

# Delving $\Upsilon$ and $\Psi$ leptonic decays with an extra gaugeboson

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#### Abstract

- We observed that lepton flavour universality can be examined through the ratio of leptonic  $\Upsilon$  =  $\Upsilon(ns)(n = 1, 2, 3)$  and  $\psi$  quarkonia decays.
- ▶ In BABAR, the neutral current transitions  $b\bar{b}$  →  $\tau \bar{\tau}$  reported that leptonic decay ratio  $R_{\Upsilon}(3s) =$  $Br(\Upsilon(3s) \to \tau \bar{\tau})/Br(\Upsilon(3s) \to \mu \bar{\mu})$  is showing an agreement with the SM at 1.8 $\sigma$  level, where the ratio  $R(J/\psi)$  expose discrepancies with resembling SM estimations.
- ► We perform an analysis by using a clarify model with an extra massive gauge boson W' that catenate substantially to left handed leptons of the third generation.

### Phenomenological Analysis

- $\blacktriangleright$   $\chi^2$  function is given by the sum of the squared pulls i.e  $\chi^2 = \sum_i pull_i^2$ .
- $\blacktriangleright pull_i = (\mathcal{O}^i_{exp} \mathcal{O}^i_{th}) / \sqrt{\sigma^{i2}_{exp} \sigma_{th}^{i2}}$ where  $\sigma^{i}_{exp,th}$  corresponds to the experimental and theoretical error.
- It is possible to obtain several models by turning on some of the couplings, while the remaining ones are set equal to zero. ► In order to adjust the experimental anomalies, any model must contain a charm-bottom interaction term in the quark sector and the corresponding  $\tau - \nu_{\tau}$  interaction term in the lepton sector, i.e it is necessary to have at least two nonzero couplings in the Lagrangian.

#### Introduction

- A general W' boson exchange leads to additional tree level effective interactions to the  $b \to c \tau \bar{\nu}_{\tau}$ transition.
- ► The total low energy effective Lagrangian is

 $\mathcal{L}_{eff}(b \rightarrow c\tau \bar{\nu}_{\tau})_{SM+W'} = \frac{4G_F}{\sqrt{2}} V_{cb} \left[ (1 + C_V^{LL})(\bar{c}\gamma_{\mu}P_Lb)(\bar{\tau}\gamma_{\mu}P_L\nu_{\tau}) + C_V^{RL}(\bar{c}\gamma_{\mu}P_Rb)(\bar{\tau}\gamma_{\mu}P_L\nu_{\tau}) + C_V^{LR}(\bar{c}\gamma_{\mu}P_R\nu_{\tau}) + C_V^{RR}(\bar{c}\gamma_{\mu}P_R\nu_{\tau}) + C_V^{RR}(\bar{c}\gamma_{\mu}P_R\nu_{\tau}$ (1)

 $\triangleright$   $C_V^{LL}$ ,  $C_V^{LR}$ ,  $C_V^{RL}$  and  $C_V^{RR}$  are the wilson coefficients, depends on the choices of chiral charges, associated with the NP operators.

$$C_{V}^{LL} = \frac{\sqrt{2}}{4G_{F}V_{cb}} \frac{\epsilon_{cb}^{L}\epsilon_{\tau\nu}^{L}}{M_{W'}^{2}} \qquad C_{V}^{RR} = \frac{\sqrt{2}}{4G_{F}V_{cb}} \frac{\epsilon_{cb}^{R}\epsilon_{\tau\nu}^{R}}{M_{W'}^{2}} \qquad (2)$$

$$C_{V}^{LR} = \frac{\sqrt{2}}{4G_{F}V_{cb}} \frac{\epsilon_{cb}^{L}\epsilon_{\tau\nu}^{R}}{M_{W'}^{2}} \qquad C_{V}^{RL} = \frac{\sqrt{2}}{4G_{F}V_{cb}} \frac{\epsilon_{cb}^{R}\epsilon_{\tau\nu}^{L}}{M_{W'}^{2}}$$

 $\blacktriangleright$  We will pay attention to the wilson coefficients and the W' boson mass which can provide an explanation for the  $b \rightarrow c \tau \bar{\nu}$  anomalies.

Table							
	Parameters on	R(D)	R(D*)	$R(J/\psi)$	$P_{\tau}(D^*)$	$F_L(D^*)$	$\chi^2$ min
	$(\epsilon_{cb}^{L}, \epsilon_{\tau\nu}^{L})$	-0.045	0.031	1.53	0.22	1.47	5.48
	$(\epsilon_{cb}^L, \epsilon_{\tau\nu}^R)$	-0.046	0.027	1.53	-0.013	1.46	5.43
	$(\epsilon_{cb}^{R}, \epsilon_{\tau\nu}^{L})$	2.57	0.46	1.56	0.19	1.52	12.27
	$(\epsilon_{cb}^R, \epsilon_{\tau\nu}^R)$	-0.047	0.027	1.53	-0.013	1.46	5.43

# LL Scenario ( $C_V^{LL} \neq 0$ )

► W' that couples only to LH quark and LH lepton currents.



**Table 1:** Turning on the parameters we analysis the  $\chi^2$  fitting, for a gauge boson mass  $M_W'=1$  Tev

- ► In order to check whether it is possible to adjust the deviations of the standard model predictions in these models, we carried out a  $\chi^2$  analysis with the five experimental observables mentioned above.
- All the models have an acceptable value for  $\chi^2_m in/dof \sim 1$  except the model with RH coupling to quarks and LH coupling to leptons.

# **RR Scenario** ( $C_V^{RR} \neq 0$ )

► W' that couples only to RH quark and RH lepton currents involving a RH sterile neutrino.



## LR and RL Scenarios

where the quark and lepton currents with different chiralities projection couple to the W' boson i.e  $C_V^{LR} \neq 0$  and  $C_V^{RL} \neq 0$ .



**Fig. 3:** Allowed parameter space in the  $(\epsilon_{cb}^{R}, \epsilon_{\tau\nu}^{L})$  plane

**Fig. 1:** Allowed parameter space in the  $(\epsilon_{cb}^{L}, \epsilon_{\tau\nu}^{L})$ plane for  $M_{W'} = 1 Tev$ .

► The purple region is obtained by considering only  $R(D^*)$  from HFLAV-2019 average, while the green one is showing the constraints for all the observables of  $b \rightarrow c \tau \bar{\nu}$  taking into account.

**Fig. 2:** Allowed parameter space in the  $(\epsilon_{cb}^{R}, \epsilon_{\tau\nu}^{R})$ plane for  $M_{W'} = 1 Tev$ .

► The purple region is obtained by considering only  $R(D^*)$  from HFLAV-2019 average, while the green one is showing the constraints for all the observables of  $b \rightarrow c \tau \bar{\nu}$  taking into account. The red red region is obtained from GMR analy-SIS.

Conclusion

- The anomalies  $R(D^*)$  related to the charged current transition  $b \to c \tau \bar{\nu}$  with in a W' boson scenario.
- Effective Lagrangian given in terms of the flavour dependent couplings  $\epsilon_{cb}^{L,R}$ ,  $\epsilon_{\tau\nu}^{L,R}$  yields to a tree level effective contribution generated by a W' boson.
- $\blacktriangleright$  we performed a  $\chi^2$  analysis by considering the cases of two non zero couplings which later show the best candidate to adjust the experimental charged current B anomalies.
- Taking into account two different datasets  $R(D^*)$  and all  $b \to c \tau \bar{\nu}$  observables we determined the regions for allowed parameter space which is favoured by those observables and our results showed that it would be interesting to study a particular NP model.





**Fig. 4:** Allowed parameter space in the  $(\epsilon_{cb}^L, \epsilon_{\tau\nu}^R)$  plane for  $M_{W'} = 1 Tev$ .