Non-perturbative QCD Dynamics at LHC and Beyond – an Experimentalist's View

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Preliminaries





This talk is not a comprehensive review...

... due to lack of time and ...

Restricted to jets in proton (*pp*) collisions at the LHC Some remarks regarding HL-LHC – high lumi operations Mostly latest findings and present day hot topics

... a bias towards ATLAS results

Better personal overview

After >25 years with this experiment

Not an intentional disregard of e.g. CMS, LHCb, ALICE, ...

They are all producing great QCD (and other) results

Focus on experimental aspects in hadronic final state reconstruction

Highly selective

Modelling issues in jet reconstruction performance and SM jet measurements

Some comments on possible future R&D

Parton shower (fragmentation & hadronization) and underlying event (multiple parton interactions, MPI) tuning



This talk



Role of soft emissions in jet reconstruction precision Small-*R* jets



Anti- k_t with R = 0.4 – standard tools for measurements & searches Underlying event (UE) contribution and pile-up (PU) mitigation

Large-*R* jets (most popular configurations)

Anti- k_t with R = 0.8 R = 1.0 - tools for tagging heavy particle decays

Constituent level pile-up suppression & jet grooming

Modeling parton emissions at all(?) scales

Tuning efforts

Selected recent results

Observations from measurements

Jet mass and substructure observables

Event shapes

Some ideas on moving on

Input to tuning

Inclusive measurements of emissions in jets at ~all scales

QCD at softer scales

Exploring energy-energy correlations in jets

Soft QCD Contribution to Jet Reconstruction





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Jet Reconstruction Sequence in ATLAS





Slide 5



Jet Reconstruction Sequence in ATLAS





*the stochastic nature of most grooming techniques including the jet-area-based pile-up suppression leads to the removal of an unknown amount of the UE contribution to the jet. August 3, 2020 Slide 6



Jet Calibration Uncertainties





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Remove pile-up from input signal

Attractive for large-R jets ...

Improve precision and resolution in measurements using structural ($p_{\rm T}$ -flow) jet information (e.g. for tagging heavy particle decays)

... and event shapes

Recovering/extracting soft signal contributions to $E_{\rm T}^{\rm miss}$ in the presence of pile-up

Non-jet context

Stochastic approach to remove diffuse emissions from pile-up - no correlation to hard scatter interactions

Local pile-up density measures with various area definitions (Voronoi-area-based, constituent subtraction with ghosts, SoftKiller, PUPPI...)

Combination of approaches -

ConstituentSubtraction + SoftKiller (CS+SK) increase efficiency

Combination with jet substructure reconstruction

Jet grooming (e.g. SoftDrop)

Reduces sensitivity to soft emissions in already biased input signals

Enhanced sensitivity to harder structure

Popular tagger inputs are jet mass & ratio of 2-point to 3-point energy-energy correlation functions D_2



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Most recent ATLAS tune: A14 (*pp* at $\sqrt{s} = 7$ TeV)

Underlying event

Evolution of transverse activity with hardest track & calorimeter jets

Parton shower

Jet structure – track jet properties, jet masses & substructure variables, jet shapes in inclusive jet and $t\bar{t}$ final states

Additional emissions beyond LO $2 \rightarrow 2$

Dijet azimuthal decorrelation Gap jet fraction in $t\bar{t}$ final states 3/2 jet ration

Z-boson $p_{\rm T}$

Performs well in Run 2

Used for all physics analysis in pp collisions at $\sqrt{s} = 13$ TeV



Dijet azimuthal decorrelations for $210 < p_{\perp}^{max}/GeV < 260$



Anti- k_T jets, R = 1.0, 300 GeV $< p_{\perp} < 400$ GeV $/\sigma d\sigma/d\sqrt{d_{23}}$ [GeV⁻¹] 0.12 ATLAS Simulation 0.1 ATLAS Data AU₂ CTEQ6L1 0.08 Monash A14-CTEQ 0.06 A14-MSTW A14-NNPDF 0.04 A14-HERA 0.02 0 1.4 MC/Data 1.2 0.8 0.6 0



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August 3, 2020

ATLAS-PHYS-PUB-2014-021







Remove soft wide angle emissions from jet

Traverses jet clustering history Drops branches reflecting wide angle soft radation Insensitive to nonglobal logarithmic corrections – removes radiation leaving the jet cone and emit particles back into the cone

SoftDrop variables

Mass, $p_{\rm T}$ balance of splitting, angular distance of emission Found to be calculable to NLL and NNLL accuracy

Excellent agreement with data in regions where calculations are accurate

Non-perturbative effects prominent theoretical uncertainty – dominate higher order effects

Free parameters z, β , jet radius RRegulate sensitivity to radiation scales







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Detector-level ...



Systematic uncertainties \rightarrow

... unfolded measurements





Soft Drop Mass – Interpretation



$\beta = 0$

Removes soft→collinear radiation



$\beta = 2$

Removes soft radiation



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 $\beta = 0$

Removes soft→collinear radiation





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Soft Drop Mass – Interpretation





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 $\beta = 0$

Removes soft→collinear radiation

 $\beta = 2$

Removes soft radiation



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 $\beta = 0$

Removes soft→collinear radiation

 $\beta = 2$

Removes soft radiation



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ATLAS Coll., Phys.Rev.D 101 052007 (2020)

Updated results ...

More SoftDrop variables

 $p_{\rm T}$ -balance of splitting $z_g = \min(p_{{\rm T},j1}, p_{{\rm T},j2})/(p_{{\rm T},j1} + p_{{\rm T},j2})$

Opening angle of splitting $r_g = \sqrt{(y_1 - y_2)^2 + (\varphi_1 - \varphi_2)^2}$

Calorimeter & track-based analyses

Updated/new systematic uncertainties

... provide a deep look in to the dynamics of jet formation

Evaluation of splitting history at various angular and momentum scales





TeV@Les Houches 2019

Use of (SoftDrop) substructure observables in tuning SoftDrop mass from studies at Les Houches 2019 Promising first investigation







Event shapes with jets

arXiv:2007.12600 [hep-ex]

Proxy for energy flow shapes in collision event

Measurement tests prediction power of fixed-order calculations, parton shower modeling, etc.

Clear expectation values for given topology

Shapes vanish for $2 \rightarrow 2$ processes with perfect forward-backward (back-to-back in transverse plane) symmetry – at maximum for uniform energy (transverse momentum) distribution

Probe for multi-jet energy flow at highest scales

 $\mathcal{O}(\text{TeV})$ for $\sqrt{s} = 13 \text{ TeV}$

Evaluated in multi-jet final states ($n^{\text{jet}} \ge 2$) as function of hardness of interaction

Representative observable for interaction activity is $H_{T2} = p_T^{lead} + p_T^{sublead}$

Measurement

Jet and event selection

Consider only fully calibrated anti- k_t jets with R = 0.4 clustered from particle flow objects with $p_{\rm T}^{\rm jet} > 100$ GeV, $\left|\eta^{\rm jet}\right| < 2.4$

Multi-jet events with $n^{\text{jet}} \ge 2$,

 $H_{\rm T2} > 1$ TeV selected

Presentation of results

Differential cross-sections as ratio to fiducial cross section $\sigma(n^{\text{jet}} \ge 2)$

 $1/\sigma(n^{\text{jet}} \ge 2) d\sigma/d\{T_{\perp}, T_m, S_{\perp}, A, C, D\} \text{ in } (H_{\text{T2}}, n^{\text{jet}}) \text{ bins}^*$

Unfolded data is compared to various Monte Carlo generators

*see additional material for description of all used event shape observables



Modeling of Jet Multiplicities





Measured as function of n^{jet} Evaluated in same three regions of H_{T2} used for event shape measurements – provides normalization Modeling $d\sigma/dn^{\text{jet}}$ shapes

Pythia 8.235

Fiducial cross section

- $2 \rightarrow 2$, LO accuracy
- Generally good agreement for all n^{jet}

Sherpa 2.2.1

 $2 \rightarrow \{2, 3\}$, LO accuracy (multi-leg) Overestimation (increasing) for $n^{\text{jet}} > 4$

Herwig 7.1.3 (angular ordered PS)

 $2 \rightarrow 2$ NLO accuracy, $2 \rightarrow 3$ LO Good description with slight underestimation for $n^{\text{jet}} > 6$

Herwig 7.1.3 (dipole PS)

 $2 \rightarrow 2$ NLO accuracy, $2 \rightarrow 3$ LO Good description for low n^{jet} underestimation for higher n^{jet}

MadGraph5 aMC 2.3.3

 $2 \rightarrow \{2, 3, 4\}$ NLO accuracy Good description for low n^{jet} . underestimation for higher n^{jet}

Modeling normalization

Well predicted at low n^{jet}

Only small differences between models

Large spread in normalization at high n^{jet}

Sherpa predicts 30% more than data Herwig (dipole PS), MadGraph predict 30% less

arXiv:2007.12600 [hep-ex]

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≥6

n^{jet}

5

4



[hep-ex]

arXiv:2007.12600

Transverse Thrust with Jets





Observations

Evolution with increasing hardness of interaction

More events with more isotropic flow at softer interactions (lower H_{T2}) Increasing H_{T2} yields increased contribution from events with close to back-to-back flow patterns

Comparisons to models

Evaluation of predictions

Generally fewer isotropic events in MC than in data at low $n^{\rm jet}$ – better agreement at higher jet multiplicities

Shapes of cross sections

None of the considered MC generators gives good description in full phase space Similar distribution shapes at high n^{jet} from all considered model

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Phys.Rev.Lett. 124, 222002

LundPlane with tracks

2-dim image of radiation pattern in jet

Wide range of emission angles and radiated energies in one measurement



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Compare emission modeling







Adjust sensitivity to scales of emission

Energy-energy correlations

Employs N-point energy correlation functions between jet constituents

 $ECF(N,\beta) = \sum_{i_1 < i_2 < \dots < i_N \in Jet} \left(\prod_{a=1}^N p_{Ti_a} \right) \left(\prod_{b=1}^{N-1} \prod_{c=b+1}^N R_{i_b i_c} \right)^{\beta}$

o-point corr. ECF(0, β) = 1 1-point corr. ECF(1, β) = $\sum_{i \in j \in I} p_{Ti}$ 2-point corr. ECF(2, β) = $\sum_{i < j \in j \in I} p_{Ti} p_{Tj} (R_{ij})^{\beta}$ 3-point corr. ECF(3, β) = $\sum_{i < j < k \in j \in I} p_{Ti} p_{Tj} p_{Tk} (R_{ij} R_{ik} R_{jk})^{\beta}$ 4-point corr. ECF(4, β) = $\sum_{i < j < k < l \in j \in I} p_{Ti} p_{Tj} p_{Tk} p_{Tl} (R_{ij} R_{ik} R_{il} R_{jk} R_{jl} R_{kl})^{\beta}$

Larkoski, Salam, Thaler, JHEP 1306(2013) 108





Adjust sensitivity to scales of emission

Energy-energy correlations

Employs N-point energy correlation functions between jet constituents

 $ECF(N,\beta) = \sum_{i_1 < i_2 < \dots < i_N \in Jet} \left(\prod_{a=1}^N p_{Ti_a} \right) \left(\prod_{b=1}^{N-1} \prod_{c=b+1}^N R_{i_b i_c} \right)^{\beta}$

Double-ratios indicate jet source

Generally $C_N^{(\beta)} = (\text{ECF}(N+1,\beta)\text{ECF}(N-1,\beta))/ECF(N,\beta)^2$

 $C_1^{(\beta)}$ is useful for quark/gluon separation for small $\beta \simeq 0.2$ – exploits different soft radiation patterns (uses 2-point correlations)

 $C_2^{(\beta)}$ helps with boosted W/Z/H identification with $\beta \simeq 0.5$ better for high mass and $\beta \simeq 2$ better for lower mass resonances – at a fixed p_T (uses 3-point correlations)

 $C_3^{(\beta)}$ distinguishes QCD jets from boosted top quarks best for $\beta \approx 1 - 2$ (uses 4-point correlations)

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Adjust sensitivity to scales of emission

Energy-energy correlations

Employs N-point energy correlation functions between jet constituents

 $ECF(N,\beta) = \sum_{i_1 < i_2 < \dots < i_N \in Jet} \left(\prod_{a=1}^N p_{Ti_a} \right) \left(\prod_{b=1}^{N-1} \prod_{c=b+1}^N R_{i_b i_c} \right)^{\beta}$

Experimentally challenging for soft QCD measurements

Track only measurement

Small effect by pile-up – tracking in dense environments at HI-LUM?

Calorimeter only measurement

Soft energy flow in jet distorted by pile-up – constituent level pile-up mitigation may hurt sensitivity to soft emissions...

Particle flow measurement

Charged response equal to track measurement – neutral response may not help much...

Could be studied with Run 2/3





Access to soft emissions in hard objects

Exploring regions of non-perturbative QCD experimentally at high precision

Looking inside of jets helps to increase understanding of radiation patterns \rightarrow input to tunes and validation of new models/calculations

Experimental limitations

Pile-up adds diffuse emissions on top of the hard scatter

Track-based measurements

Constituent-level pile-up mitigation in particle flow measurements

Hard emission modeling not too well controlled

ATLAS analysis of event shapes

More work needed on the side of calculations – or matching calculations with tuned parton showers and MPI?

May be an issue in multi-jet searches...

Soft QCD effects

Several new observables available

Testing of tuning, evaluation of achievable precision and biases?

Tuning using 1-dim distributions of jet shapes and structure is already established

Tuning to $\geq 2 \text{ dim "images"}$

Machine learning in soft QCD tuning seems a very interesting route to follow – exploration of otherwise had to determine correlations, ranking of observables by Frameworks like Professor are already performing multi-dimensional fits, as far as I know...

Additional Material





Event shapes with jets

Proxy for energy flow (shapes) in collision event

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Event shape	Name	Comments
T_{\perp}	Transverse thrust	$\tau_{\perp} = 1 - T_{\perp}, 0 \le \tau_{\perp} < 1 - 2/\pi, \tau_{\perp} \nearrow $ back-to-back topology
T _m	Transverse thrust, minor component	$0 \le T_m < 2/\pi$, $T_m \nearrow$ increased energy flow outside of plane spanned by thrust and beam axes
S_{\perp}	Transverse sphericity	from eigenvalues { μ_k } of transverse sphericity tensor \mathcal{M}_{xy} , $S_{\perp} = 2\mu_2/(\mu_1 + \mu_2), 0 \le S_{\perp} \le 1, \downarrow$ back-to-back, \uparrow isotropic
Α	Aplanarity	from eigenvalues { λ_k } of sphericity tensor \mathcal{M}_{xyz} , $A = \frac{3}{2}\lambda_3$, $0 \le A \le 1$, $A \nearrow \Rightarrow$ event less planar
С	3-jet observable	$C=3(\lambda_1\lambda_1+\lambda_1\lambda_3+\lambda_2\lambda_3),C=0$ for $n^{\rm jet}<3,0< C\leq 1$ for $n^{\rm jet}>2$
D	4-jet observable	$D = 27(\lambda_1 \lambda_2 \lambda_3), 0 \le D \le 1, D = 0$ if all jets are in same plane

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Linearized Sphericity Tensor

$$\mathcal{M}_{xyz} = \frac{1}{\sum_{i} |\vec{p}_{i}|} \sum_{i} \frac{1}{|\vec{p}_{i}|} \begin{pmatrix} p_{x,i}^{2} & p_{x,i}p_{y,i} & p_{x,i}p_{z,i} \\ p_{y,i}p_{x,i} & p_{y,i}^{2} & p_{y,i}p_{z,i} \\ p_{z,i}p_{x,i} & p_{z,i}p_{y,i} & p_{z,i}^{2} \end{pmatrix}$$

Transverse Linearized Sphericity Tensor

$$\mathcal{M}_{xy} = \frac{1}{\sum_{i} |\vec{p}_{i}|} \sum_{i} \frac{1}{|\vec{p}_{i}|} \begin{pmatrix} p_{x,i}^{2} & p_{x,i}p_{y,i} \\ p_{y,i}p_{x,i} & p_{y,i}^{2} \end{pmatrix}$$





Examples: transverse thrust & transverse sphericity





Hadronic Event Shapes: $au_{\perp}(H_{\text{T2}}, n^{\text{jet}})$, $S_{\perp}(H_{\text{T2}}, n^{\text{jet}})$





Hadronic Event Shapes: $A(H_{T2}, n^{jet})$, $D(H_{T2}, n^{jet})$



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A multi-purpose detector system **Muon Detectors Tile Calorimeter** Liquid Argon Calorimeter Total weight : 7000 t **Overall length: 46 m Overall diameter: 23 m Toroid Magnets** Solenoid Magnet SCT Tracker Pixel Detector TRT Tracker Magnetic field: 2T solenoid + (varying) toroid field Peter Loch – Snowmass EF05 Slide 34

Detectors for Hadronic Final State Reconstruction

Calorimeters

Provides principal signals for e^{\pm}/τ^{\pm} and jet kinematics – and other measurements Full coverage within $|\eta| < 4.9$ with depth $\geq 10 \lambda_{int}$ Highly segmented for energy flow measurements (~188,000 cells)

High granularity in $\Delta \eta \times \Delta \varphi = 0.025 \times \pi/128$ (central EM)

Up to seven depth layers (*samplings*)

Inner detector

Provides charged particle tracks and vertices

Coverage $|\eta| < 2.5$

Jet energy calibration refinement

Provides vertex for jet origin correction/jet vertex association/jet vertex tagging (JVT) Flavor/fragmentation sensitive response measures – mitigation of jet flavor response dependencies

Particle flow

Replace charged response in calorimeter with kinematics from well-measured tracks Missing transverse momentum soft contributions

Tracks not used or associated with (hard) reconstructed particles and jets

Muon spectrometer

Reconstructed muons

Contribution to missing transverse momentum reconstruction

Track segments

Proxy for energy leakage behind a jet