Constraining the Higgs potential at the Multi-TeV Muon Collider

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(partially) based on arXiv:2003.13628

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H self-couplings measurement: future colliders (HHHH)

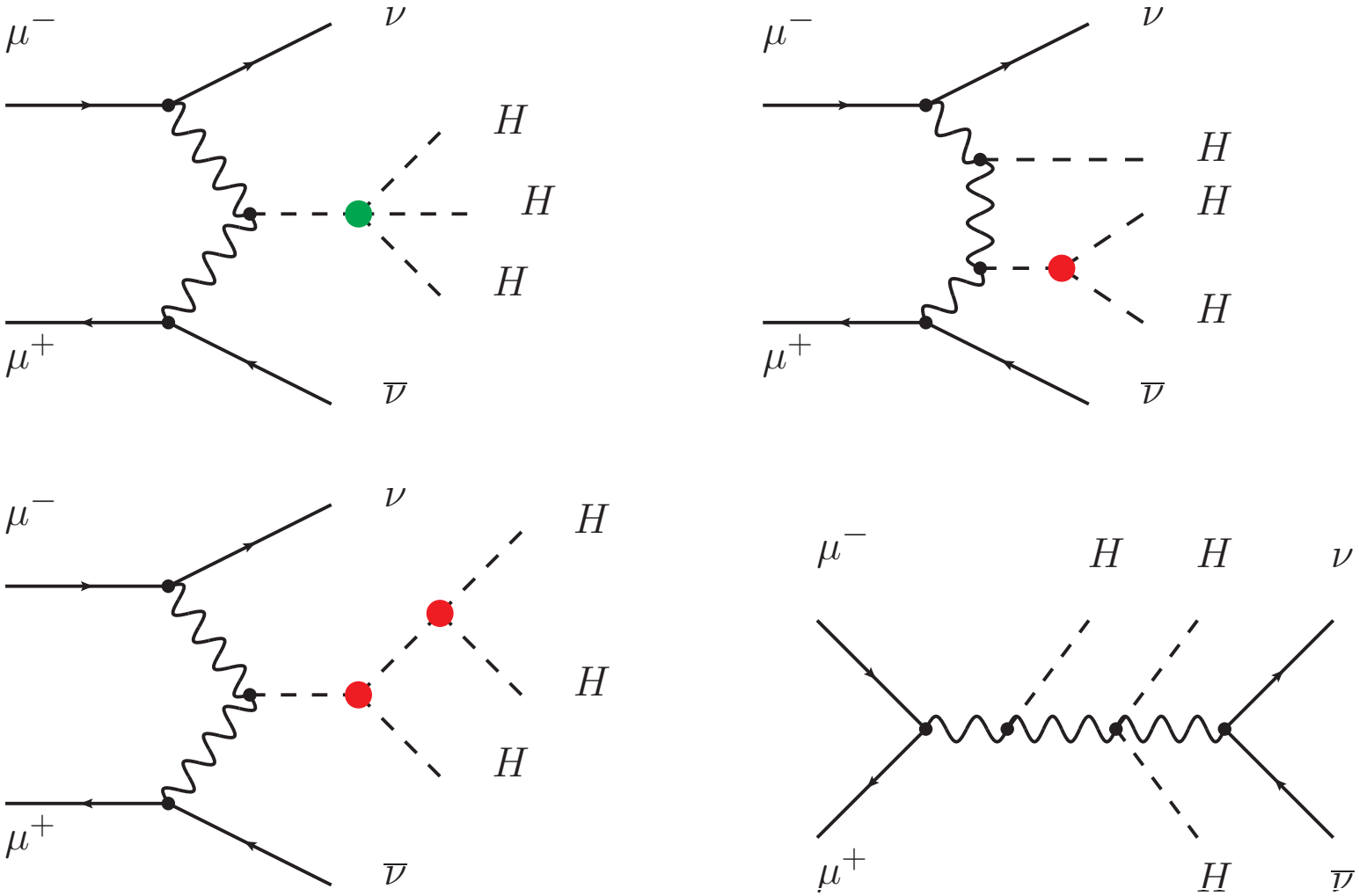
- the proposed future colliders can put strong constraints on the triple Higgs coupling δ_3: $\pm 10\%$ 1-σ bound at CLIC and ILC, $\pm 5\%$ at FCC
- the bounds on the quartic couplings δ_4 are very loose (68% CL)
  - ILC: $\sim [-10, +10]$ ($\pm 1000\%$!)
  - CLIC: $\sim [-5, +5]$
  - FCC: $\sim [-5, +15]$, from $pp \to HHH$
  - FCC: $\sim [-2, +4]$, from $pp \to HH$

I will focus on the sensitivity of the muon collider to the quartic coupling

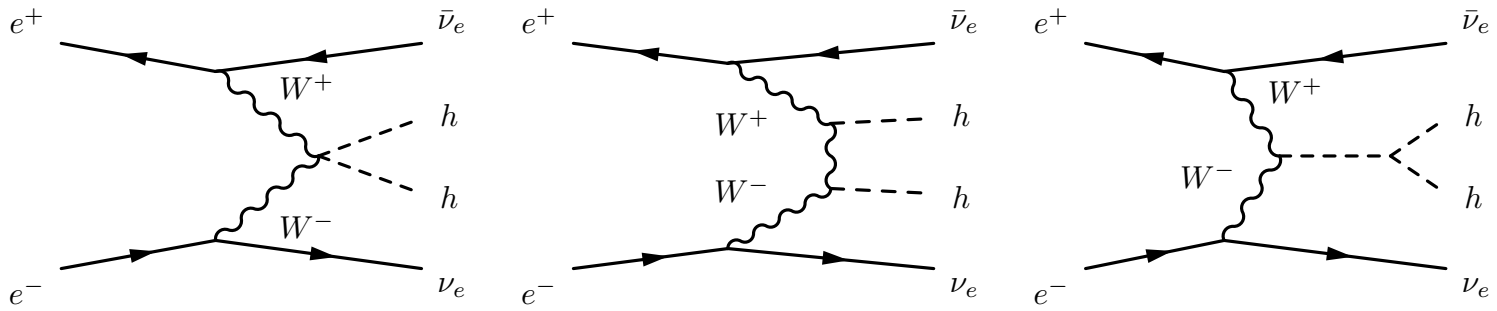
Spoiler:

under (reasonable) assumptions on the energy and the luminosity, the muon collider can do a pretty good job in constraining the quartic Higgs coupling

$\mu^+\mu^- \to HHH\nu\overline{\nu}$



Remark: HHH coupling at the muon collider



- Though $\mu^+\mu^- \to HHH\nu\bar{\nu}$ is sensitive to the HHH coupling, it is more natural to measure κ_3 from double Higgs production
- Sensitivity to δ_3 at the μ-coll. studied in:
  - A. Constantini et al. 2005.10289
  - T. Han et al. 2008.12204
  - D. Buttazzo et al. 2012.11555
- roughly speaking (under reasonable assumptions on the luminosity):
  - $\delta_3 \in \pm 6\%$ for $\sqrt{s} = 10$ TeV at 2σ
  - $\delta_3 \in \pm 2\%$ for $\sqrt{s} = 30$ TeV at 2σ

$\mu^+\mu^- \to HHH\nu\bar{\nu}$: SM Higgs couplings (energy)

$\sqrt{s}$ (TeV) / L (ab^{-1})	1.5 / 1.2	3 / 4.4	6 / 12
σ_{SM} (ab) $[N_{\rm ev}]$			
$\sigma^{\rm tot}$	0.03 [0]	0.31 [1]	1.65 [20]
$\sigma(M_{HHH} < 3$TeV$)$	0.03 [0]	0.31 [1]	1.47 [18]
$\sigma(M_{HHH} < 1$TeV$)$	0.02 [0]	0.12 [1]	0.26 [3]

$\sqrt{s}$ (TeV) / L (ab^{-1})	10 / 20	14 / 33	30 / 100
σ_{SM} (ab) $[N_{\rm ev}]$			
$\sigma^{\rm tot}$	4.18 [84]	7.02 [232]	18.51 [1851]
$\sigma(M_{HHH} < 3$TeV$)$	2.89 [58]	3.98 [131]	6.69 [669]
$\sigma(M_{HHH} < 1$TeV$)$	0.37 [7]	0.45 [15]	0.64 [64]

σ increases with $\sqrt{s}$

$\mu^+\mu^- \to HHH\nu\bar{\nu}$: SM Higgs couplings (luminosity)

$\sqrt{s}$ (TeV) / L (ab^{-1})	1.5 / 1.2	3 / 4.4	6 / 12
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The lower energy setups (1.5 and 3 TeV) do not have enough events to study the quartic Higgs self coupling

$\mu^+\mu^- \to HHH\nu\bar{\nu}$: SM Higgs couplings (luminosity)

- the luminosities assumed for $\sqrt{s} = 1.5, 3, 6, 14$ TeV are based on MAP studies

V. Shiltsev FERMILAB-FN_1083-AD-APC,

talks by D. Shulte and M. Palmer https://indico.cern.ch/event/847002/

- at $\sqrt{s} = 10, 30$ TeV, the luminosity is fixed by (see arXiv:1910.06150)

Luminosity:

$$L \gtrsim \frac{5\,\text{years}}{\text{time}} \left(\frac{\sqrt{s}_\mu}{10\,\text{TeV}} \right)^2 2 \cdot 10^{35} \text{cm}^{-2}\text{s}^{-1}$$

- for the 10 and 30 TeV setups, it might be that higher luminosity could be achieved

Deviations from SM Higgs couplings

$$\mathcal{L} = -\frac{1}{2} M_H^2 H^2 - \left(1+\delta_3\right)\frac{M_H^2}{2v} H^3 - \left(1+\delta_4\right)\frac{M_H^2}{8v^2} H^4$$

We consider 3 different scenarios:

1. $\delta_3 = 0$, δ_4 arbitrary
2. δ_3 arbitrary, $\delta_4 = 6\delta_3$ (well behaved SMEFT)

S. Borowka et al. arXiv:1811.12366

3. δ_3 arbitrary and δ_4 arbitrary

Sensitivity to δ_3 and δ_4

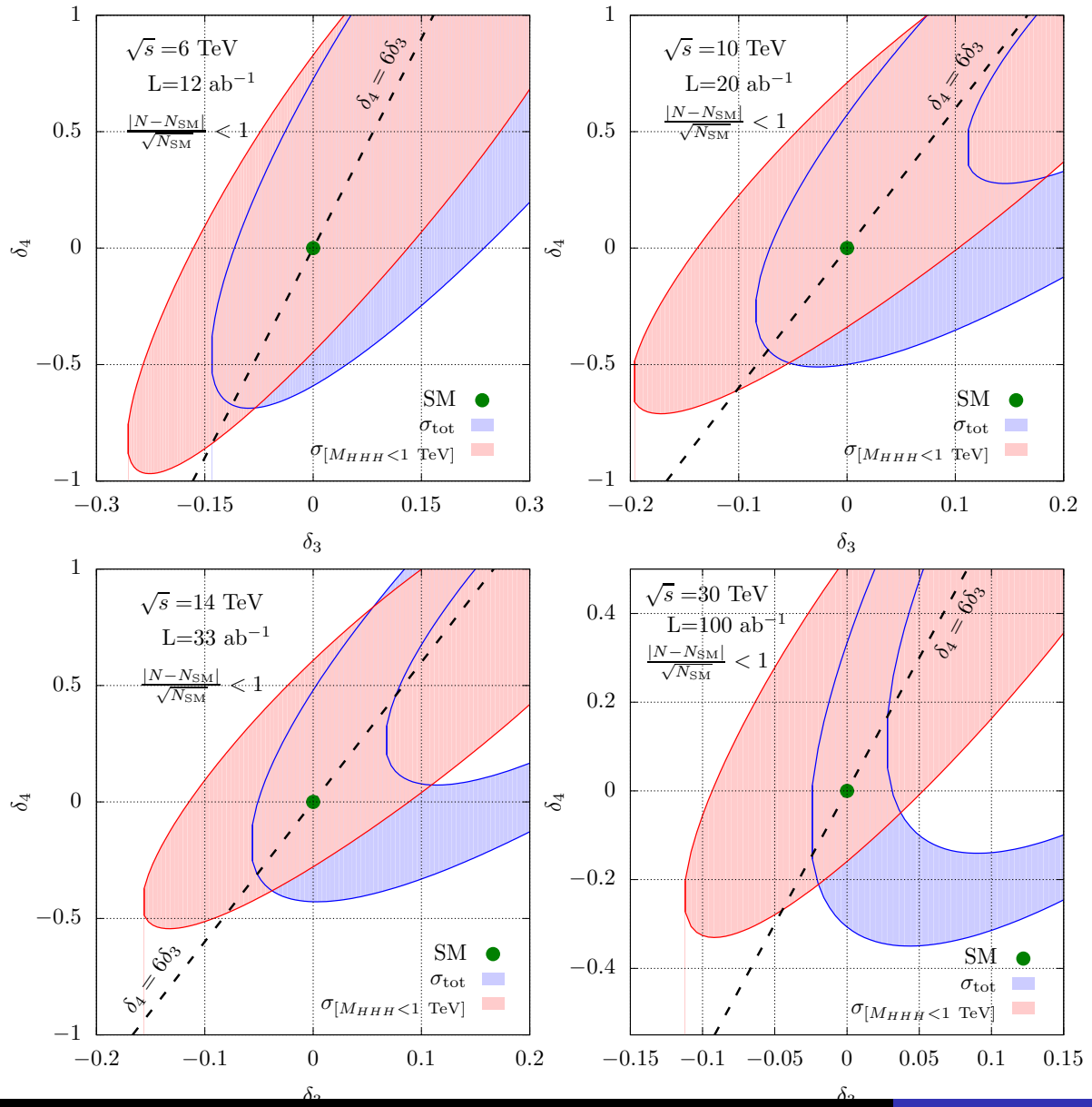
No background process considered:
we quantify the sensitivity in terms of standard deviations from the SM expectation

$$\frac{|N - N_{\mathrm{SM}}|}{\sqrt{N_{\mathrm{SM}}}}$$

Remarks

- no background is considered, but the environment should be rather clean
- no branching ratio is applied, but if the environment is clean enough all the main decay channels should be visible
- (almost) no optimization based on kinematics is performed, so there is room for improvement

Sensitivity to δ_3 and δ_4 (small δ_3)



- no cuts
- $M_{HHH} < 1$ TeV

$$\delta_3 = 0$$

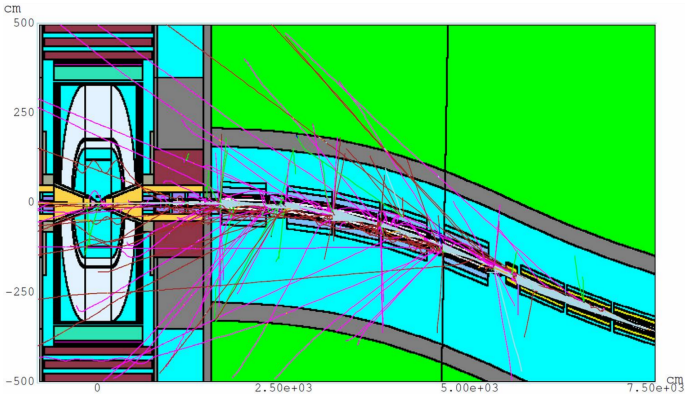
6 TeV $\delta_4 \sim [-0.45, 0.8]$

10 TeV $\delta_4 \sim [-0.4, 0.7]$

14 TeV $\delta_4 \sim [-0.35, 0.6]$

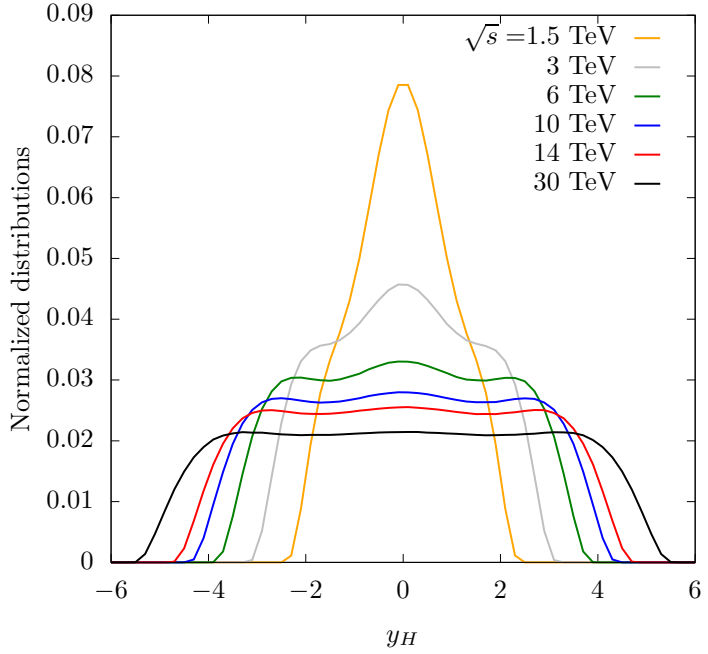
30 TeV $\delta_4 \sim [-0.2, 0.5]$

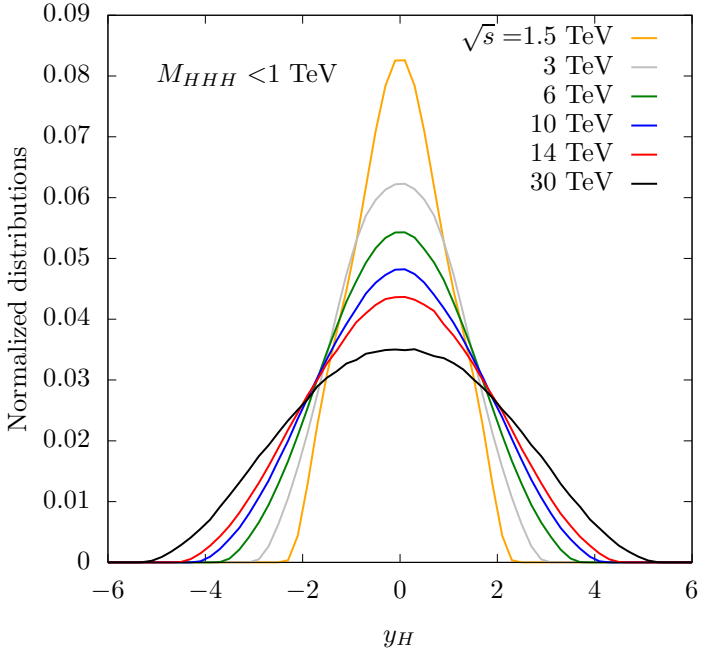
Remark on detector acceptance (1)



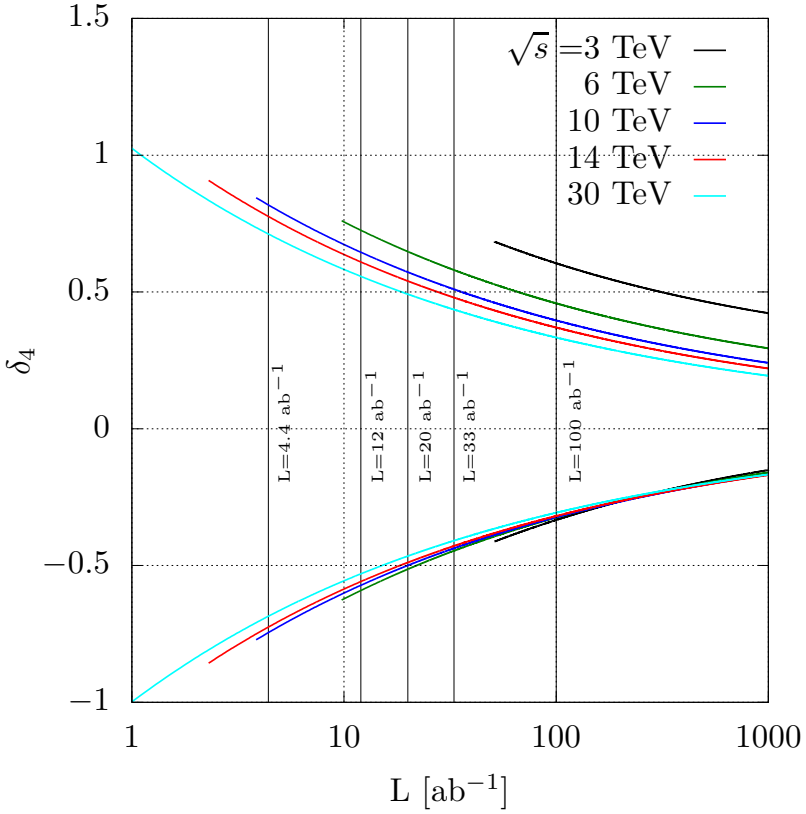
The detector must be shielded from the beam radiation

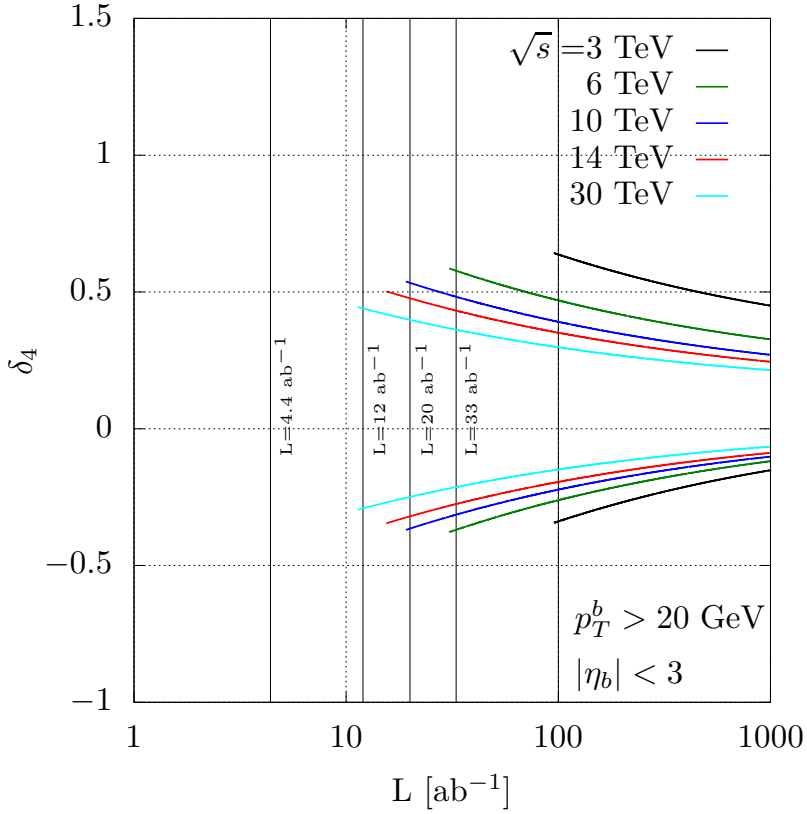
- 5-10 degrees blind spot in the forward region for $\sqrt{s} = 3$ TeV
- angle could be reduced at higher energies





Remark on detector acceptance (2)





- only geometric acceptance considered (no BR applied)
- sensitivity increases because the SM production is forward, the BSM central

Conclusions (1)

- we studied the sensitivity of the muon collider to the Higgs quartic coupling by considering the process $\mu^+\mu^- \to HHH\nu\overline{\nu}$
- no background was considered
- (almost) no optimization based on kinematics was performed
- the sensitivity increases with $\sqrt{s}$ and/or the luminosity

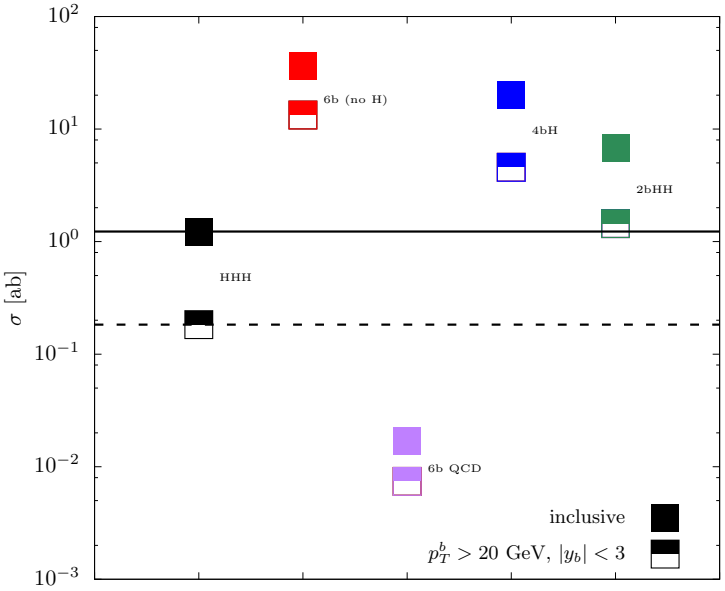
$\sqrt{s}$ [TeV]	L [ab^{-1}]	δ_4 (arbitrary δ_3)	δ_4 ($\delta_3 = 0$)
6	12	[-1,1.7]	[-0.45,0.8]
10	20	[-0.7,1.55]	[-0.4,0.7]
14	33	[-0.55,1.4]	[-0.35,0.6]
30	100	[-0.35,1.2]	[-0.2,0.5]

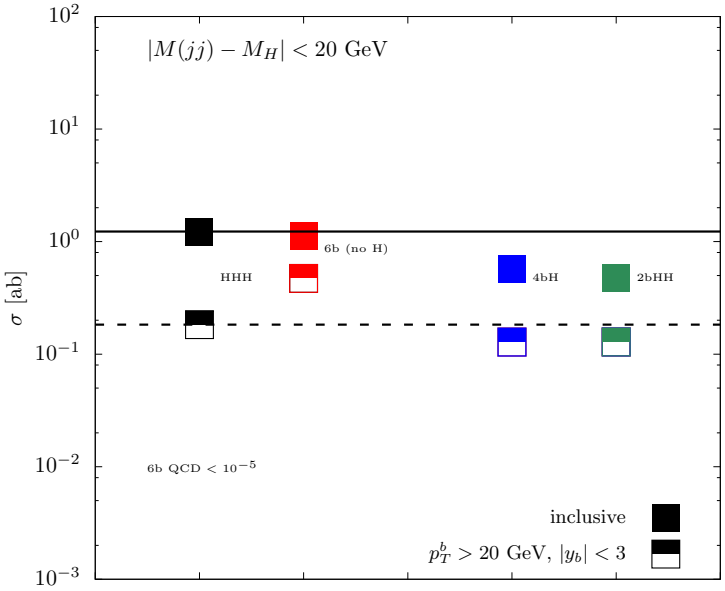
- under (reasonable) assumptions on the energy and the luminosity, the muon collider can do a pretty good job in constraining the quartic Higgs coupling

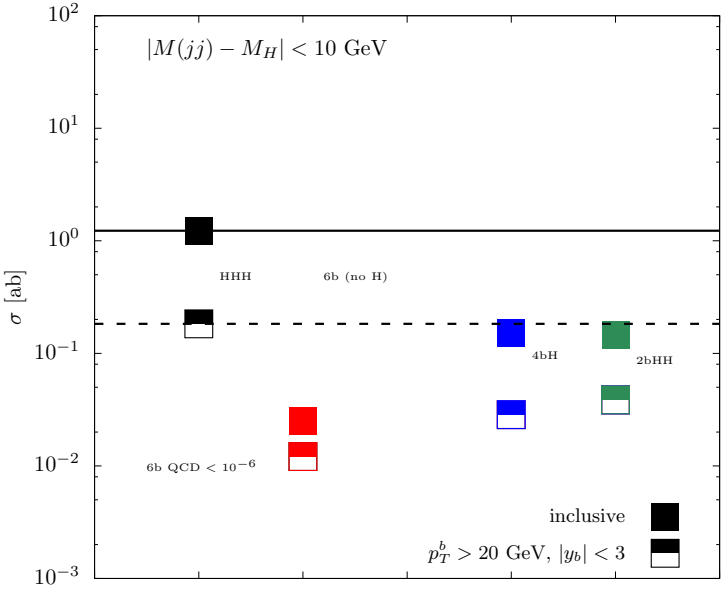
First estimate of possible backgrounds

preliminary, $\sqrt{s} = 14$ TeV only

1st strategy: search for 6b quarks

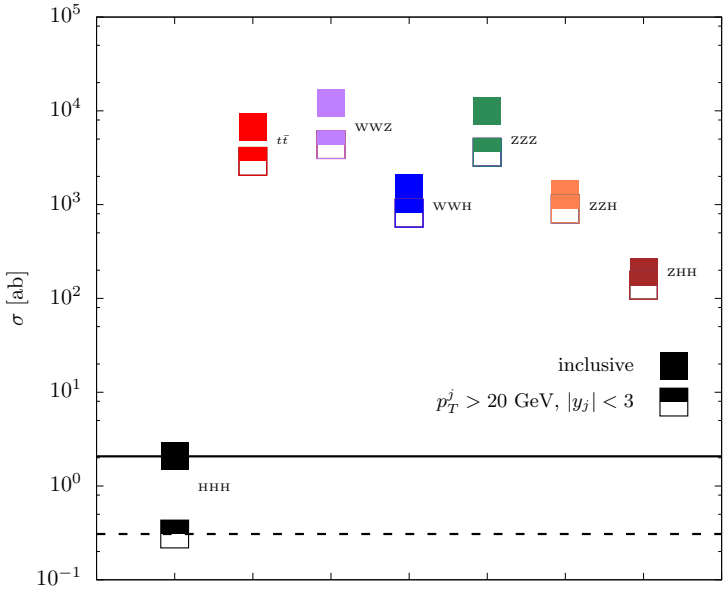


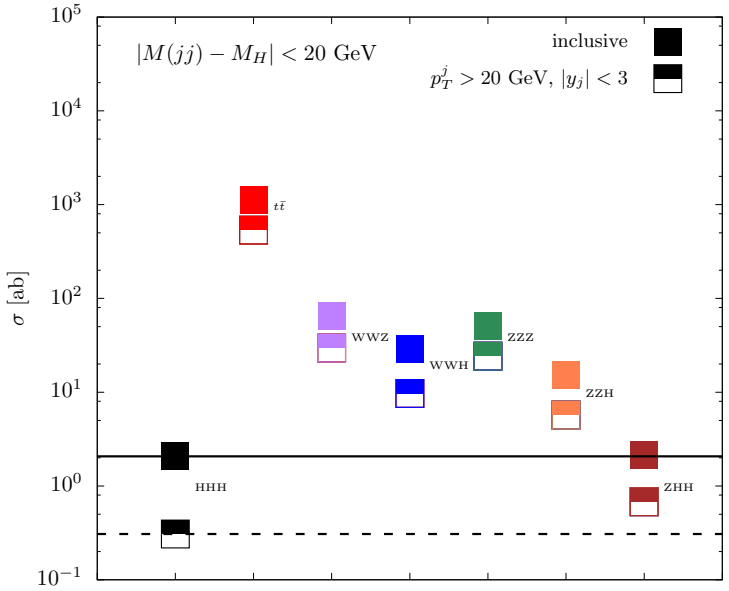


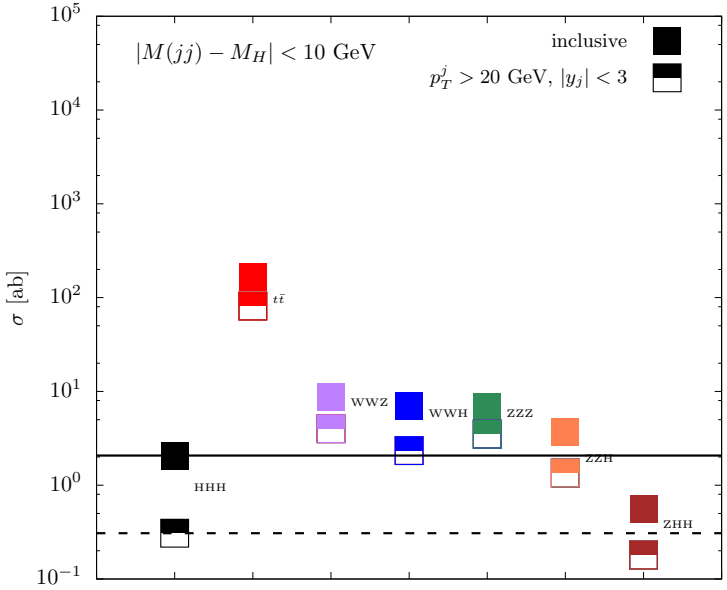


- only 1/5 of the signal survives after selecting $H \to bb$
- Higgs reconstruction needed (at least with 20 GeV resolution, better with 10 GeV res)
- full simulation for 6b quarks (no Higgses) computed with a modified version of `ALPGEN`

2nd strategy: look for 6 jets (no b-tagging)

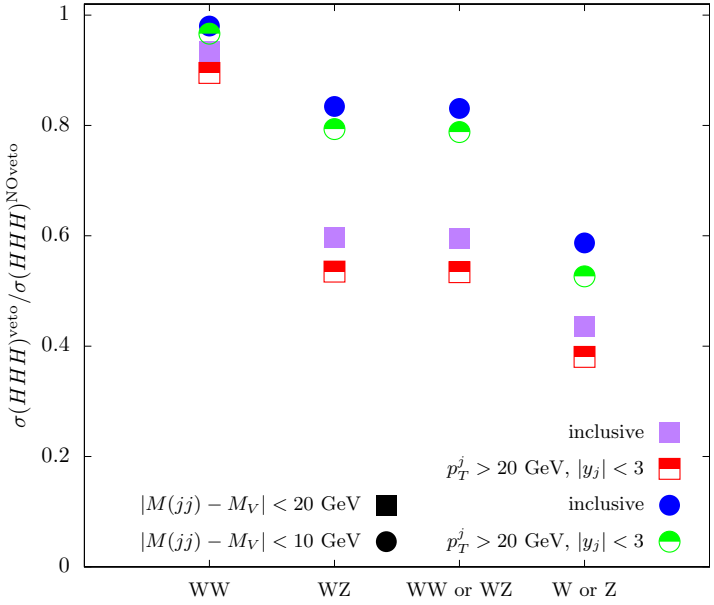


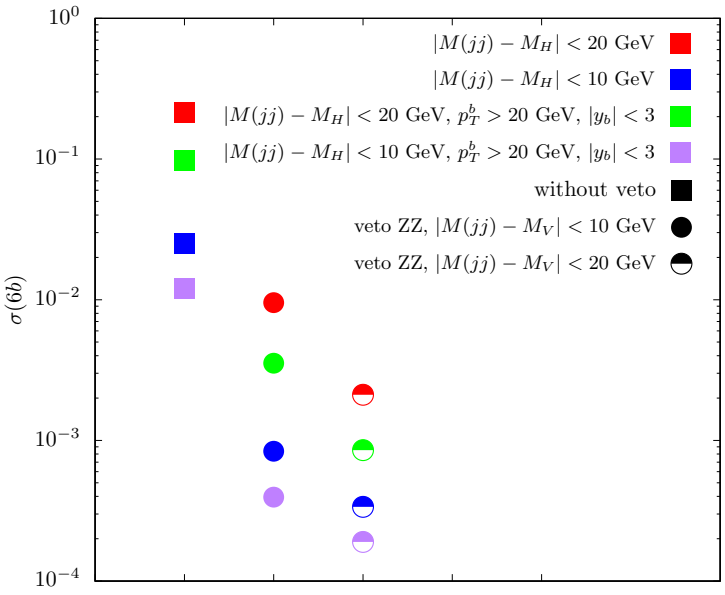




- only 1/3 of the signal survives after selecting $H \to jj$
- Higgs reconstruction needed (at least with 20 GeV resolution, better with 10 GeV res)
- full simulation for 6j still missing, however the largest contribution should be from 6q (gluons have larger combinatorics, but pure QCD is suppressed)

3rd strategy: veto W and/or Z bosons





- vetoing Ws and/or Zs does not kill too much signal (10 GeV resolution)
- vetoing Ws and/or Zs kills the corresponding background
- best background suppression strategy should involve vetoes on Ws and/or Zs

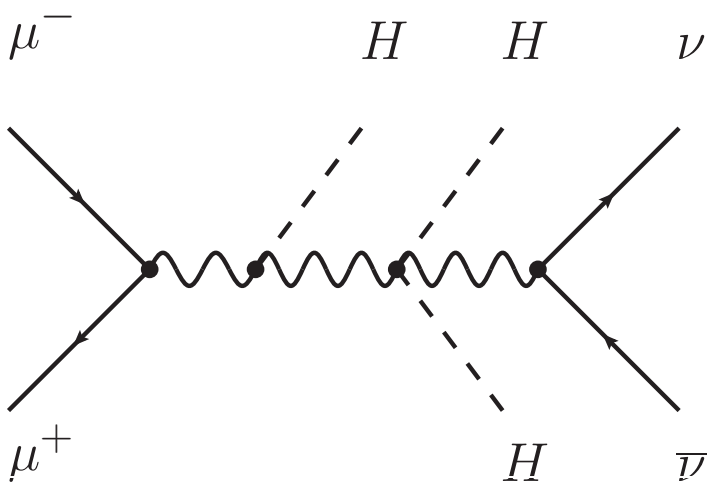
Conclusions (2)

background processes for $\mu^+\mu^- \to HHH\nu\bar{\nu}$

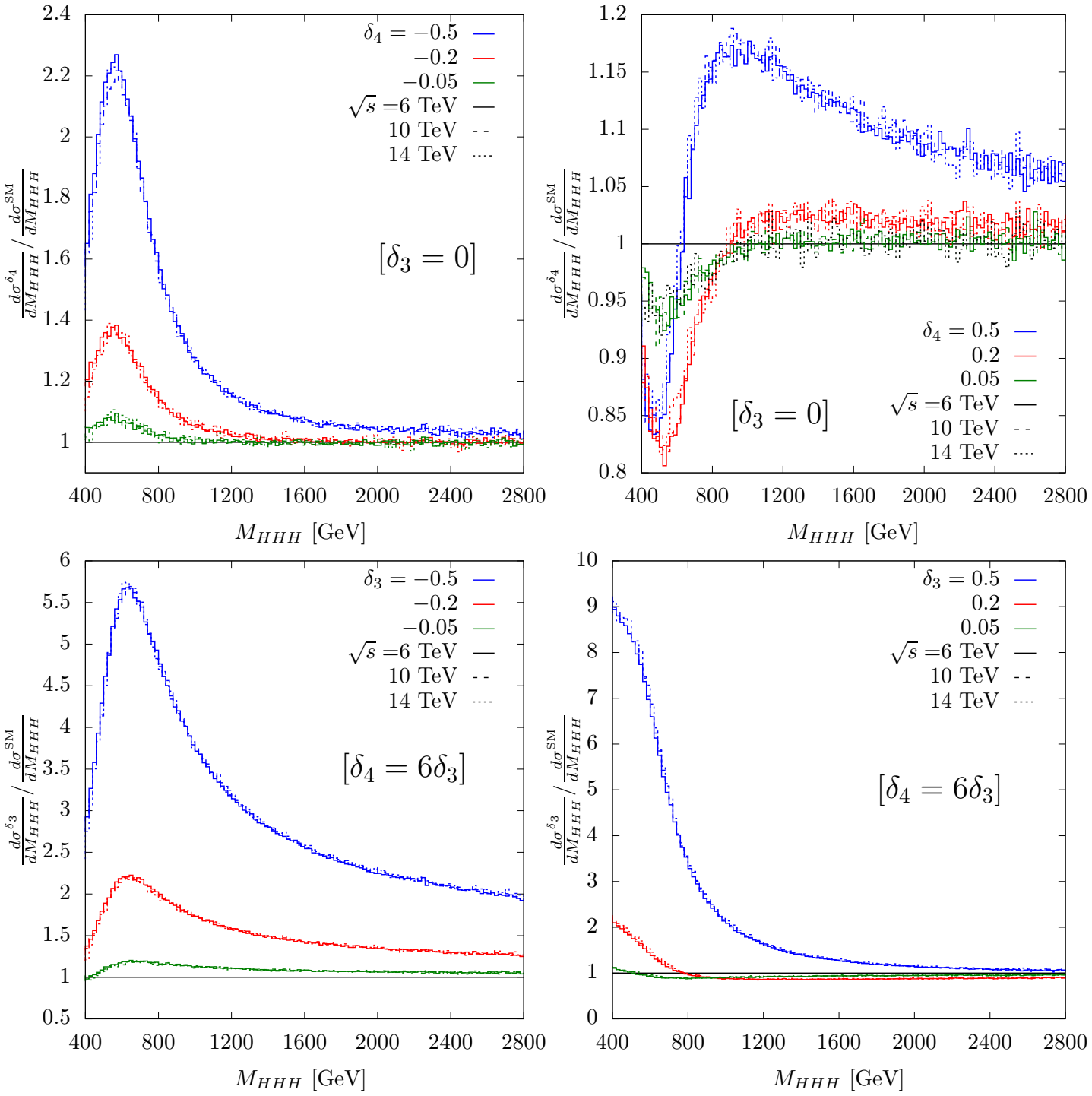
- are definitively there
- might be large
- they can be largely suppressed with some combination of:
  - b-tagging (soft?)
  - Higgs reconstruction
  - Z/W vetoes
- optimal suppression strategy seems to require good resolution in dijet invariant mass reconstruction

Backup slides

Details of the calculations



- H produced on shell
- $H \to b\bar{b}$ (on-shell) decays added at the LHE level
- $\Gamma_W = \Gamma_Z = \Gamma_H = 0$ to avoid issues with gauge invariance
- technical cut $M(\nu\overline{\nu}) > 150$ GeV
- σ and $d\sigma$ computed with `WHIZARD` at LO
- all results cross-checked with `MadGraph` and an independent calculation by X. Zhao


$\delta_4 = -0.5$
-0.2
-0.05
$\sqrt{s}$ =6 TeV
10 TeV
14 TeV
$[\delta_3 = 0]$
$\frac{d\sigma^{\delta_4}}{dM_{HHH}} / \frac{d\sigma^{SM}}{dM_{HHH}}$
M_{HHH} [GeV]
$\delta_4 = 0.5$
0.2
0.05
$\sqrt{s}$ =6 TeV
10 TeV
14 TeV
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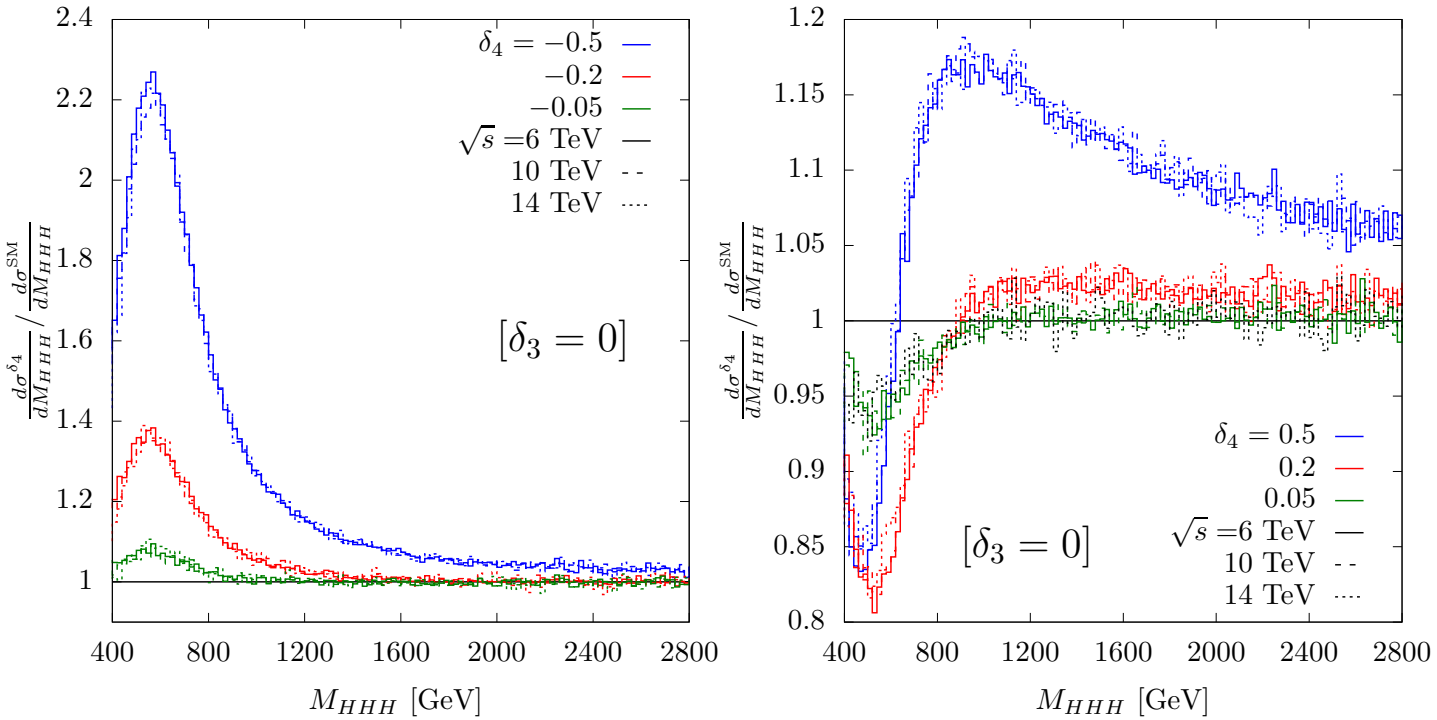
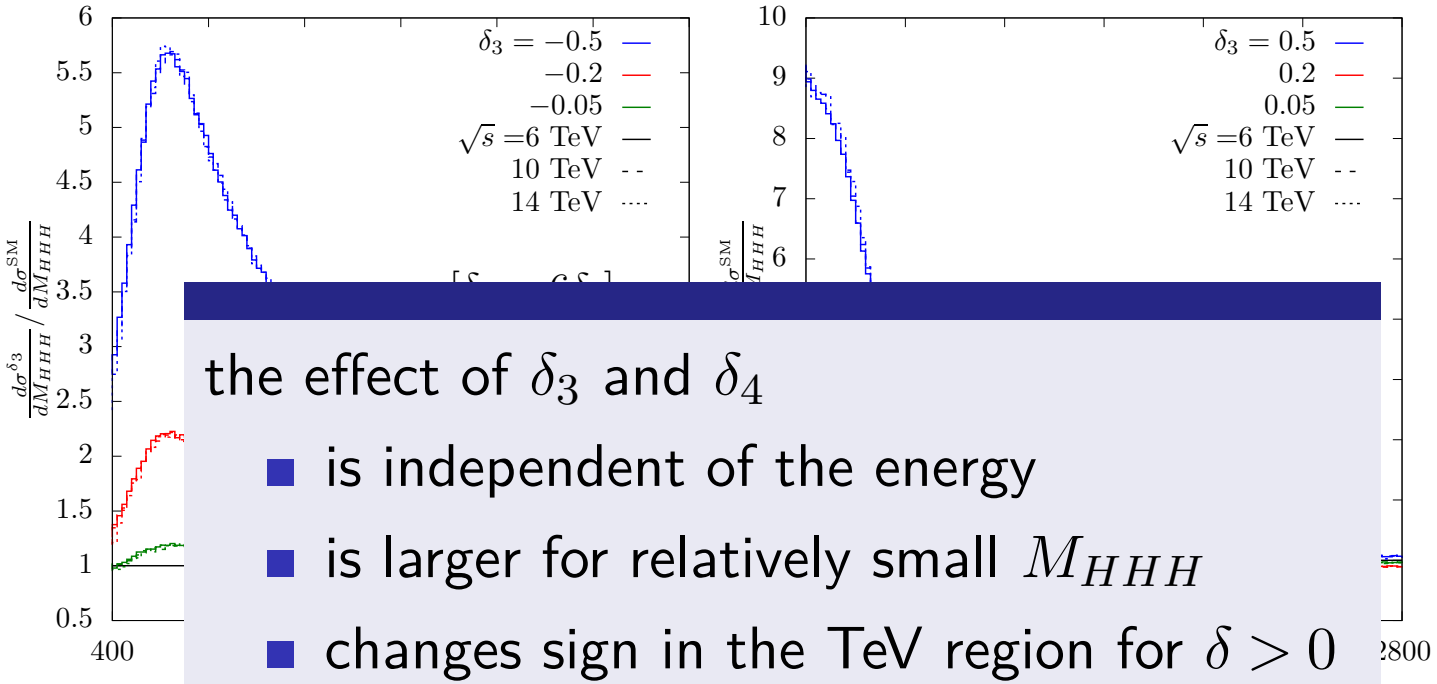
M_{HHH} [GeV]

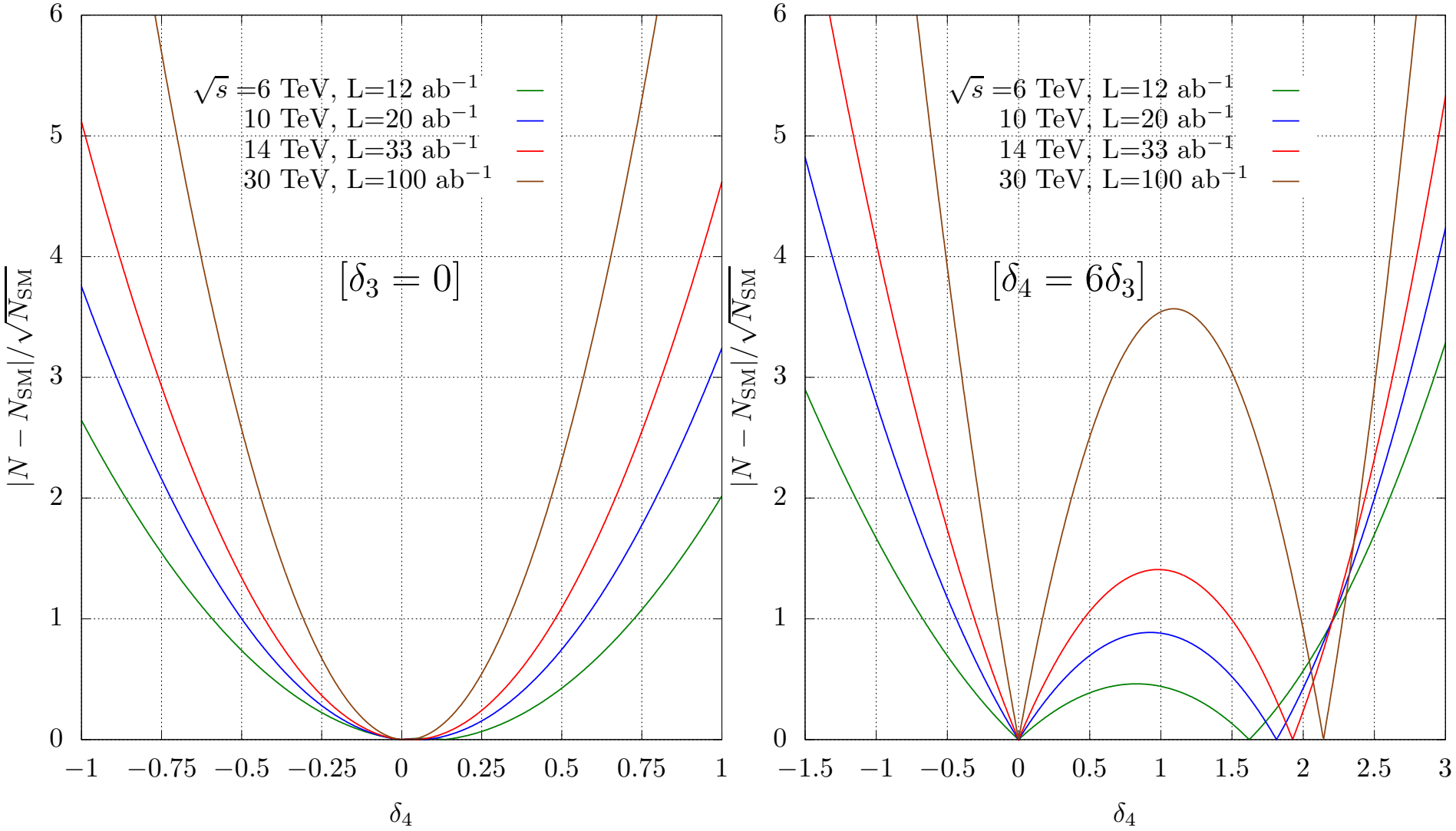




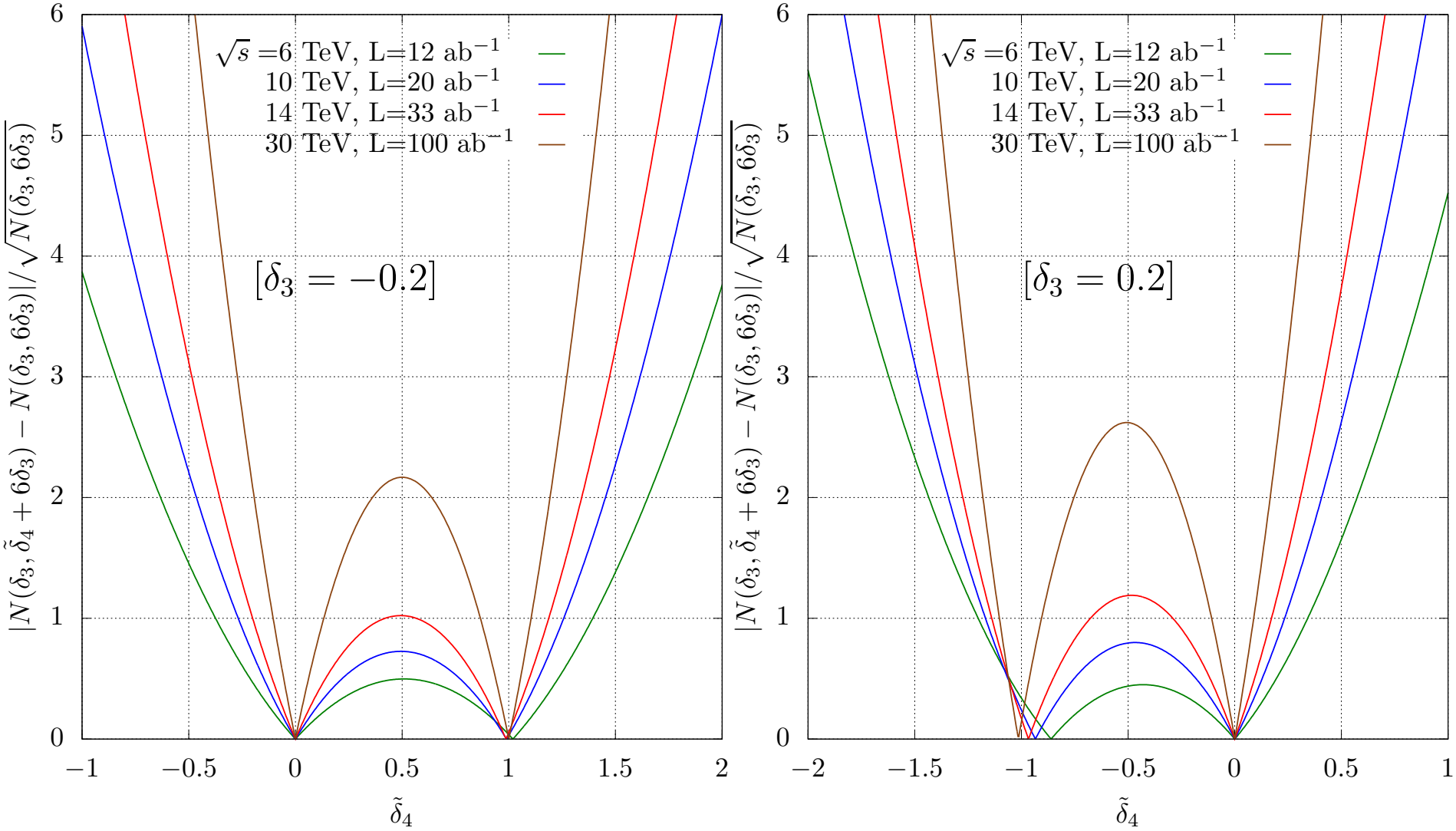
the effect of δ_3 and δ_4

- is independent of the energy
- is larger for relatively small M_{HHH}
- changes sign in the TeV region for $\delta > 0$

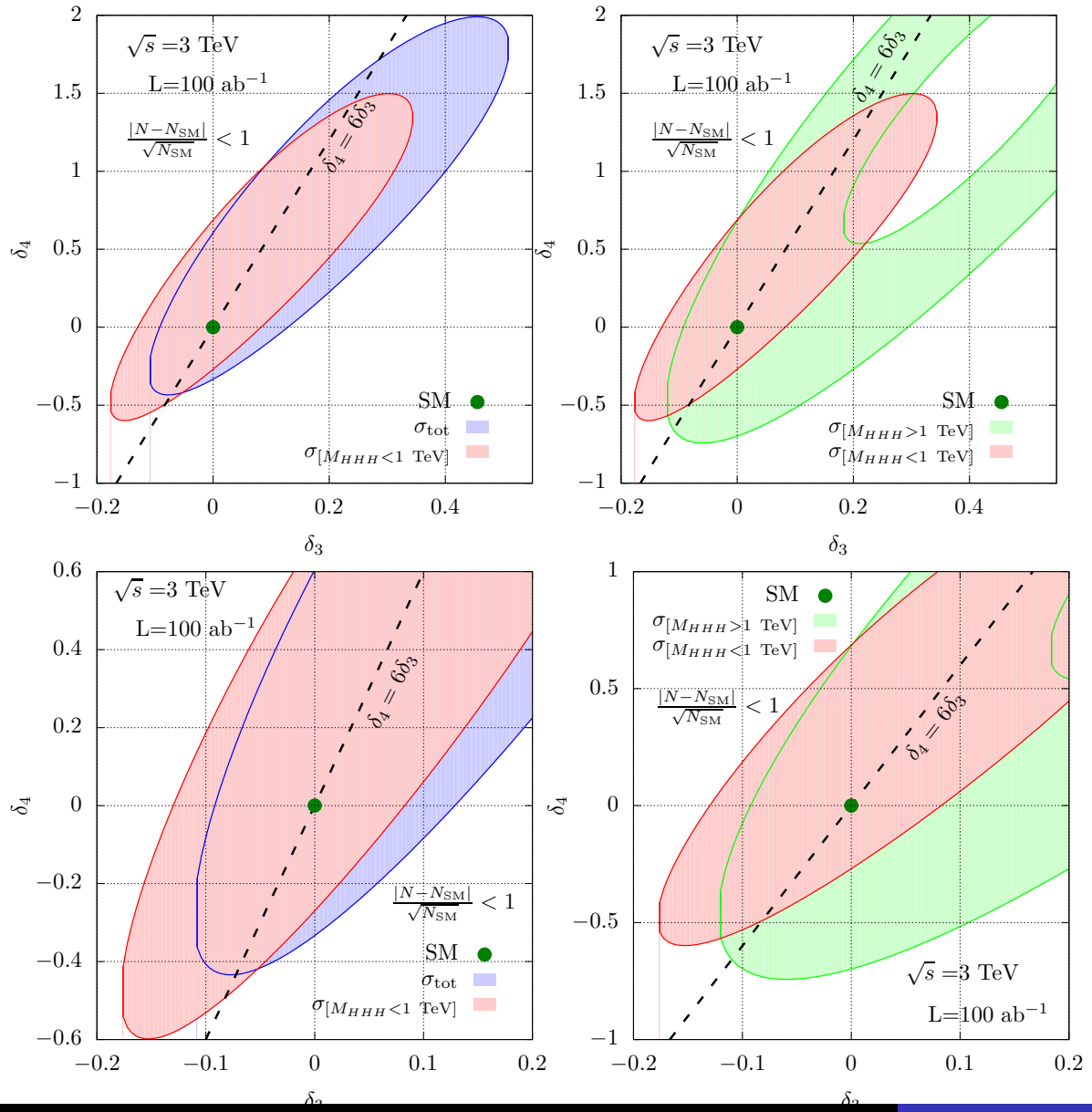
Sensitivity to δ_3 and δ_4



Sensitivity to $\tilde{\delta}_4$ (deviation from SMEFT)



Sensitivity to δ_3 and δ_4 ($\sqrt{s} = 3$ TeV, $L = 100$ ab^{-1})



- no cuts
- $M_{HHH} < 1$ TeV
- $M_{HHH} > 1$ TeV

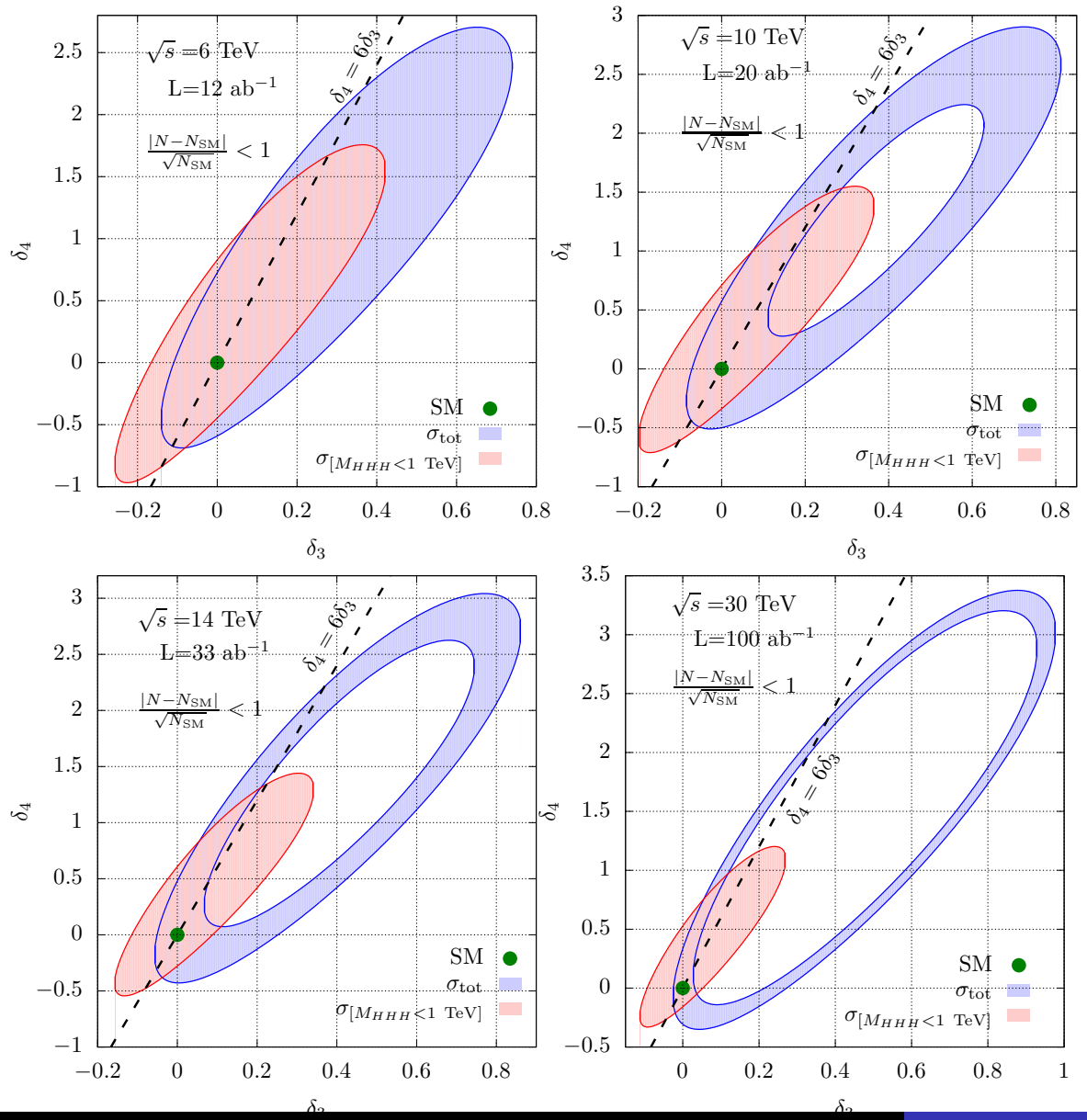
$$\delta_4 \sim [-0.6, 1.5]$$

if $\delta_3 = 0$

$$\delta_4 \sim [-0.3, 0.65]$$

Using 20 times the expected luminosity!

Sensitivity to δ_3 and δ_4 (arbitrary δ_3)



- no cuts
- $M_{HHH} < 1$ TeV

$$\delta_3 = 0$$

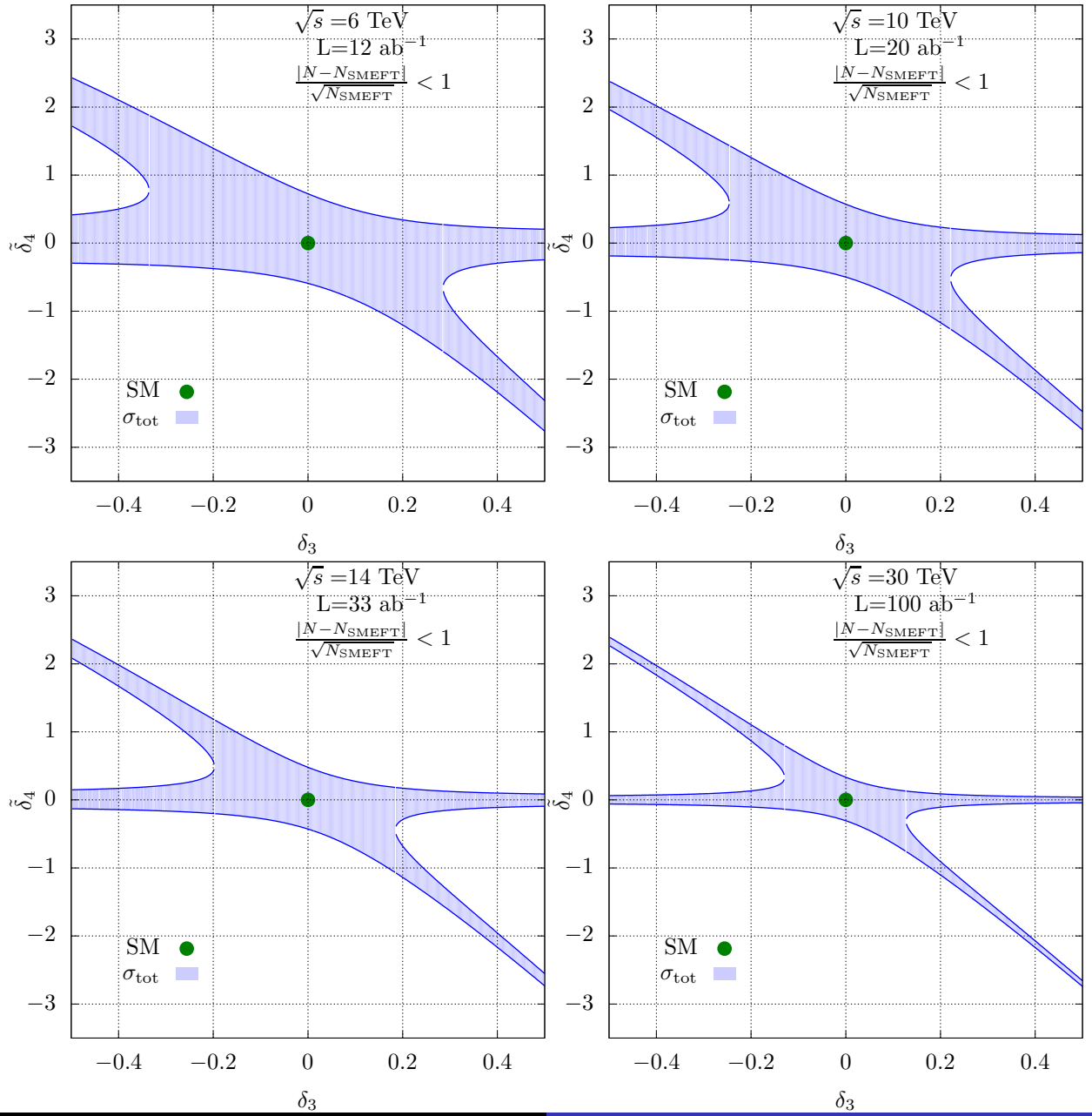
6 TeV $\delta_4 \sim [-0.1, 1.7]$

10 TeV $\delta_4 \sim [-0.7, 1.55]$

14 TeV $\delta_4 \sim [-0.55, 1.4]$

30 TeV $\delta_4 \sim [-0.35, 1.2]$

Sensitivity to δ_3 and δ_4: comments

- stronger constraints on negative δs
- constraints on positive δs improve with the cut $M_{HHH} < 1$ TeV (provided that the cross section after the cut is large enough)
- the bounds improve at large $\sqrt{s}$ because the cross section increases
- the most interesting region is $\delta_3 \sim 0$, as bounds on δ_3 can be obtained form other processes (i.e. $\mu^+\mu^- \to HH\nu\overline{\nu}$). It is reasonable to assume that such bounds will be competitive or stronger than the ones form linear colliders
- if $\delta_3 \neq 0$, one can constrain possible deviations from the SMEFT expectation for δ_4: $\tilde{\delta}_4 = \delta_4 - 6\delta_3$

Sensitivity to $\tilde{\delta}_4$ (deviation from SMEFT)