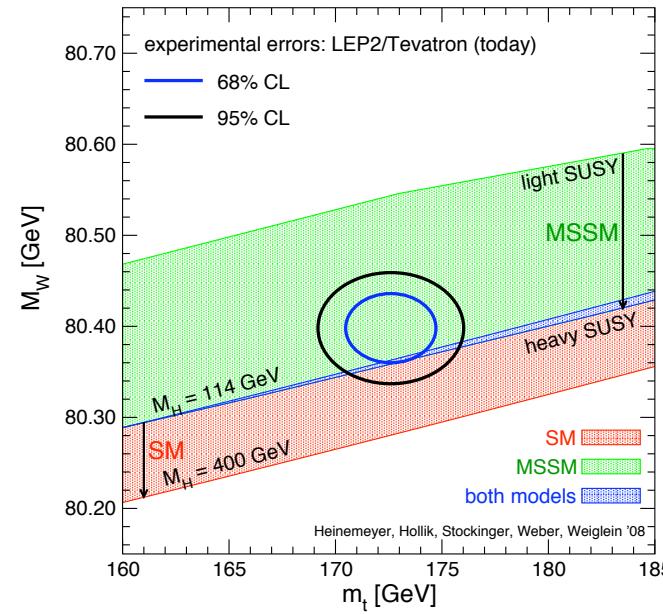
Lykken's talk at the "Muon Collider Physics, Detectors and Backgrounds Workshop"



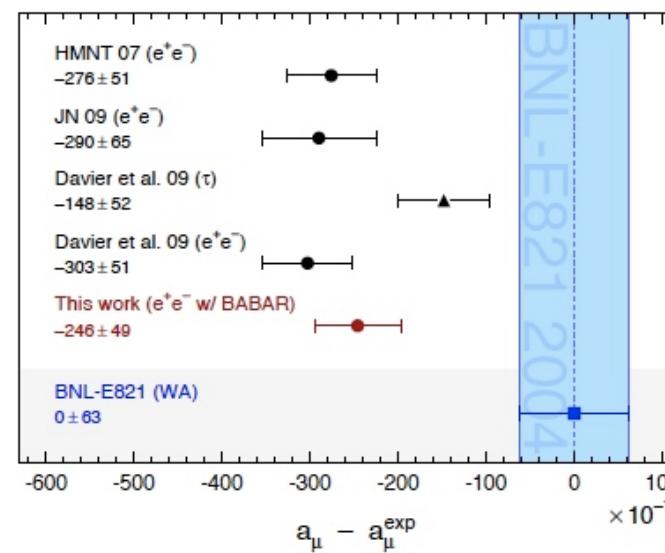
Physics Landscape

- Experimental hints for new physics

Higgs



muon ($g-2$)





Crossroad In Theoretical Physics

- Energy Frontier Experiments

- Tevatron - Operating well

$\sqrt{s} = 1.96 \text{ TeV}$ pbar p

Luminosity - $3.52 \times 10^{32} \text{ cm}^{-2} \text{ sec}^{-1}$ (peak)

9.1 fb^{-1} (to date Run II)

CDF, DO

- LHC - Online

$\sqrt{s} = 7.0 \text{ TeV}$ p p

Luminosity - 1 fb^{-1} in 2011

$\sqrt{s} = 14 \text{ TeV}$ p p

Luminosity - $10^{34} \text{ cm}^{-2} \text{ sec}^{-1}$

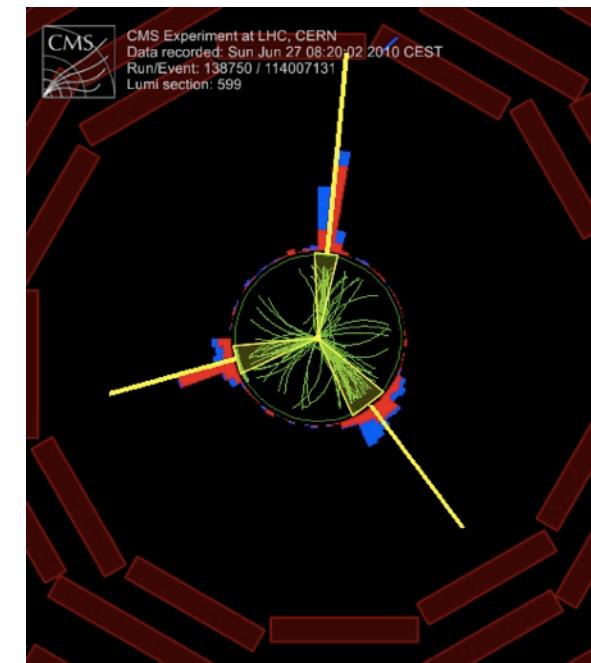
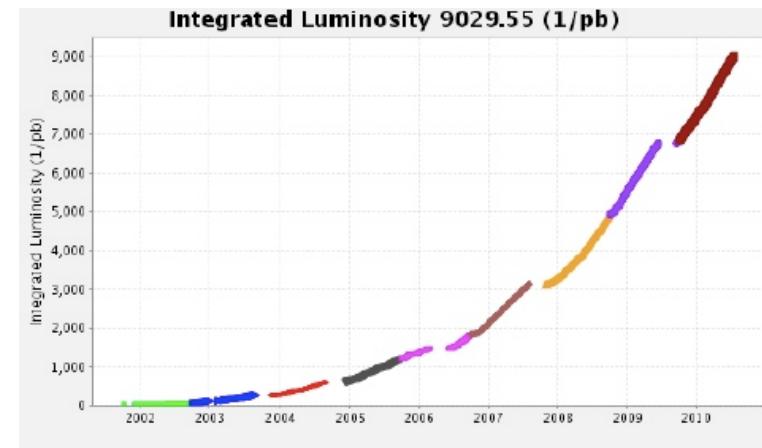
ATLAS, CMS, LHCb, ALICE

- Neutrino Experiments

Accelerators: MiniBooNE, SciBooNE, MINOS, OPERA, NOvA, LBNE, Project X, ...

Reactors: Double CHOOZ, Daya Bay, ...

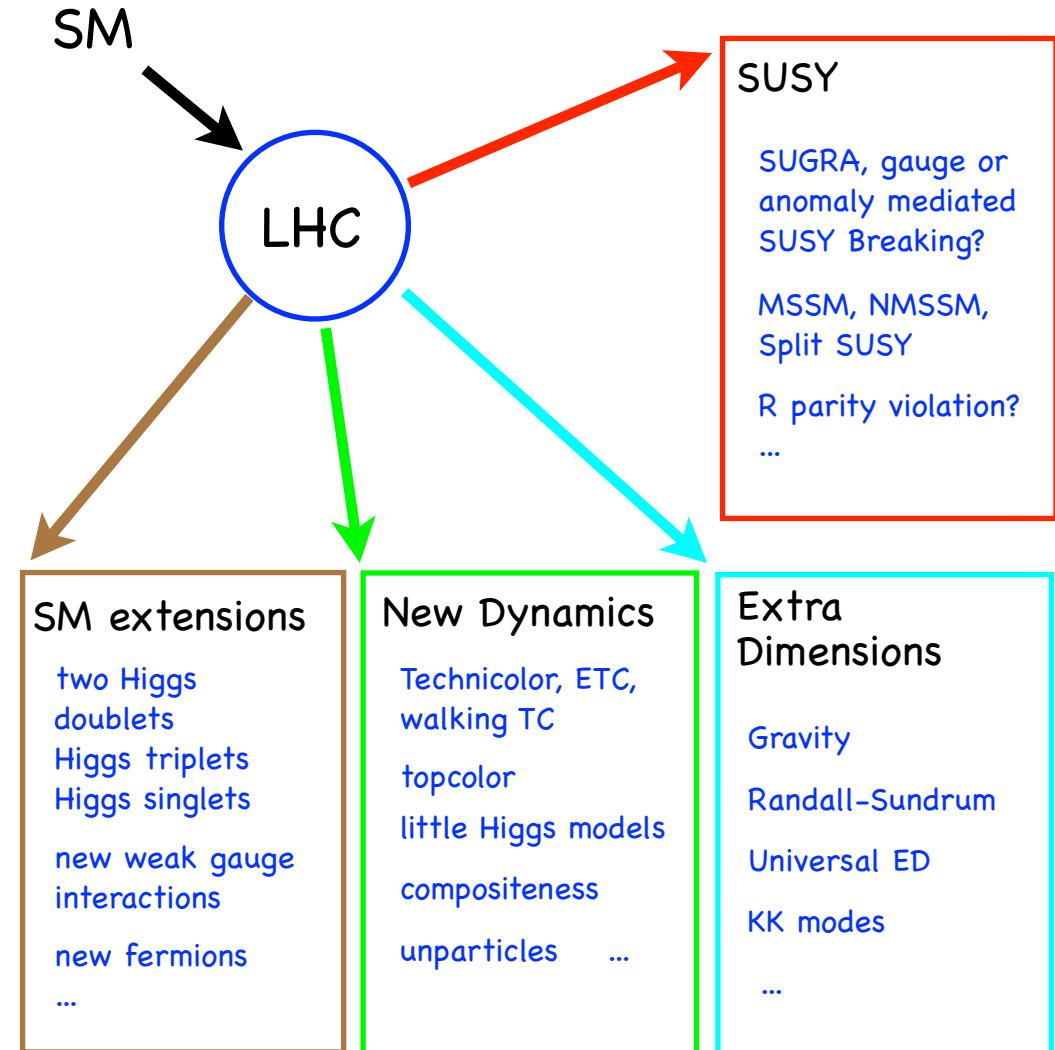
Double Beta Decay, Super Beams, Beta Beams, Astrophysical Sources





Crossroad In Theoretical Physics

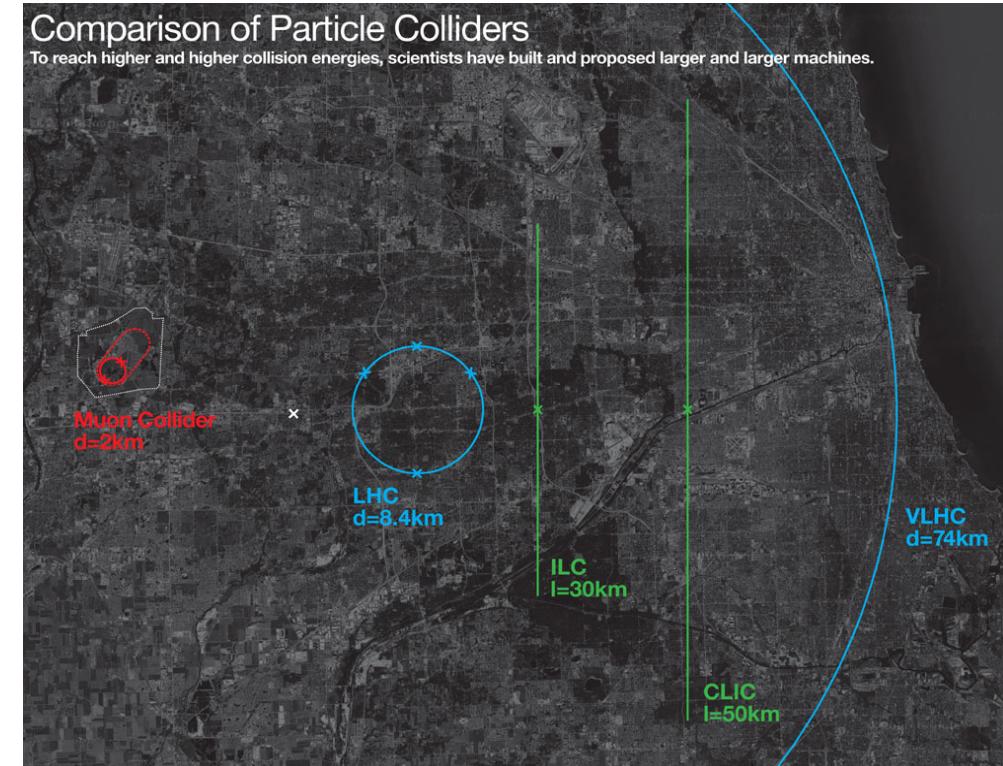
- Each path (SM-like, SUSY, New Dynamics and Extra Dimensions) represents a different mechanism for Electroweak Symmetry Breaking
- There is great excitement in the field, that after more than 25 years of theoretical speculation we will have an answer.
- We expect the results from the LHC will determine the mechanism of EWSB.
- However it will likely be necessary to build a further collider to fully explore physics at the Terascale.





Options for the Next Collider

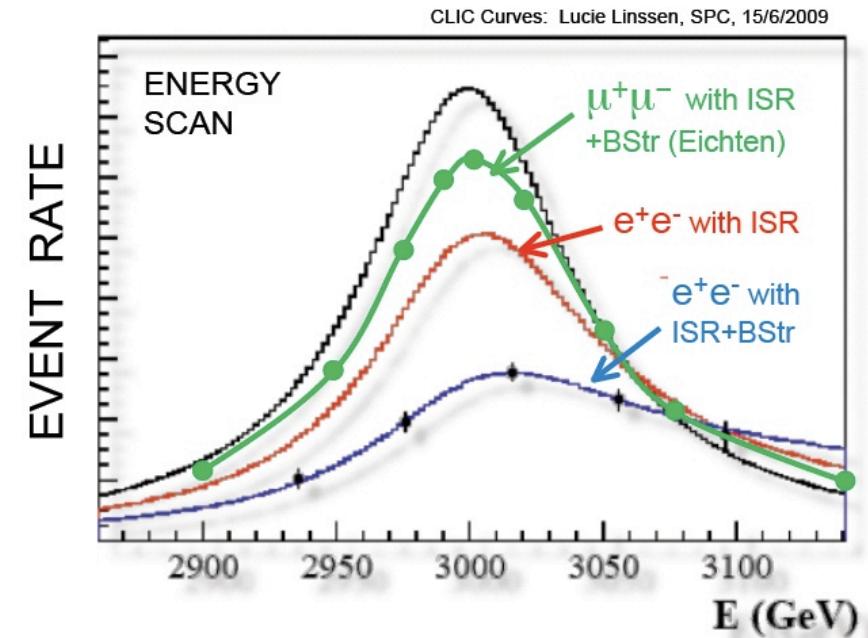
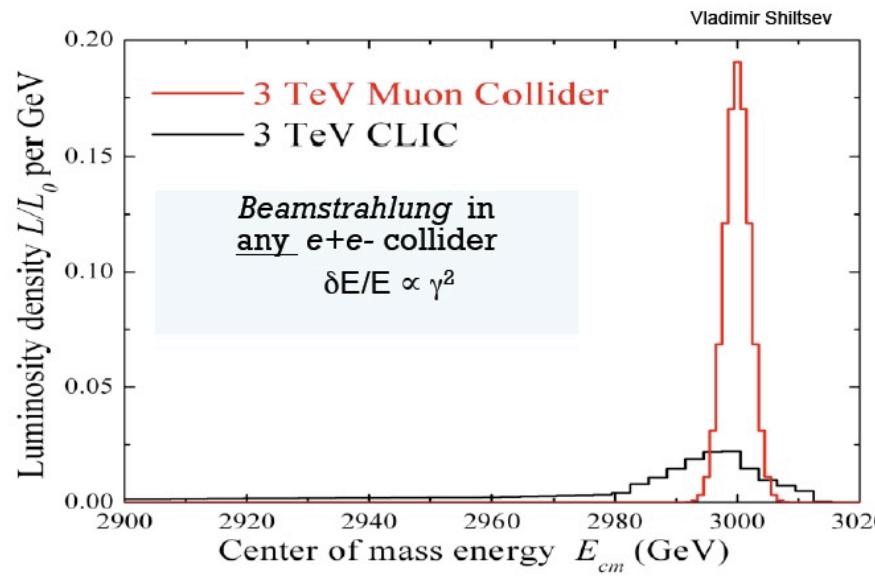
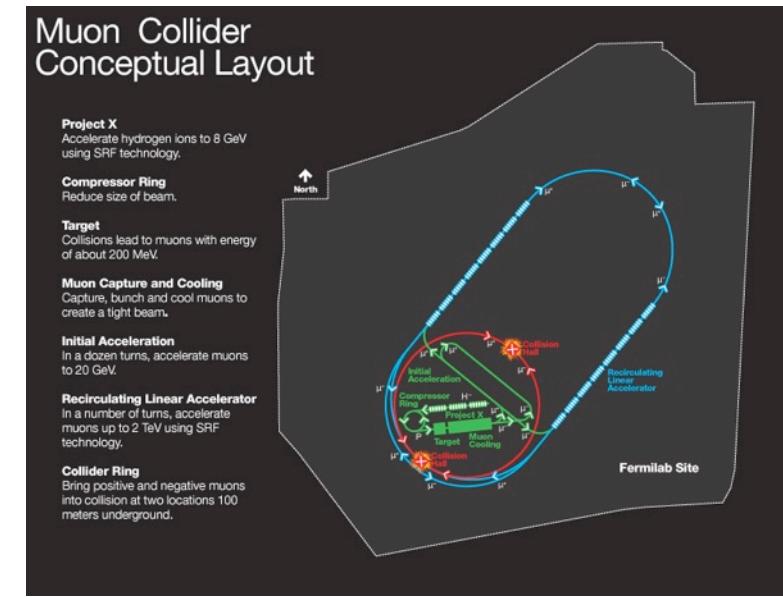
- Existing facilities in 2025:
 - LHC with luminosity or energy upgrade
- Options for next facility:
 - low energy lepton collider:
ILC (500 GeV) (upgradable) or
muon collider - Higgs Factory
 - lepton collider in the multi-TeV range:
CLIC or muon collider
 - hadron collider in hundred TeV range:
VLHC
- I will argue that a high energy lepton collider will likely be required for full study of Tevасcale physics.
- Will explore the Muon Collider option.





A Muon Collider

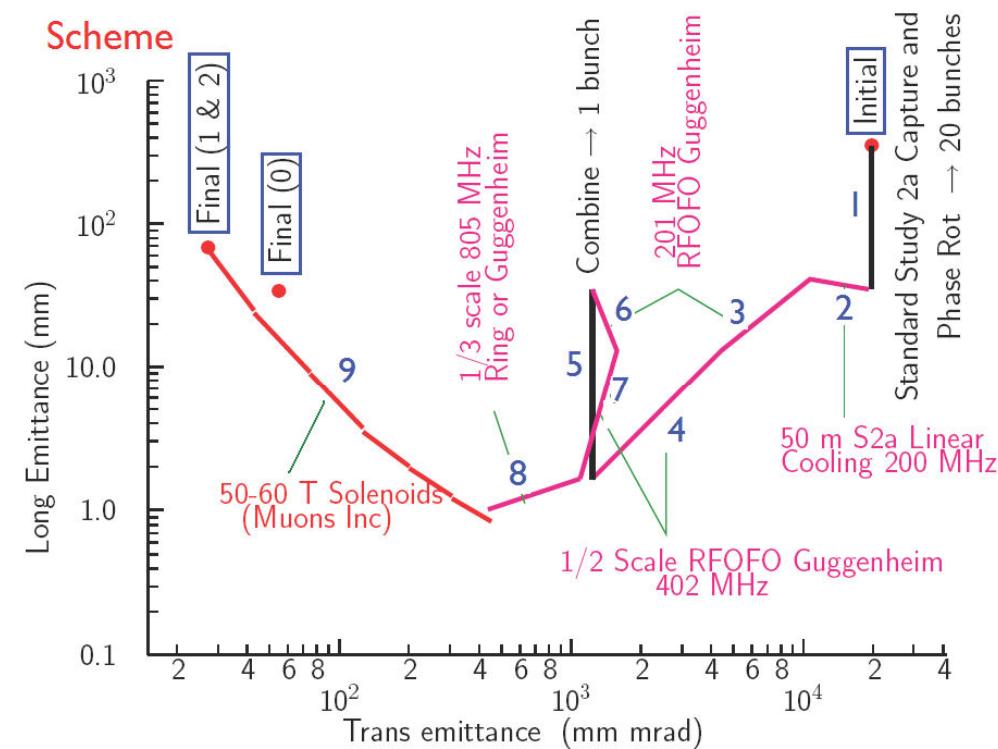
- $\mu^+\mu^-$ Collider:
 - Center of Mass energy: 1.5 - 5 TeV (focus 3 TeV)
 - Luminosity $> 10^{34} \text{ cm}^{-2} \text{ sec}^{-1}$ (focus 400 fb^{-1} per year)
- Compact facility
- Superb Energy Resolution
 - MC: 95% luminosity in $dE/E \sim 0.1\%$
 - CLIC: 35% luminosity in $dE/E \sim 1\%$





A Muon Collider

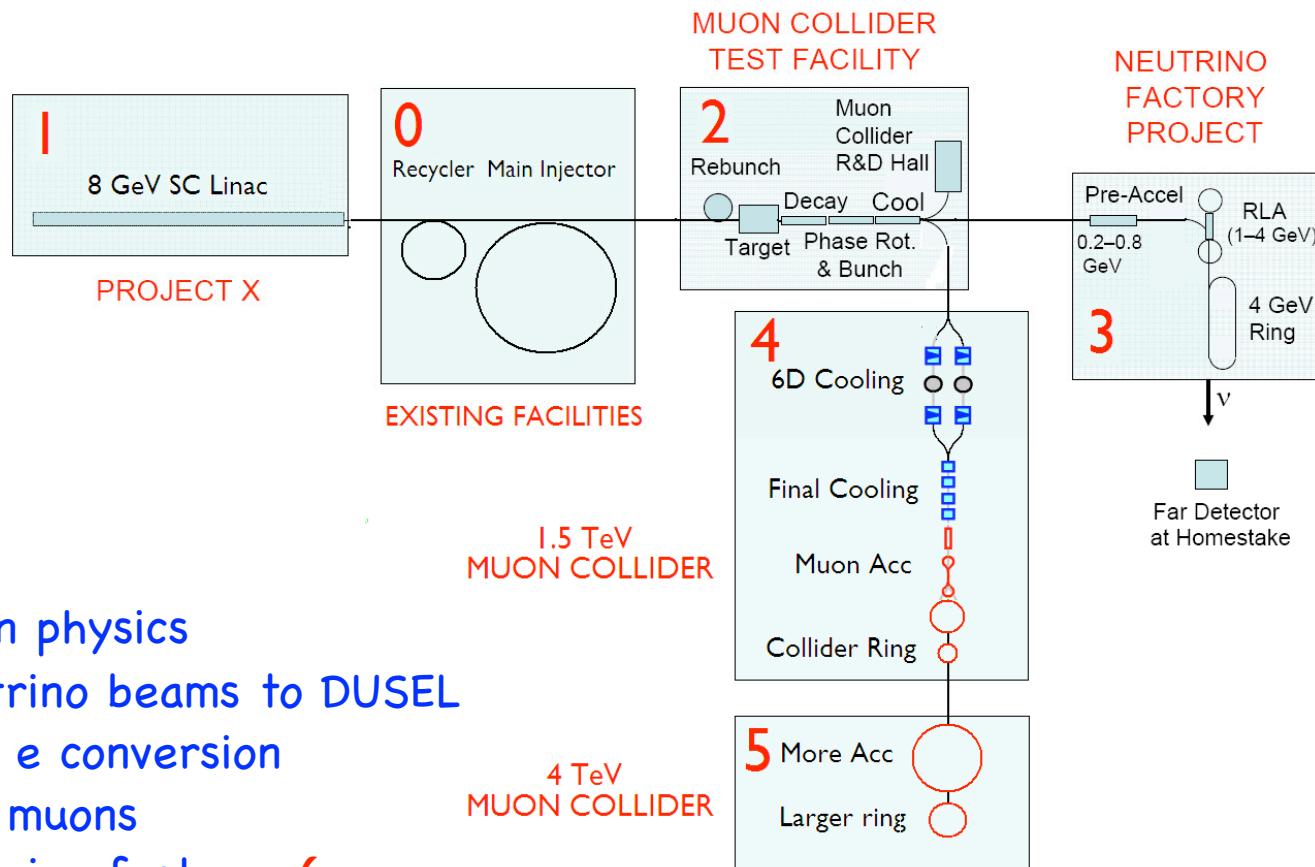
- Muons decay:
 - muon lifetime: $(2.197034 \pm 0.000021) \times 10^{-6}$ sec
 - A 3 GeV muon travels 18.7 km in one lifetime
 - A 1.5 TeV muon travels 9,300 km in this time -> More than 2000 turns in final collider ring.
 - The muon beams must be accelerated and cooled in phase space (factor $\approx 10^6$) rapidly -> ionization cooling
 - requires a complex cooling scheme
 - The decay products ($\mu^- \rightarrow \nu_\mu \bar{\nu}_e e^-$) have high energies. Serious issue for $E_{cm} \geq 4$ TeV





Path to Muon Collider Facility

- A flexible scenario with physics at each stage:

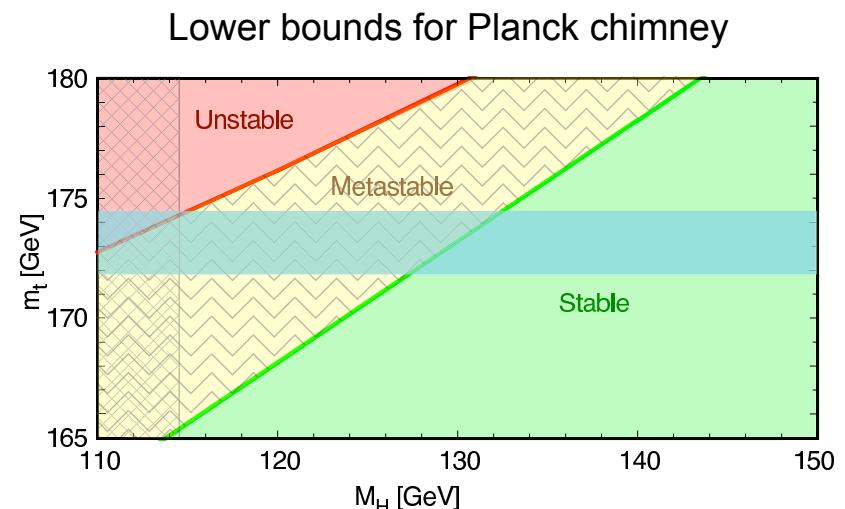
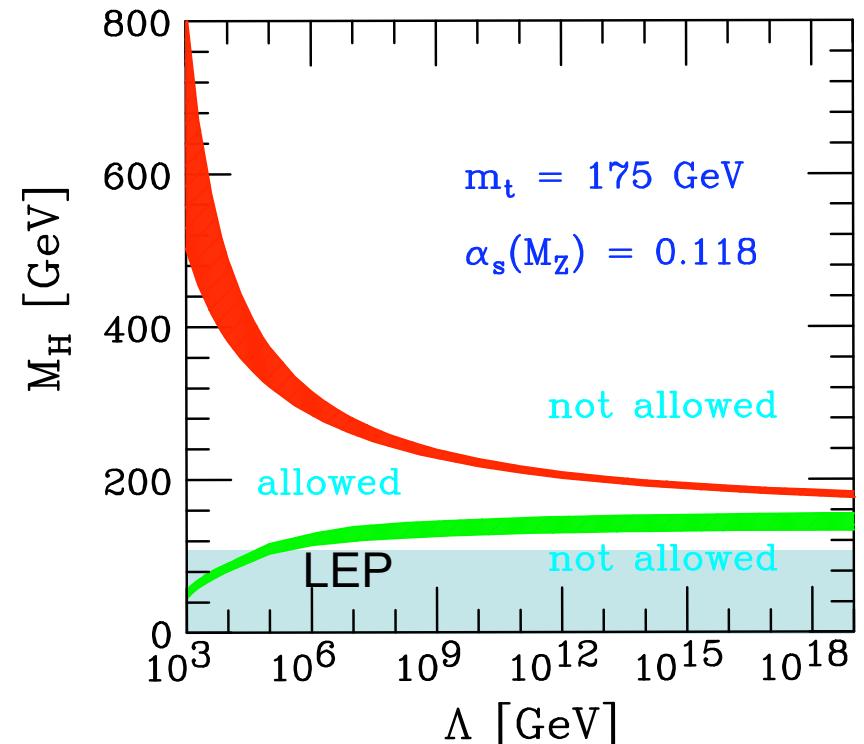


- Kaon physics
- Neutrino beams to DUSEL
- $\mu \rightarrow e$ conversion
- cold muons
- Neutrino factory ✓
- Muon collider - Higgs factory ✓
- Multi-TeV Muon Collider ✓



The Standard Model Higgs Boson

- Theoretical Constraints:
 - The standard model with an elementary Higgs scalar is only self-consistent up to some maximum energy scale (Λ).
 - Upper bound - A large Higgs mass requires a large higgs self-coupling term. This coupling increases with the scale Λ until perturbative theory breaks down.
 - Lower bound - For small Higgs mass, the quantum corrections can lead to vacuum instability.
 - Planck Chimney: SM self-consistent to Planck scale ($\approx 10^{19}$ GeV)

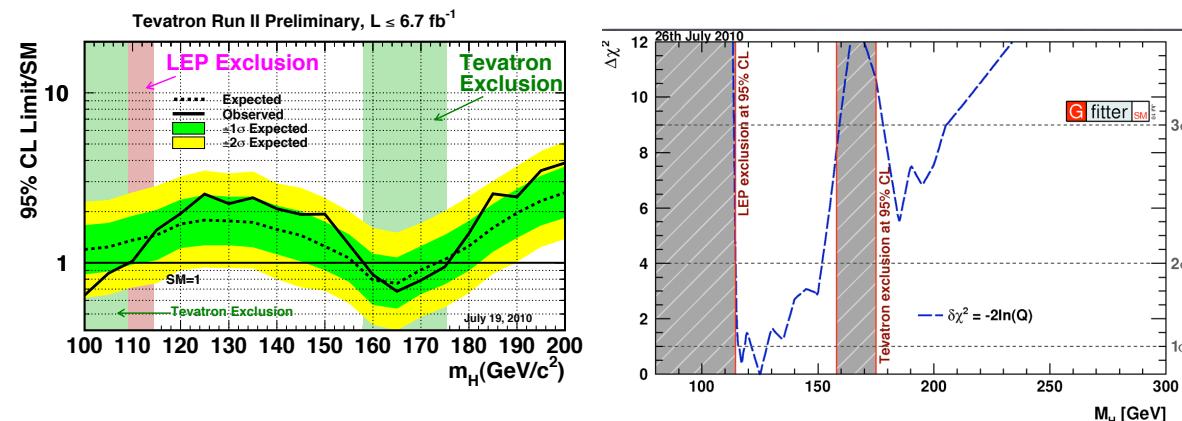




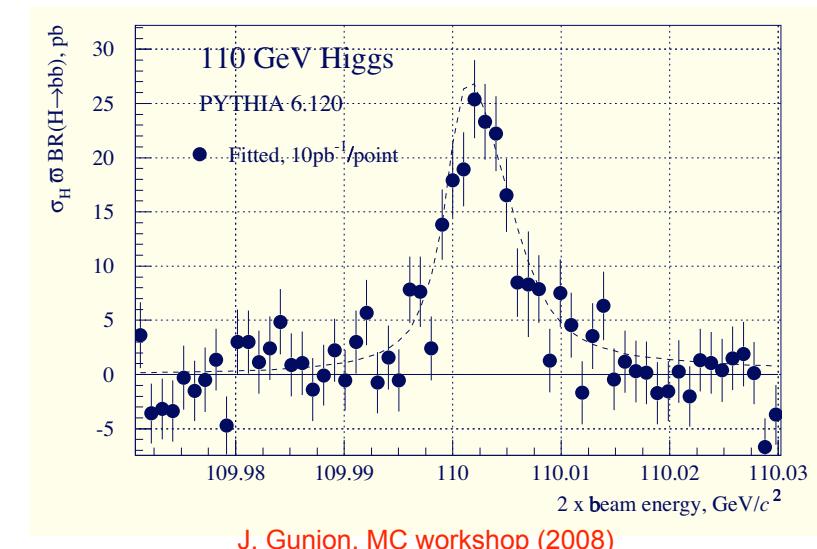
The Standard Model Higgs Boson

□ Experimental Constraints:

- Direct: LEP $m_H > 114.7 \text{ GeV}$ (95% CL)
CDF/D0 $m_H < 158$ or $> 175 \text{ GeV}$ (95% CL)
- Indirect: LEP/SLC $m_H < 190 \text{ GeV}$ (95% CL)
- Combined all information: Gfitter
 $114.6 < m_H < 151.8 \text{ GeV}$ (2σ)



- LHC will discover the SM Higgs. If Higgs mass is not in the Planck chimney (130–190), new physics “nearby”.
- Large Higgs mass implies a strong Higgs self interaction and presumably a nearby strong interaction.
- For a low mass Higgs, the new physics can be perturbative. This case is favored by the present indirect Higgs bounds. Many of the Higgs couplings could be measured at the LHC.
- The ILC(500) allows detailed study of the light Higgs properties.
- Only a low energy Muon Collider can directly measure Higgs width.



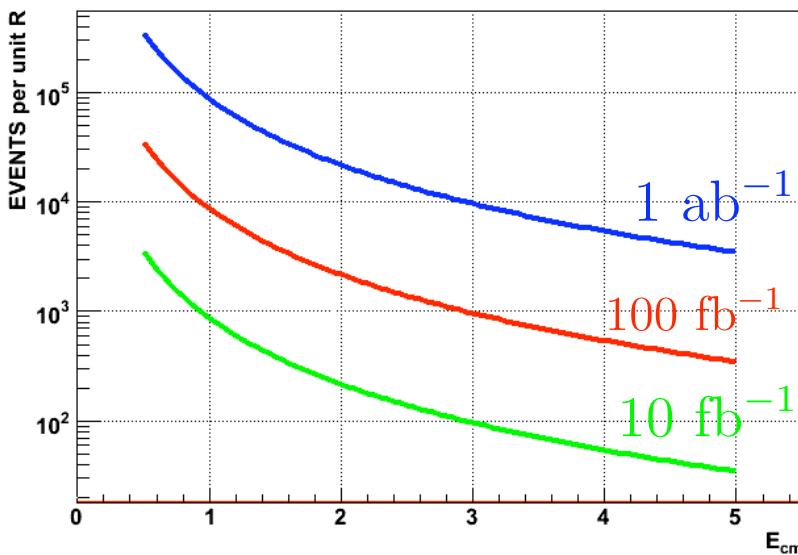


MultiTeV Muon Collider Basics

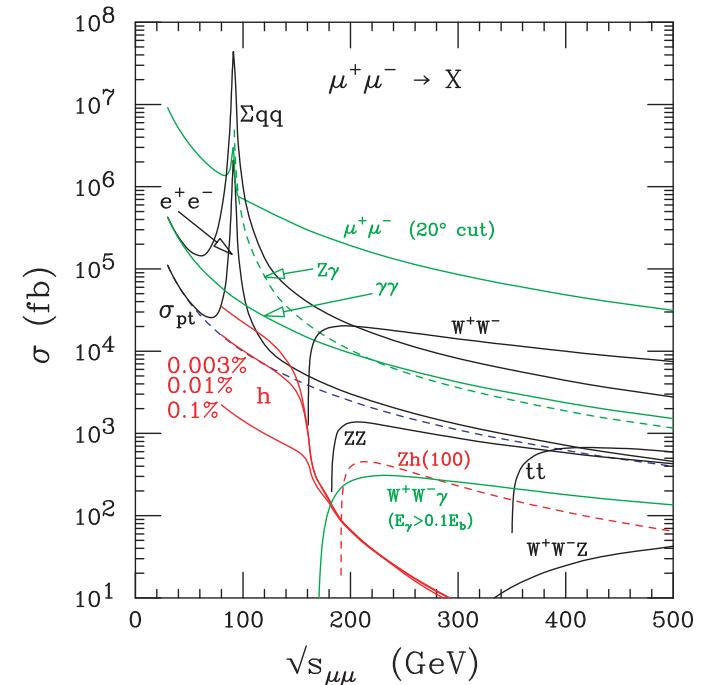
- For $\sqrt{s} < 500 \text{ GeV}$
 - SM threshold region: top pairs; W^+W^- ; Z^0Z^0 ; Z^0h ; ...
- For $\sqrt{s} > 500 \text{ GeV}$
 - For SM pair production ($|\theta| > 10^\circ$)

$$R = \sigma / \sigma_{\text{QED}}(\mu^+\mu^- \rightarrow e^+e^-) \sim \text{flat}$$

$$\sigma_{\text{QED}}(\mu^+\mu^- \rightarrow e^+e^-) = \frac{4\pi\alpha^2}{3s} = \frac{86.8 \text{ fb}}{s(\text{TeV}^2)}$$
 - High luminosity required



Standard Model Cross Sections



$$\begin{aligned} \sqrt{s} = 3.0 \text{ TeV} \quad \mathcal{L} &= 10^{34} \text{ cm}^{-2}\text{sec}^{-1} \\ &\rightarrow 100 \text{ fb}^{-1}\text{year}^{-1} \end{aligned}$$

$\Rightarrow 965 \text{ events/unit of } R$

Processes with $R \geq 0.1$ can be studied

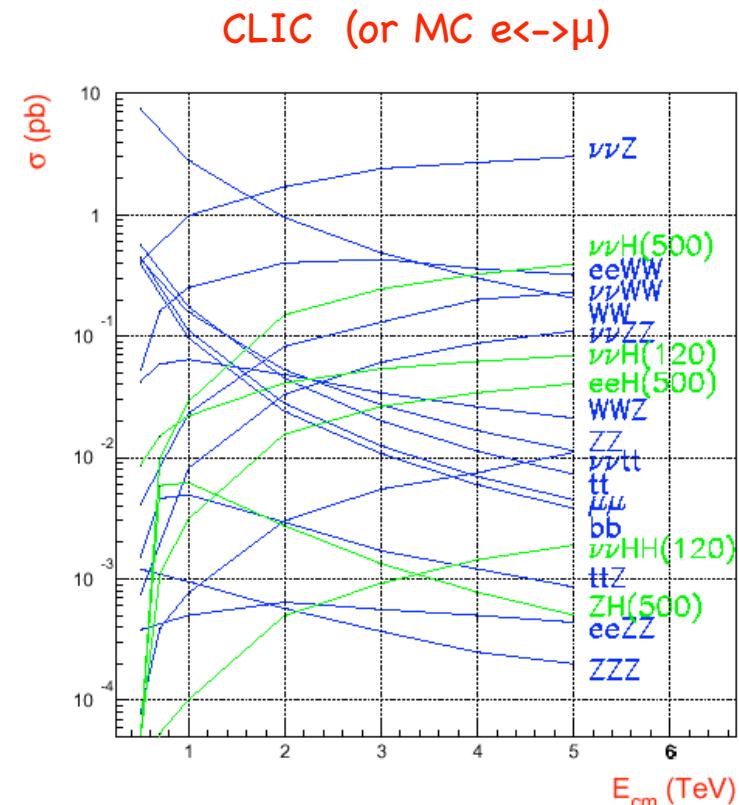
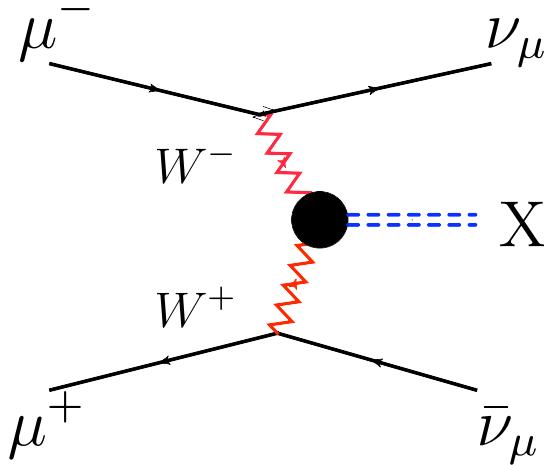
Total - 540 K SM events per year



Fusion Process

- For $\sqrt{s} > 1$ TeV - Fusion Processes

- Large cross sections
- Increase with s .
- Important at multi-Tev energies
- $M_X^2 < s$
- Backgrounds for SUSY processes
- t-channel processes sensitive to angular cuts



$$\sigma(s) = C \ln\left(\frac{s}{M_X^2}\right) + \dots$$

- An Electroweak Boson Collider



Minimum Luminosity for Muon Collider

Universal behavior for s-channel resonance

$$\sigma(E) = \frac{2J+1}{(2S_1+1)(2S_2+1)} \frac{4\pi}{k^2} \left[\frac{\Gamma^2/4}{(E - E_0)^2 + \Gamma^2/4} \right] B_{in} B_{out}$$

Convolute with beam resolution ΔE .

If $\Delta E \ll \Gamma$

$$R_{\text{peak}} = (2J+1)3 \frac{B(\mu^+\mu^-)B(\text{visible})}{\alpha_{\text{EM}}^2}$$

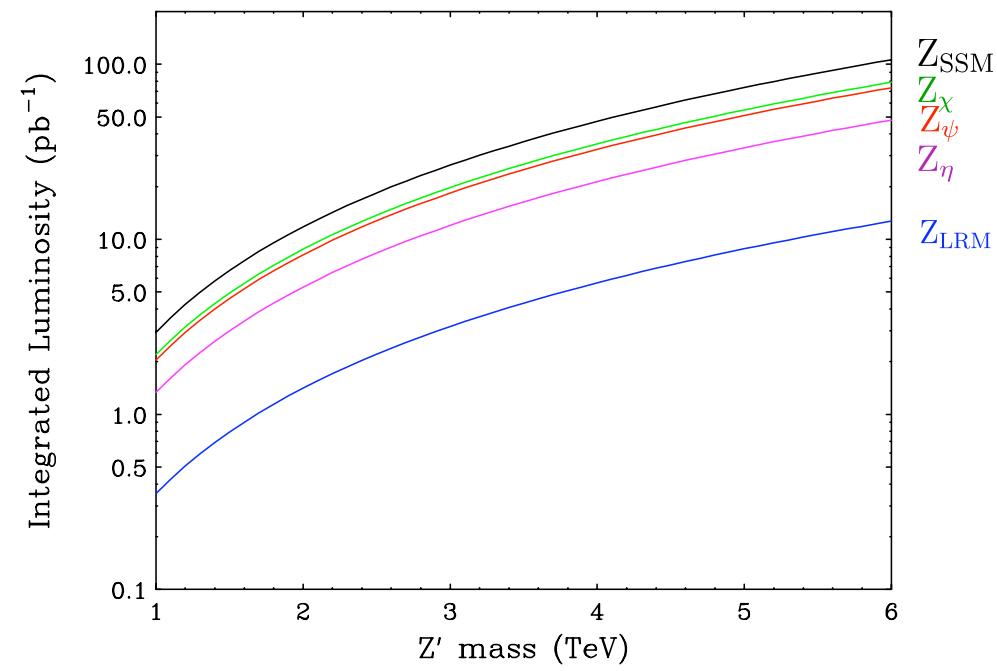
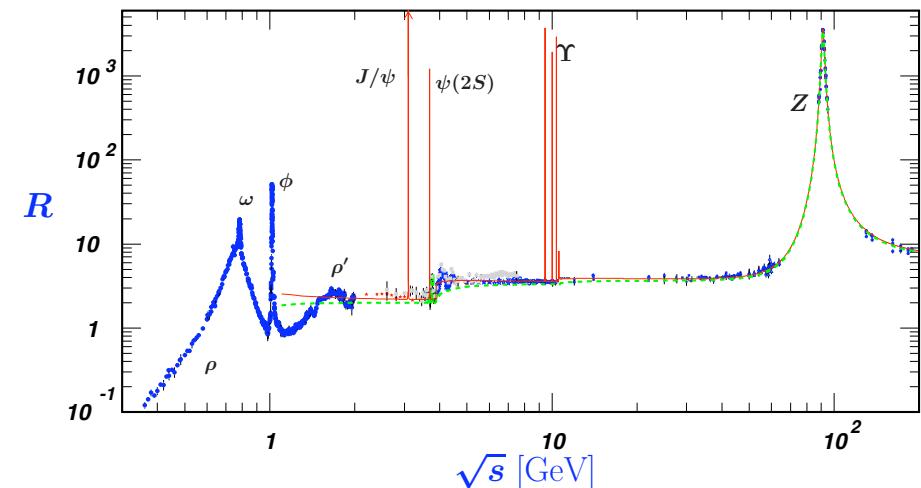
Can use to set minimum required luminosity

- Likely new physics candidates:
 - scalars: h , H^0 , A^0 , ...
 - gauge bosons: Z'
 - new dynamics: bound states
 - ED: KK modes
- Example - new gauge boson: Z'
 - SSM, E6, LRM
 - 5σ discovery limits: 4-5 TeV at LHC (@ 300 fb^{-1})

Minimum luminosity at Z' peak:

$$\mathcal{L} = 0.5-5.0 \times 10^{30} \text{ cm}^{-2} \text{ sec}^{-1}$$

for $M(Z') \rightarrow 1.5-5.0 \text{ TeV}$



The integrated luminosity required to produce 1000 $\mu^+\mu^- \rightarrow Z'$ events on the peak



Studying the Higgs Boson

- Theoretical Issues:

- Higgs boson couplings SM?
- Scalar interaction self-coupling SM?
- Any additional scalars? EW doublets, triplets or singlets?
- Where's the next scale? GUT?

- Various processes available for studying the Higgs at a multi-TeV muon collider

- Associated production: $Z h^0$ MC or CLIC:

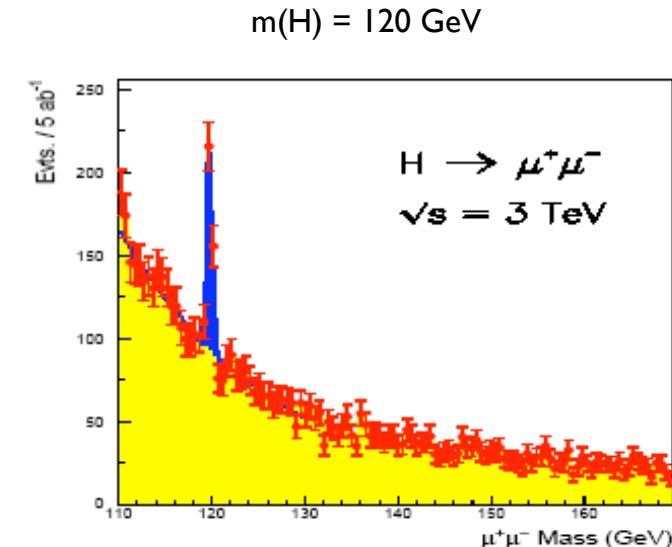
- $R \sim 0.12$
 - search for invisible h^0 decays

- Higgsstrahlung: $t\bar{t} h^0$ needs 10 ab^{-1}
- $R \sim 0.01$
 - measure top coupling

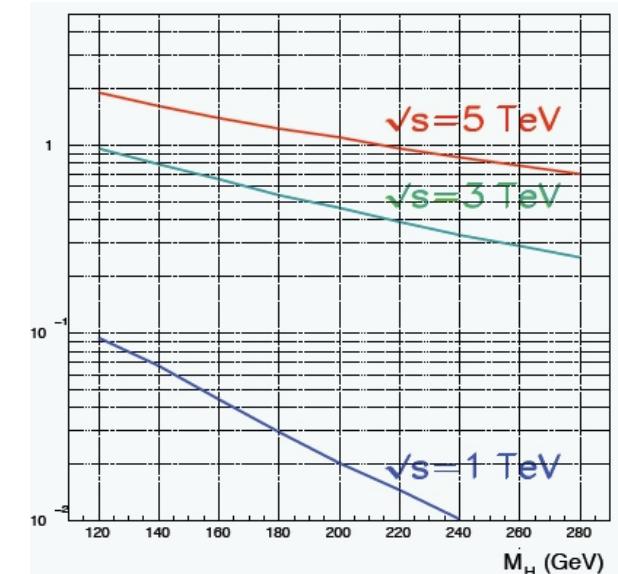
- W^*W^* fusion ($m_h = 120 \text{ GeV}$)
 - $\nu_\mu \bar{\nu}_\mu h^0$: $R \sim 1.1 s \ln(s)$ (s in TeV^2)
 - $\nu_\mu \bar{\nu}_\mu h^0 h^0$: measure Higgs self couplings

MC or CLIC:

good benchmark process



$$\sigma(\mu^+ \mu^- \rightarrow \nu \bar{\nu} h^0 h^0) (\text{fb}^{-1})$$





Two Higgs Doublets (MSSM)

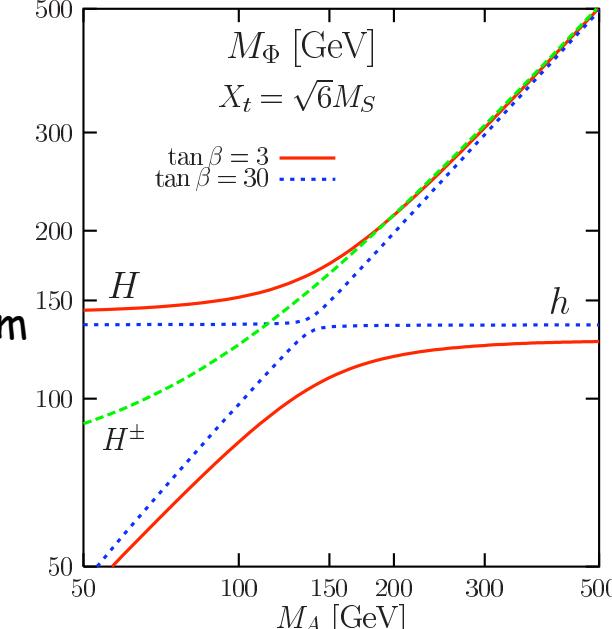
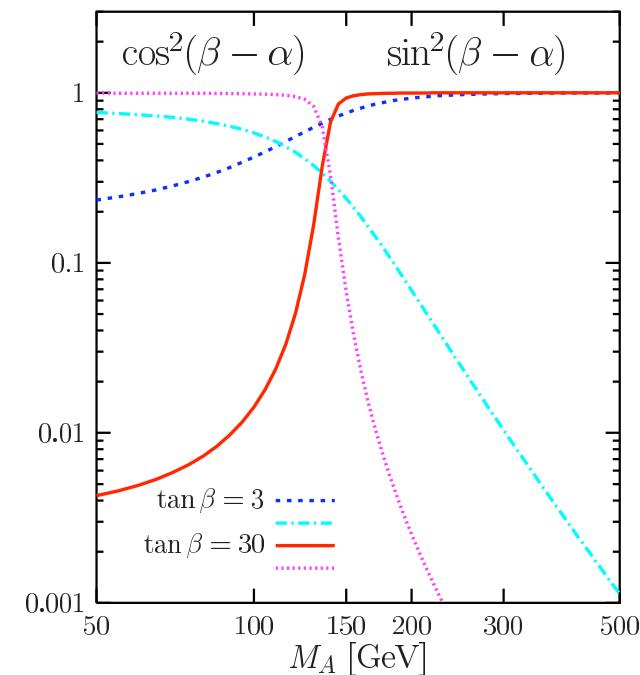
- Five scalar particles: h^0, H^0, A^0, H^\pm
- Decay amplitudes depend on two parameters: (α, β)

	$\mu^+ \mu^-, b\bar{b}$	$t\bar{t}$	ZZ, W^+W^-	$Z A^0$
h^0	$-\sin \alpha / \cos \beta$	$\cos \alpha / \sin \beta$	$\sin(\beta - \alpha)$	$\cos(\beta - \alpha)$
H^0	$\cos \alpha / \cos \beta$	$\sin \alpha / \sin \beta$	$\cos(\beta - \alpha)$	$-\sin(\beta - \alpha)$
A^0	$-i\gamma_5 \tan \beta$	$-i\gamma_5 / \tan \beta$	0	0

$$\tan 2\alpha = \frac{M_A^2 + M_Z^2}{M_A^2 - M_Z^2} \tan 2\beta.$$

- decoupling limit $m_{A^0} \gg m_{Z^0}$:

- h^0 couplings close to SM values
- H^0, H^\pm and A^0 nearly degenerate in mass
- H^0 small couplings to VV , large couplings to $Z A^0$
- For large $\tan \beta$, H^0 and A^0 couplings to charged leptons and bottom quarks enhanced by $\tan \beta$. Couplings to top quarks suppressed by $1/\tan \beta$ factor.



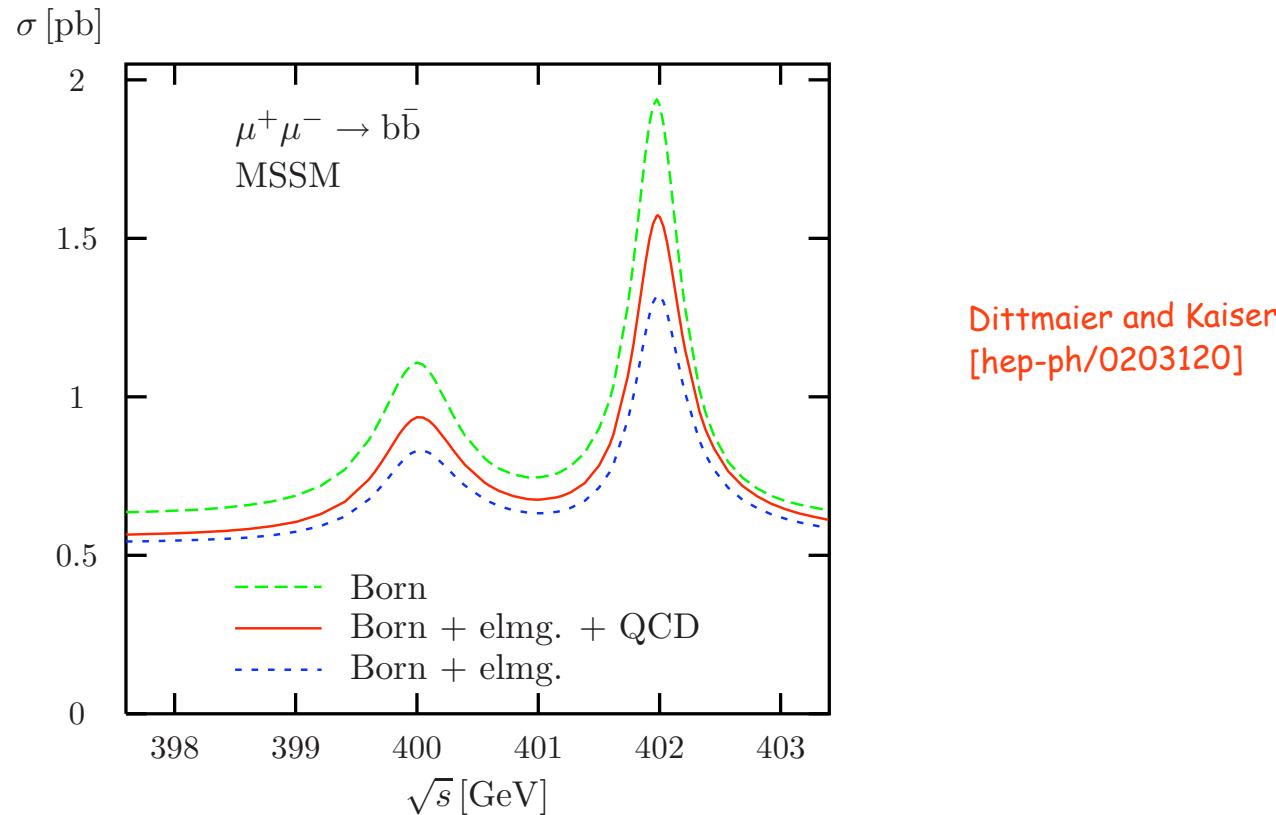


Two Higgs Doublets (MSSM)

- good energy resolution is needed for H^0 and A^0 studies:

- for s-channel production of H^0 : $\Gamma/M \approx 1\%$ at $\tan\beta = 20$.
- nearby in mass need good energy resolution to separate H and A .
- can use bremsstrahlung tail to see states using bb decay mode.

good benchmark process





New Fermions and Gauge Bosons

- Present CDF/D0 bounds on W' , Z' , and new quarks effectively rule out production at ILC(500).

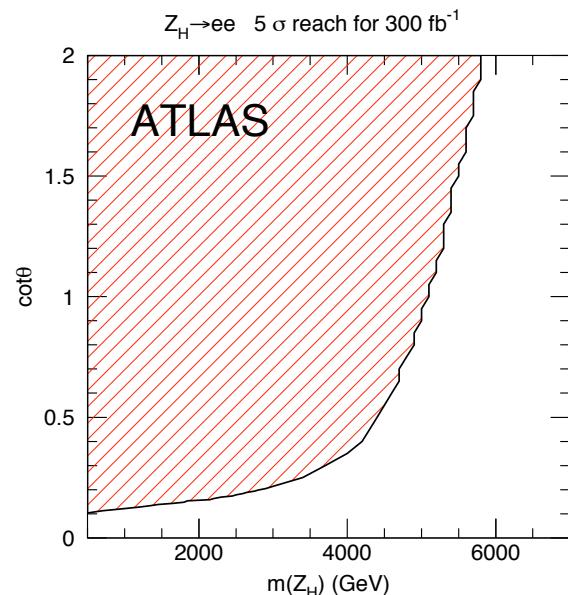
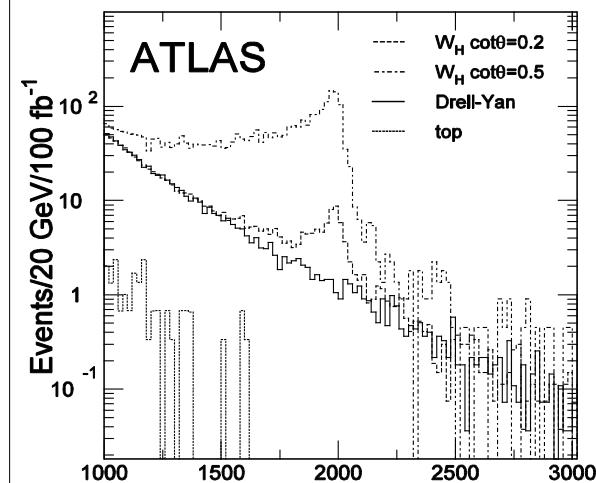
State	CDF/D0 Limit (GeV)
Quark: $(W, Z, h) + \text{jet}$	325
Z' (SM)	923
W' (SM)	860

- Littlest Higgs Model:
charge (2/3) quark T (EW singlet),
new W , Z , and A gauge bosons, Higgs triplet

At the LHC, T observable for $m(T) < 2.5$ TeV
For W , Z , and A dependent on mixing parameters

- Muon collider will allow detailed study.
Requires high luminosity 1 ab^{-1} for T

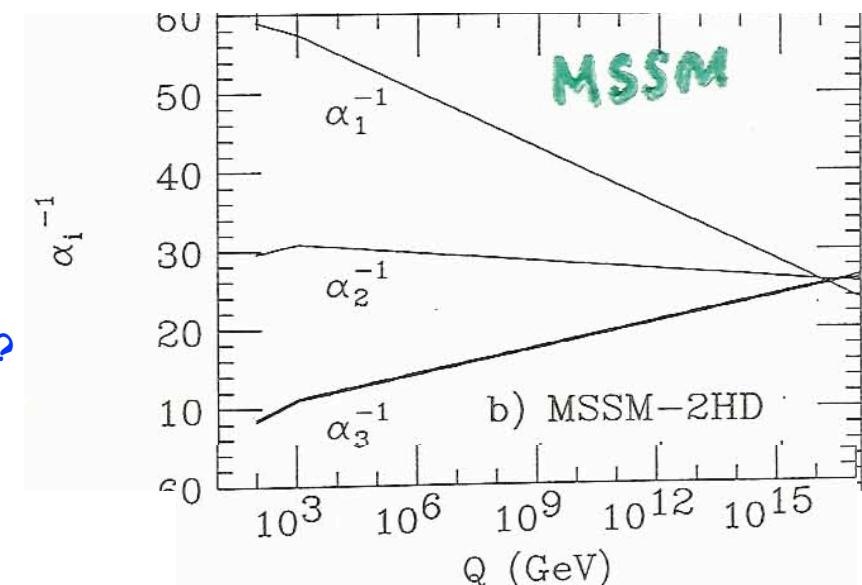
ATLAS study LHC [hep-ph/0402037]





Supersymmetry

- Solves the Naturalness Problem: Scalars associated with fermions. Higgs mass associated with SUSY breaking scale.
- Couplings of sparticles determined by symmetry. Masses depend on SUSY breaking mechanism.
- If discovered at LHC →
 - What is the spectrum of superpartner masses?
 - Dark matter candidates?
 - Are all the couplings correct?
 - What is the structure of flavor mixing interactions?
 - Are there additional CP violating interactions?
 - Is R parity violated?
 - What is the mass scale at which SUSY is restored?
 - What is the mechanism of SUSY breaking?



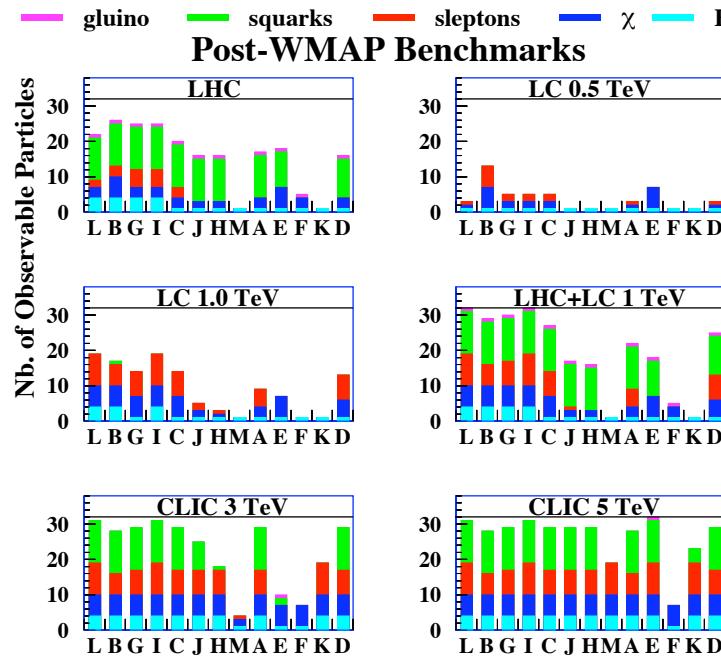
- cMSSM [Constrained Minimal Supersymmetric Standard Model]
 - Five parameters: m_0 , $m_{1/2}$, $\tan\beta$, A/m_0 , $\text{sign}(\mu)$
 - Experimental constraints
 - Direct limit (LEP, CDF, Dzero): $m_{h^0}, m_{\chi^\pm}, m_{\tilde{t}}, \dots$
 - Electroweak precision observables (EWPO): $M_W^2, \sin^2\theta_{sw}, (g-2)_\mu, \dots$
 - B physics observables (BPO): $b \rightarrow s + \gamma$, $\text{BR}(B_s \rightarrow \mu^+ \mu^-), \dots$
 - Cold dark matter (CDM): $\Omega_{DM} = .23 \pm .04$
 - Allowed regions are narrow filaments in parameter space



Supersymmetry

- The combination of the LHC and a multiTeV lepton collider is required to fully study the SUSY spectrum.

- cMSSM 2004 CLIC study SUSY reach

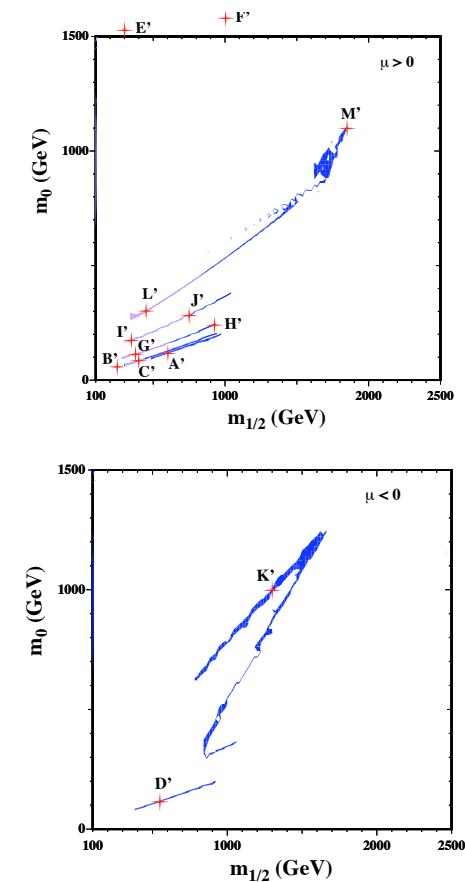


Similar Conclusion for MC

Anupama Atre, Low Emittance Muon Collider Workshop, Fermilab, April 2008

- Alternate supersymmetry breaking schemes (mGMSB, mAMSB) also require multiTeV lepton collider. S. Heinemeyer, X. Miao, S. Su, G. Wieglein [arXiv:0805.2359]
- Supersymmetry provides a strong case for a multiTeV muon collider

Allowed regions and sample points

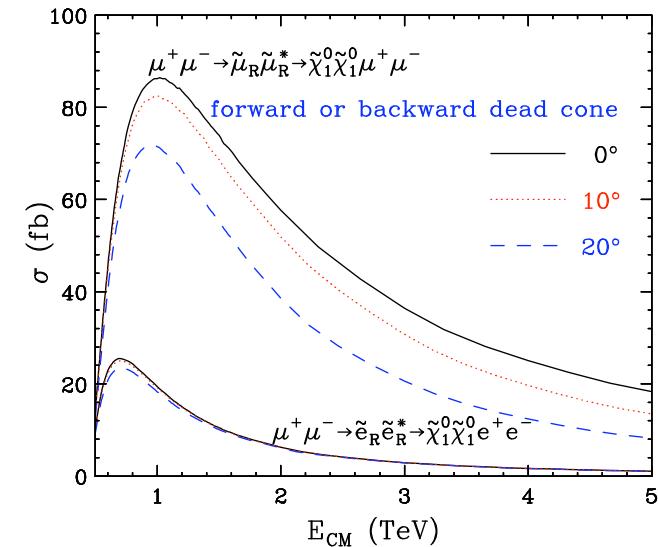




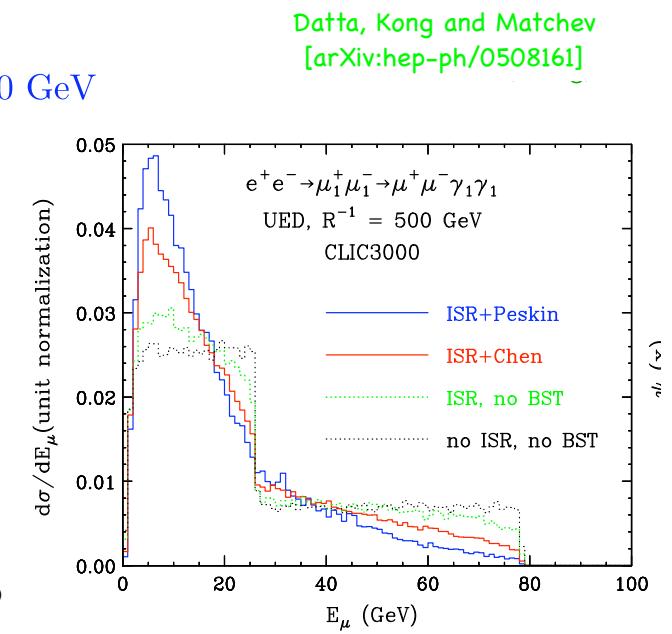
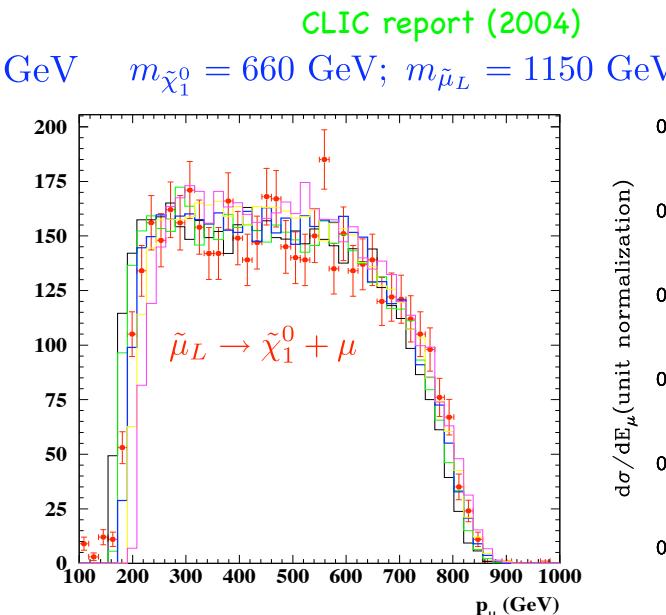
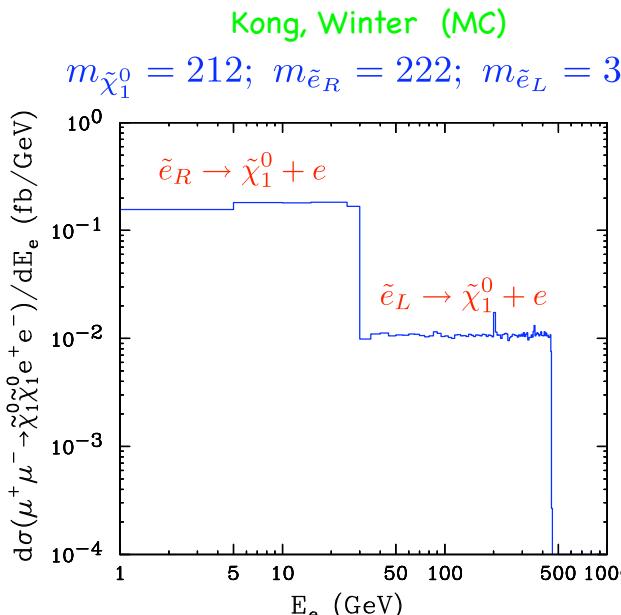
Example Process at Muon Collider

- $\mu^+ \mu^- \rightarrow \tilde{e}_1^+ \tilde{e}_1^- \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0 e^+ e^-$
 - Angular cut at 20° from beam direction:
 - 50% reduction for smuon pairs
 - 20% reduction for selectron pairs
 - Mass measurements using edge method better for MC than CLIC:

$$E_{\max/\min} = \frac{1}{2} M_{\tilde{e}} \left[1 - \frac{M_{\tilde{\chi}_1^0}^2}{M_{\tilde{e}}^2} \right] \gamma(1 \pm \beta)$$



Effect of beamstrahlung





New Strong Dynamics

- Solves the Naturalness Problem: Electroweak Symmetry Breaking is generated dynamically at a nearby scale. May or may not be a light Higgs boson.

Theoretical issues

- What is the spectrum of low-lying states?
- What is the ultraviolet completion? Gauge group? Fermion representations?
- What is the energy scale of the new dynamics?
- Any new insight into quark and/or lepton flavor mixing and CP violation?
- ...

Technicolor, ETC, Walking TC, Topcolor , ...

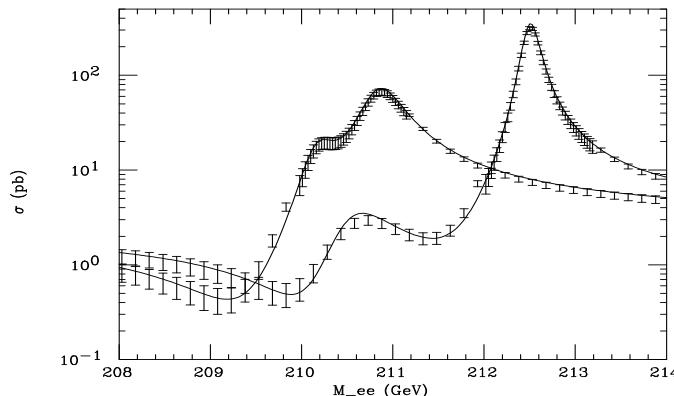
For example with a new strong interaction at TeV scale expect:

- Technipions - s channel production (Higgs like)
- Technirhos - Nearby resonances (p_T, ω_T) - need fine energy resolution of muon collider.

Eichten, Lane, Womersley PRL 80, 5489 (1998)

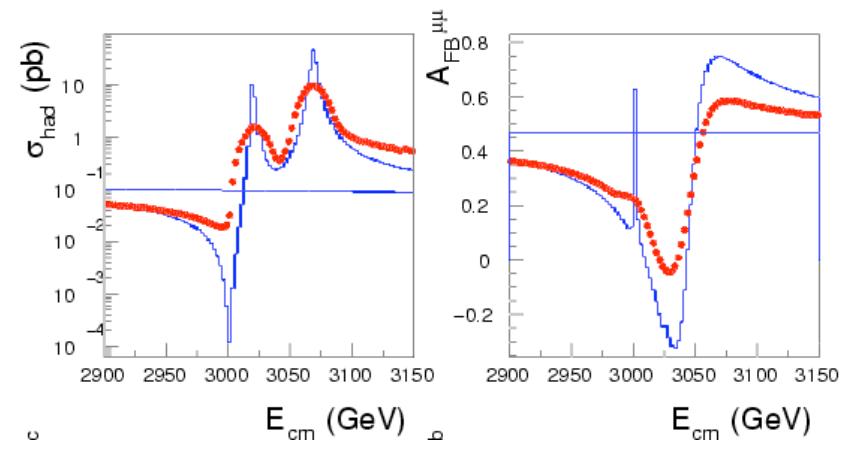
$M(p_T) = 210 \text{ GeV}$ $M(\omega_T) = 211, 209 \text{ GeV}$

MC 40 steps (total 1 fb^{-1})



good benchmark
processes

CLIC - D-BESS model (resolution 13 GeV)





Contact Interactions

- Solves the Naturalness Problem: The SM is only an effective theory valid below the compositeness scale.

- New interactions (at scales not directly accessible) give rise to contact interactions.

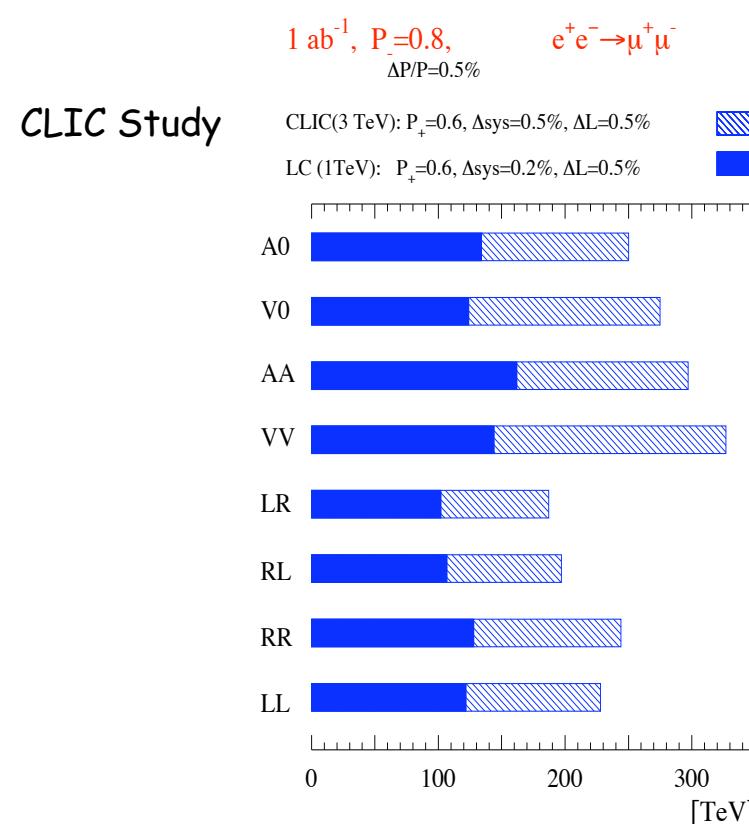
$$\mathcal{L} = \frac{g^2}{\Lambda^2} (\bar{\Psi} \Gamma \Psi)(\bar{\Psi} \Gamma' \Psi)$$

- Muon collider is sensitive to contact interaction scales over **200 TeV** as is CLIC.
- Cuts on forward angles for a muon collider not an issue.
- Polarization useful to disentangle the chiral structure of the interaction.
(CLIC)

good benchmark process

Muon Collider Study

E.Eichten, S.~Keller, [arXiv:hep-ph/9801258]





Extra Dimensions

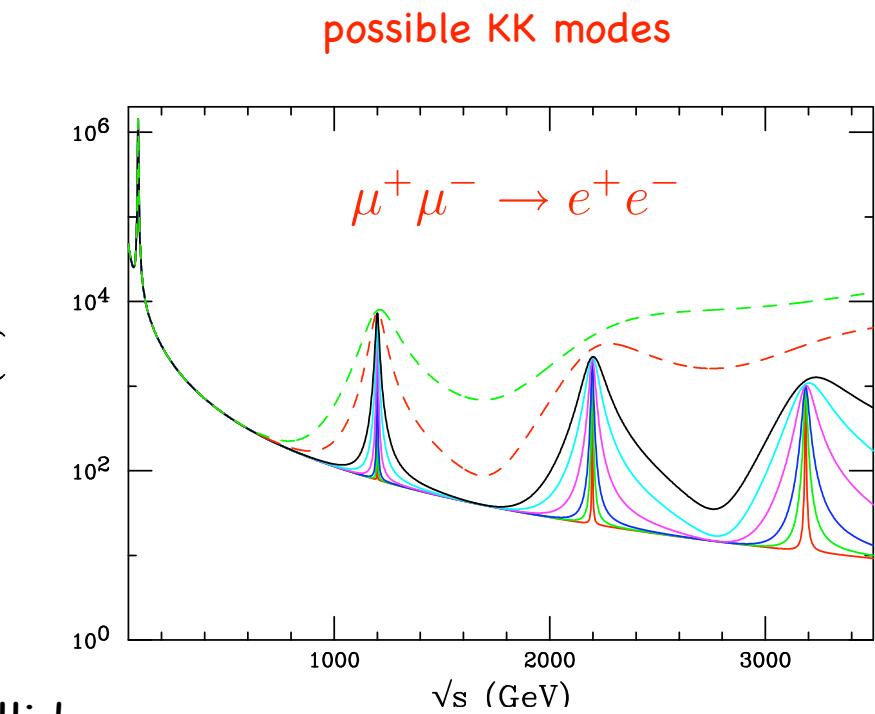
- Solves the Naturalness Problem: The effective GUT scale is moved closer.

Theoretical issues

- How many dimensions?
- Which interactions (other than gravity) extend into the extra dimensions?
- At what scale does gravity become a strong interaction?
- What happens above that scale?
- ...

Randall-Sundrum model:
warped extra dimensions

- two parameters:
 - mass scale \propto first KK mode;
 - width \propto 5D curvature / effective 4D Planck scale.



LHC discovery - Detailed study at a muon collider



Summary of Physics Potential

- A multiTeV lepton collider is required for full coverage of Terascale physics.
- The physics potential for a muon collider at $\sqrt{s} \sim 3$ TeV and integrated luminosity of 1 ab^{-1} is outstanding. Particularly strong case for SUSY and new strong dynamics.
- Narrow s-channel states played an important role in past lepton colliders. If such states exist in the multi-TeV region, they will play a similar role in precision studies for new physics. Sets the minimum luminosity scale.
- A staged Muon Collider can provide a Neutrino Factory to fully disentangle neutrino physics.
- A detailed study of physics case for 1.5-4.0 TeV muon collider has begun. Goals:
 - Identify benchmark processes: pair production (slepton; new fermion), Z' pole studies, h^0 plus missing energy, resolving nearby states (H^0-A^0 ; $\rho_T-\omega^0_T$), ...
 - Dependence on initial beam [electron/muon, polarization and beam energy spread] as well as luminosity to be considered.
 - Estimates of collision point environment and detector parameters needed.
 - Must present a compelling case even after ten years of running at the LHC.



Backup Slides



○ Fine tuning problems in the cMSSM

$M(h^0) > 113.8 \text{ GeV}$ (95% cl) LEP combined bound]

tree **1-loop**

$$M_{h^0}^2 = m_Z^2 \cos^2(2\beta) + \frac{3}{4\pi^2} \sin^2 \beta y_t^2 \left[m_t^2 \ln \left(m_{\tilde{t}_1} m_{\tilde{t}_2} / m_t^2 \right) + c_{\tilde{t}}^2 s_{\tilde{t}}^2 (m_{\tilde{t}_2}^2 - m_{\tilde{t}_1}^2) \ln(m_{\tilde{t}_2}^2 / m_{\tilde{t}_1}^2) \right. \\ \left. + c_{\tilde{t}}^4 s_{\tilde{t}}^4 \left\{ (m_{\tilde{t}_2}^2 - m_{\tilde{t}_1}^2)^2 - \frac{1}{2} (m_{\tilde{t}_2}^4 - m_{\tilde{t}_1}^4) \ln(m_{\tilde{t}_2}^2 / m_{\tilde{t}_1}^2) \right\} / m_t^2 \right].$$

$\tan \beta = v_u/v_d$

top squark

masses: $m_{\tilde{t}_1}, m_{\tilde{t}_2}$

mixing: $c_{\tilde{t}}, s_{\tilde{t}}$

+ ...

with measured top mass and $\tan \beta$ constraints,
need large top squark mass. BUT

$$m_Z^2 = -2(|\mu|^2 + m_{H_u}^2) - \frac{1}{v_u} \frac{\partial}{\partial v_u} \Delta V + \mathcal{O}(1/\tan^2 \beta).$$

soft SUSY breaking mass term
in higgs field coupling to top

loop part of effective potential

the largeness the soft SUSY breaking mass term
means a fine tuned cancellation between the μ^2 and
 m_H^2 terms to more than a few percent.

Relax the soft breaking restrictions at the GUT scale ?



cMSSM, mGMSB, mAMSB Studies

- More generally, full coverage likely requires a multi TeV lepton collider

S. Heinemeyer, X. Miao, S. Su, G. Wieglein [arXiv:0805.2359]
(using only EWPO, BPO and LEP)

Second lightest neutralino:

$m(\tilde{\chi}_2^0) < 900$ GeV for $\Delta\chi^2 < 4$

Heavy for LHC - possibly in decay chain ?

Lepton collider: $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 + X$

Lightest chargino:

$m(\tilde{\chi}_1^\pm) < 800, 900, 300$ GeV for $\Delta\chi^2 < 4$

Heavy for LHC - possibly in decay chain ?

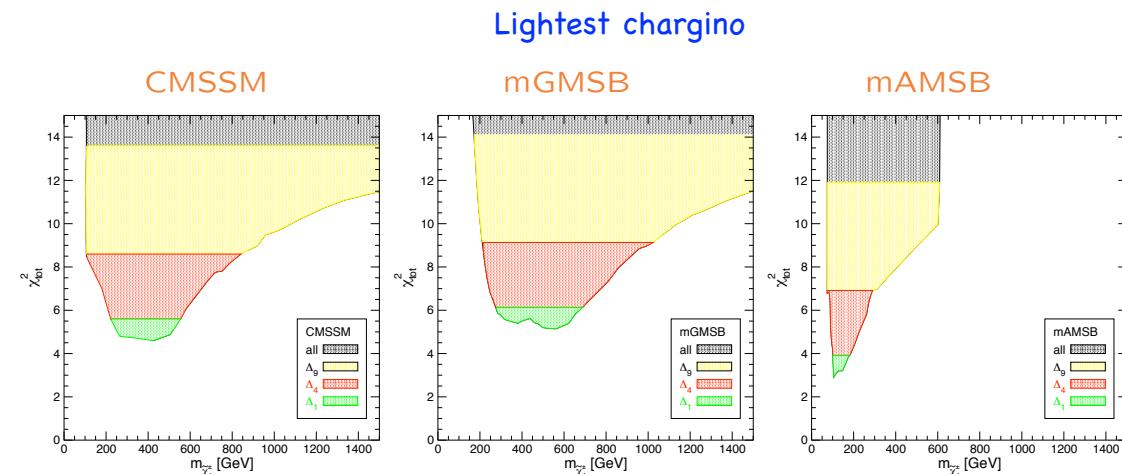
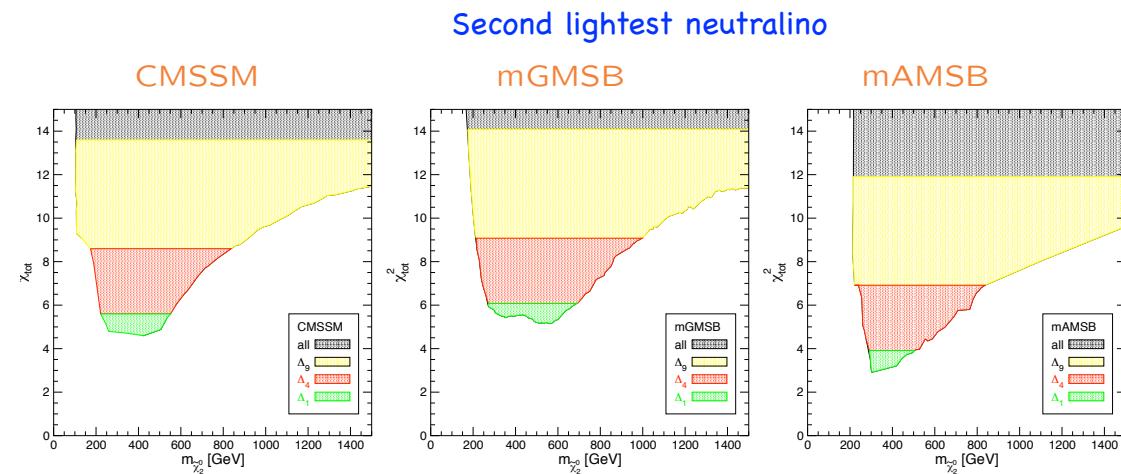
Lepton collider: Observable at ILC for mAMSB

Lightest stop, sbottom and gluino:

$m(\tilde{t}_1) > 500$ for $\Delta\chi^2 < 4$

Easy for LHC up to 2 TeV

Lepton collider: Detailed study?

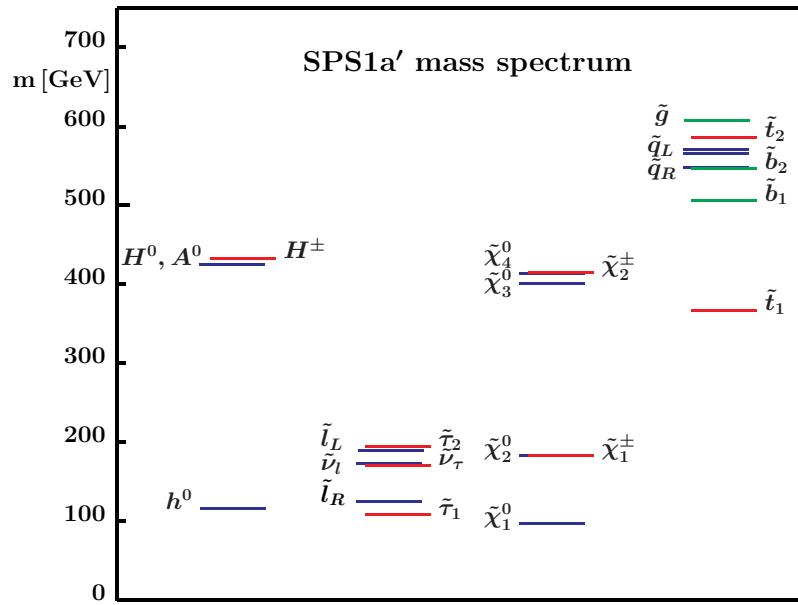




Modifying cMSSM

- Fine tuning problems in the cMSSM - Allow non universal $m_{1/2}$

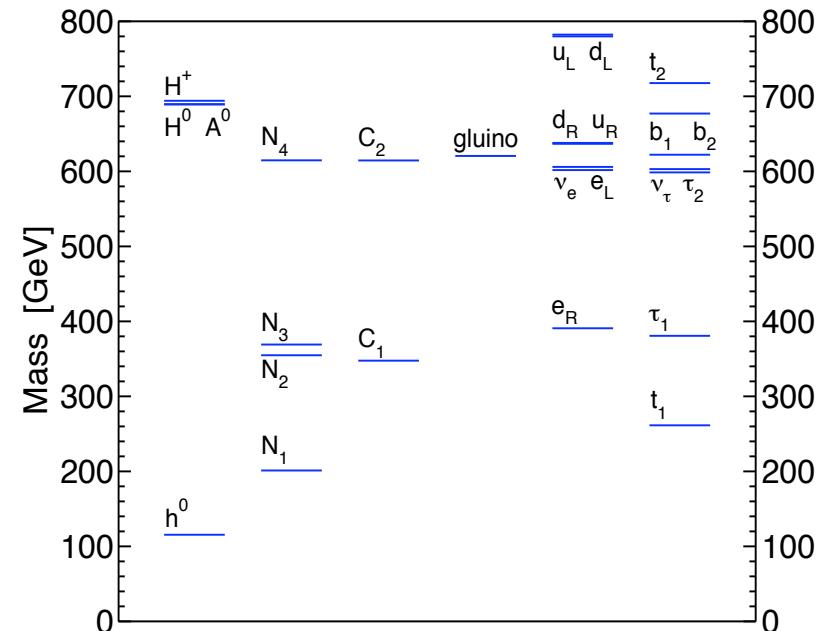
cMSSM ILC Benchmark



Many visible superpartners within reach of the ILC (500 GeV).
All pair production thresholds are below 1.2 TeV.

Compressed SUSY

S. Martin [PR D75:115005,2007]



No visible superpartners within reach of the ILC (500 GeV).
All pair production thresholds are below 1.6 TeV.

Supersymmetry provides strong case for a multi-TeV lepton collider



Challenges Ahead

- Many technical challenges exist. Two of the most difficult are:
 - 6D cooling - needed to obtain sufficient luminosity for physics.
 - The interaction region and detectors have to be designed to do physics with the background environment generated by nearby muon decays.
- Many practical issues also need consideration.
 - Cost of building a multi-TeV muon collider? - Staging likely necessary and desirable
 - How to deal with the high energy neutrinos from muon decay? - Limiting factor for energy reach of the Muon Collider?
- However a multiTeV Muon Collider would address the most fundamental issues in our field. It would allow a detailed look at the mechanism of EW symmetry breaking and likely provide clues to the origin of fermion masses.
- This extraordinary opportunity justifies the serious research efforts presently underway into the feasibility of such a collider.