
Top Mass Determinations at Hadron Colliders

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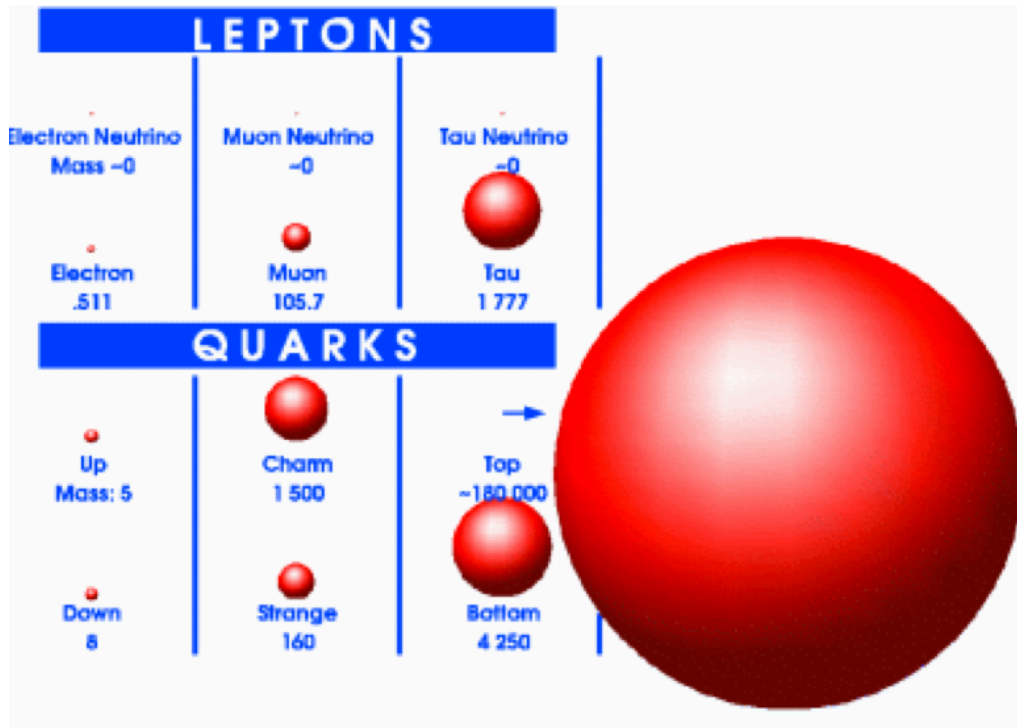
Please refer to EF03 Meeting September 10 2020 and
arXiv:2004.12915 for more details.

fdk Π Doktoratskolleg
Particles and Interactions



FWF
Der Wissenschaftsfonds.

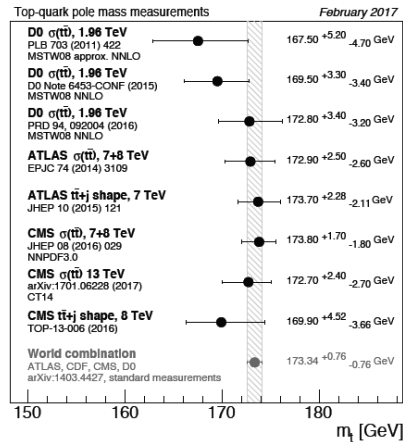
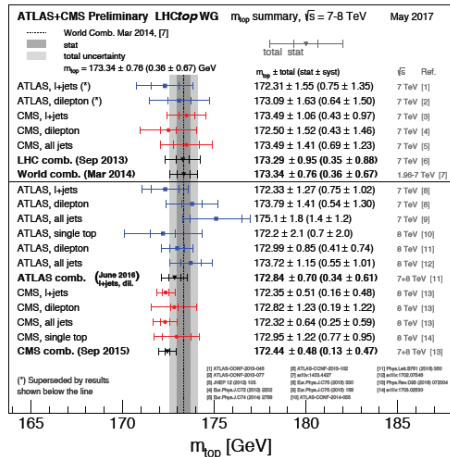
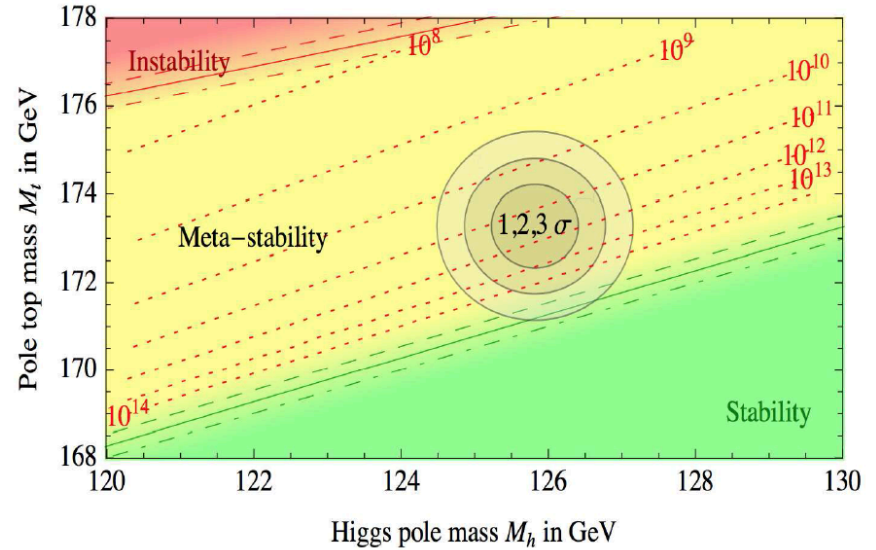
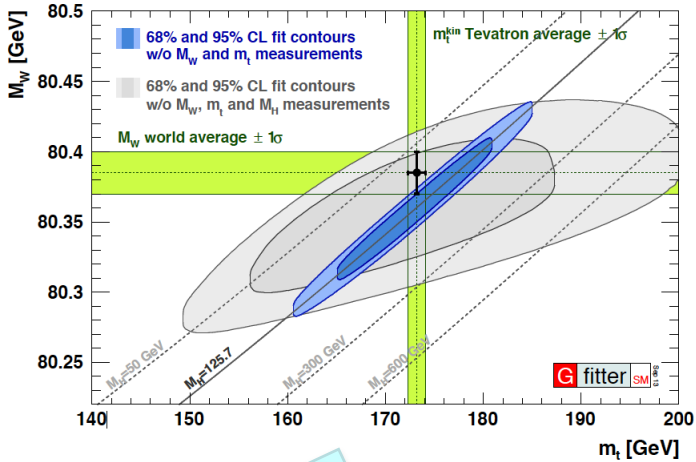
.. not just the heaviest SM particle



- Top quark: heaviest known particle
- Most sensitive to the mechanism of mass generation
- Peculiar role in the generation of flavor.
- Top might not be the SM-Top, but have a non-SM component.
- Top as calibration tool for new physics particles (SUSY and other exotics)
- Top production major background in new physics searches
- One of crucial motivations for New Physics

- Very special physics laboratory: $\Gamma_t \gg \Lambda_{\text{QCD}}$
 - Top treated as a particle: p_T , spin, σ_{tot} , $\sigma(\text{single top})$, $\sigma(\text{tt}+X)$,... $\rightarrow q \gg \Gamma_t$
 - Quantum state sensitive to low-E QCD and unstable particle effects: m_t , endpoint regions $\rightarrow q \sim \Gamma_t$
 - Multiscale problem: $p_T, m_t \gg \Gamma_t \gg \Lambda_{\text{QCD}}, \dots$ (depends on resolution scale of observable)

Why a Precision Top Mass is Important



Aims: m_{top} wanted!
 M_{top} is a renormalized QCD parameter!

- Reduce error in $m_{\text{top}}^{\text{MC}}$
- Improve / understand better MC
- Clarify mass scheme $m_{\text{top}}^{\text{MC}}$!

Top Mass Measurements

Most precise method: Direct Reconstruction

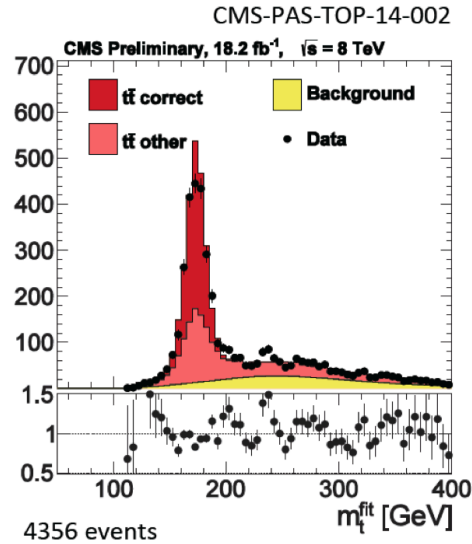
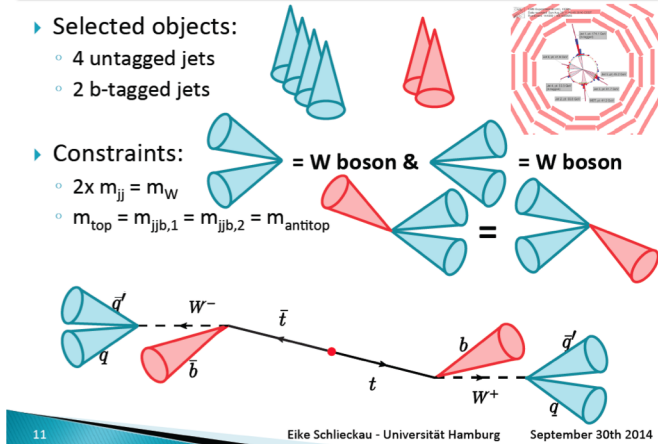
Kinematic Fit

Selected objects:

- 4 untagged jets
- 2 b-tagged jets

Constraints:

- $2 \times m_{jj} = m_W$
- $m_{top} = m_{jib,1} = m_{jib,2} = m_{antitop}$



kinematic mass determination

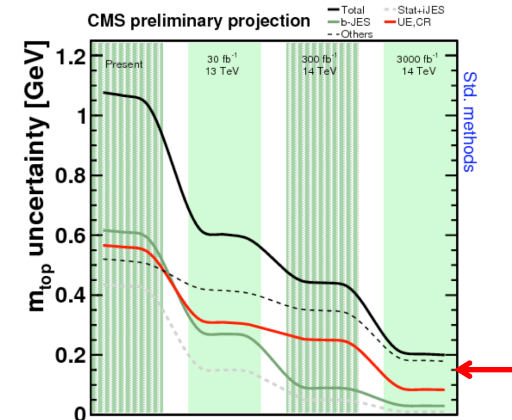
Determination of the best-fit value of the Monte-Carlo top quark mass parameter

⊕ High top mass sensitivity

⊖ Precision of MC ?

⊖ Meaning of m_t^{MC} ?

VALUE (GeV)	DOCUMENT ID	TECN	COMMENT
172.76 ± 0.30 OUR AVERAGE	Error includes scale factor of 1.2.		
172.69 ± 0.25 ± 0.41	1 AABOUD	19AC ATLS	7, 8 TeV ATLAS combination
172.26 ± 0.07 ± 0.61	2 SIRUNYAN	19AP CMS	lepton+jets, all-jets channels
172.33 ± 0.14 ⁺ ₋ 0.66 0.72	3 SIRUNYAN	19AR CMS	dilepton channel ($e\mu$, $2e$, 2μ)
172.95 ± 0.77 ⁺ ₋ 0.93	4 SIRUNYAN	17L CMS	t-channel single top production
172.44 ± 0.13 ± 0.47	5 KHACHATRY...16AK	CMS	7, 8 TeV CMS combination
174.30 ± 0.35 ± 0.54	6 TEVEWWG	16 TEVA	Tevatron combination



Mass Extraction and Renormalization Schemes

The Principle of Top Mass Determinations

- Top quark is not a physical particle (“colored parton”)
- Top mass defined from theoretical prescriptions (renormalization schemes)
- Different schemes are related by a perturbative series.

$$m_t^A - m_t^B = \sum_{n=1} c_n \alpha_s^n(\mu)$$

Parton level cross section formally scheme-invariant,
but can be practically scheme-dependent due to truncation

$$\hat{\sigma}(Q, m_t^A, \alpha_s(\mu), \mu; \delta m^A) = \hat{\sigma}(Q, m_t^B, \alpha_s(\mu), \mu; \delta m^B)$$

- For comparison with exp. data one has to account for non-perturbative corrections

$$\sigma^{\text{exp}} = \hat{\sigma}(Q, m_t^X, \alpha_s(\mu), \mu; \delta m^X) + \sigma^{\text{NP}}(Q, \Lambda_{\text{QCD}})$$

Typically at LHC: $\sigma^{\text{NP}} \sim \left(\frac{\Lambda_{\text{QCD}}}{Q}\right)^n, \quad n = 1$

Linear effects always arise from color neutralization processes.

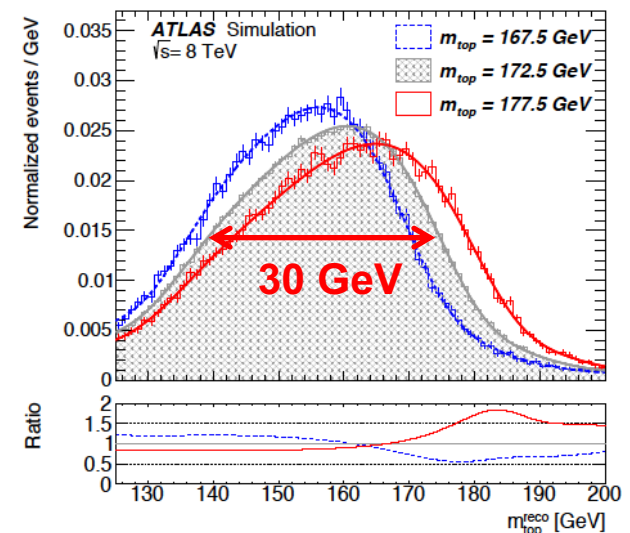
→ High precision control over soft partonic and NP effects needed when
mass sensitivity generated by small dynamical scales

Mass Extraction and Renormalization Schemes

- Parton level cross section and NP corrections MUST be separately consistent with QCD so that the top quark mass (as well as $\alpha_s(Q)$) can be determined reliably!
→ otherwise systematic bias: model instead of field theory parameters
- Which mass scheme is best? → Consider analogy to strong coupling α_s
 - Relevant dynamical scale $Q \Rightarrow \alpha_s(Q)$ frequently best choice (MSbar)
 - All quantum corrections to quark-gluon interactions from scales above Q are absorbed into $\alpha_s(Q)$ → IR-safe definition of strong coupling
 - Multiple scale problems: factorization allows to make adequate scale choices

We seek for a scale-dependent mass scheme $m_t(Q)$ with properties similar to the strong coupling $\alpha_s(Q)$.

- Multi-scale issue:
In general high mass sensitivity is associated with QCD dynamics at a low scale
→ typically: scale \sim width of distribution



Mass Extraction and Renormalization Schemes

Top Pole mass

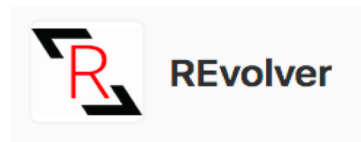
- Theoretical precision limit: $\sim 120 - 250$ MeV (pole mass renormalon)
- Most codes naturally in this scheme
- Scale independent

Top Mass Renormalization Schemes (renormalon-free running masses)

- Theoretical precision limit: $\sim 10 - 20$ MeV
- Theoretical work needed to implement scheme change
- Scale-dependent

\overline{MS} mass: • Adequate for total cross sections, production rates (scales above m_t)

MSR mass: • Adequate for thresholds, resonances, kinks (scales below m_t)



- C++ / Mathematica / Python package
- All common mass schemes supported
- All known corrections implemented

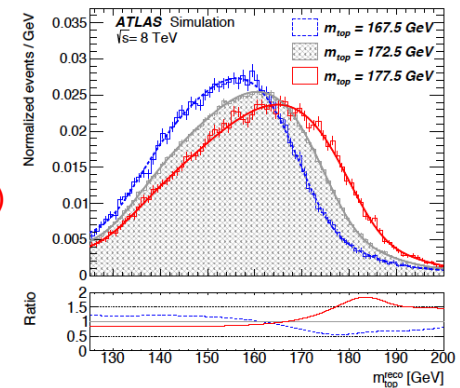
Release shortly arXiv: 2101:xxx

Status of Top Mass Determinations at the LHC

Direct Measurements:

→ 1st path to make progress

- Template method (ATLAS), matrix element/ideogram method (CMS)
- Based on highly top mass sensitive distributions (M_{lb-jet} , m_t^{reco} , etc) that are dominated by parton shower and hadronization model and **cannot be systematically improved by NLO or NNLO matching.**
(Mazitelli etal. arXiv:2012.14267)
- Problem: How is m_t^{MC} related to field theory mass schemes?
(Top mass interpretation problem)



Better theoretical understanding of MC event generators

needed! → work in progress (will not be resolved quickly,

comparable in complexity to the task to develop NLL precise MC generators)

See talk at EF03, Sept 10, 2020 and arXiv:2004.12915

Status of Top Mass Determinations at the LHC

“Pole Mass Measurements”:

- Based on total and differential cross section for which the parton level calculation can be done reliably at NLO or NNLO/NNLL → mass scheme under control
- Called “pole mass measurements” only because theorists used pole mass scheme for their calculations. → misleading!

Better: Measurements of m_t in well-defined scheme

- Total inclusive cross section:

$$m_t^{\text{pole}} = 172.9^{+2.5}_{-2.6} \text{ GeV (ATLAS, 7 and 8 TeV data)}$$

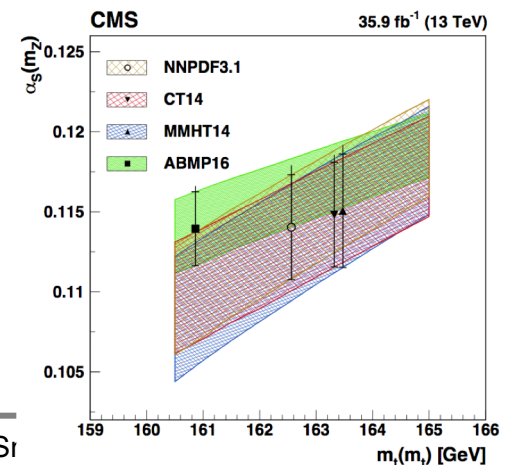
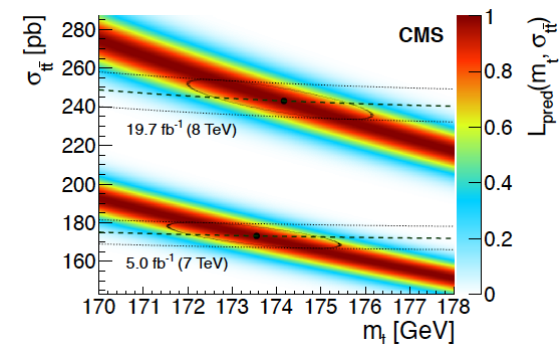
$$m_t^{\text{pole}} = 173.8^{+1.7}_{-1.8} \text{ GeV (CMS, 7 and 8 TeV data)}$$

$$m_t^{\text{pole}} = 169.9^{+2.0}_{-2.2} \text{ GeV (CMS, 13 TeV data)}$$

lower precision due to impact of norm uncertainties

(strong additional correlation to pdfs, α_S)

→ reliable mass interpretation, but imprecise



CMS arXiv:1812.10505

Status of Top Mass Determinations at the LHC

Differential Cross Section Measurements:

→ 2nd path to make progress

- Recently also differential cross sections: $M_{t\bar{t}+jet}$, $M_{t\bar{t}} + y(tt)$, lepton energies
→ based on concrete theory improvable (FO) calculations (with mass scheme control)
→ distributions elevate top mass sensitivity due to structures

$$M_{t\bar{t}} + y(t\bar{t}) : m_t^{\text{pole}} = 170.5 \pm 0.8 \text{ GeV} \quad (\text{CMS})$$

$$M_{t\bar{t}+jet} : m_t^{\text{pole}} = 171.1_{-1.1}^{+1.2} \text{ GeV} \quad (\text{ATLAS})$$

$$\text{leptons} : m_t^{\text{pole}} = 173.2 \pm 1.6 \text{ GeV} \quad (\text{ATLAS})$$

Important questions to address:

- Reliability of FO parton level differential cross sections
 - Test pole mass versus running masses → Garzelli, Kemmler, Moch, Zenaiev 2009.07763
 - Much more difficult (theory + experiment) than inclusive cross sections
(Hard work needed: Do not expect easy competition with direct measurement)
- Recent studies:
 - Soft-dropped boosted top jet masses → AHH, Mantry, Pathak, Stewart 1708.02586
 - Lepton energy distribution (t-channel single top) → Yuan, Gao, Gao 2007.15527