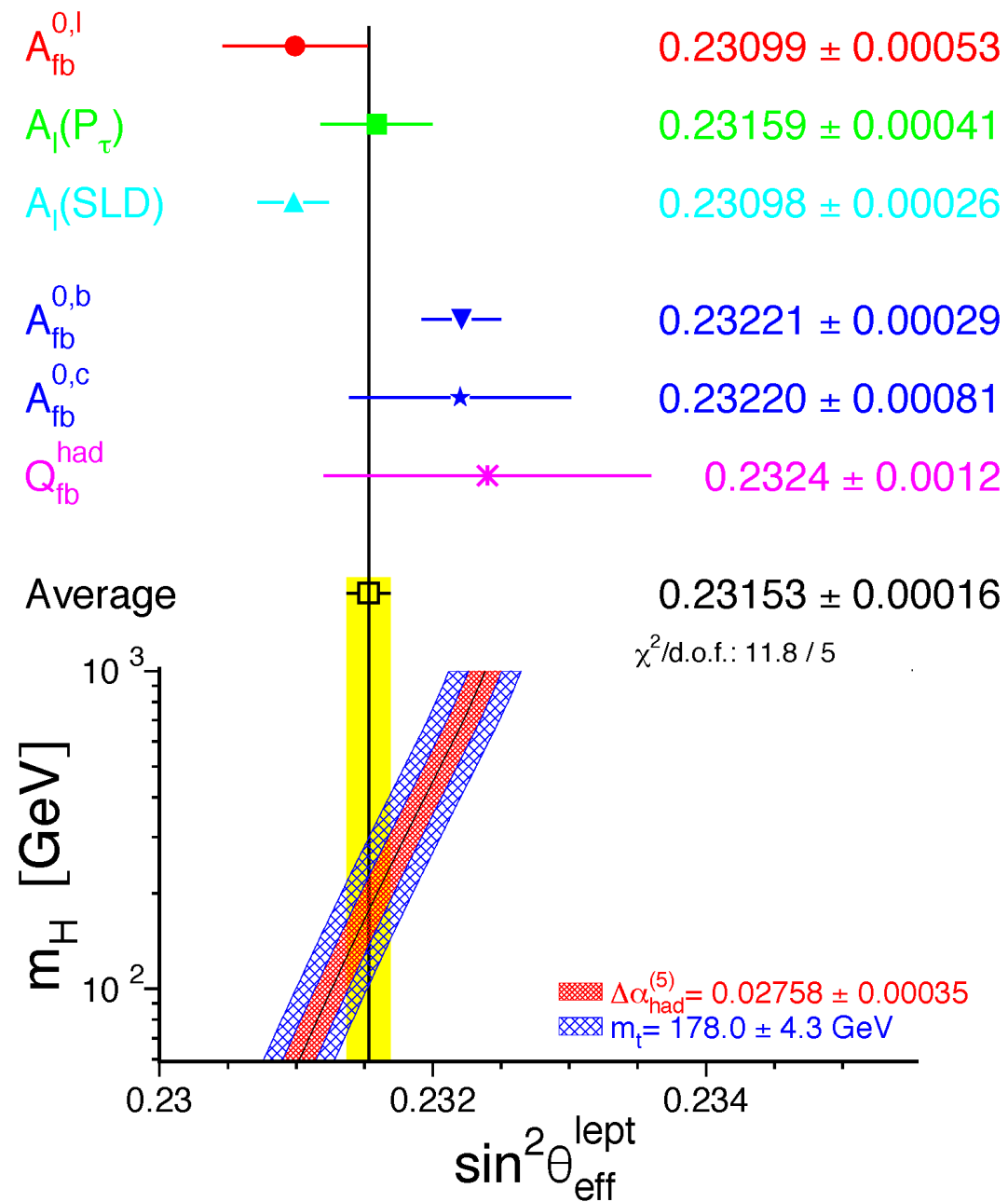


# Heavy quark studies at linear collider other than top – Quick review and outlook

Roman Pöschl



Snowmass Process EF03 Kickoff-Meeting 28/5/20



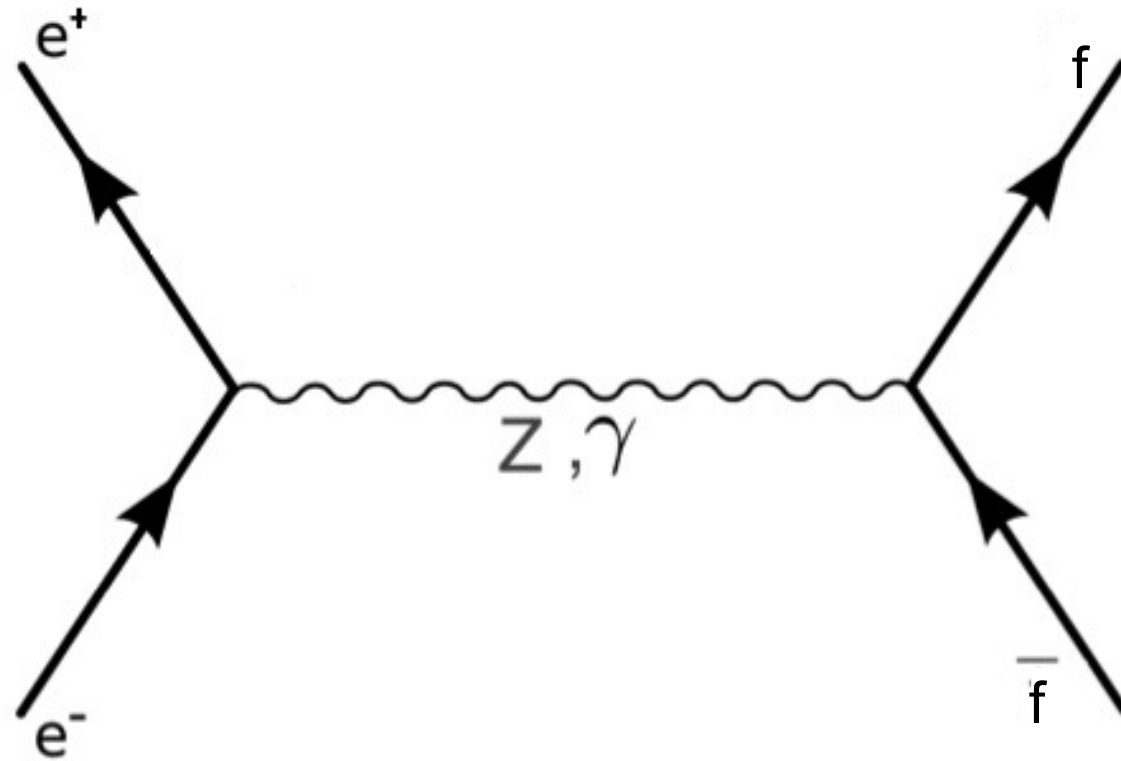
Most precise single Individual determination of  $\sin^2 \theta_{\text{eff}}^{\ell}$  from SLC

- Left-right asymmetry of leptons

- Most precise measurement of  $\sin^2 \theta_{\text{eff}}^{\ell}$  from forward backward asymmetry  $A_{FB}^b$  in  $ee \rightarrow bb$  at LEP

:

- Most precise determinations of  $\sin^2 \theta_{\text{eff}}^{\ell}$  differ significantly
  - Requires verification
  - **Heavy quark effect, effect on all quarks/fermions, no effect at all?**



Differential cross sections for (relativistic) di-fermion production\*:

$$\frac{d\sigma}{d\cos\theta}(e_L^- e_R^+ \rightarrow f \bar{f}) = \Sigma_{LL}(1 + \cos\theta)^2 + \Sigma_{LR}(1 - \cos\theta)^2$$

$$\frac{d\sigma}{d\cos\theta}(e_R^- e_L^+ \rightarrow f \bar{f}) = \Sigma_{RL}(1 + \cos\theta)^2 + \Sigma_{RR}(1 - \cos\theta)^2$$

\*add term  $\sim \sin^2\theta$  in case of non-relativistic fermions e.g. top close to threshold

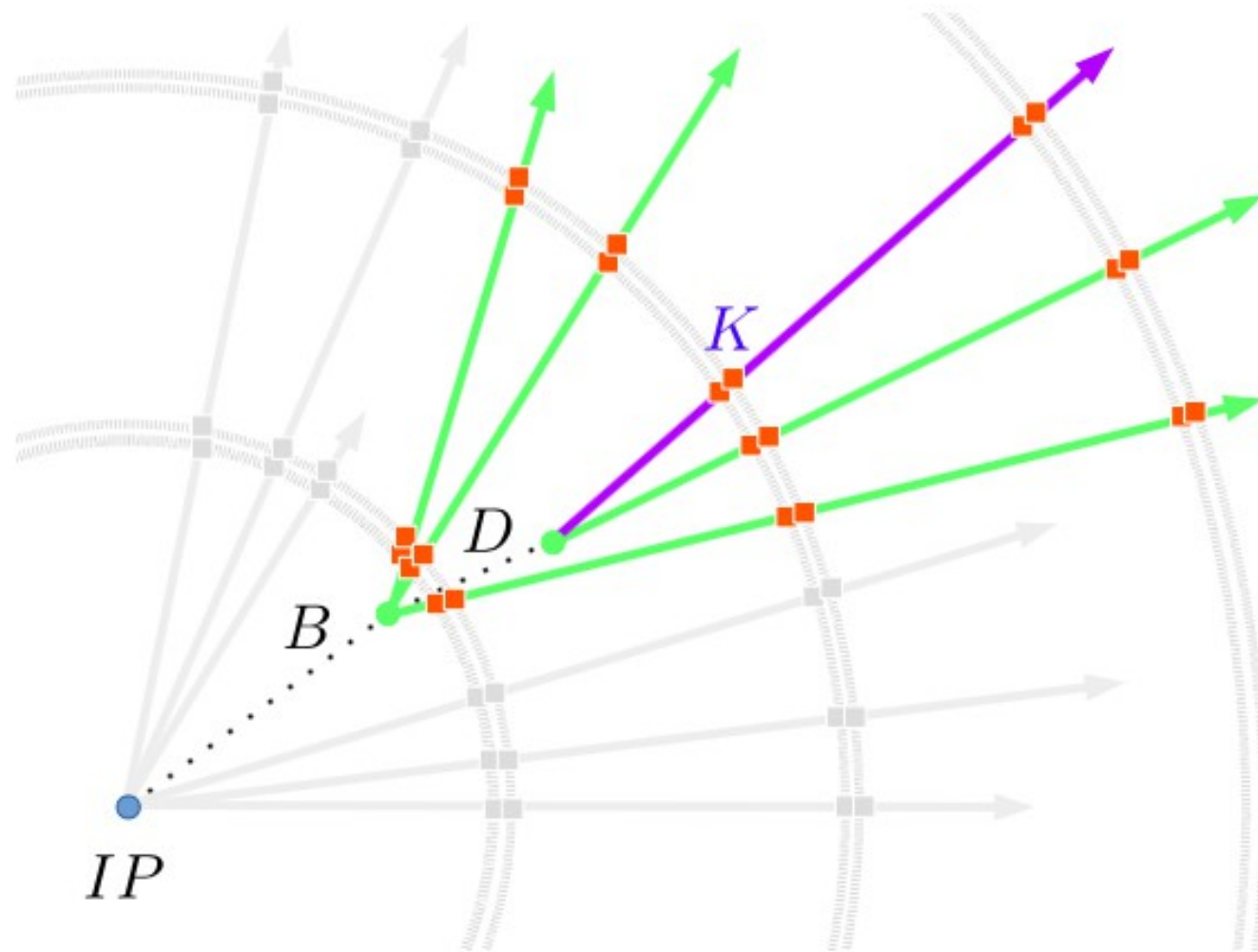
$\Sigma_{IJ}$  are helicity amplitudes that contain couplings  $g_L, g_R$  (or  $g_V, g_A$ )

$\Sigma_{IJ} \neq \Sigma_{I'J'} \Rightarrow$  (characteristic) asymmetries for each fermion

Forward-backward in angle, general left-right in cross section

**All four helicity amplitudes for all fermions only available with polarised beams**

**Here we focus first on (tt), bb and cc pair production**

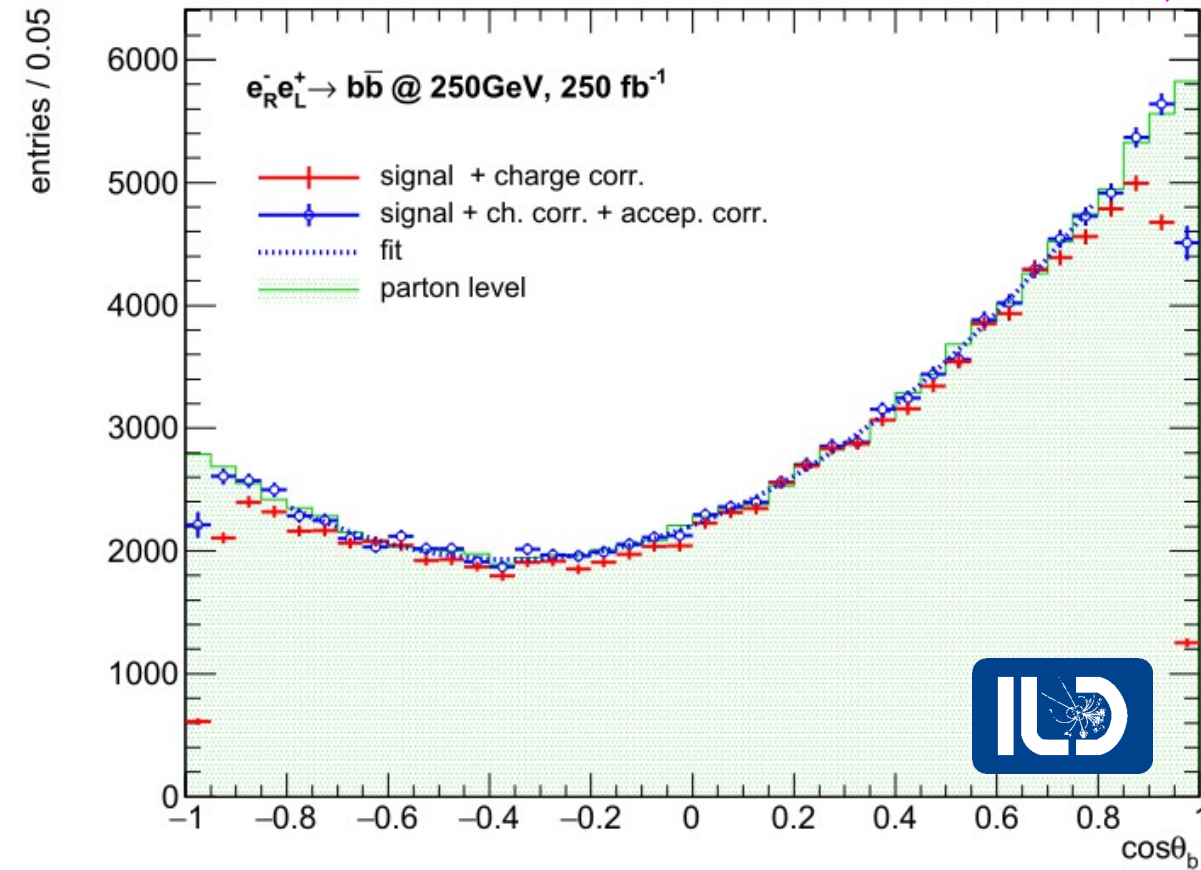
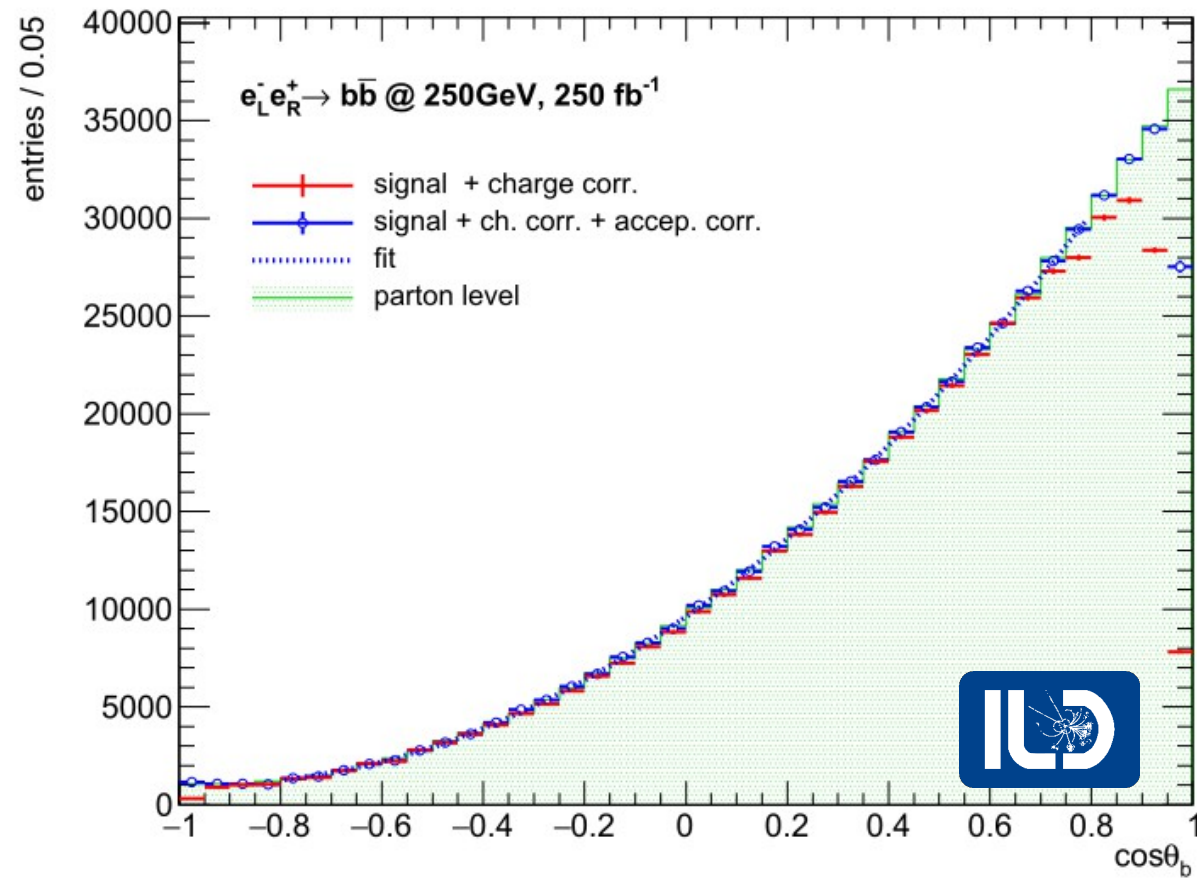


- Flavor tagging
  - Indispensable for analyses with final state quarks
- Quark charge measurement
  - Important for top quark studies,
  - indispensable for  $ee \rightarrow bb, cc, ss, \dots$
- Control of migrations:
  - Correct measurement of vertex charge
  - Kaon identification by  $dE/dx$  (and more)
- Future detectors can base the entire measurements on double Tagging and vertex charge
  - LEP/SLC had to include single tags and Semi-leptonic events

PhD thesis: S. Bilokin  
A. Irles



Arxiv:1709.04289, ILD Paper in progress

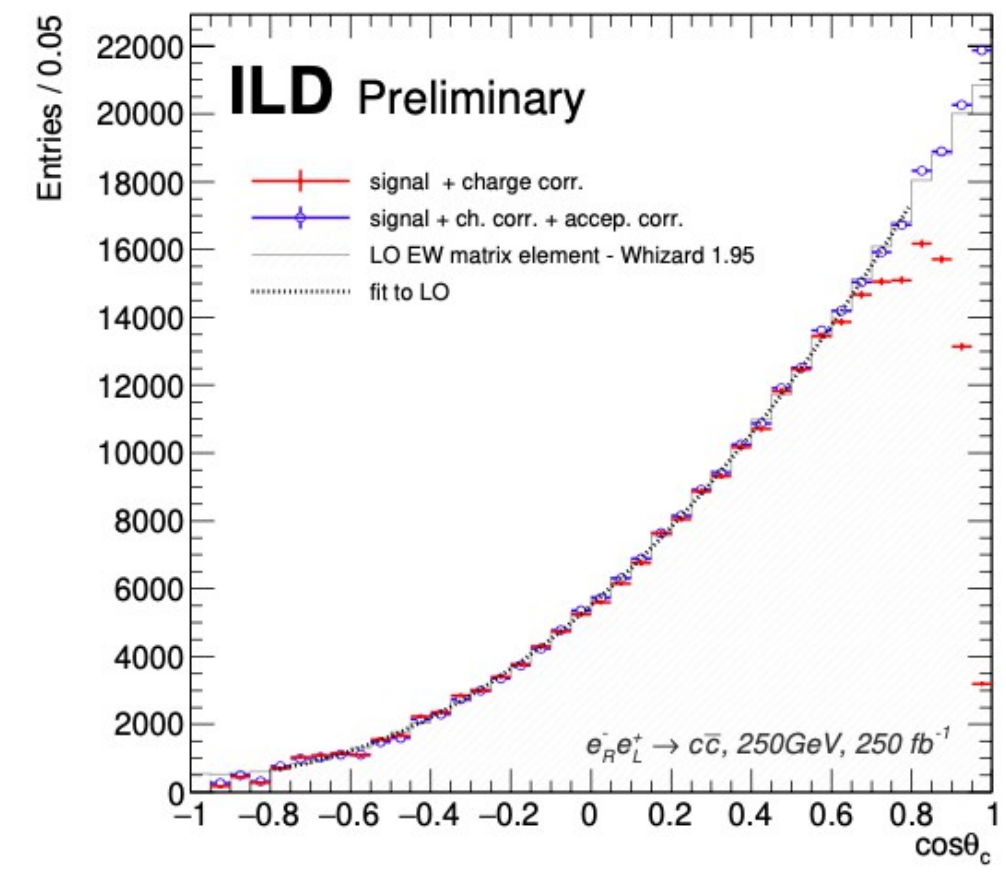
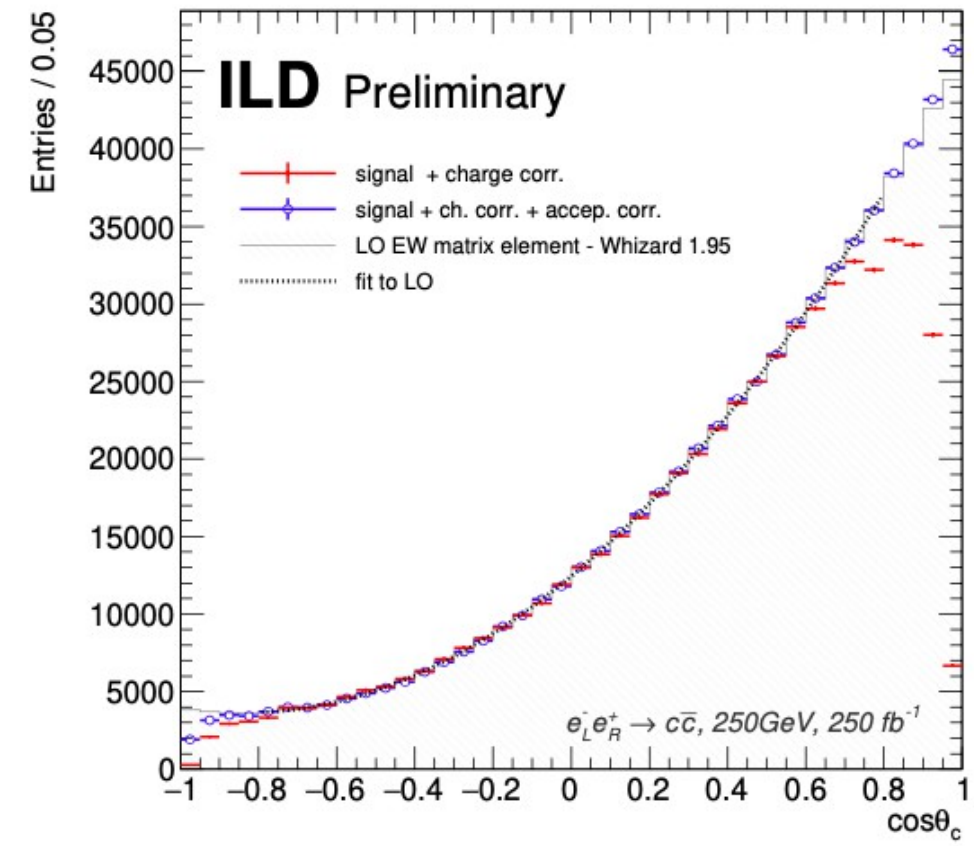


- Full simulation study (with ILD concept), Benchmark reaction
- Long lever arm in  $\cos \theta_b$  to extract from factors or couplings

Status and plans in coming year:

- Preliminary results exist
- Plan to publish a paper still in 2020 with full assessment of potential systematic errors
- Note that the precision will reach the per-mill level -> requires full control over detector performance

arxiv:2002.05805



Full simulation study (with ILD concept)  
 Long lever arm in  $\cos \theta_c$  to extract from factors or couplings

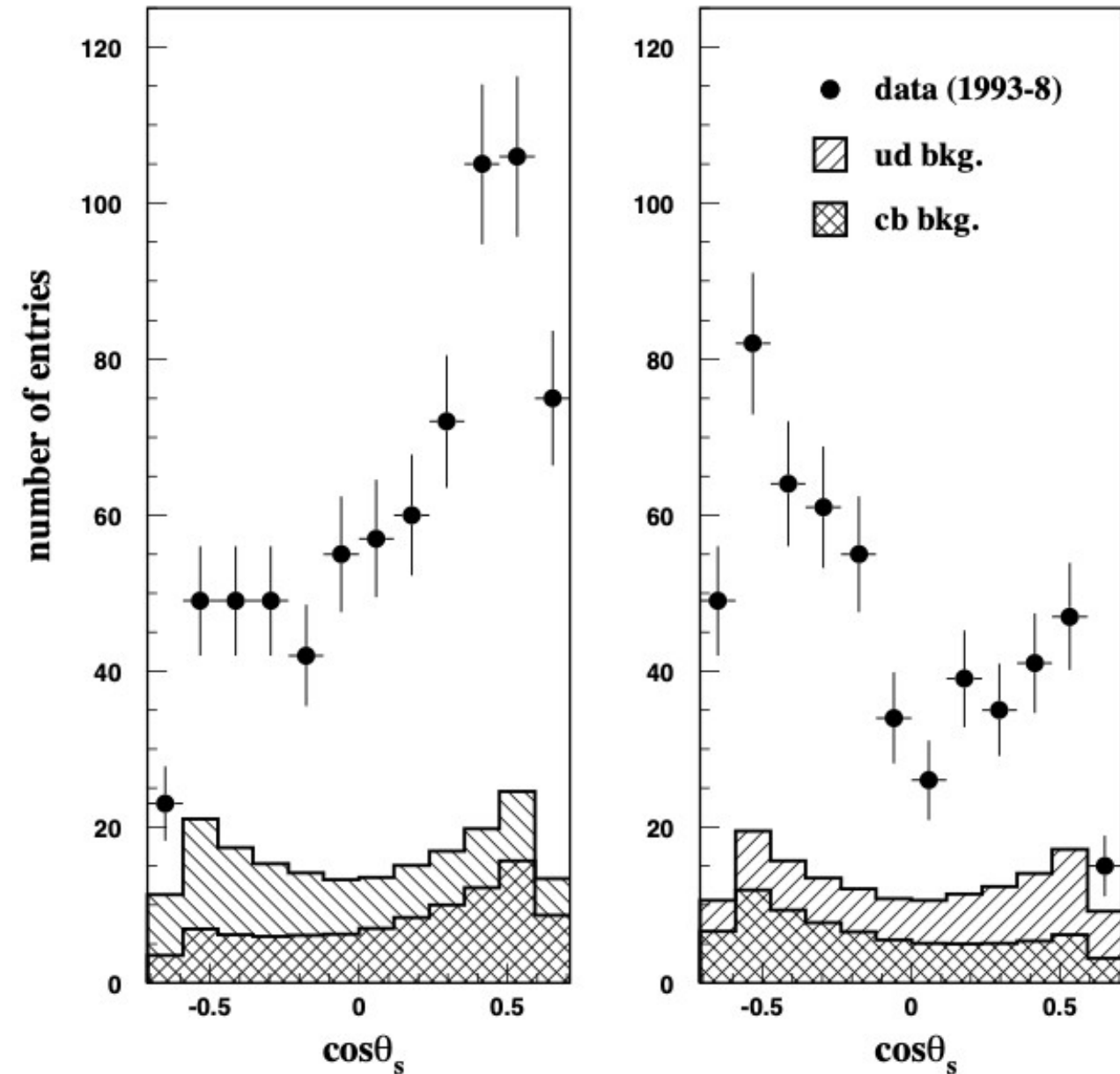
Status and plans in coming year:

- Preliminary results exist
- Plan to publish a paper in early 2021 with full assessment of potential systematic errors
- Note that the precision will reach the per-mill level -> requires full control over detector performance

## ee -->ss: SLD Analysis at Z Pole

neg. polarization

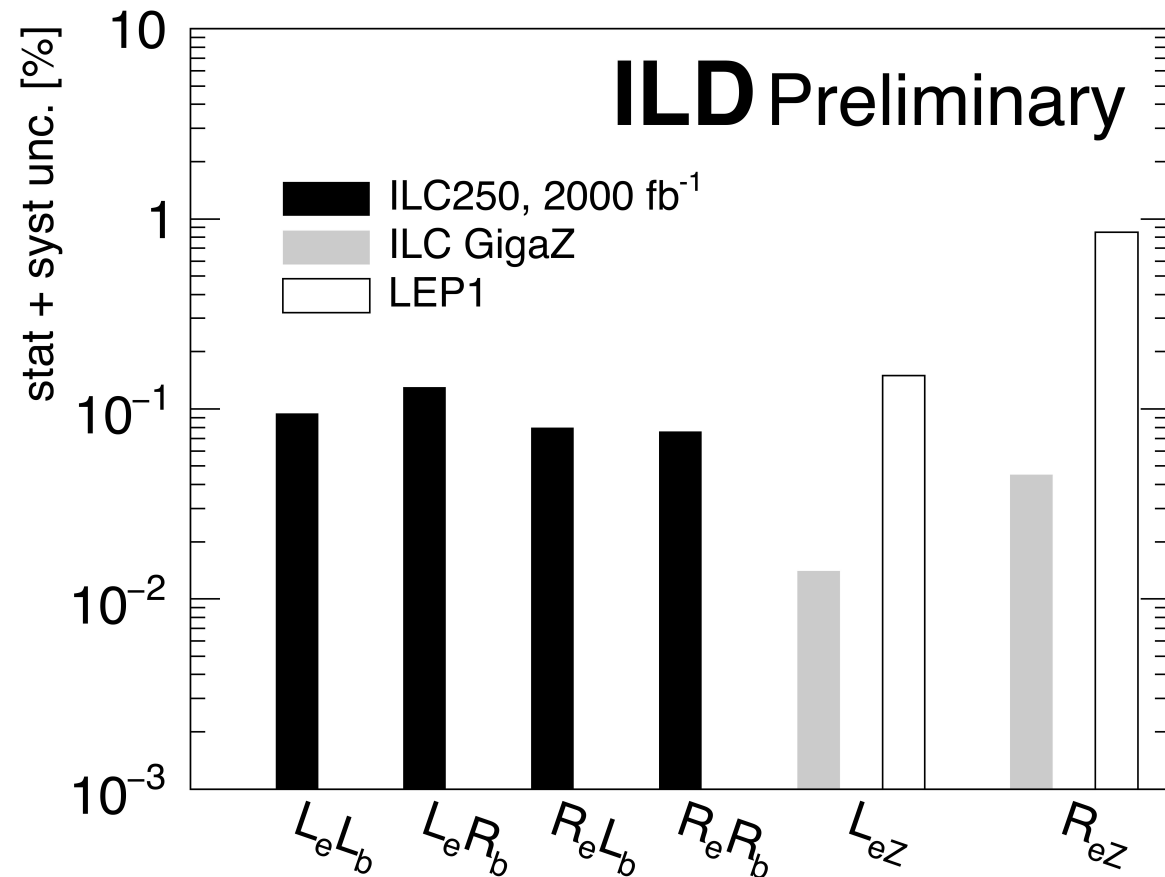
pos. polarization



- Extend the heavy quark analyses to light quarks to get full picture
- Optimise vertexing and particle ID (i.e .Kaon ID with full simulation studies)



## Example b-couplings (same observation for c-couplings, arxiv:2002.05805)



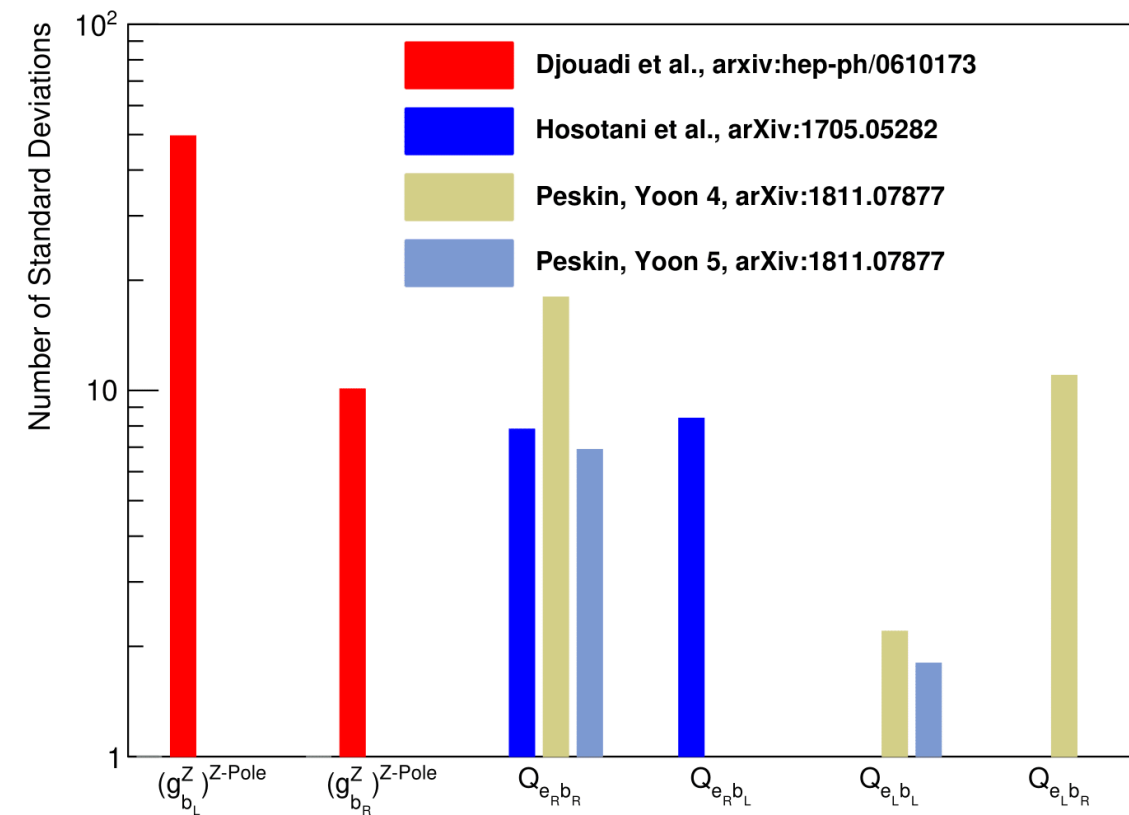
Couplings are order of magnitude better than at LEP

- In particular right handed couplings are much better constrained

New physics can also influence the Zee vertex

- in 'non top-philic' models

Full disentangling of helicity structure for all fermions only possible with polarised beams!!



Impressive sensitivity to new physics in Randall Sundrum Models with warped extra dimensions

- **Complete tests only possible at LC**
- **Discovery reach O(10 TeV)@250 GeV and O(20 TeV)@500 GeV**

Pole measurements critical input

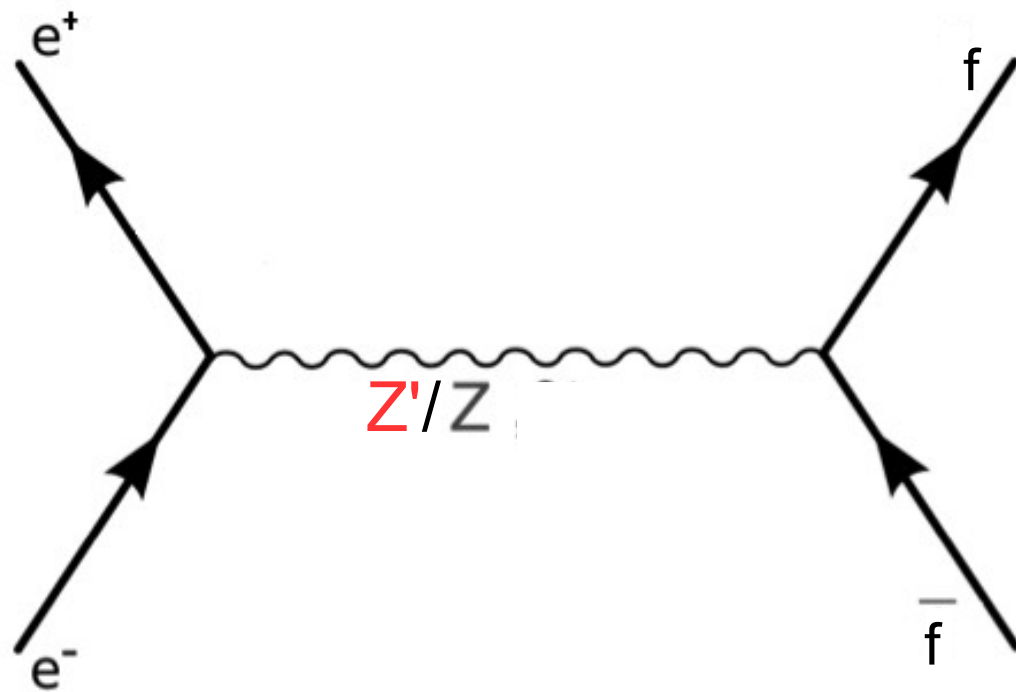
EF03 Kickoff • Only poorly constrained by LEP



- **Future e+e- machines are more than Higgs factories**
  - ee→qq (here b and c) processes will reach the per-mill precision
  - These processes are important to establish full patterns of electroweak couplings that may lead to discoveries
    - Note also that in some current models Higgs couplings will be agnostic to new physics
    - Marcel and Gauthier will also address issues like contact terms
  - Precision on Z-Pole as input to search for new physics at higher energies (see backup)
    - This may imply complementarity between linear and circular colliders if one can afford two machines
    - Detector optimisation to be studied (and to be brought in phase with the detector optimisation at higher energies)
      - One may exchange several times a vertex detector but not a calorimeter
- **Main challenge at future machines will be the control of systematic errors**
  - Experimentally
    - Vertex charge and particle ID
    - PFO for final state jets
  - Theoretically (not discussed)
    - Need at least NLO electroweak predictions (and MC programs) for correct interpretation of results

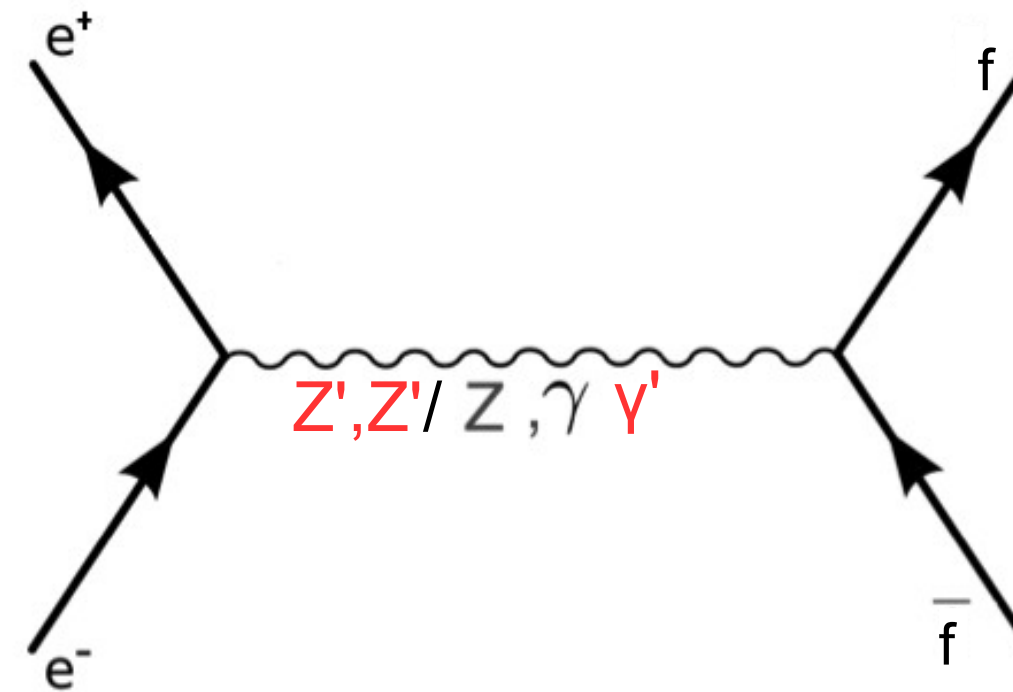
Backup

## On the Z-pole



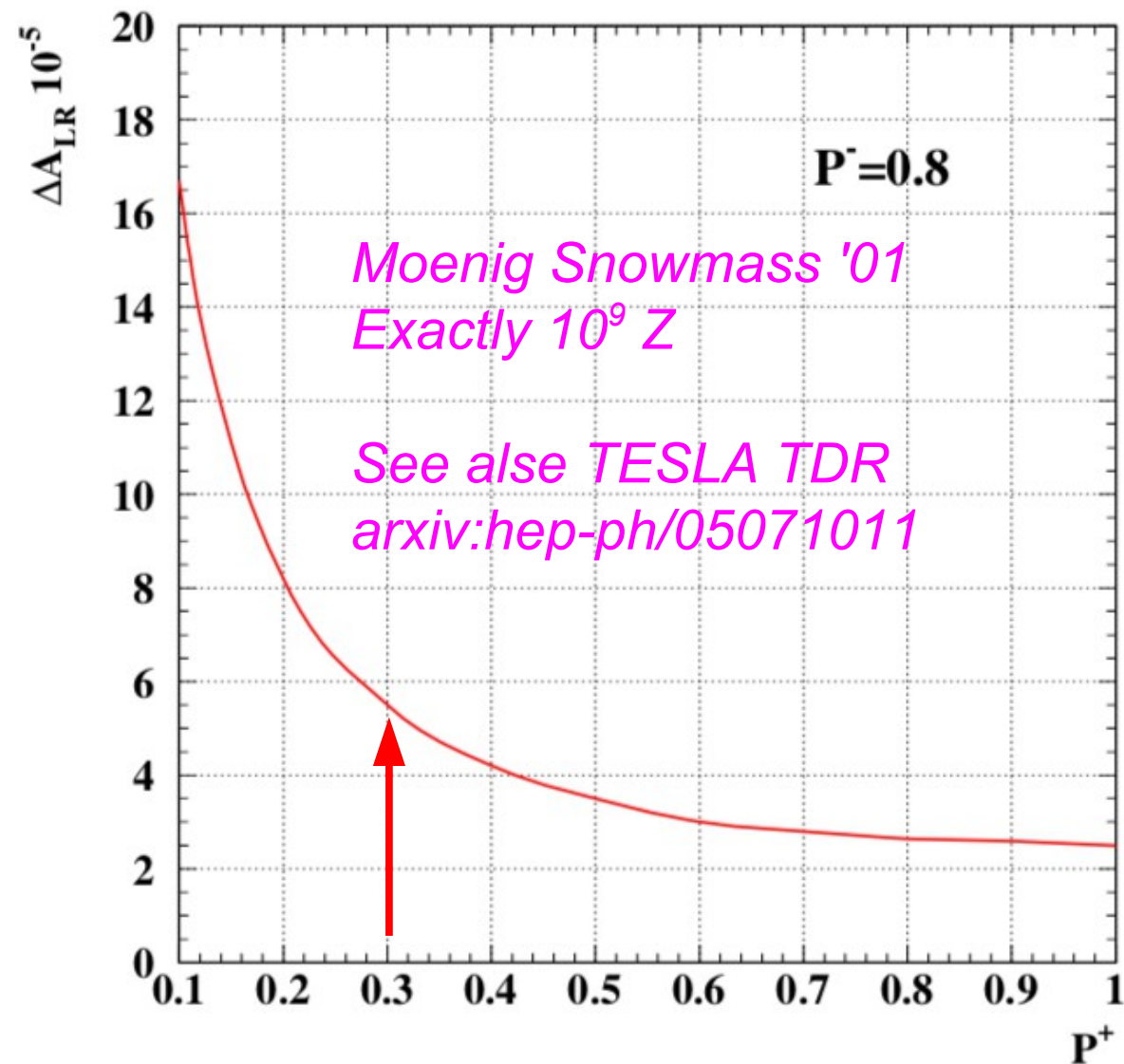
Sensitivity to  $Z/Z'$  mixing  
 Sensitivity to vector (and tensor) couplings of the  $Z$   
 •the photon does not “disturb”

## Above the Z-pole



Sensitivity to interference effects of  $Z$  and photon!!  
 Measured couplings of photon and  $Z$  can be influenced by new physics effects  
 Interpretation of result is greatly supported by precise input from  $Z$  pole

Blondel scheme: 
$$A_{LR} = \sqrt{\frac{(\sigma_{++} + \sigma_{-+} - \sigma_{+-} - \sigma_{--})(-\sigma_{++} + \sigma_{-+} - \sigma_{+-} + \sigma_{--})}{(\sigma_{++} + \sigma_{-+} + \sigma_{+-} + \sigma_{--})(-\sigma_{++} + \sigma_{-+} + \sigma_{+-} - \sigma_{--})}}$$



Blondel scheme independent of polarimeter precision

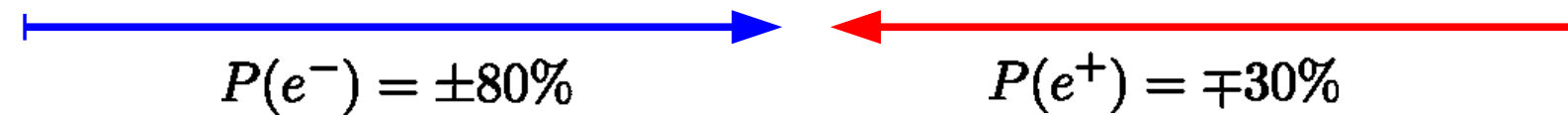
- Assumes perfect spin flip for polarised beams
- Residuals must be monitored by polarimeter
- Residual uncertainty of  $\Delta A_{LR} = 0.5 \times 10^{-4}$  seems possible
- The more positron polarisation the better (see backup)
- Don't forget energy dependency ( $dA_{LR}/d\sqrt{s} \sim 2 \times 10^{-5}/\text{MeV}$ )

Precision  $\Delta A_{LR} = 1 \times 10^{-4}$  is a realistic assumption for GigaZ

$\Rightarrow \delta \sin^2 \theta_{\text{eff}}^l \sim 1.3 \cdot 10^{-5}$



## With two beam polarisation configurations



There exist a number of observables sensitive to chiral structure, e.g.

$\sigma_I$	$A_{FB,I}^t = \frac{N(\cos\theta > 0) - N(\cos\theta < 0)}{N(\cos\theta > 0) + N(\cos\theta < 0)}$	$(F_R)_I = \frac{(\sigma_{tR})_I}{\sigma_I}$
x-section	Forward backward asymmetry	Fraction of right handed top quarks



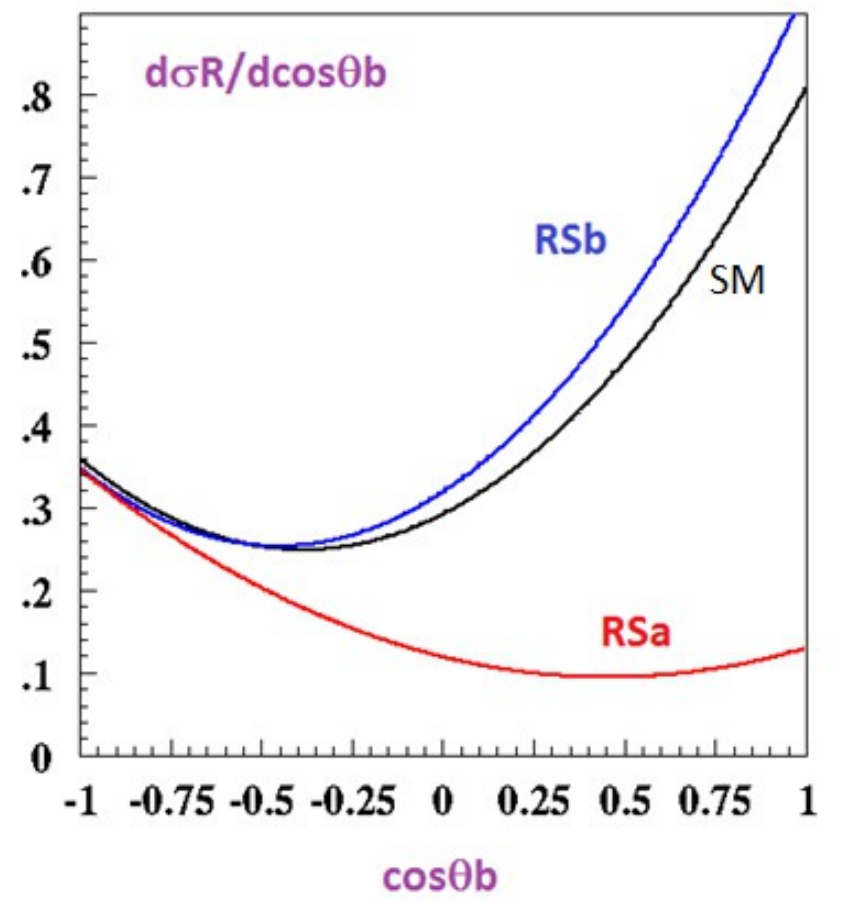
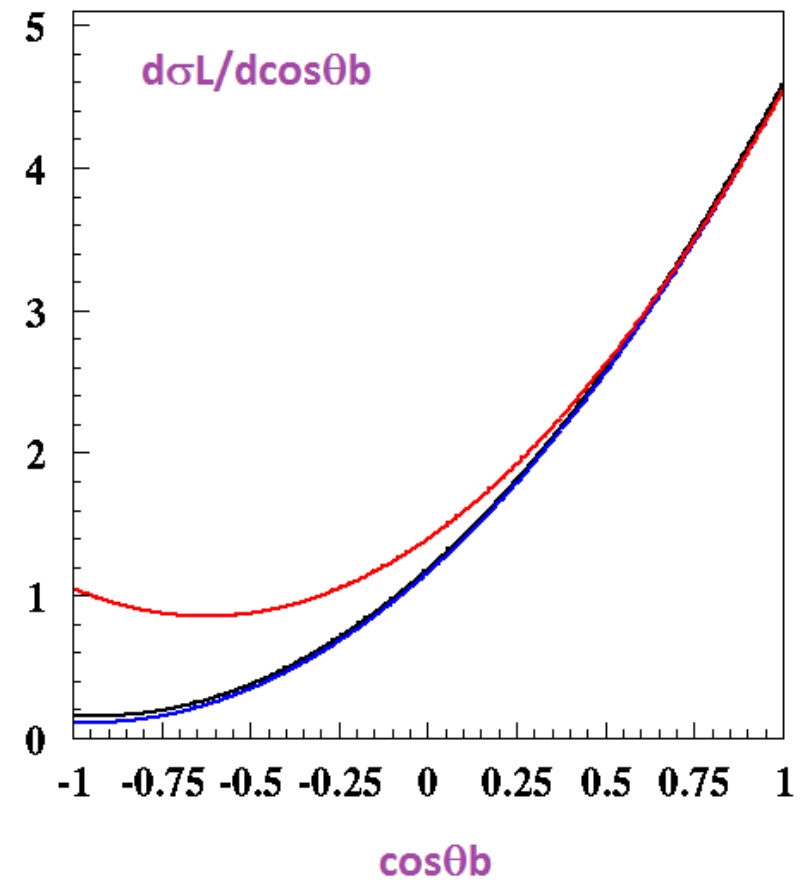
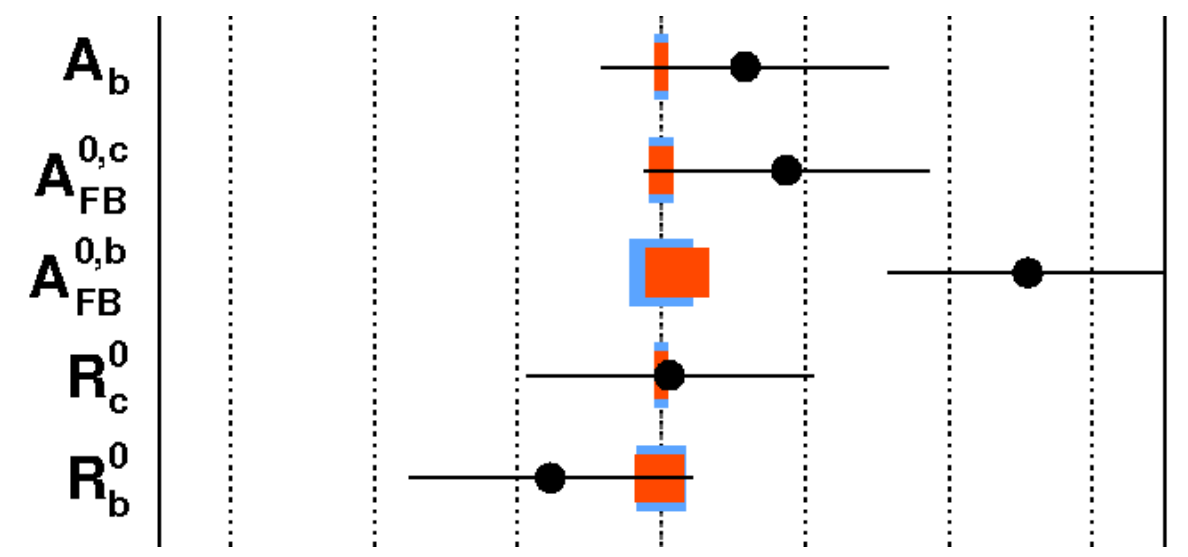
### Extraction of relevant unknowns

$$F_{1V}^\gamma, F_{1V}^Z, F_{1A}^\gamma = 0, F_{1A}^Z \quad \text{or equivalently} \quad g_L^\gamma, g_R^\gamma, g_L^Z, g_R^Z$$

$$F_{2V}^\gamma, F_{2V}^Z$$

$\sim 3\sigma$  in heavy quark observable  $A_{FB}^b$

$ee \rightarrow bb @ 250 \text{ GeV}$



- Is tension due to underestimation of errors or due to new physics?

- High precision  $e^+e^-$  collider will give final word on anomaly

Randall Sundrum Models Djouadi/Richard '06

- In case it will persist polarised beams will allow for discrimination between effects on left and right handed couplings (Remember  $Zb_l b_l$  is protected by cross section)

- Note that also B-Factories report on anomalies EF03 Kickoff

- There is a strong motivation to measure electroweak heavy quark couplings at the ILC
- New physics models predict deviations and b and c quarks are at the cross roads between 'top-philic' and 'non-top-philic' models
- Remember also LEP anomaly on  $A_{FB}^b$
- ILC with GigaZ is a unique opportunity for a complete set of measurements and an unambiguous interpretation of the results
- Relevant observables at GigaZ are  $A_b$  (see above) and

$$R_q = \frac{N_q}{N_{had}} = \frac{\Gamma_q}{\Gamma_{had}} = \frac{(g_q^L)^2 + (g_q^R)^2}{\sum_{i=1}^{n_q} [(g_i^L)^2 + (g_i^R)^2]}$$

- Here  $\Gamma_{had}$  is constrained by the fact that all hadrons are produced from the known quark species i.e.  $R_b + R_c + R_{uds} = 1$  and has therefore no error, but the  $g_i$  are correlated to fulfill this constraint
- The measured  $\Gamma_{had}$ , which is sensitive to the experimental Z mass resolution has to be considered as a consistency check

## Beam spot size



	FCCee	ILC	SLC	LEP
$\sigma_x$ [nm]	13700	516	1500	200000
$\sigma_y$ [nm]	36	7.7	500	2500

Source SLC, LEP, PDG

LEP

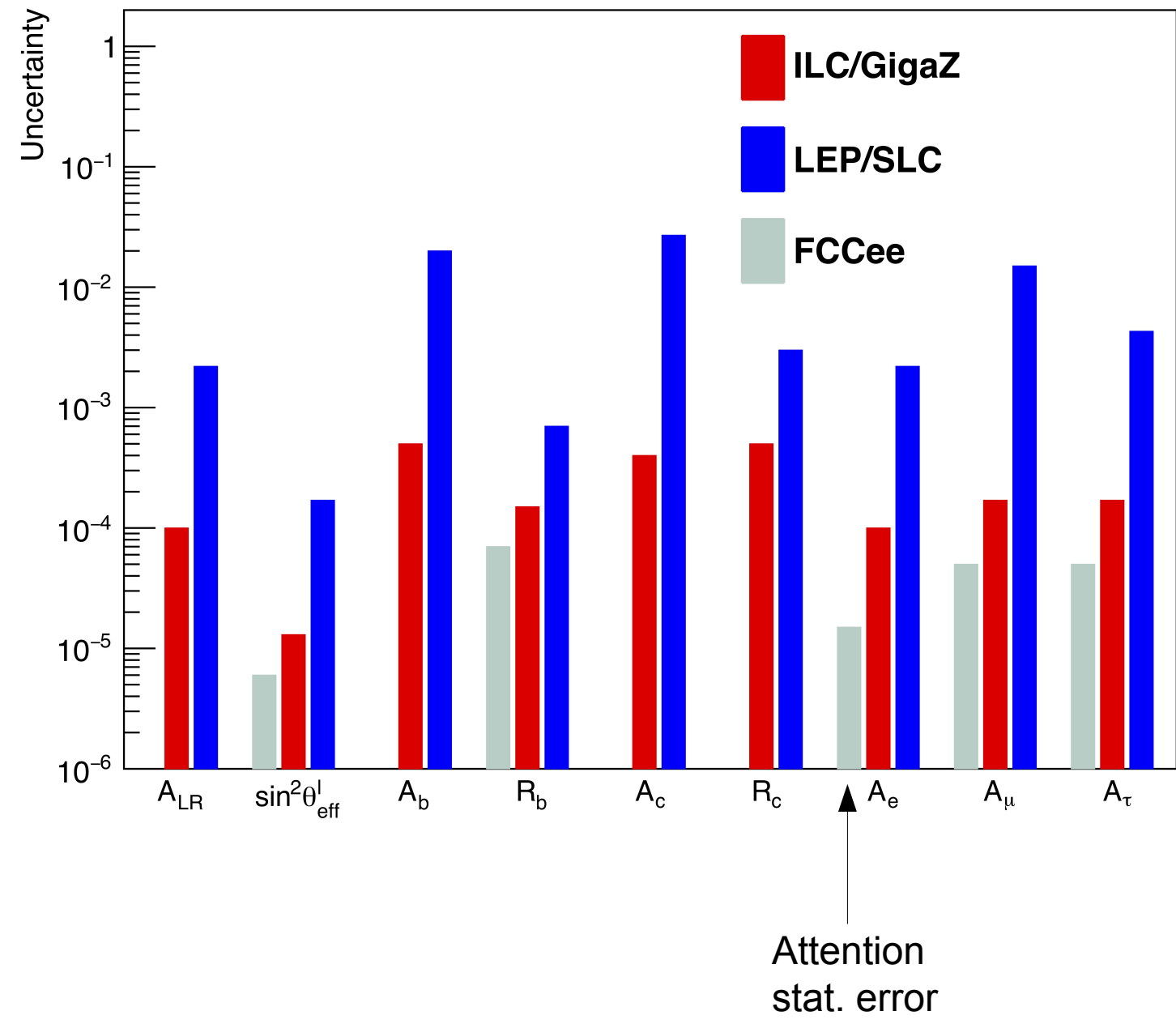
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SLC

>>

ILC





## Precise measurement of $\sin^2 \theta_{eff}^l$ .

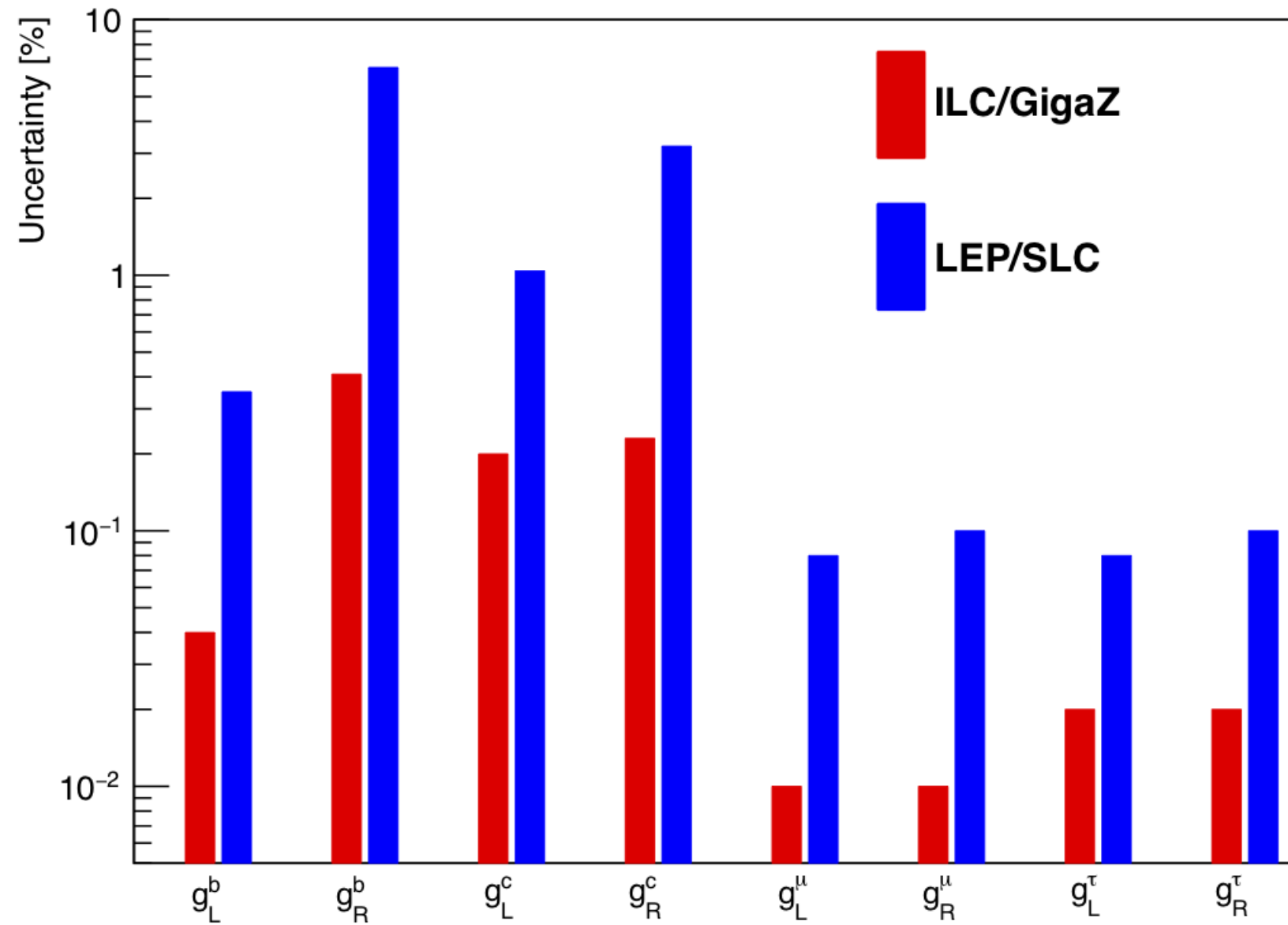
- Ten times better than LEP/SLD and competitive with FCC
- **Polarisation compensates for ~30 times luminosity**
- ... and  $A_{LR}$  at LC can benefit from hadronic Z decays
- **No assumption on lepton universality at LC**

## Complete test of lepton universality

- Precisions of order 0.05%

## Note excellent measurement of quark asymmetries

- See above for  $ee \rightarrow bb$  at 250 GeV



Partial fermion width:

$$R_f = \frac{N_f}{N_{had}} = \frac{(g_f^L)^2 + (g_f^R)^2}{\sum_{i=1}^{n_q} [(g_i^L)^2 + (g_i^R)^2]}$$

Sensitive to sum of coupling constants  
Available at linear and circular colliders

Left-right asymmetry:

$$A_{LR} = \frac{1}{|\mathcal{P}_{eff.}|} \frac{\sigma_L - \sigma_R}{\sigma_L + \sigma_R} = \mathcal{A}_e = \frac{(g_f^L)^2 - (g_f^R)^2}{(g_i^L)^2 + (g_i^R)^2} \sim 1 - 4\sin^2 \theta_{eff.}^l$$

Direct sensitivity to Zee vertex  
Only available at linear colliders due to beam polarisation  
Circular colliders need auxiliary measurement  
•e.g.  $P_\tau \sim A_e$

Forward-backward asymmetry:

$$A_{FB}^f = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B} = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f \text{ for } \mathcal{P}_e = 0.$$

“Classical” observable to study P-violating effects in ee->ff  
Available at circular and linear colliders  
**Without beam polarisation interpretation is always model dependent**

Left-right-forward-backward asymmetry:

$$A_{FB,LR}^f = \frac{(\sigma_F - \sigma_B)_L - (\sigma_F - \sigma_B)_R}{(\sigma_F + \sigma_B)_L + (\sigma_F + \sigma_B)_R} = -\frac{3}{4} \mathcal{A}_f$$

Combination of asymmetries above  
Only available linear colliders due to beam polarisation  
Direct and model independent measurement of  $A_f$

Helicity amplitudes can be analysed in several ways (not mutually exclusive):

Oblique Parameters W, Z:

$$Q_{eifj} = Q_e^\gamma Q_f^\gamma + \frac{g_{e_i}^Z g_{f_j}^Z}{\sin^2 \theta_W \cos^2 \theta_W} \frac{s}{s - M_Z^2 + i\Gamma_Z M_Z} + \frac{s}{m_W^2} f_{i,j}(W, Y)$$

Contact interactions with e.g. compositeness scale  $\Lambda$ :

$$Q_{eifj} = Q_e^\gamma Q_f^\gamma + \frac{g_{e_i}^Z g_{f_j}^Z}{\sin^2 \theta_W \cos^2 \theta_W} \frac{s}{s - M_Z^2 + i\Gamma_Z M_Z} + \frac{g_{\text{contact}}^2}{2\Lambda^2} \eta_{eifj}$$

New propagators in concrete models of new physics:

$$Q_{eifj} = Q_e^\gamma Q_f^\gamma + \frac{g_{e_i}^Z g_{f_j}^Z}{\sin^2 \theta_W \cos^2 \theta_W} \frac{s}{s - M_Z^2 + i\Gamma_Z M_Z} + \sum \frac{g_{e_i}^{Z'} g_{f_j}^{Z'}}{\sin^2 \theta_W \cos^2 \theta_W} \frac{s}{s - M_{Z'}^2 + i\Gamma_{Z'} M_{Z'}}$$

Always with I,j being the helicities of the initial state electron e and the final state fermion f