# Analaysis of X-arapuca

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In the following, we focus on the shifting and evolution of the LAr slow scintillation component only. We assume that the fast LAr scintillation is not affected by the N2 quenching and the Xenon doping and it can anyway be treated analogously. Under these conditions, the quenching, the shifting and the scintillation processes can be described with these three equations:



#### What happens in the vacuum ultraviolet:



Jirich-Ar-Xe-mix.t  $OQ$ 3 / 14

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- $\bullet$  the ArAr<sup>\*</sup> excited dimer (AA) disappears by scintillation (at 128 nm), quenching through nitrogen, and shifting to ArXe<sup>\*</sup> (AX) through the process:
	- $ArAr^* + Xe \rightarrow ArXe^* + Ar$
- **■** the ArXe<sup>\*</sup> dimer is formed from ArAr<sup>\*</sup> dimer and disappears through scintillation  $(at 150 nm)$  and shifting to  $XeXe^*$   $(XX)$ through the process

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- $\bullet$   $ArXe^* + Xe \rightarrow XeXe^* + Ar$
- **3** the XeXe\* dimer is formed from the ArXe\* dimer and disappears through scintillation (at 175 nm)
	- $\bullet$  XeXe<sup> $* \rightarrow Xe + Xe$ </sup>

[\(arXiv:1511.07721\)](https://arxiv.org/abs/1511.07721) [\(arXiv:1403.0525\)](https://arxiv.org/ftp/arxiv/papers/1403/1403.0525.pdf)

### $A r A r^* + X e \rightarrow A r X e^* + A r$

$$
\frac{dAA}{dt} = -\frac{AA}{\tau_{128}} - \frac{AA}{\tau_{N2}} - \frac{AA}{\tau_{AX}}
$$
\n
$$
= -\frac{AA}{\tau_{TA}}
$$
\n(1)

#### $ArXe^* + Xe \rightarrow XeXe^* + Ar$



### $XeXe^* \rightarrow Xe + Xe$



These equations can be solved analytically resulting in the following solutions.

(2)

#### $ArAr^* + Xe \rightarrow ArXe^* + Ar$

$$
AA(t) = Ke^{-t/\tau_{TA}} \tag{4}
$$

 $ArXe^* + Xe \rightarrow XeXe^* + Ar$ 

$$
AX(t) = K\tau_{TA} \frac{\tau_{TX}}{\tau_{AX}} \left( \frac{e^{-t/\tau_{TA}} - e^{-t/\tau_{TX}}}{\tau_{TA} - \tau_{TX}} \right)
$$
(5)

#### $XeXe^* \rightarrow Xe + Xe$

$$
XX(t) = K \frac{\tau_{TA}}{\tau_{XX}} \frac{\tau_{TX}}{\tau_{AX}} \left( \frac{\tau_{175}}{\tau_{TA} - \tau_{TX}} \right)
$$

$$
\left[ \left( \frac{\tau_{TA}}{\tau_{TA} - \tau_{175}} \right) e^{-t/\tau_{TA}} - \left( \frac{\tau_{TX}}{\tau_{TX} - \tau_{175}} \right) e^{-t/\tau_{TX}} \right]
$$

$$
- \left( \frac{\tau_{TA}}{\tau_{TA} - \tau_{175}} - \frac{\tau_{TX}}{\tau_{TX} - \tau_{175}} \right) e^{-t/\tau_{175}} \right]
$$
(6)

Where  $K^*$  is initial amount of excited dimers  $ArAr^*$  in the triplet state. Note that the extended expression for XX(t) is valid also in case of two different  $\tau_{175}$  (a fast and a slow one), with two additional coefficients to take into account the relative splitting between the two components.

### $XeXe^* \rightarrow Xe + Xe$

If  $\tau_{175}$  is much smaller than  $\tau_{TA}$  and  $\tau_{TX}$ , it is easily seen that the expression of XX(t) can be approximately simplified as follows:

$$
XX(t) = K \frac{\tau_{TA}}{\tau_{XX}} \frac{\tau_{TX}}{\tau_{AX}} \left( \frac{\tau_{175}}{\tau_{TA} - \tau_{TX}} \right) \left( e^{-t/\tau_{TA}} - e^{-t/\tau_{TX}} \right)
$$
(7)

## scintillation

 $\bullet$ 

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From the above expressions, we can directly derive the expected rate of scintillation for the three expected decay processes, in the assumption that  $\tau_{175}$  is much faster than the other  $\tau$ 's :

$$
\frac{dAA}{dt}(\text{scint}) = K \frac{\tau_{TA}}{\tau_{128}} \frac{e^{-t/\tau_{TA}}}{\tau_{TA}}
$$
(8)

$$
\frac{dAX}{dt}(\text{scint}) = K \frac{\tau_{TA}}{\tau_{150}} \frac{\tau_{TX}}{\tau_{TAX}} \left( \frac{e^{-t/\tau_{TA}} - e^{-t/\tau_{TX}}}{\tau_{TA} - \tau_{TX}} \right) \tag{9}
$$

$$
\frac{dXX}{dt}(\text{scint}) = K \frac{\tau_{TA}}{\tau_{XX}} \frac{\tau_{TX}}{\tau_{AX}} \left( \frac{e^{-t/\tau_{TA}} - e^{-t/\tau_{TX}}}{\tau_{TA} - \tau_{TX}} \right) \qquad (10)
$$

# scintillation

- The X-Arapuca with no Quartz window (XN) would see the sum of the three spectra. The X-Arapuca with the Quartz window (XQ) would only be sensitive to the third spectrum (the one from XeXe\*).
- At times  $>> 3\tau_{TA}$  the first spectrum should be negligible, hence the ratio XQ/XN should be constant and its value should be  $\frac{\tau_{150}}{\tau_{XX}+\tau_{150}}$ . For  $\tau_{XX} < \tau_{150}$  (higher xenon concentration) the ratio should tend to 1. The deviation from 1, at high Xenon concentration, gives the inter-calibration (ic) factor between XN and XQ.
- The following plot shows that the experimental data can actually be nicely described by the above calculation with  $\tau_{XX} \sim \tau_{150}$  3/[Xe ppm] and ic = 1/0.95.

# inter-calibration

The difference  $XN - i c * XQ$  should be sensitive to the sum of the first and second spectra only. A fit to the data will allow to determine in the exponential functions.



A ratio of the integrals of the first and second spectra will then proportional to  $\frac{\tau_{128}\tau_{\mathcal{T}X}}{\tau_{150}\tau_{\mathcal{A}X}}$  The  $\tau_{\mathcal{T}150}$  can then be evaluated.

## Each Run deconvoluted and averaged for each Xe injections



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## Each Run deconvoluted and averaged for each Xe injections



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## $nonQ - Q * (lcalibration)$



 $Icalibration = 1.2$ 



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