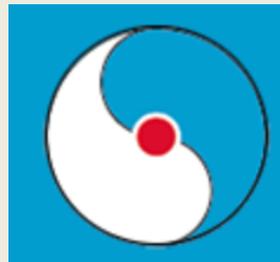


# The Standard Model effective field theory at low energies

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# SM-EFT at low energy

- **Tautology:** *the saying of the same thing twice in different words*
- By construction: SM-EFT approximates some UV-complete theory at a high-energy scale  $\sim \Lambda$
- **All experiments are low-energy experiments.**
- Of course: SM-EFT might break down @LHC if  $\Lambda$  close to EW scale
- That being said: in general a split is made



**Large Hadron Collider  
and beyond**

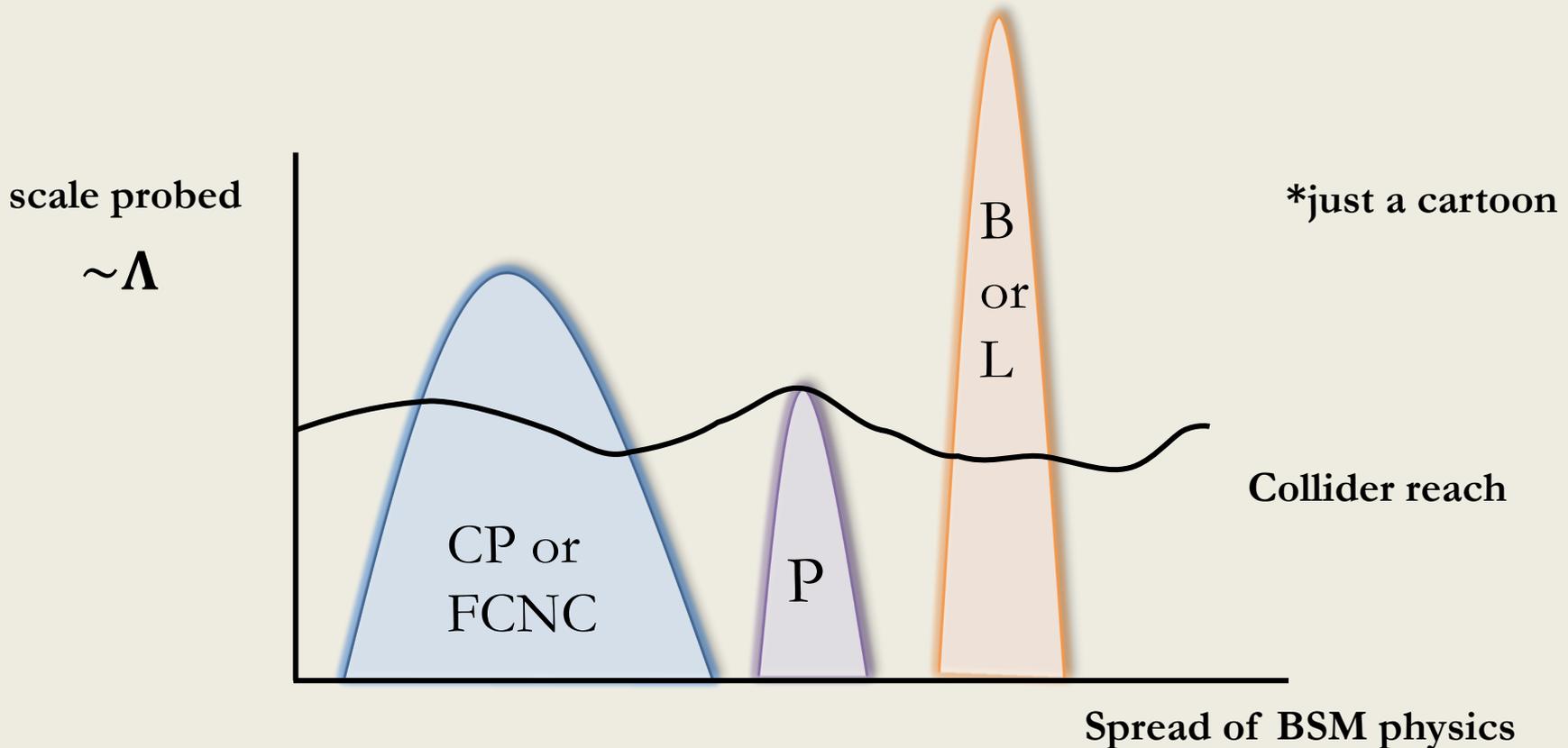


**Everything else at lower  
energies (might include LEP)**

Split is understandable from community point of view but of course artificial

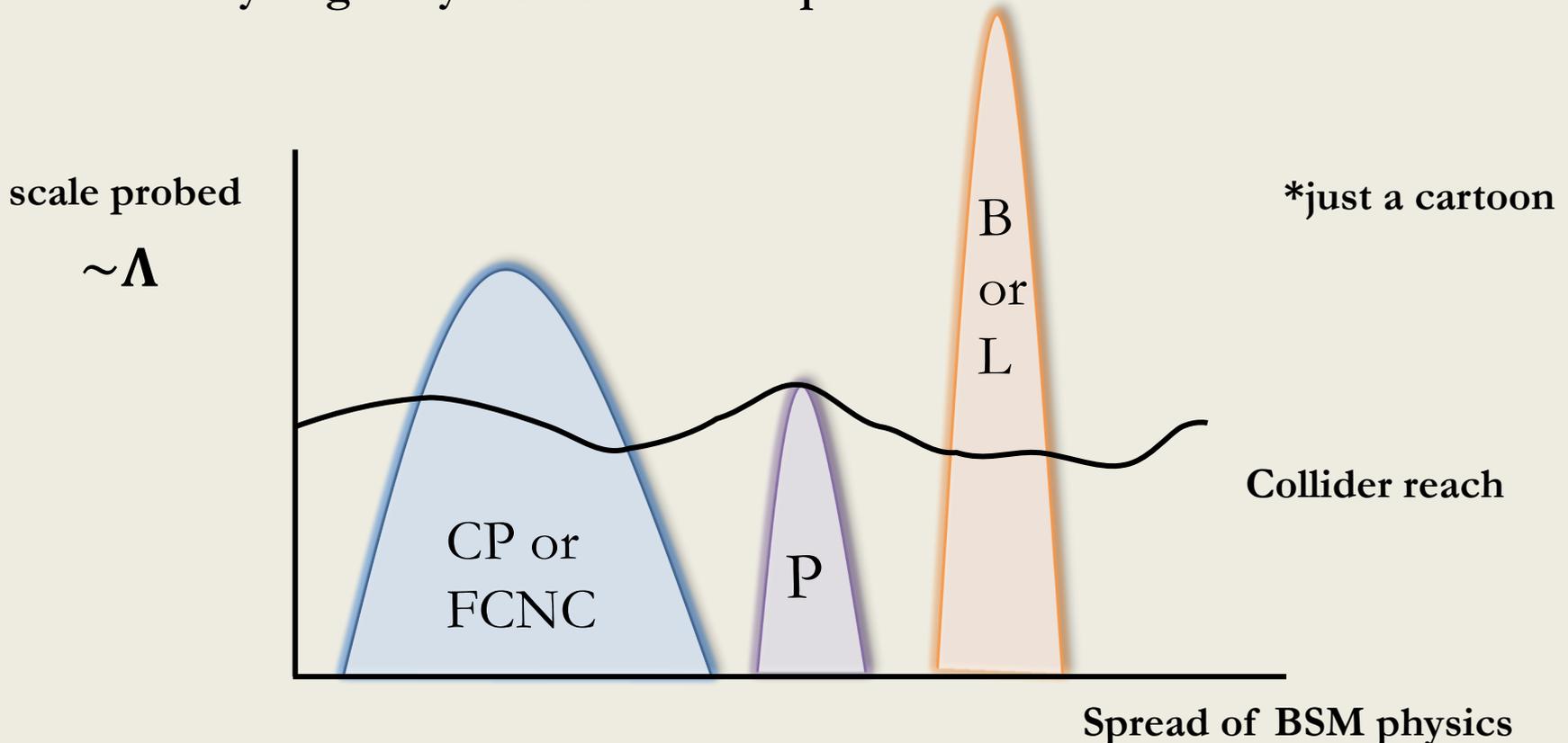
# Large number of low-energy SM tests

- Proton decay ( $\text{dim } 6^2$ )
- neutron-antineutron oscillations ( $\text{dim } 9$ )
- Neutrinoless double beta decay ( $\text{dim } 5^2$ )
- Lepton flavor violation ( $\text{dim } 6^2$ )
- Electric dipole moments ( $\text{dim } 6$ )
- Pion, neutron, nuclear beta decay ( $\text{dim } 6$ )
- Electron/muon  $g-2$  ( $\text{dim } 6$ )
- Flavor physics ( $\text{dim } 6$ )
- Coherent neutrino scattering ( $\text{dim } 6$ )
- .....

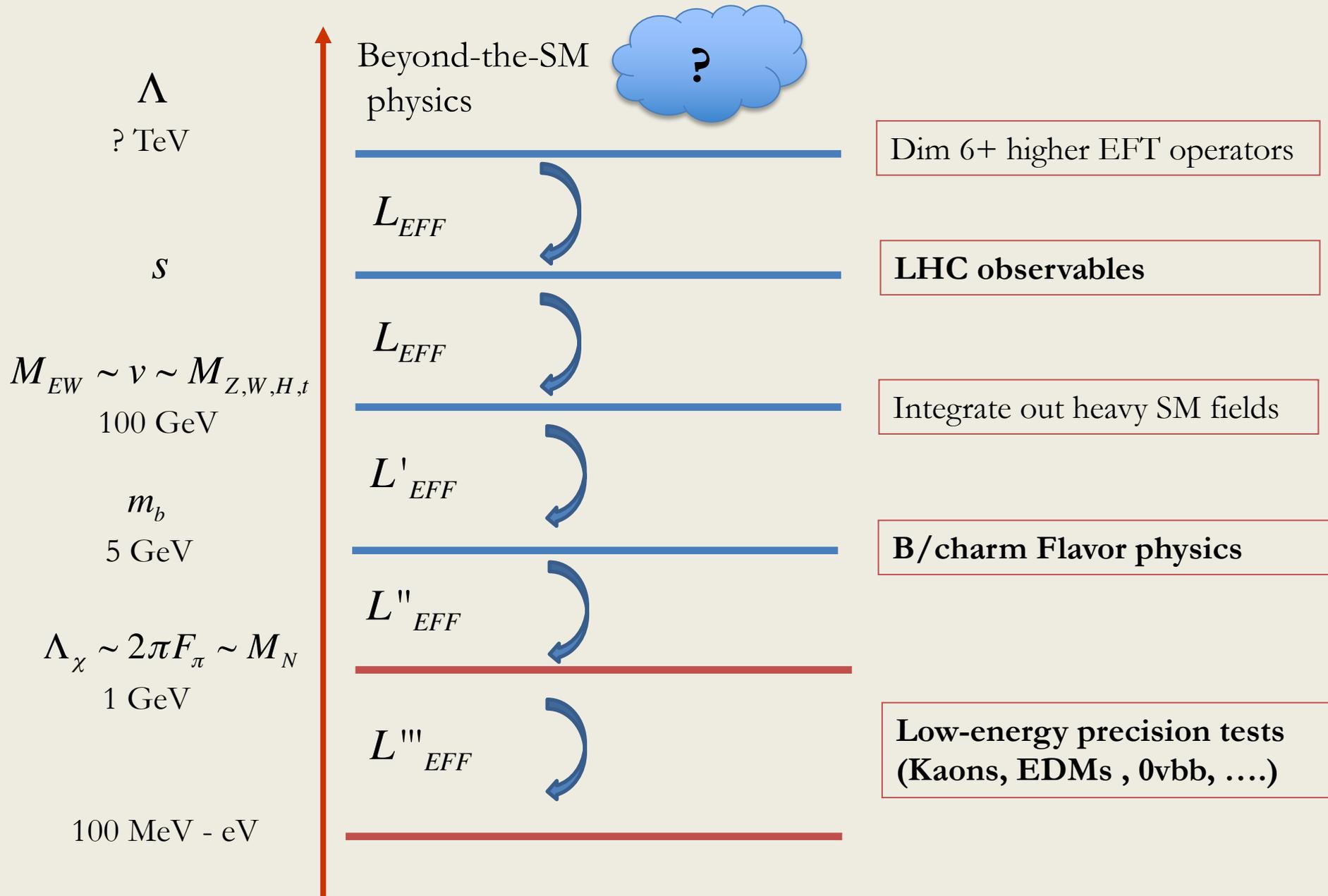


# Large number of low-energy SM tests

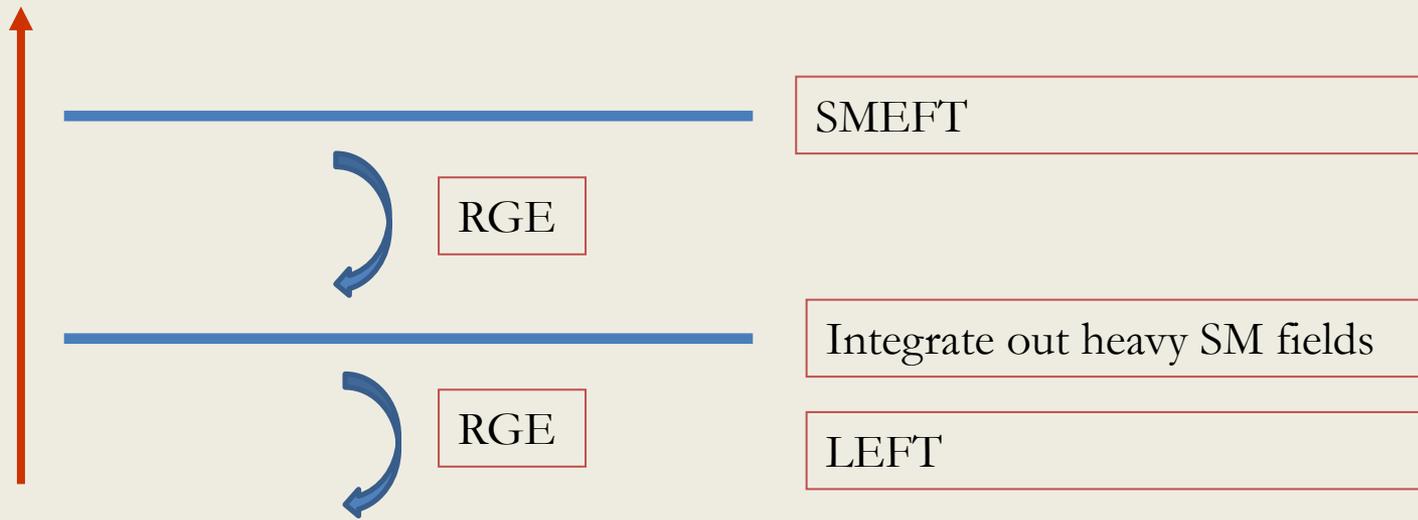
- For BSM physics to be observable at the LHC, we need to avoid the ‘peaks’
  1. Avoid FCNC and CP by hand or via Minimal Flavor Violation
  2. Assume B+L conservation (accidental in SM)
- Reasonable: but of course strong assumptions on BSM models
- **Best way to go beyond these assumptions ?**



# The road to incorporate low-energy tests

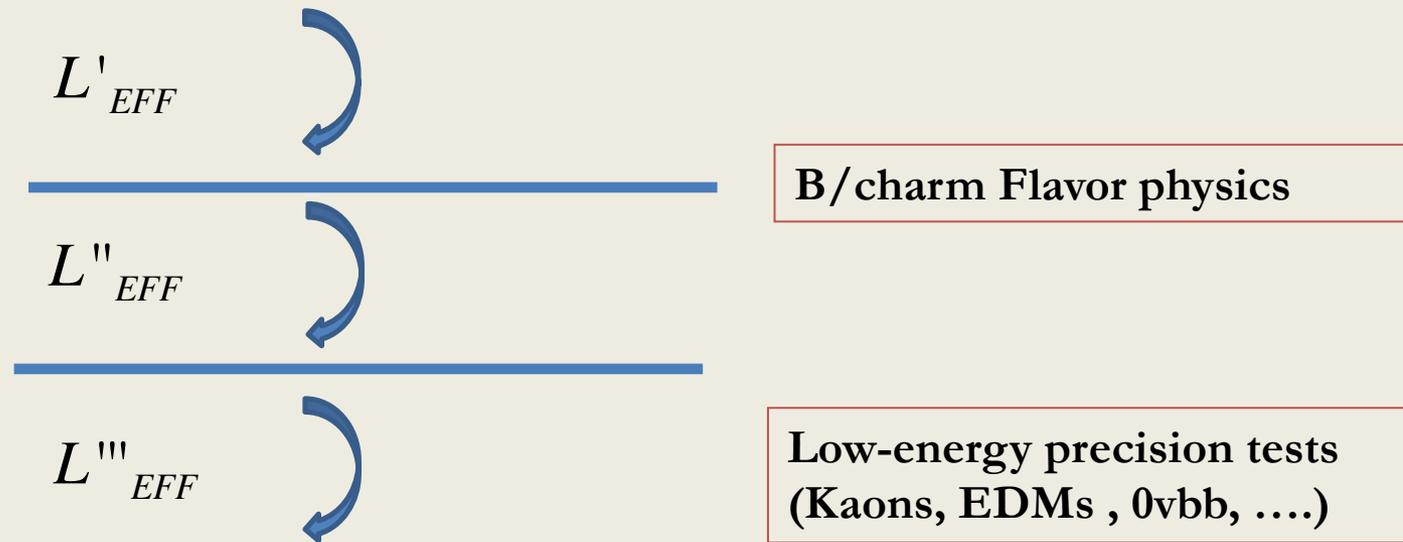


# Connecting to low-energy scales



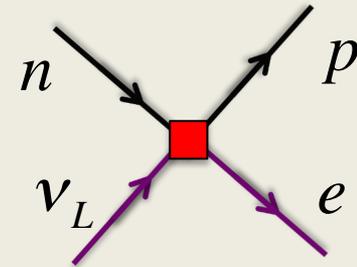
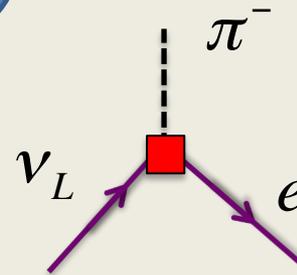
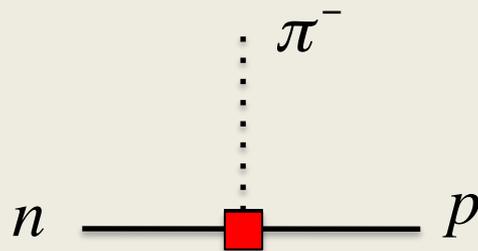
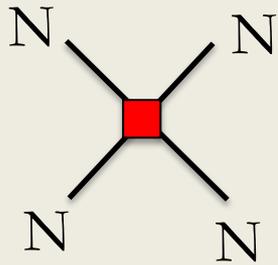
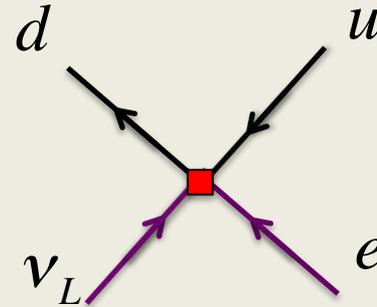
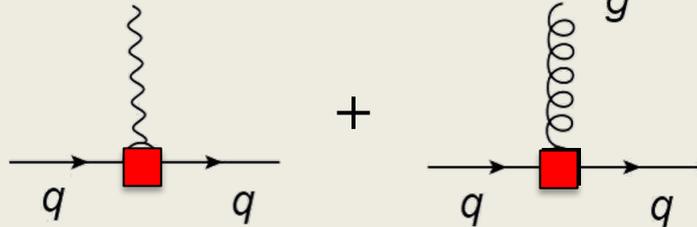
- **Renormalization-group equations** to evolve to low energies
- Full 1 loop below and above EW scale known (Jenkins et al '13 '14 '17)
- 2- or 3-loop known for subset of operators (e.g. 1907.04923)
  
- Matching to **LEFT** below EW scales (integrate out top, Higgs, W-Z )
- **Full one-loop matching by Dekens/Stoffer: 1908.05295**
- Some cases 2-loop matters (e.g. Barr-Zee diagrams)

# Challenges of non-perturbative QCD



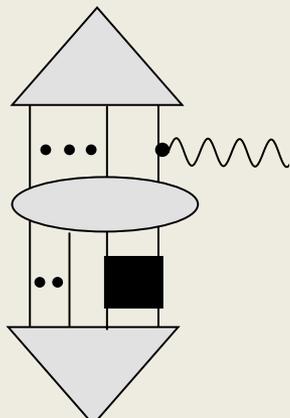
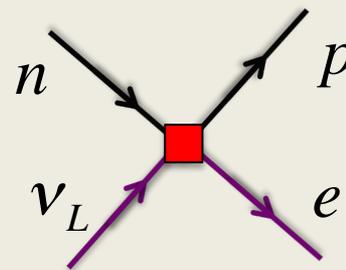
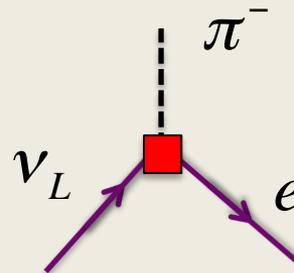
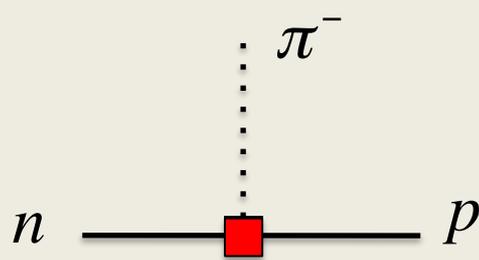
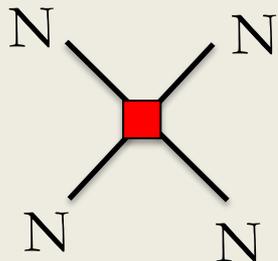
- SM-EFT can be matched to EFTs for non-perturbative QCD
  1. Heavy Quark EFT (bottom, charm)
  2. Chiral Perturbation Theory (nucleons, pions, kaons)
  3. Chiral EFT (nuclei)

# From LEFT to ChPT



- ChPT to translate LEFT to hadrons. **Power counting identifies leading terms.**
- Matrix elements ideally from lattice QCD
- Great progress for nucleon form factors, LNV, CPV, g-2 QCD corrections
- But still 100% uncertainties for many cases (e.g.  $0\nu\beta\beta$ , EDMs)

# From LEFT to ChEFT

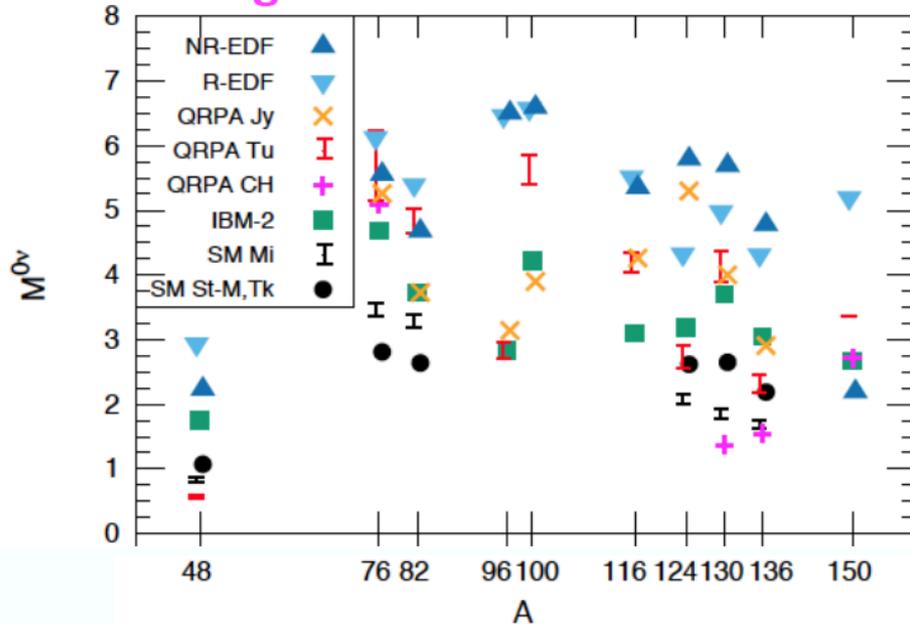


ChEFT to systematically calculate nuclear BSM observables  
EDMs,  $0\nu\beta\beta$ ,  $\mu$ -to- $e$  conversion, DM direct detection, ....

- Chiral EFT for both **wave functions** and **BSM currents**.
- **First-principle** for light nuclei (tremendous progress last few years)
- Nuclear approximations for heavier systems (e.g. Shell Model)
- **Works great in some cases, for others still big problem**

# Open problems

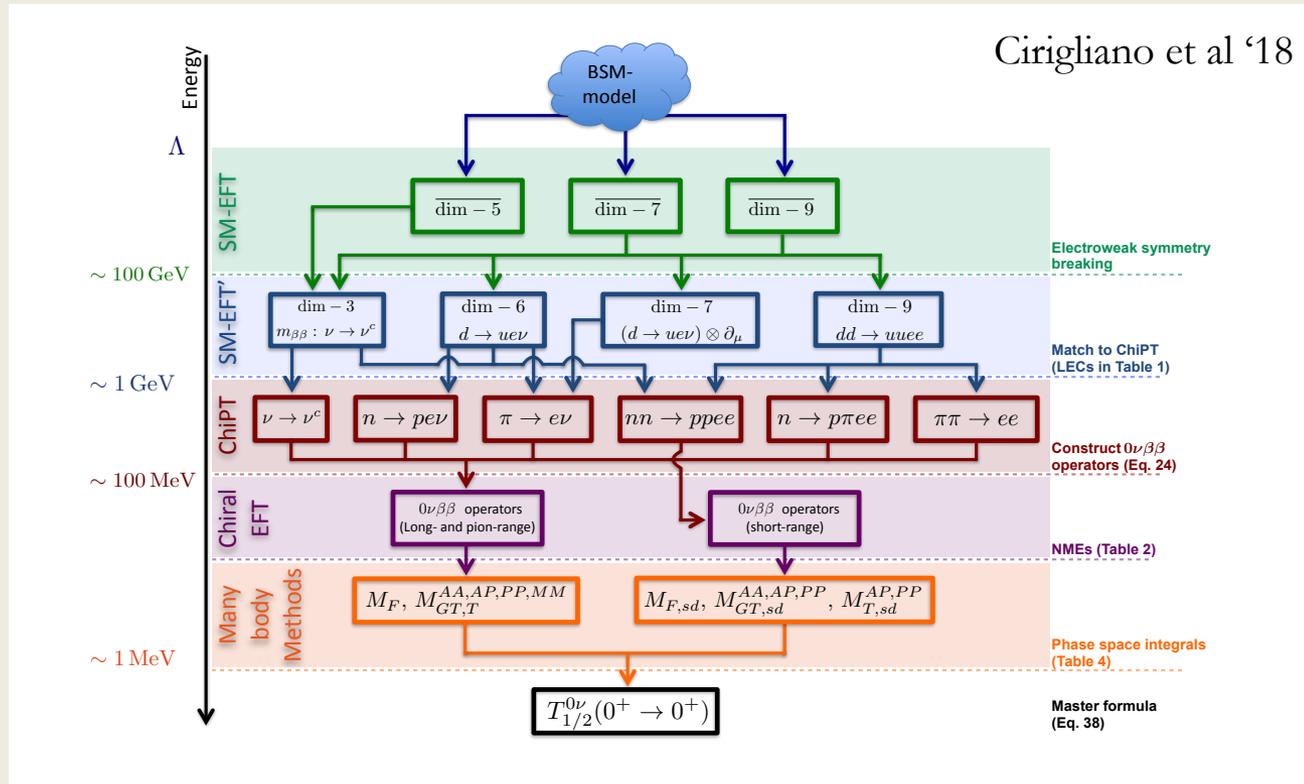
## Engel-Menendez 1610.06548



EDMs	$a_1$ range (best)
$^{199}\text{Hg}$	$0.030 \pm 0.060$
$^{225}\text{Ra}$	$14 \pm 12$
$^{129}\text{Xe}$	$-0.03 \pm 0.025$

- Nuclear observables that depend on two-nucleon processes (neutrinoless double beta decay, CPV nuclear force, Hadronic Parity Violation) are **problematic**
- **Power counting** not always understood (*do we even know what is leading order?*)
- **Chiral EFT is rich topic: benefits from closer particle-nuclear connection and the EFT experts in the hep-ph community**
- See for instance 2020 INT program: BSM physics with nucleons and nuclei

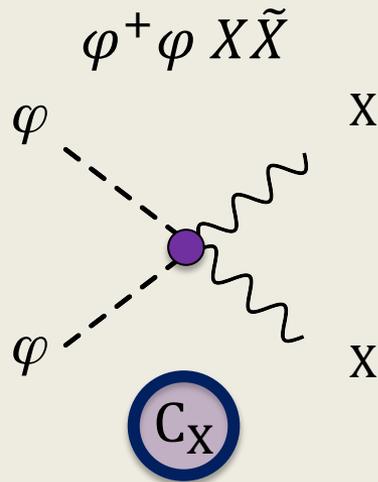
# End-to-End EFT framework



- Cascade of EFTs to connect low-energy observables to SMEFT at high scales
- Exists e.g. for  $0\nu\beta\beta$ , EDMs, beta decay, coherent scattering, many flavor observables
- **How to deal with hadronic/nuclear uncertainties in fits ?**
- **How to truncate operator basis to make fit manageable ?**
- Automization tools would be very welcomes (a la **flavio** for flavor physics)

# An explicit example: CPV in Higgs sector

- CP violation typically ignored in fits of Higgs, top, EW global fits
- Lot of interest in CPV in higgs and top physics (e.g. for EW baryogenesis)
- *See also: Snowmass-2021 CPV Higgs couplings by Andrei Gritsan*
- 4 gauge-Higgs operators exist (B, W, BW, G)



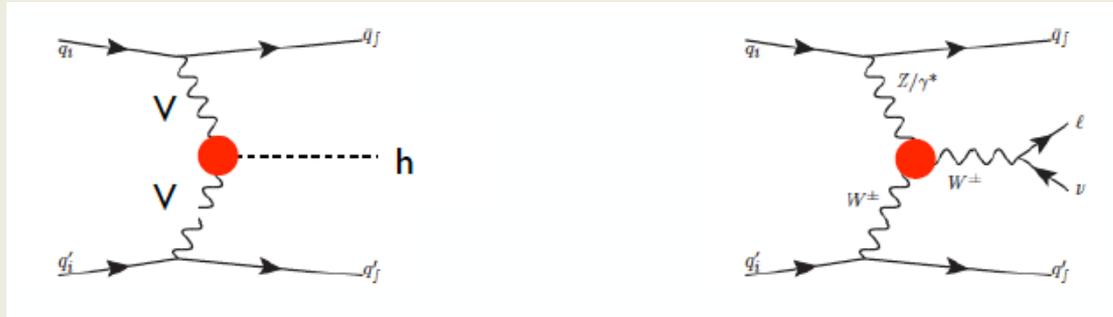
→  
After EW  
symmetry  
breaking

- **h-gluon-gluon**
- **h-gamma-gamma**
- **h-gamma-Z**
- **W-W-gamma**
- **h-Z-Z (not independent)**

- Evades flavor constraints (MFV automatic). Scale can be relatively low
- Motivated by universal theories (BSM couples to SM bosons/fermions through SM currents)

# Collider and low-energy probes

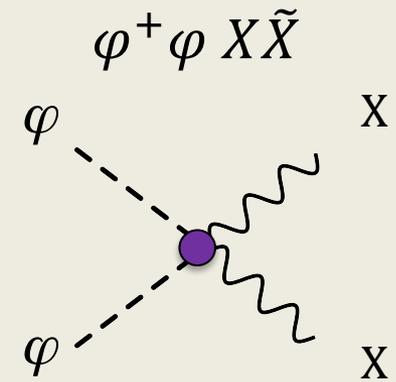
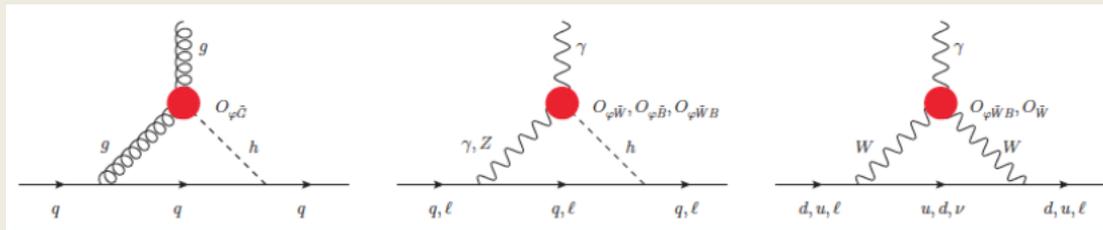
- Induce CPV angular distribution in  $pp \rightarrow h/V + 2 \text{ jets}$



e.g. ATLAS 2006.15458  $0.23 < \tilde{C}_{HWB}/\Lambda^2 < 2.34 \text{ (TeV}^{-2}\text{)}$

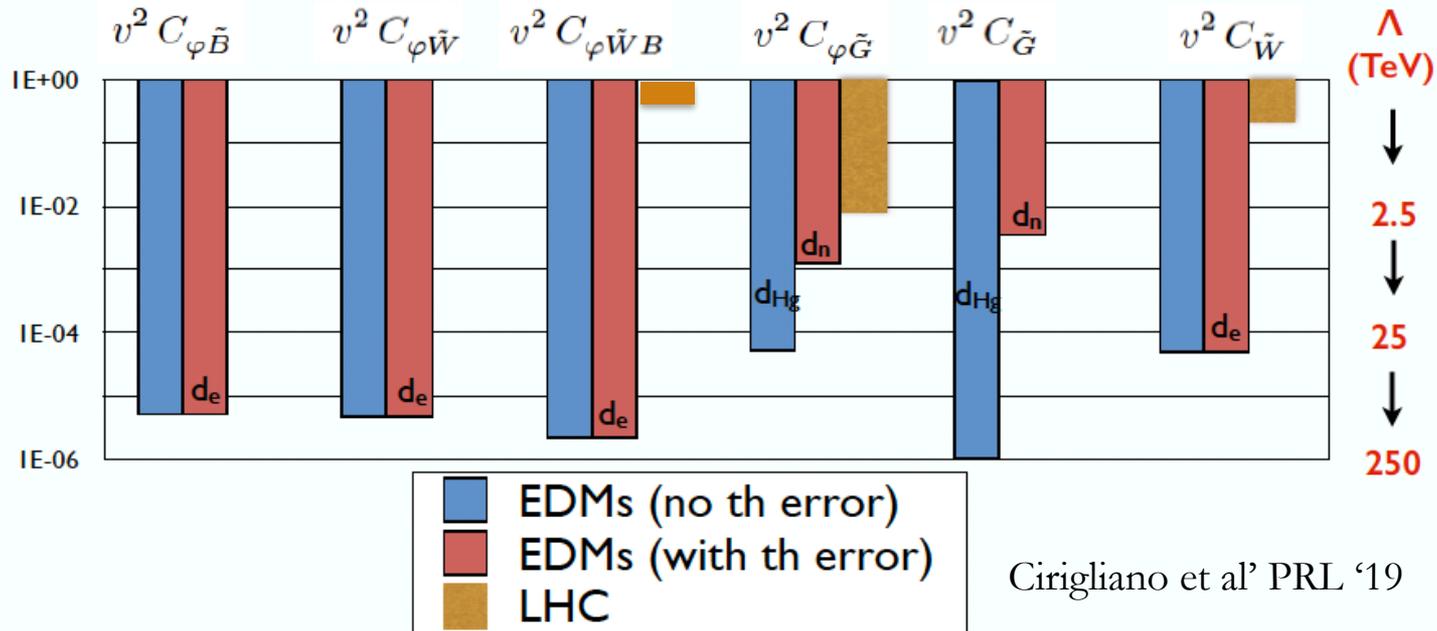
Bernlochner et al '19  $-0.19 < \tilde{C}_{HGG}/\Lambda^2 < 0.03 \text{ (TeV}^{-2}\text{)}$

- Same couplings induce CPV fermionic operators at one loop
- Also induce CPV in  $B \rightarrow s \text{ gamma}$  transitions



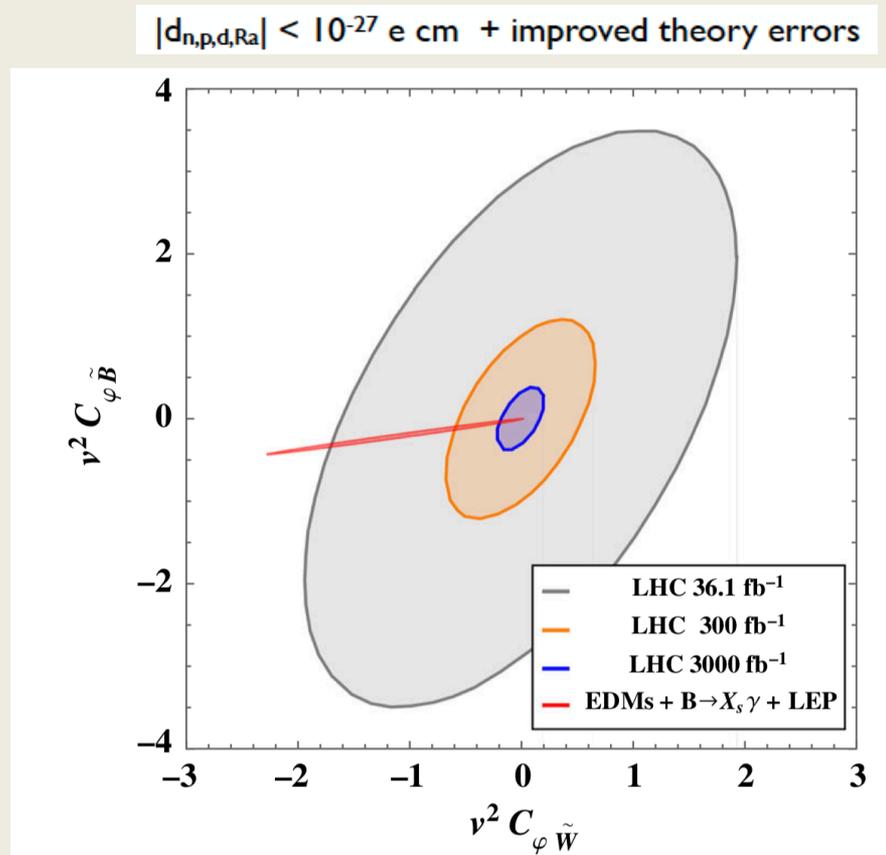
# Low-energy constraints are stringent

- Current constraints, “turning on” one coupling at a time: EDMs vs LHC



- EDM constraints are very stringent for single couplings
- Theory errors are relevant for nuclear CP violation (very severe!)
- **But EDMs suffer from ‘free directions’**  $0.17 C_{\varphi\tilde{B}} + 0.86 C_{\varphi\tilde{W}} + 0.48 C_{\varphi\tilde{W}B}$
- **Unconstrained direction: appears in global fits**

# CP violation in ‘universal theories’



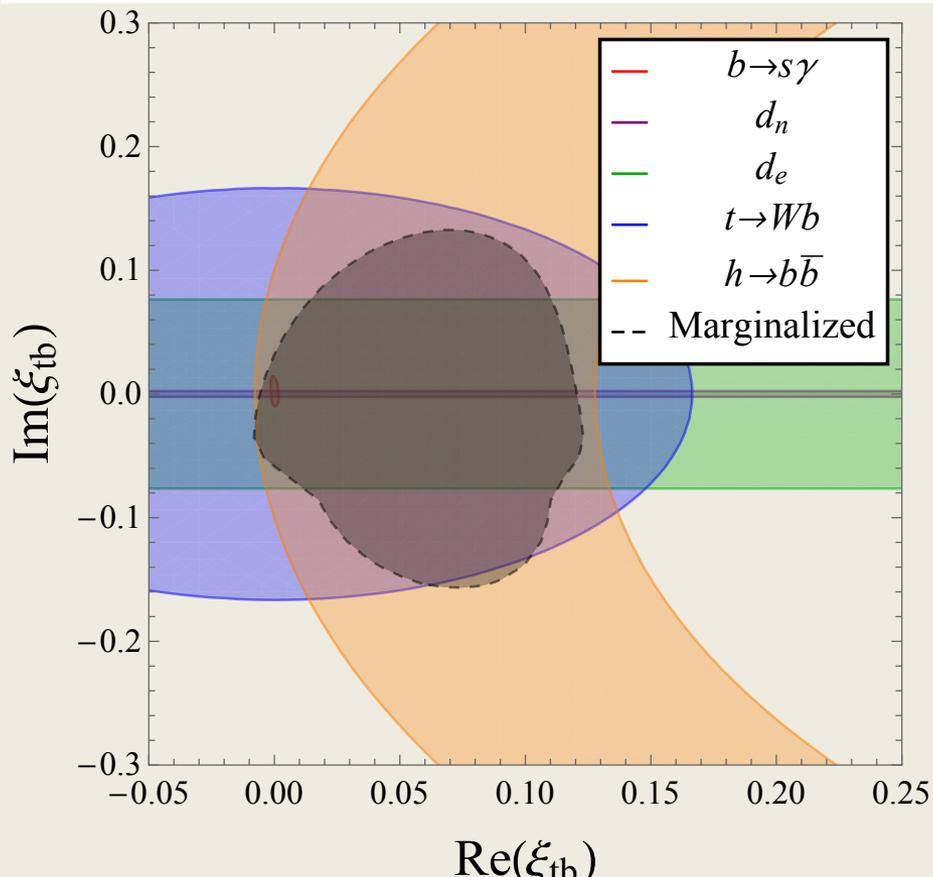
Cirigliano et al PRL ‘19

HL-LHC projections from  
Bernlochner et al ‘18

- Low-energy limits avoided in global fits (free directions)
- Strong constraints on possible UV completion (must point in free direction or appear at high scales only)
- **Low- and high-energy data together will close parameter space**

# Another example: modified $Wtb$ vertices

Real	Individual	Marginalized	Imaginary	Individual	Marginalized
$\xi_{tb}$	$[-1.4, 1.5] \cdot 10^{-3}$	$[-0.01, 0.12]$	$\xi_{tb}$	$[-2.4, 2.4] \cdot 10^{-3}$	$[-0.16, 0.13]$
$\delta_{LL}$	$[-0.03, 0.04]$	$[-0.03, 0.04]$	—		
$v^2 C_{Wt}$	$[-0.09, 0.05]$	$[-0.10, 0.04]$	$v^2 C_{Wt}$	$[-6.0, 6.0] \cdot 10^{-4}$	$[-1.0, 0.9] \cdot 10^{-3}$
$v^2 C_{Wb}$	$[-0.04, 0.05]$	$[-3.5, 0.4]$	$v^2 C_{Wb}$	$[-3.5, 3.5] \cdot 10^{-2}$	$[-1.9, 2.2]$



Alioli, Cirigliano, Dekens, JdV, Mereghetti '17  
de Beurs, Laenen, Vreeswijk, Vryonidou '18

# These are just examples

- **All experiments are low-energy experiments in SMEFT**
- Even if low-energy constraints ‘disappear’ from global limits they still enforce non-trivial constraints on UV models
- We should include as much data as we can
- Benefit from **closer connection between communities** (e.g collider versus fundamental symmetries) who are looking for the same thing!
- **Personal interests for future of SMEFT@low energies**
  1. Hopefully more global analyses including low-energy data
  2. Better understanding of chiral EFT power counting for BSM
  3. Go beyond SM-EFT by adding light d.o.f.: e.g. sterile neutrinos or scalars
  4. Theta term not protected in SMEFT (unlike SM): new hierarchy problem. CP-odd observables could provide indirect evidence for Peccei-Quinn mechanism in nature *Draper, Kozaczuk et al' 18*
  5. Potential of new experiments (e.g. EIC) (see talk by Daniel Wiegand)
  6. ....