

# *Extending subtraction schemes to $N^3LO$*

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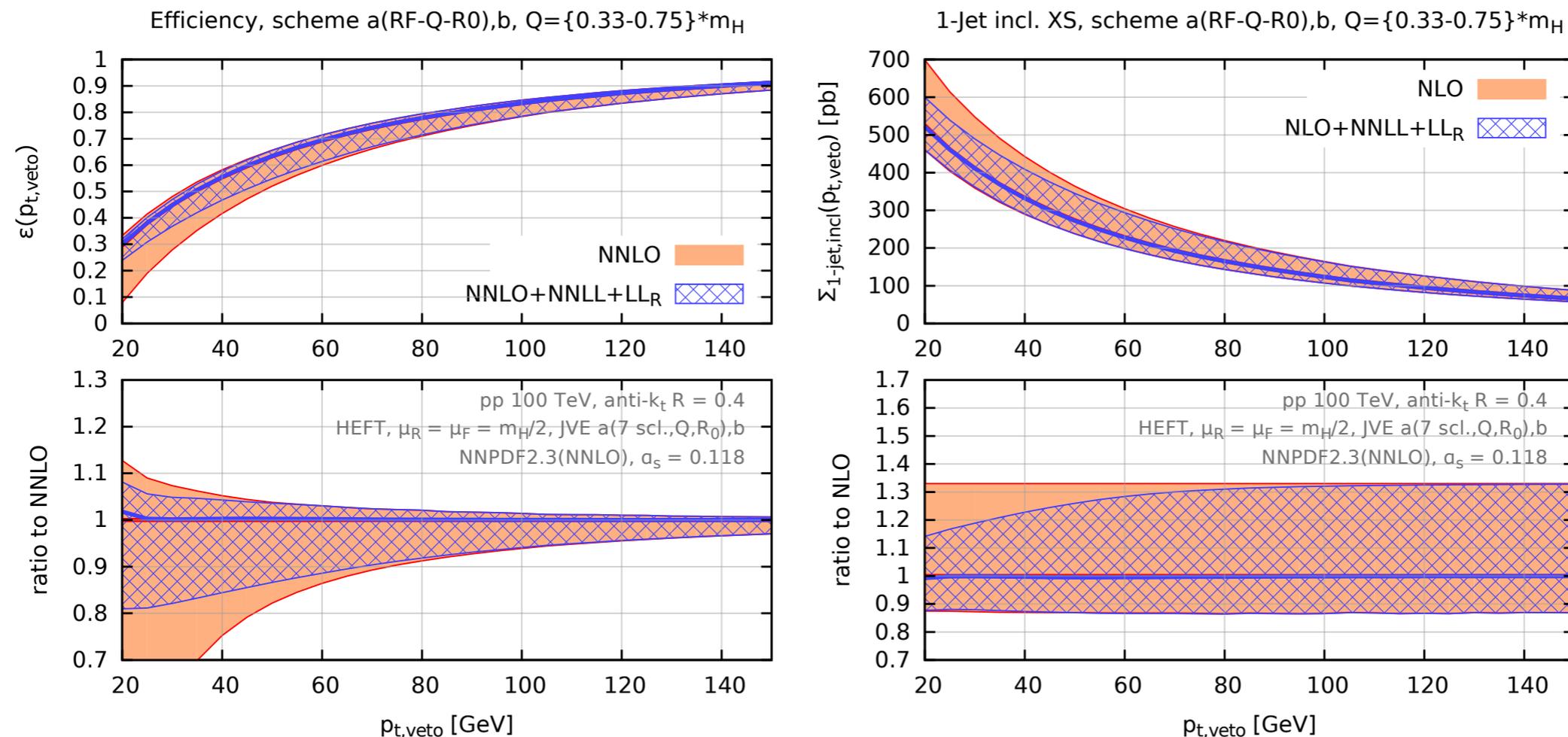
Snowmass Community Planning Meeting, Oct. 6th 2020



# Why N<sup>3</sup>LO?

LHC: fixed-order calculation essential for precision@colliders\*

- a clean, systematically improvable [*see later*] framework for first principle calculations → solid, unambiguous predictions → new physics from small deviations
- applicable whenever we are far from soft/collinear regions → likely to play a crucial role for high energy colliders (*less issues with acceptances etc*)



Higgs jet veto efficiency at 100 TeV  
[FCC report (2016)]

\* NB: this does not mean that f.o. is the only important ingredient

# Why N<sup>3</sup>LO?

## 1) PRECISION

$$\alpha_s \sim 0.1, \quad \alpha_s C_A \sim 0.3 \quad \longrightarrow \quad \text{few percent} \Leftrightarrow \text{N}^{(3)}\text{LO}$$

“Few percent”

- Experimentally achievable, at least for standard candles (V, H, top...)
- $\Lambda_{\text{NP}}/Q^2 \sim 1\% \rightarrow$  sensitive to new physics at the (multi)-TeV scale, may not be accessible directly
- $\alpha \sim 0.01 \rightarrow$  study EW theory at the quantum level, without worrying about QCD contaminations

# Why N<sup>3</sup>LO?

## 1) PRECISION: SOME CAVEATS

[Mistlberger, QCD@LHC2016]

$$\sigma = 48.58 \text{ pb}^{+2.22 \text{ pb} (+4.56\%)}_{-3.27 \text{ pb} (-6.72\%)} \text{ (theory)} \pm 1.56 \text{ pb} (3.20\%) \text{ (PDF} + \alpha_s).$$

48.58 pb =	16.00 pb	(+32.9%)	(LO, rEFT)
	+ 20.84 pb	(+42.9%)	(NLO, rEFT)
	- 2.05 pb	(-4.2%)	(( <i>t, b, c</i> ), exact NLO)
	+ 9.56 pb	(+19.7%)	(NNLO, rEFT)
	+ 0.34 pb	(+0.7%)	(NNLO, 1/ <i>m<sub>t</sub></i> )
	+ 2.40 pb	(+4.9%)	(EW, QCD-EW)
	+ 1.49 pb	(+3.1%)	(N <sup>3</sup> LO, rEFT)

At the few-percent level: many non-trivial subtle effects

# Why N<sup>3</sup>LO?

## 1) PRECISION: SOME CAVEATS

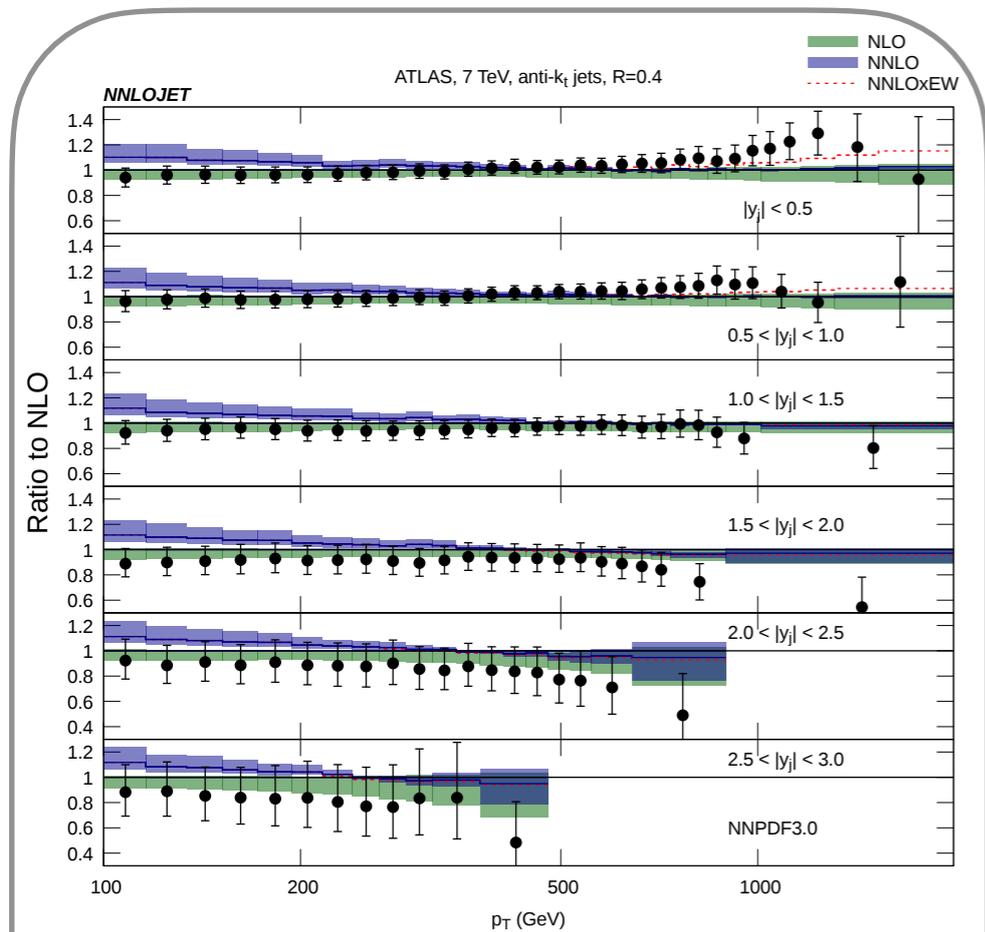
$$d\sigma = \int dx_1 dx_2 f(x_1) f(x_2) d\sigma_{\text{part}}(x_1, x_2) F_J \left[ 1 + \mathcal{O} \left( \frac{\Lambda_{\text{QCD}}^n}{Q^n} \right) \right]$$

In many cases, few-percent could be a theoretical limit (given our current understanding of QCD)

- $Q \sim \text{EW scale}$ , linear power corrections (e.g. jets)  $\rightarrow$  percent
- At least a partial understanding would be as important as developing higher order calculations (already now, e.g. *top mass measurements*)

# Why N<sup>3</sup>LO?

## 2) MULTI-PARTON DYNAMICS

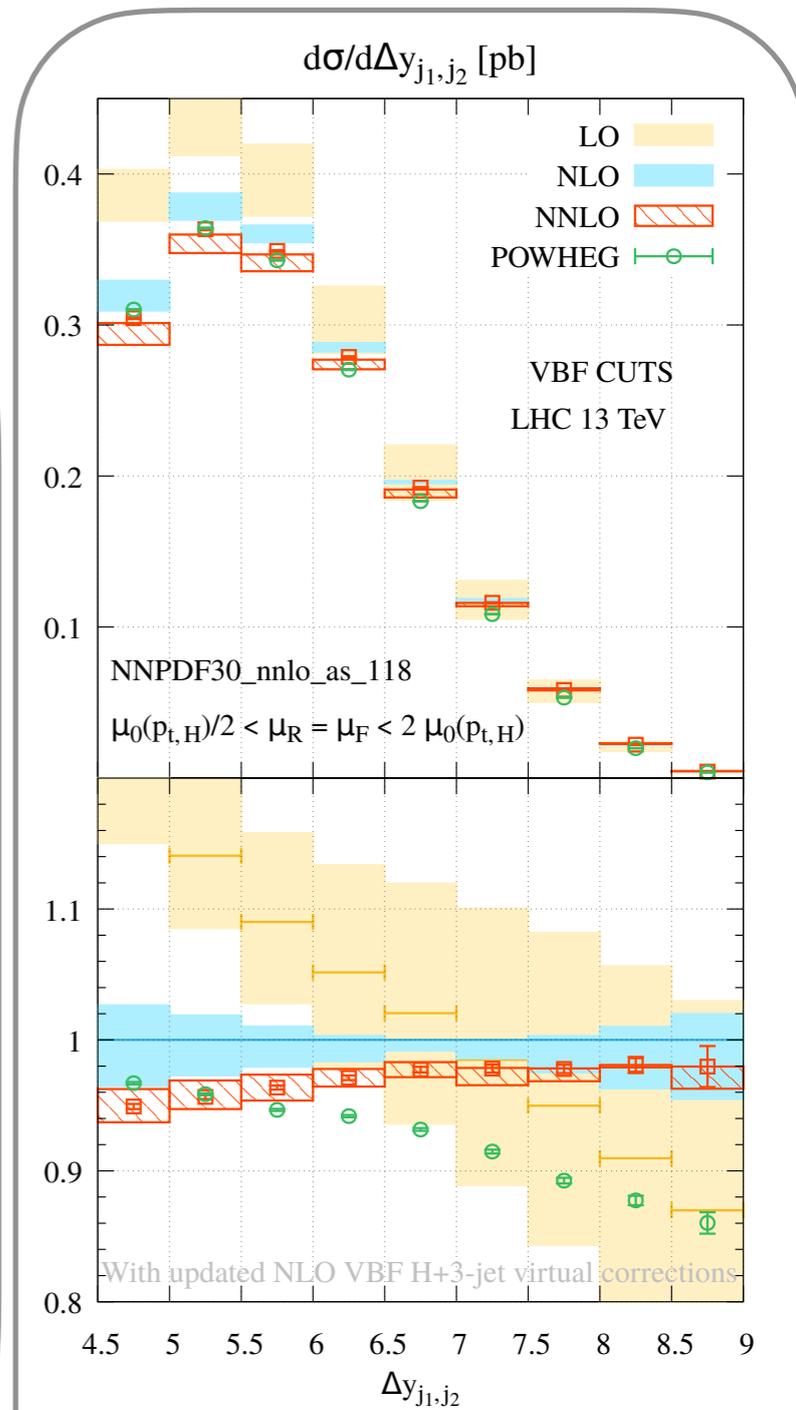
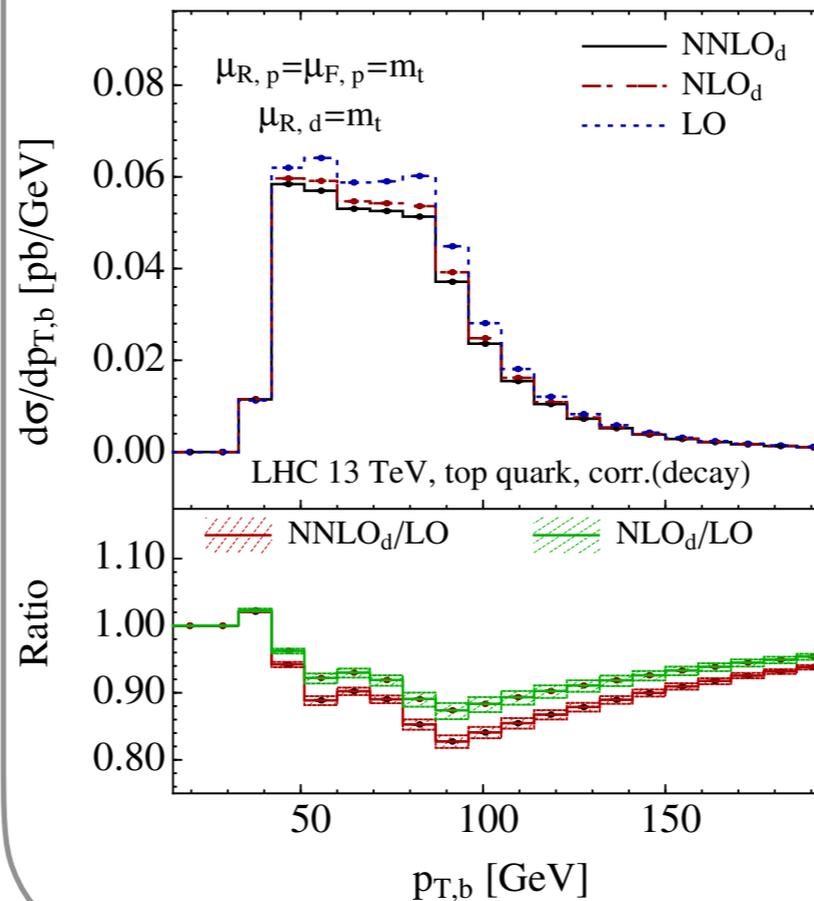


### Dijet

[Currie et al., NNLO<sub>JET</sub> (2016-...);  
Czakon, van Hameren, Mitov,  
Poncelet (2020)]

### Single-top

[Berger, Gao, Yuan, Zhu (2016-...)]



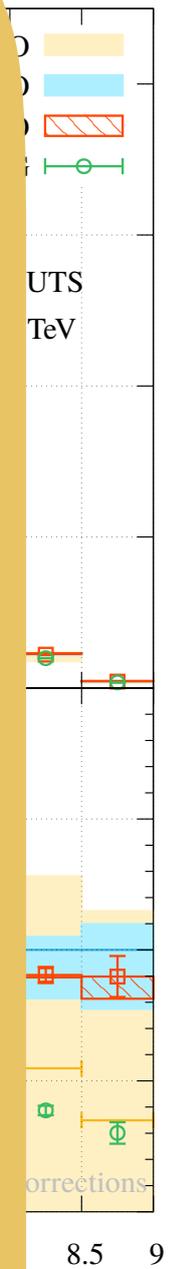
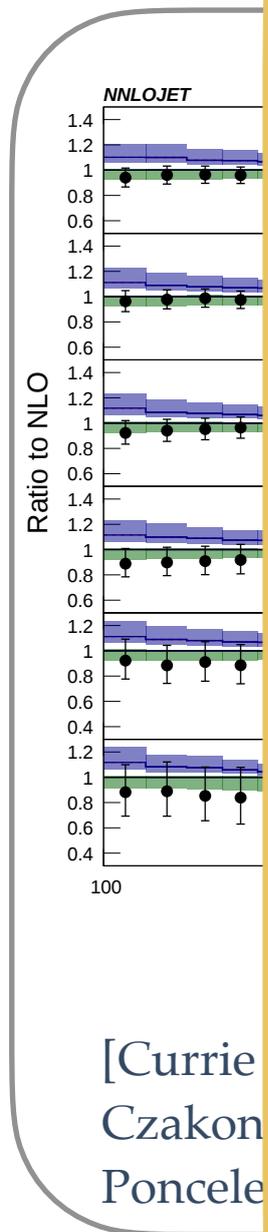
### VBF

[Cacciari et al. (2015);  
Cruz-Martinez, Gehrmann,  
Glover, Huss (2-18)]

New handles on non-trivial multi-particle dynamics, complementary to standard ones

2)

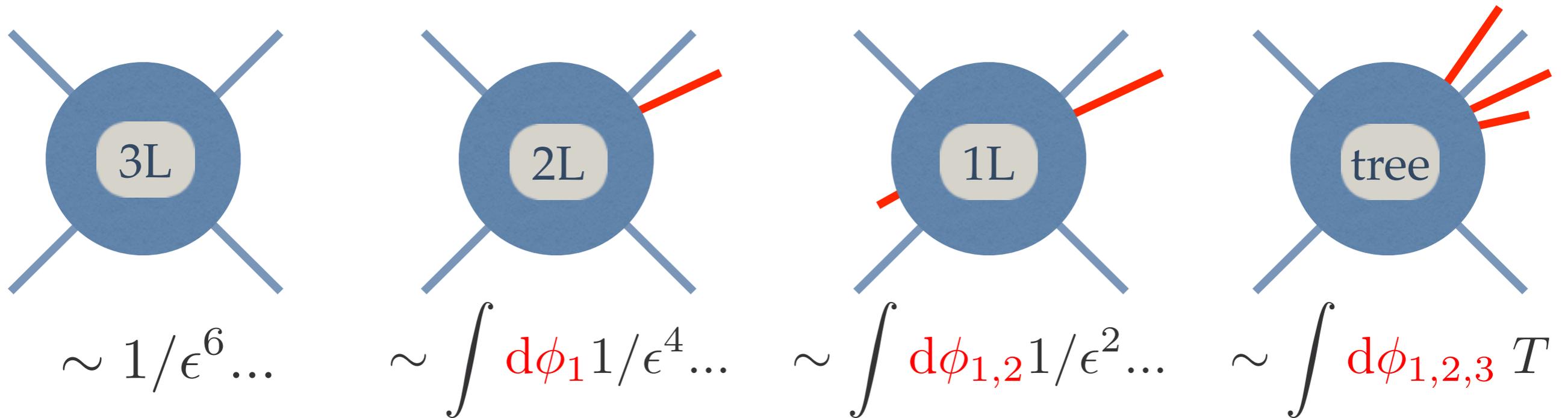
- Good phenomenological reasons to explore  $N^3\text{LO}$ , at least for standard candles
- If result not (purely) result-driven ( $\rightarrow$  *brute force*): interesting opportunity to explore soft/collinear structure of gauge theories



New  
dynam

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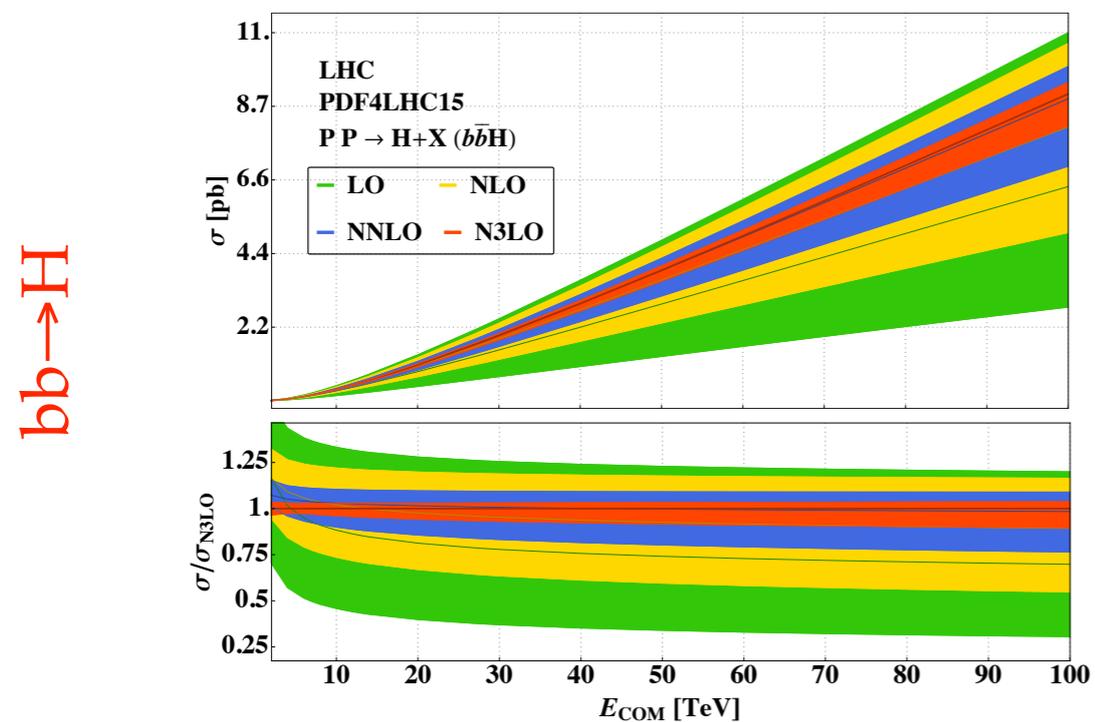
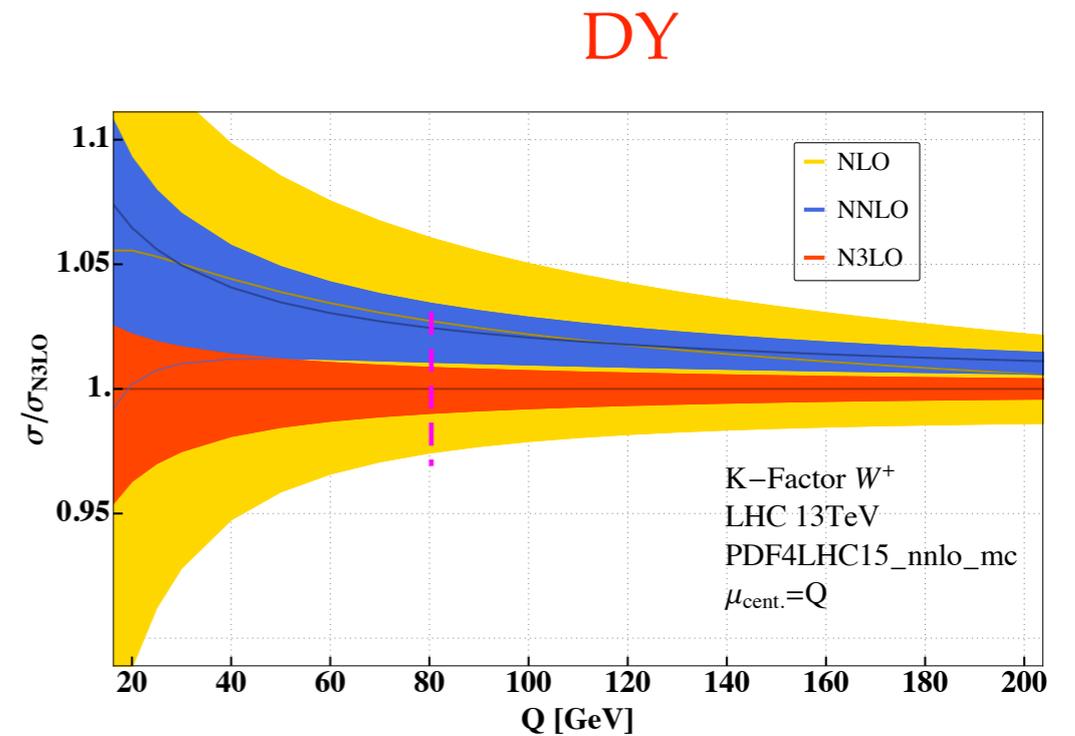
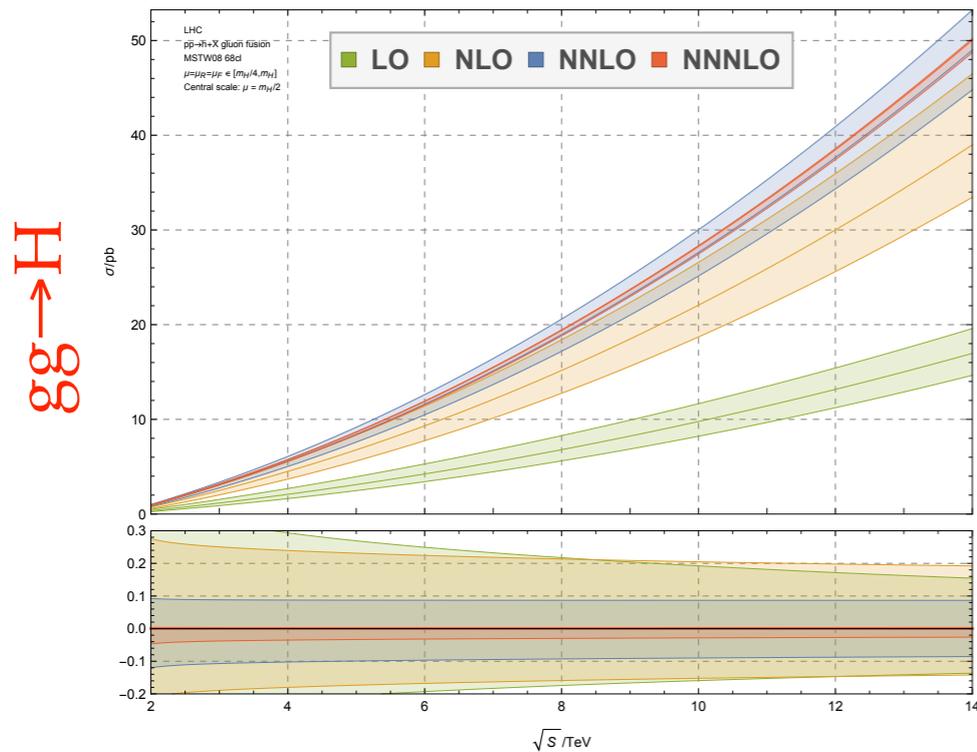
# N<sup>3</sup>LO: how to get there



## Two main issues

- Multi-loop amplitudes (*especially in unresolved configurations*) → see Harald's talk
- How to properly extract (and regulate) IR from loop integrals. Highly non-trivial for fully exclusive calculations → **focus here**. An optimal solution at NNLO is not yet available (*my opinion*), we can get results thanks to large computing resources

# N<sup>3</sup>LO right now: fully inclusive



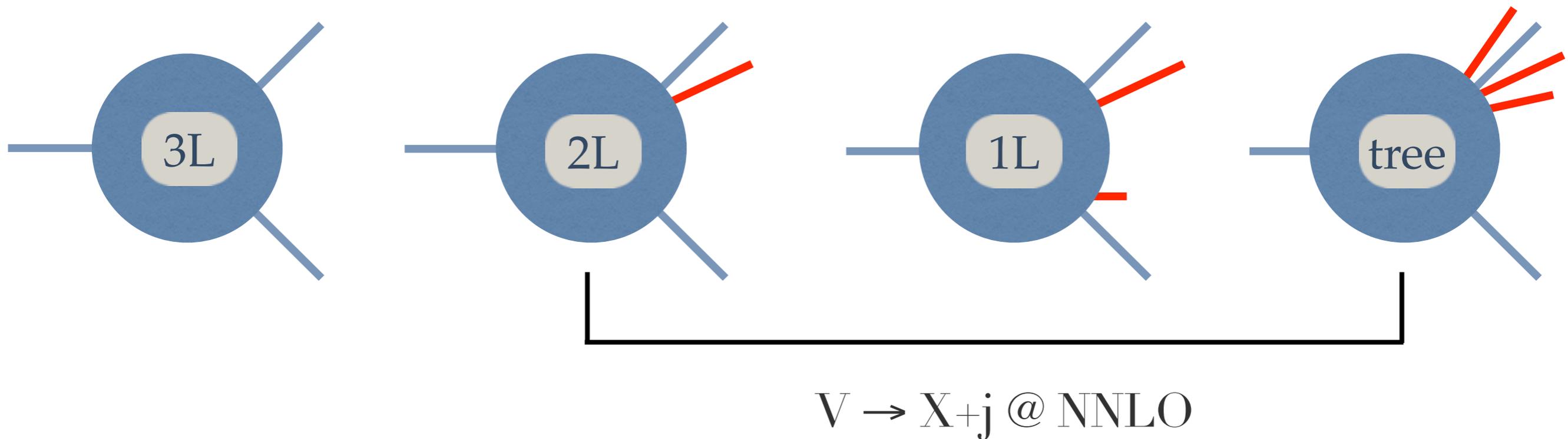
Fully inclusive N<sup>3</sup>LO  
corrections for 2 → 1 color  
singlet production is possible

[Dulat, Duhr, Mistlberger +  
Anastasiou et al (2015-...)]

[VBF: Dreyer, Karlberg (2018)]

# Colour singlet: differential results

Consider the decay of a colour-singlet state  $V$

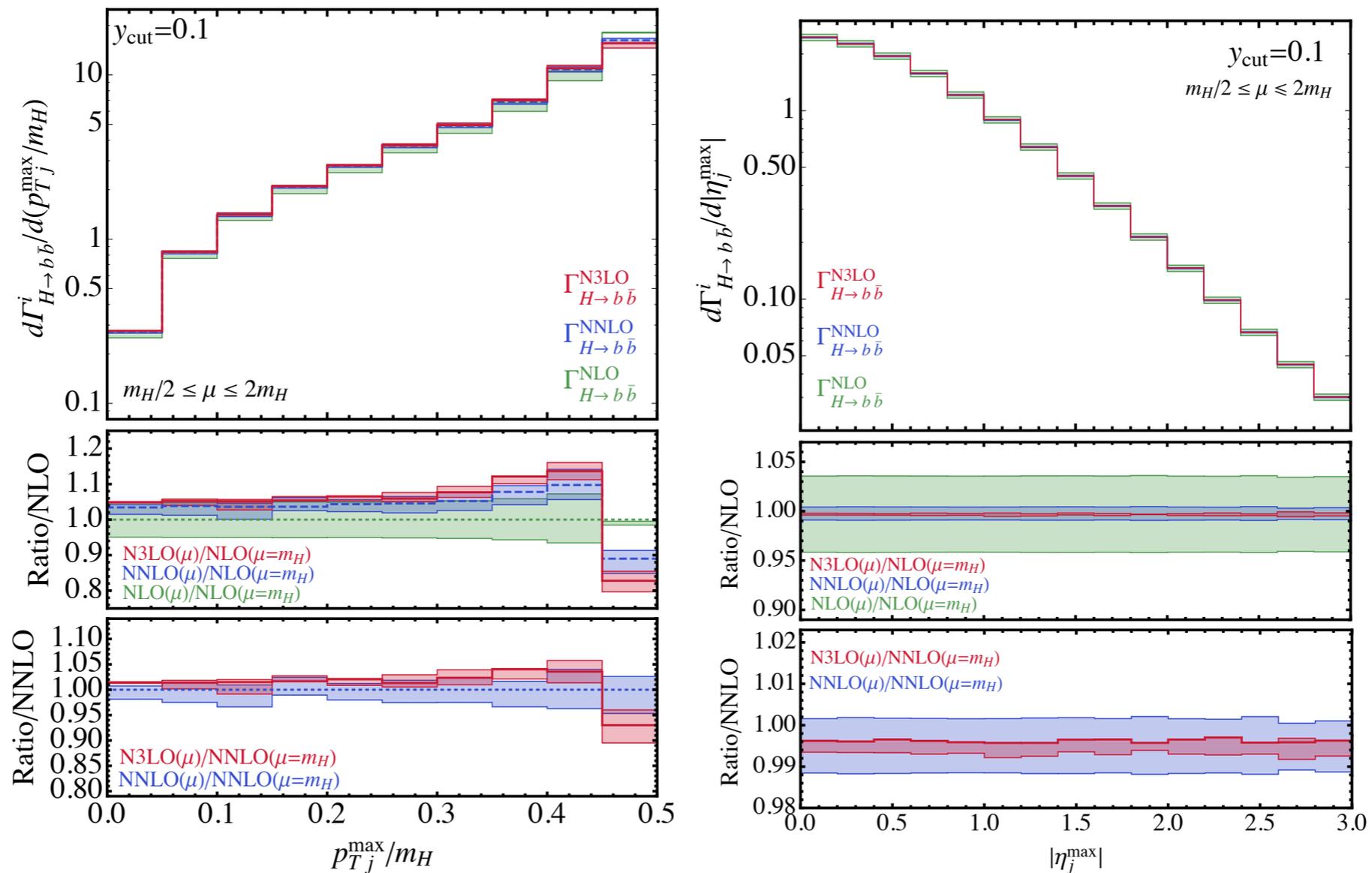


- Most of the information contained in  $V \rightarrow X_{+j} @ \text{NNLO}$
- Born-like configuration: no kinematical dependence
- As a consequence:  $V \rightarrow X_{+j} @ \text{NNLO}$  & total rate @  $\text{N}^3\text{LO}$  enough for differential rate at  $\text{N}^3\text{LO}$

# Colour singlet: differential results

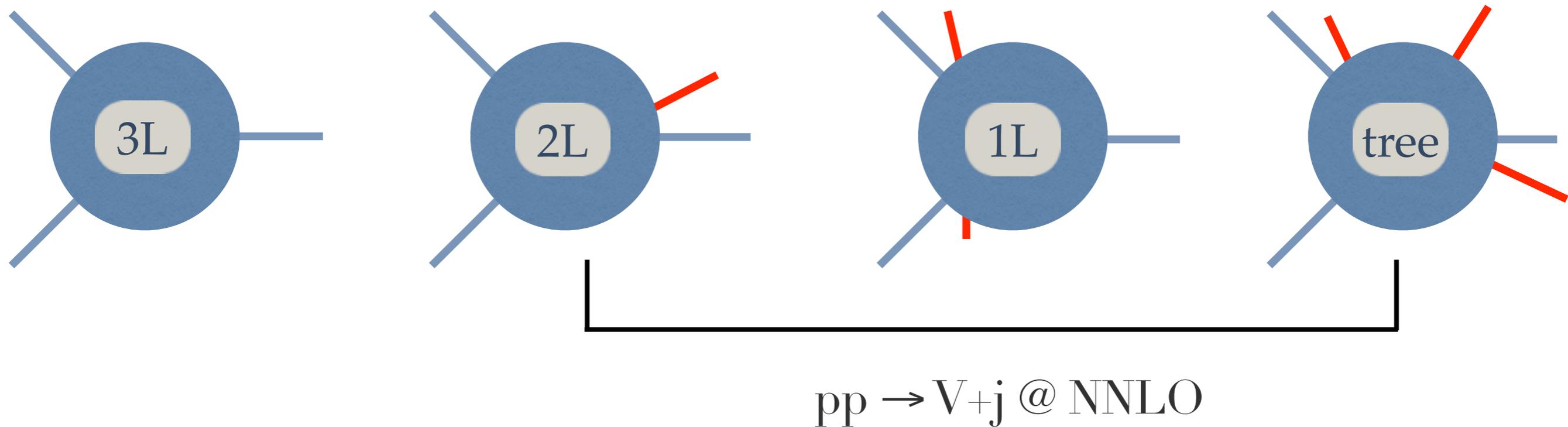
$H \rightarrow b \bar{b}$  @ N<sup>3</sup>LO known

[Mondini, Schiavi, Williams (2019)]



# Colour singlet: differential results

Consider the ~~decay~~ production of a colour-singlet state  $V$



- Most of the information contained in  $V \rightarrow X+j @ NNLO$
- Born-like configuration: ~~no kinematical~~ **non-trivial rapidity** dependence
- As a consequence:  $pp \rightarrow V+j @ NNLO$  &  $(0-p_t)$  **rapidity dependence** @N<sup>3</sup>LO enough for differential rate at N<sup>3</sup>LO

Cons

- In principle: we can compute fully differential N<sup>3</sup>LO predictions for colour singlet production with current technology
- In practice: need to “loose a jet” in  $pp \rightarrow V+j @ NNLO$ , computationally intensive

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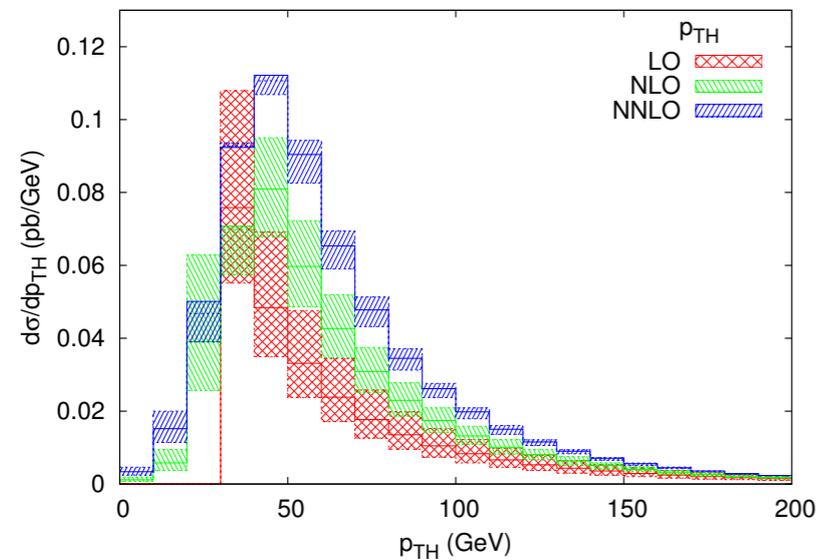
*The bad part of the story: this trick is not (immediately) generalisable to arbitrary processes*

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# Colour singlet: differential results

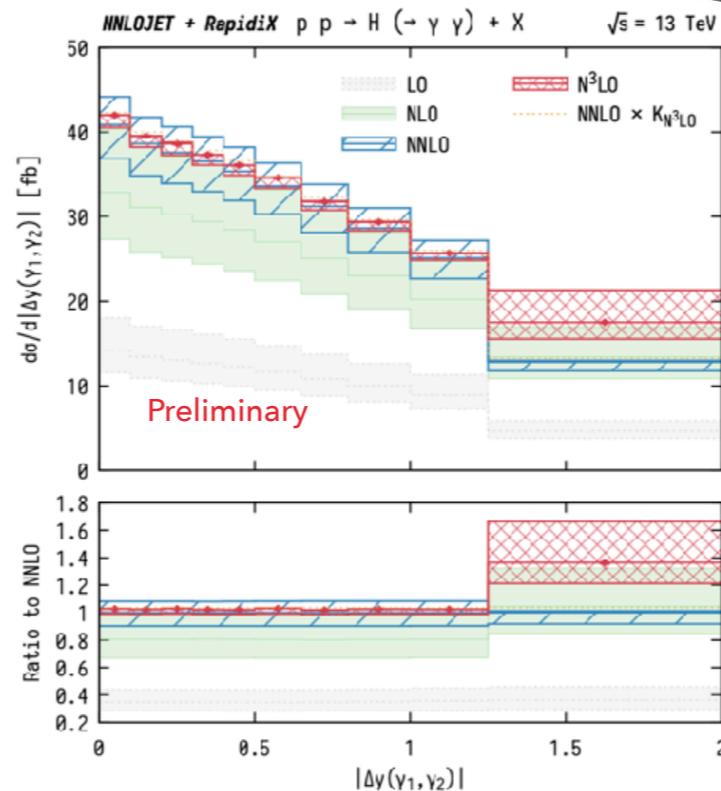
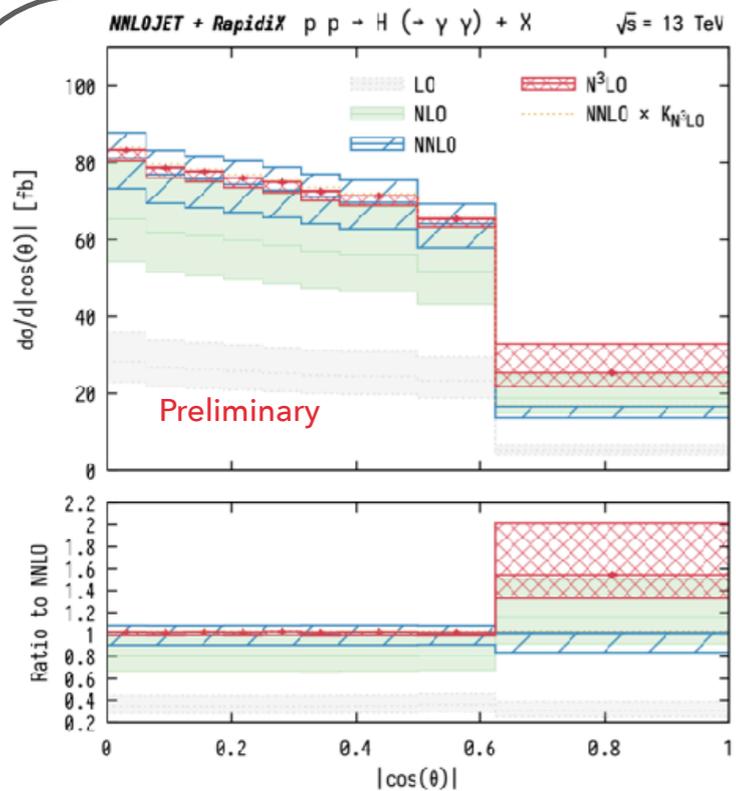
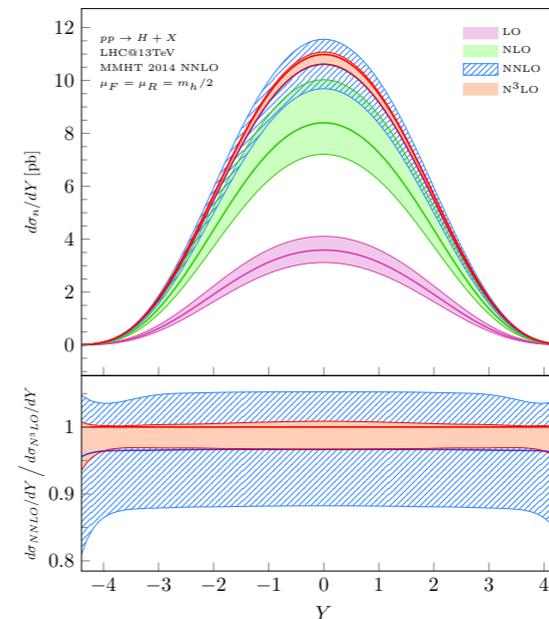
H+J@NNLO

[Chen, Cruz-Martinez, Gehrmann, Glover, Jaquier (2016)]



$d\sigma/dy$  @N<sup>3</sup>LO

[Dulat, Mistlberger, Pelloni (2018)]



[Chen, Dulat, Gehrmann, Glover, Huss, **BM**, Pelloni, to appear]

[Mistlberger, talk at QCD@LHC-X]

(See also Cieri et al. (2018) for earlier work assuming trivial rapidity dependence)

# Colour singlet: beyond the Higgs

- For colour-singlet production: only require rapidity-dependence at zero transverse momentum
- Factorisation  $\rightarrow$  universality. “**Beam functions**”
- Over the past year or so: all ingredients computed
  - $p_t$  beam function [Luo, Zang, Zhu, Zhu (2020); Ebert, Mistlberger, Vita (2020)]
  - N-jettiness beam function\*: [Behring, Melnikov, Rietkerk, Tancredi, Wever (2019); Ebert, Mistlberger, Vita (2020)]
- Provided that we have relevant loop amplitudes and **enough computing power** this would allow for a brute-force calculation of N<sup>3</sup>LO QCD corrections for any colour-singlet state

*\*Not enough even for colour-singlet, but has the potential to be extended to arbitrary processes*

# Beyond colour-singlet

- Beyond colour-singlet, new frameworks are required → better understanding of soft/collinear emission (and beyond)
- *Slicing techniques*
  - similar to what discussed so far. Introduce resolution variable that separates Born-like vs +jet
  - -: typically, this introduce some non-locality in IR cancellations → large numerical cancellation, huge computing power required
  - +: giving up locality makes everything much simpler → may be a shortcut to the result, if enough computing power
  - +: several ingredients already known / under computation *N-jettiness beam function, towards N-jettiness soft function* [Baranowski (2020)]

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  - +: several ingredients already known / under computation *N-jettiness beam function, towards N-jettiness soft function [Baranowski (2020)]*
  - **The way forward** (*my opinion*): sub-leading power corrections. Would allow for much better numerics / very interesting from a conceptual point of view [Moult et al. (2017-...); Boughezal, Isgro, Petriello (2018); Beneke et al (2018-...); Laenen et al. (2018-...); Neubert et al (2020-...)]

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- *Subtraction techniques*
  - allow for local cancellation of IR singularities via universal subtraction terms → interesting *per se*
  - +: theoretically appealing, much higher efficiency (slicing eventually abandoned @ NLO)
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  - +: theoretically appealing, much higher efficiency (slicing eventually abandoned @ NLO)
  - -: the price of locality: much more challenging to devise a generic framework
  - +: in the recent past, several of the universal ingredients computed [Catani, Colferai, Torrini (2019); Del Duca, Duhr, Haindl, Lazopoulos, Michel (2019-20); Bern, Dixon, Kosower (2004); Badger, Glover (2004); Duhr, Gehrmann, Jaquier (2014); Duhr, Gehrmann (2014); Li, Zhu (2013); Dixon, Hermann, Kai, Zhu (2019); Catani, de Florian, Rodrigo (2019); Badger, Buciuni, Peraro (2015); **Zhu (2020)**]

*(Personal opinion: regulating soft/collinear emission would proceed faster than for NNLO. Big conceptual gap was between NLO and NNLO)*

# N<sup>3</sup>LO: the “real” challenges?

- Thanks to the recent progress in NNLO: good position to start thinking at N<sup>3</sup>LO generalisations
- However, there may be surprises...
  - Top pair production:  $\sigma^{\text{N}^3\text{LO}} \neq \text{“R} + \text{V”}$ .
    - Non-trivial threshold effects, starting at  $\mathcal{O}(\alpha_s^3)$  [Beneke, Ruiz-Femenia (2016), see also Melnikov, Vainshtein, Voloshin (2014)]
    - Not captured by “standard” perturbation theory
  - Factorization breaking effects for jet production
    - Non-trivial absorptive part of loop integrals starts playing a role at N<sup>3</sup>LO. Could spoil universality of collinear singularity in processes with non-trivial color structure (e.g. di-jet...) [Forshaw, Siemour et al (2006-...); Catani, de Florian, Rodrigo (2011)]
    - Analysis confirmed in [Dixon, Hermann, Kai, Zhu (2019)]

*This sort of issues are likely to be the real challenges for N<sup>3</sup>LO computations*

# Conclusion

- N<sup>3</sup>LO likely to be relevant for current and future precision programs at the energy frontier
- There is something we can do already now, but not immediately generalisable (either conceptually or practically)
- For “standard” soft/collinear regularisation, the path is clear → (massively) more complicated than NNLO, but at least in principle we know where to start
- In this talk, I focused on “standard” techniques. Recently: old ideas (LT duality...) powered by new insight and better computing → this may turn out to be a good way forward
- There may however be unexpected effects, and this may require a better understanding of the structure of gauge theories

*Interesting times ahead!*

*Thank you very much  
for your attention!*