

# Time evolution of lepton number of Majorana neutrinos: the Schrödinger picture

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## Motivation

In our past work we studied the time evolution of lepton number (L) of Majorana neutrinos by constructing the Heisenberg operator for it. The initial state is a neutrino state with a definite lepton number but is not a mass eigenstate. The expectation value of L oscillates. we study the same phenomena by adopting the Schrödinger representation.

### Questions:

- 1) Time evolution of the state will be clarified
- 2) The understanding of the difference of two vacua one which the operator with a definite lepton number operates and the other on which the massive Majorana operator acts on.

## Goals

Starting with an initial state with a definite lepton number and let it evolve with respect to time and construct the state at arbitrary time. With the state, one can compute the expectation value of lepton number.

$$\langle \Psi(0) | L_H(t) | \Psi(0) \rangle = \langle \Psi(t) | L_S | \Psi(t) \rangle$$

One can also check the previous result in Heisenberg rep..

## Method

- (1) Relation bet. neutrino and anti-neutrinos creation and annihilation operators, (a(p), b(p)) and massive Majorana particles with different helicities (h) (a<sub>M</sub>(p, h=+1), a<sub>M</sub>(p, h=-1)) is obtained in the form Bogoliubov transformation.
- (2) Rewrite the lepton number operator which is originally written with a and b in terms of operators of massive Majorana neutrinos.

$$L = \int'_{p \in A} \frac{d^3 p}{(2\pi)^3 2|p|} \{ (a^\dagger(p)a(p) - b^\dagger(p)b(p)) + (a^\dagger(-p)a(-p) - b^\dagger(-p)b(-p)) \}$$

$$L = \int'_{p \in A} \frac{d^3 p}{(2\pi)^3 2E_p} p \{ a_M^\dagger(p, -)a_M(p, -) + a_M^\dagger(-p, -)a_M(-p, -) - a_M^\dagger(p, +)a_M(p, +) - a_M^\dagger(-p, +)a_M(-p, +) \}$$

$$+ \int'_{p \in A} \frac{d^3 p}{(2\pi)^3 2E_p} im \{ a_M^\dagger(p, -)a_M^\dagger(-p, -) - a_M(-p, -)a_M(p, -) - a_M^\dagger(p, +)a_M^\dagger(-p, +) + a_M(-p, +)a_M(p, +) \}$$

$$|0\rangle = \prod_{p \in A} (\cos^2 \phi_p |0_M\rangle_p - \sin^2 \phi_p |4_M\rangle_p + i \sin \phi_p \cos \phi_p |2_M\rangle_p)$$

The vacuum with 0 lepton number is expressed by the state in which one and two pairs of the Majorana neutrinos are condensed.

Bogoliubov transformation.

$$T_p(\phi_p) = \text{Exp}[i(g_+(p) + g_-(p))\phi_p]$$

$$g_\lambda(p) = \alpha_M(p, \lambda)\alpha_M(-p, \lambda) + \alpha_M^\dagger(-p, \lambda)\alpha_M^\dagger(p, \lambda)$$

$$\alpha_M(p, \lambda) = \frac{a_M(p, \lambda)}{\sqrt{2E_p(2\pi)^3\delta^{(3)}(0)}} \quad (\text{Dimensionless operators})$$

$$B_{M\lambda}^\dagger(p) = \alpha_M^\dagger(-p, \lambda)\alpha_M^\dagger(p, \lambda)$$

$$\cos \phi_p = \frac{\sqrt{1+v_p}}{\sqrt{2}}$$

$$\sin \phi_p = \frac{\sqrt{1-v_p}}{\sqrt{2}}$$

## Results

The state at time t is derived.

$$|\psi(t)\rangle_q = N_q \sqrt{\cos 2\phi_q} e^{-iE_q t} \{ \cos \phi_q + i \sin \phi_q e^{-i2E_q t} B_{M+}^\dagger(q) \} a_M^\dagger(q, -) |0_M\rangle_q$$

$$\hat{L} = \int'_{p \in A} \frac{V d^3 p}{(2\pi)^3} \hat{l}(p)$$

The Fock state with a momentum q is a superposition of a massive Majorana state and a state with a Majorana particle accommodated with a pair of Majorana particles. The mixing of the latter state is maximum for the velocity zero limit.

$$\langle \psi(t) | \hat{l}_q | \psi(t) \rangle = v^2 + (1 - v^2) \cos \left[ \frac{2mt}{\sqrt{1-v^2}} \right]$$

Results of the time evolution of L is the same as that obtained by Heisenberg picture.

