

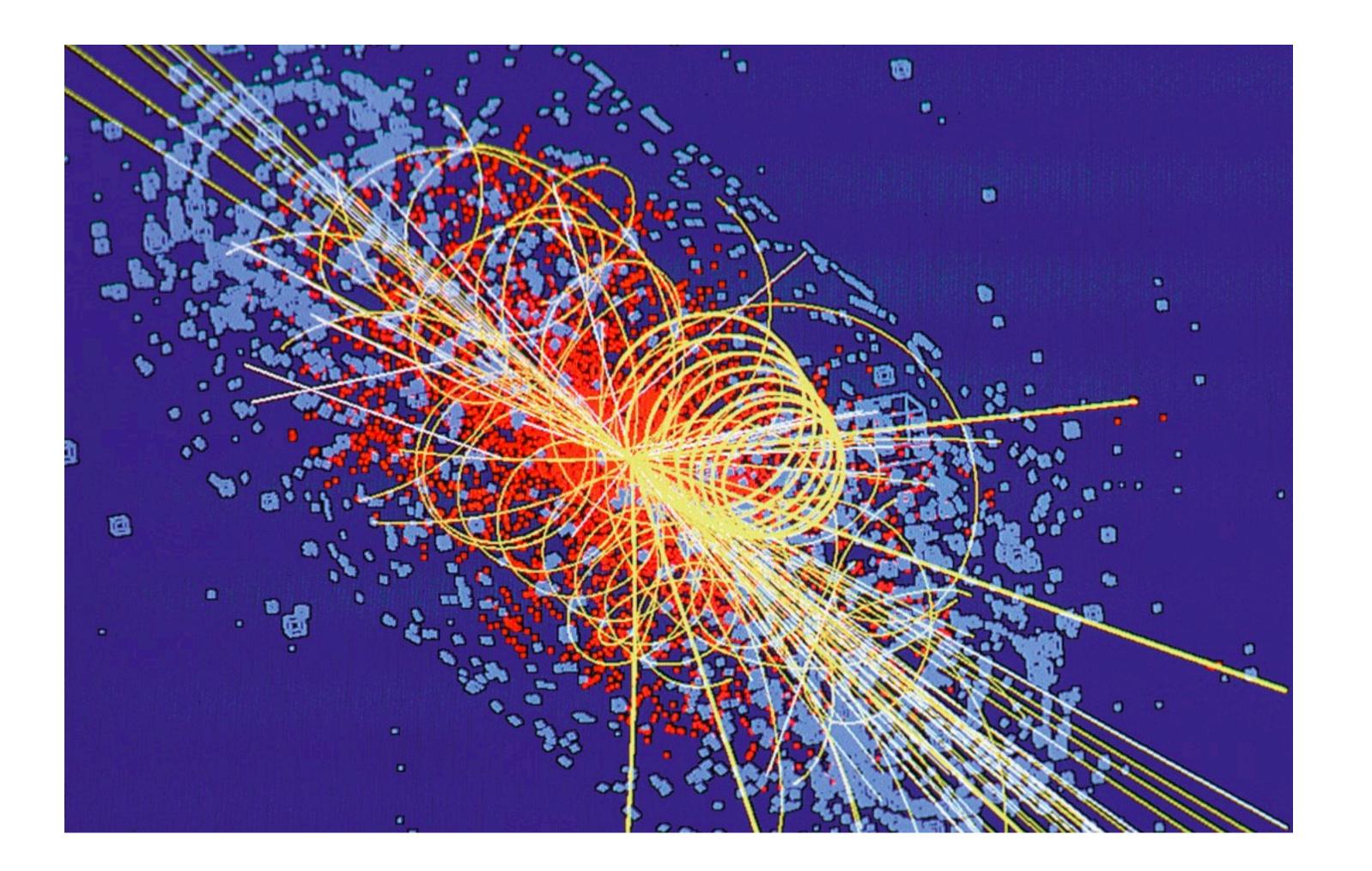


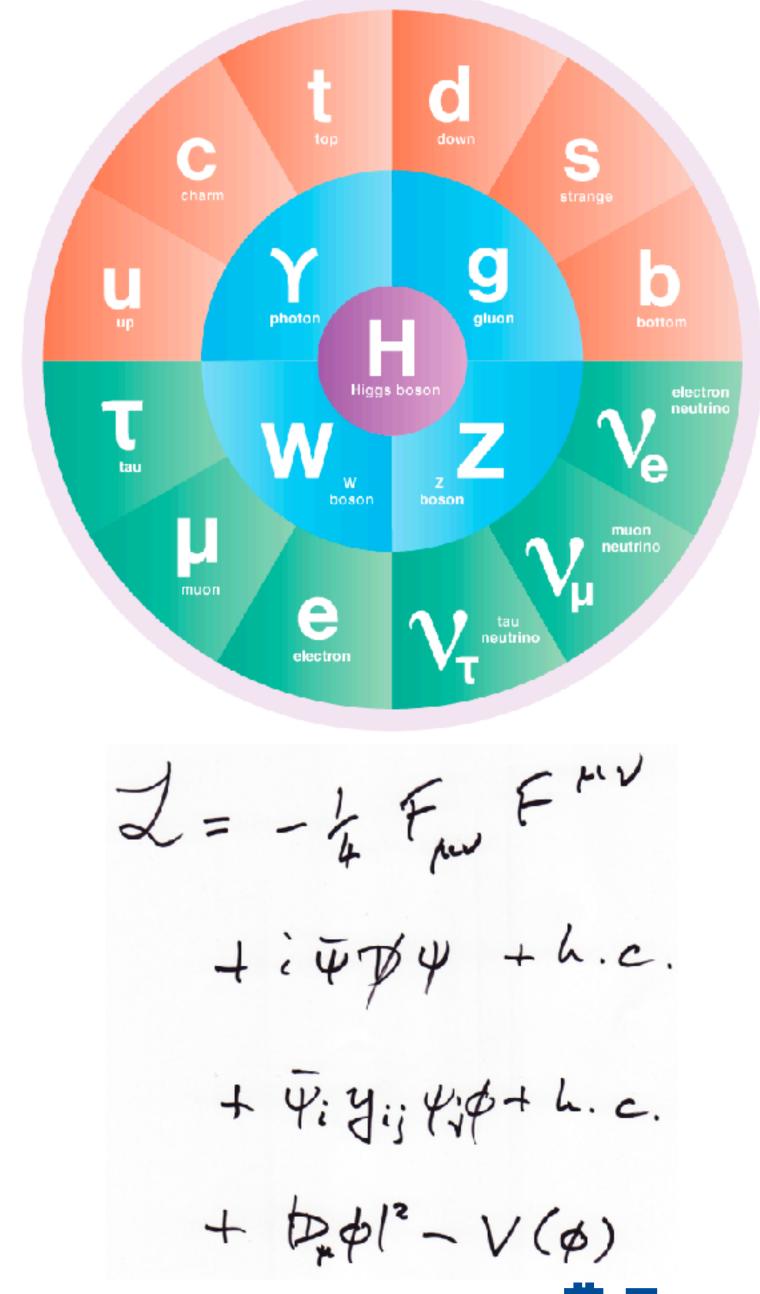


Introduction to Particle Physics

Kevin J. Kelly (kkelly12 at fnal dot gov) Summer Lecture Series, June 1st, 2021

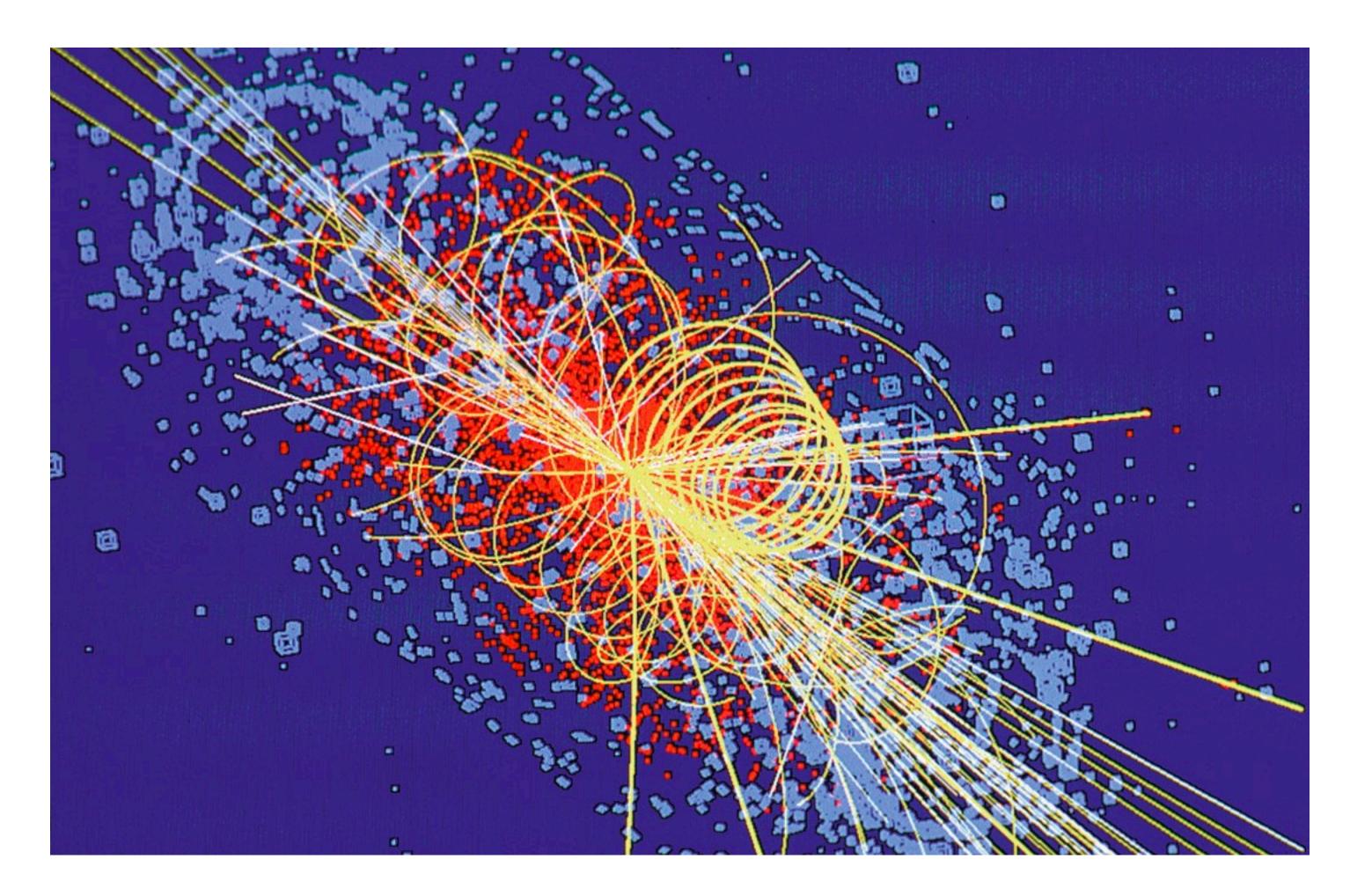
What is Particle Physics?



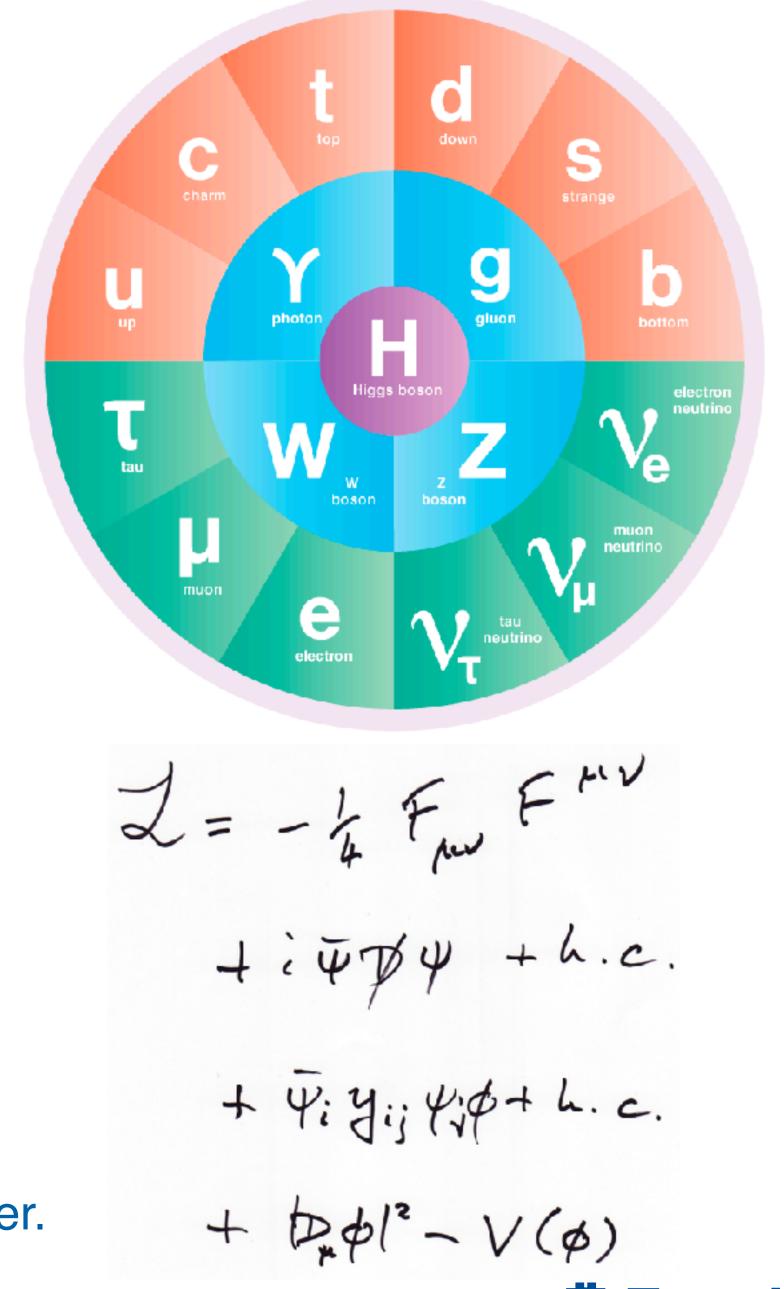




What is Particle Physics?



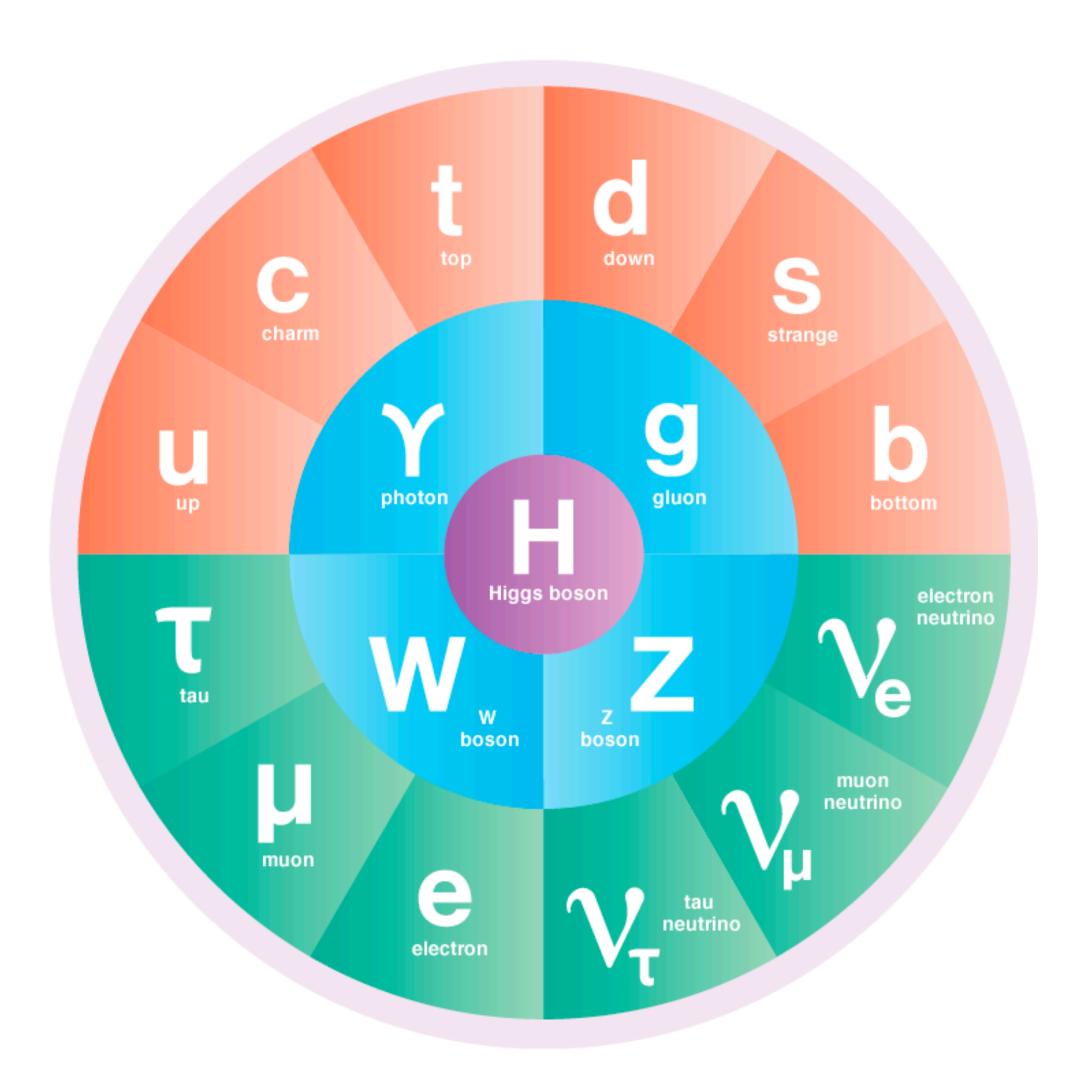
Broader applications than "just particles" — we'll get to those later.





A bit of history...

 Physicists today claim to know what "everything" is made up of — the Standard Model. How did we get here?

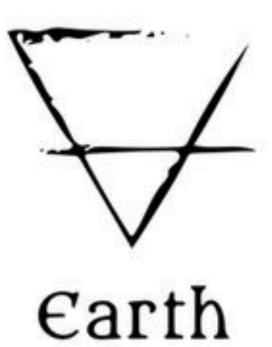








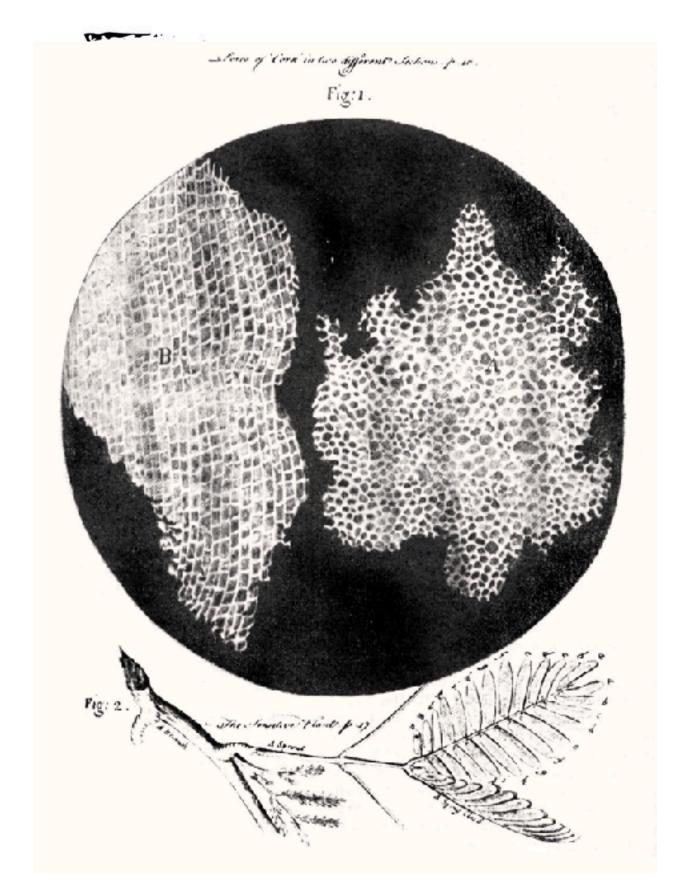




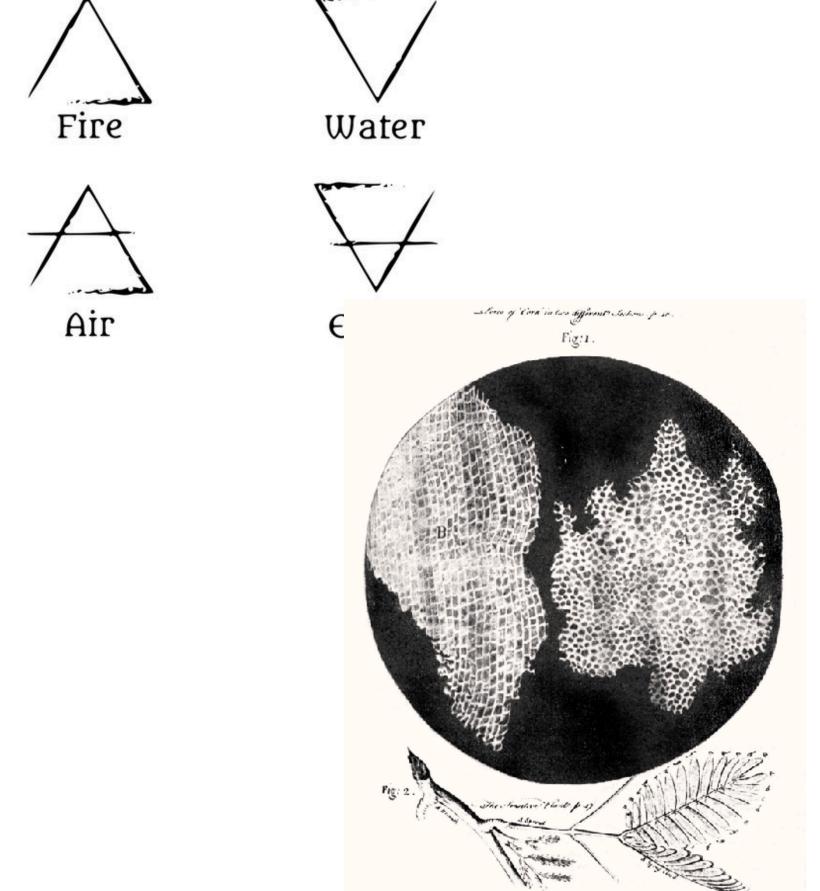


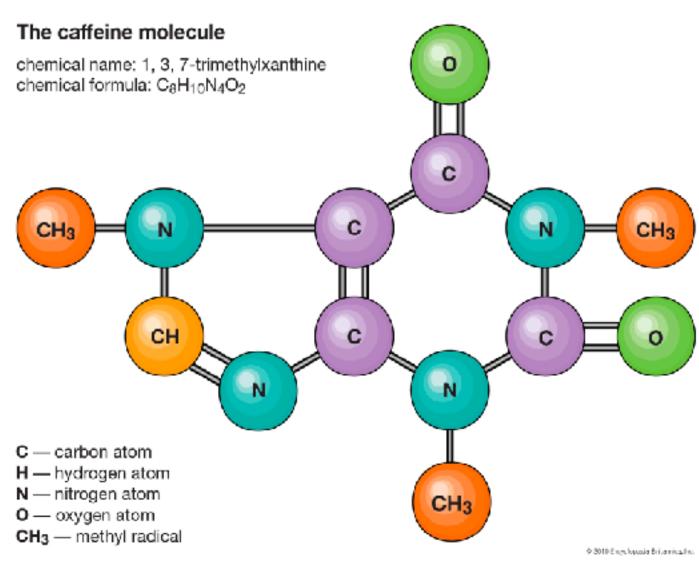












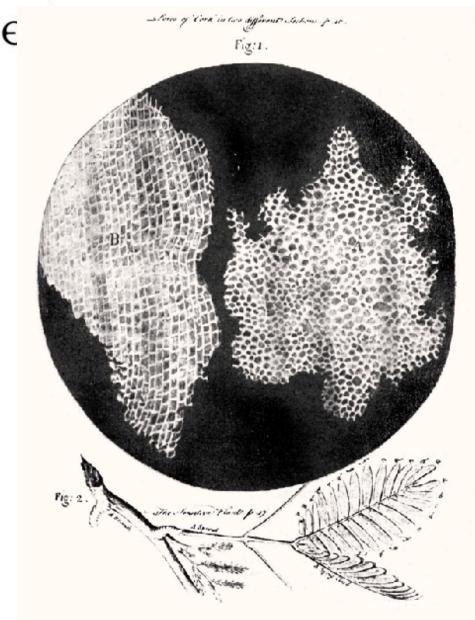


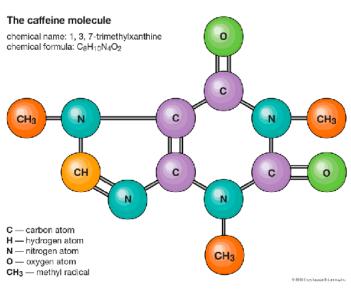


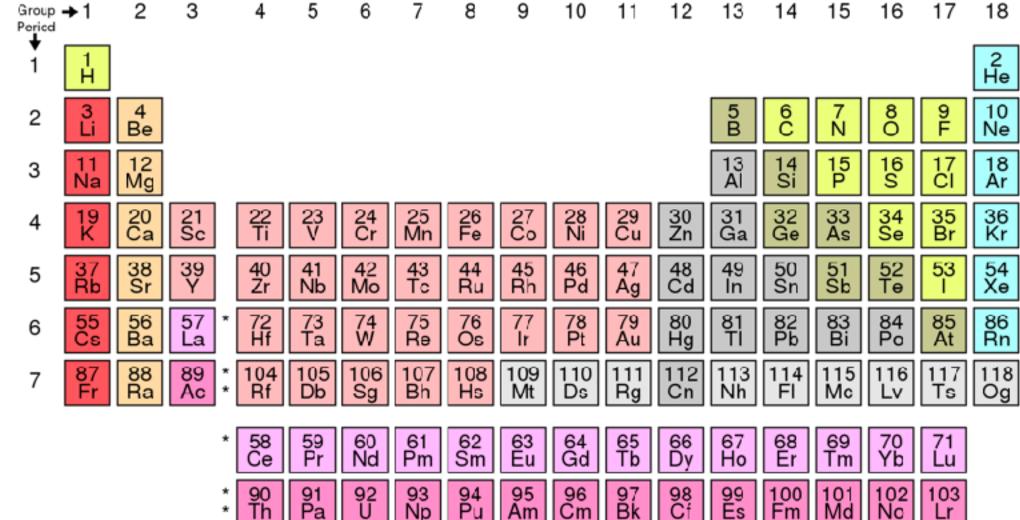












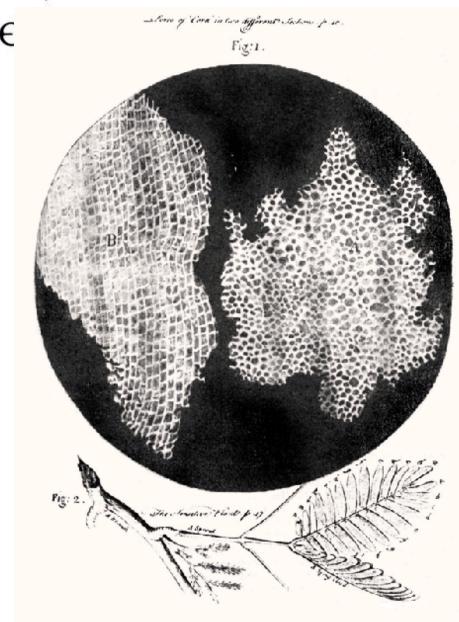


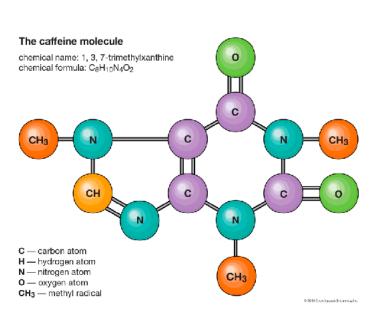


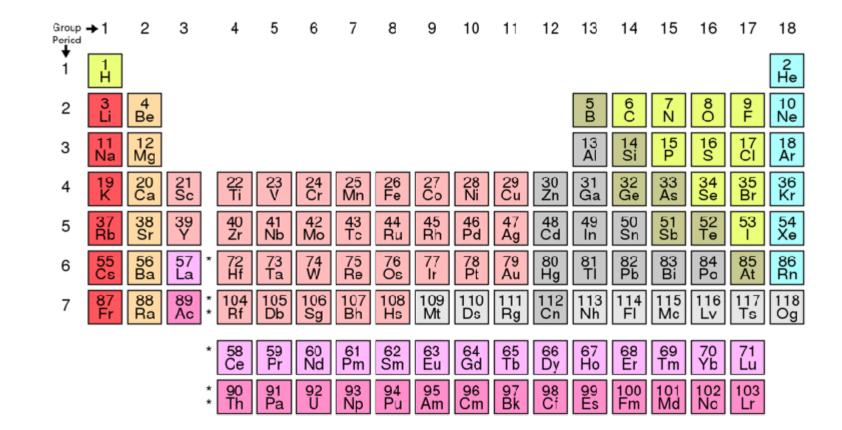






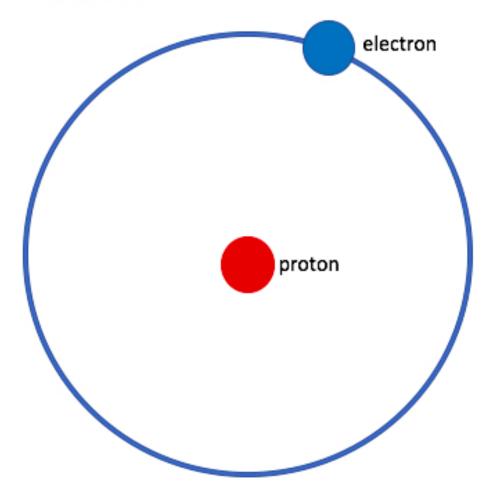






Hydrogen Atom

*Not to scale.



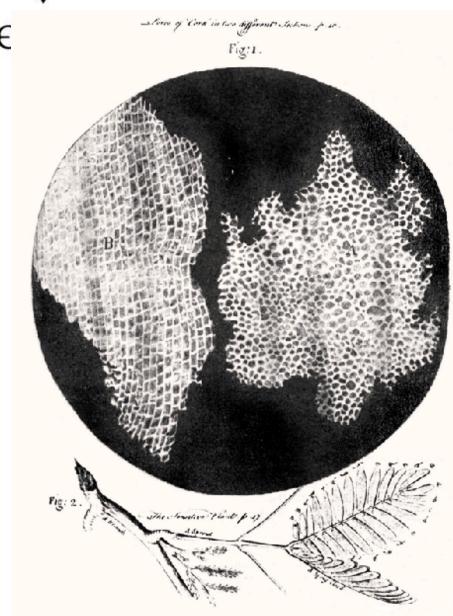


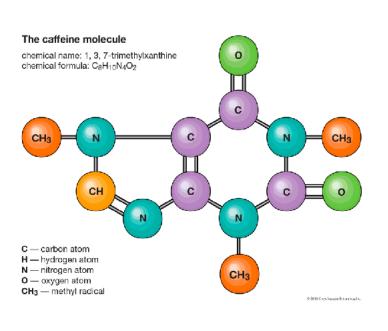


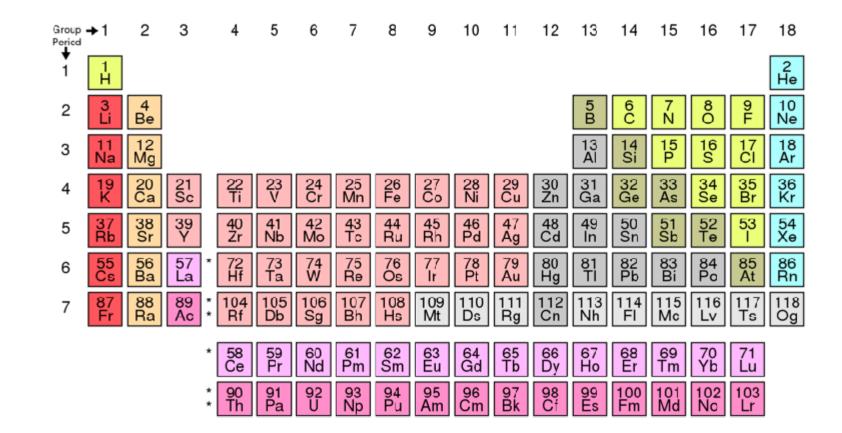


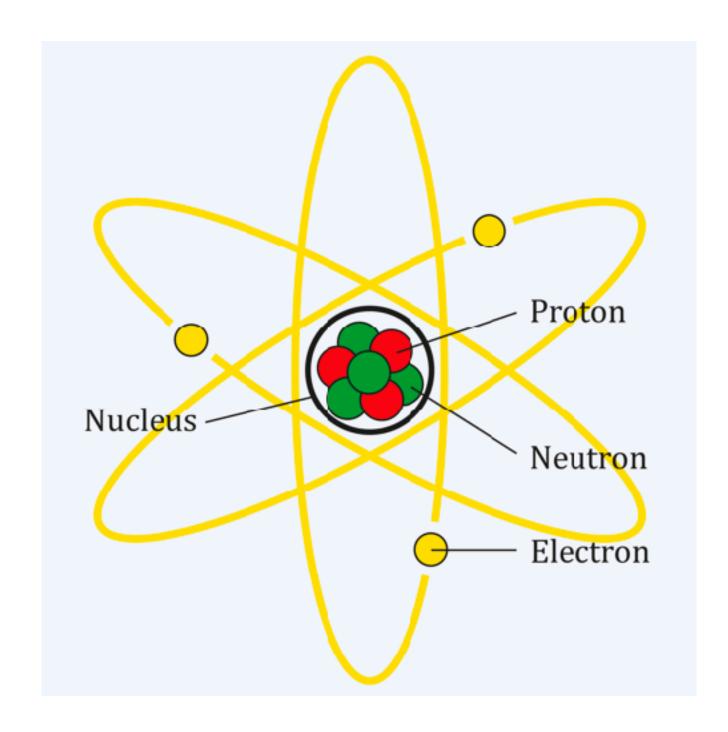




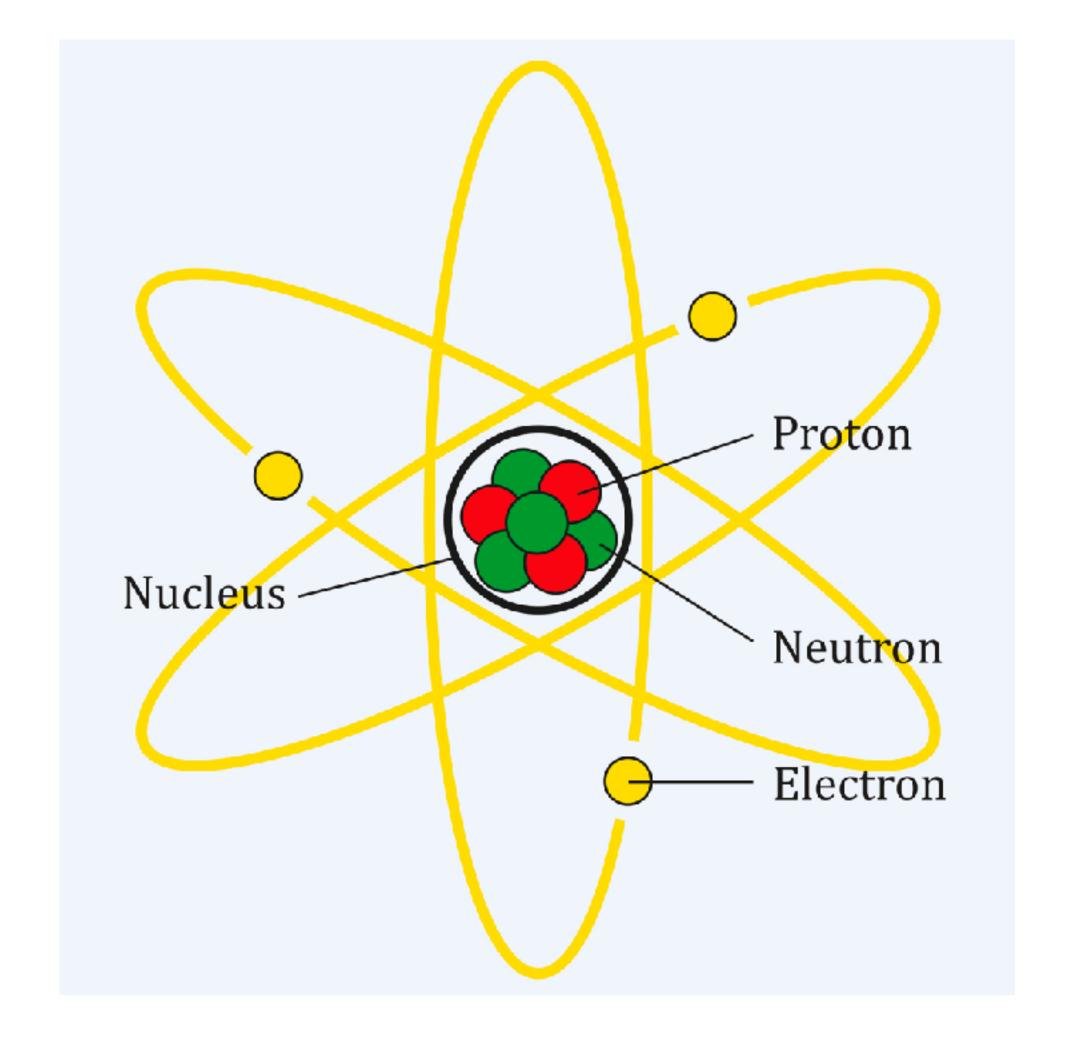




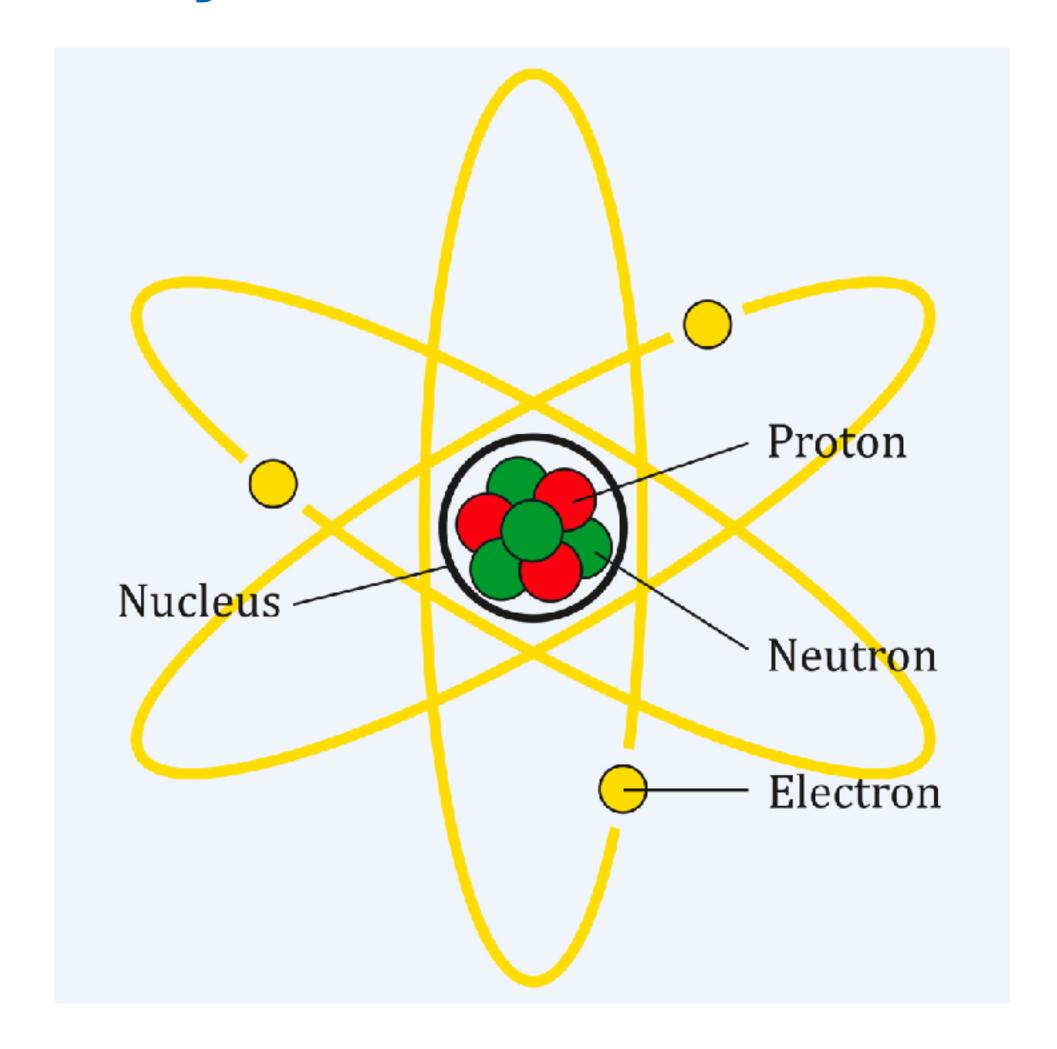






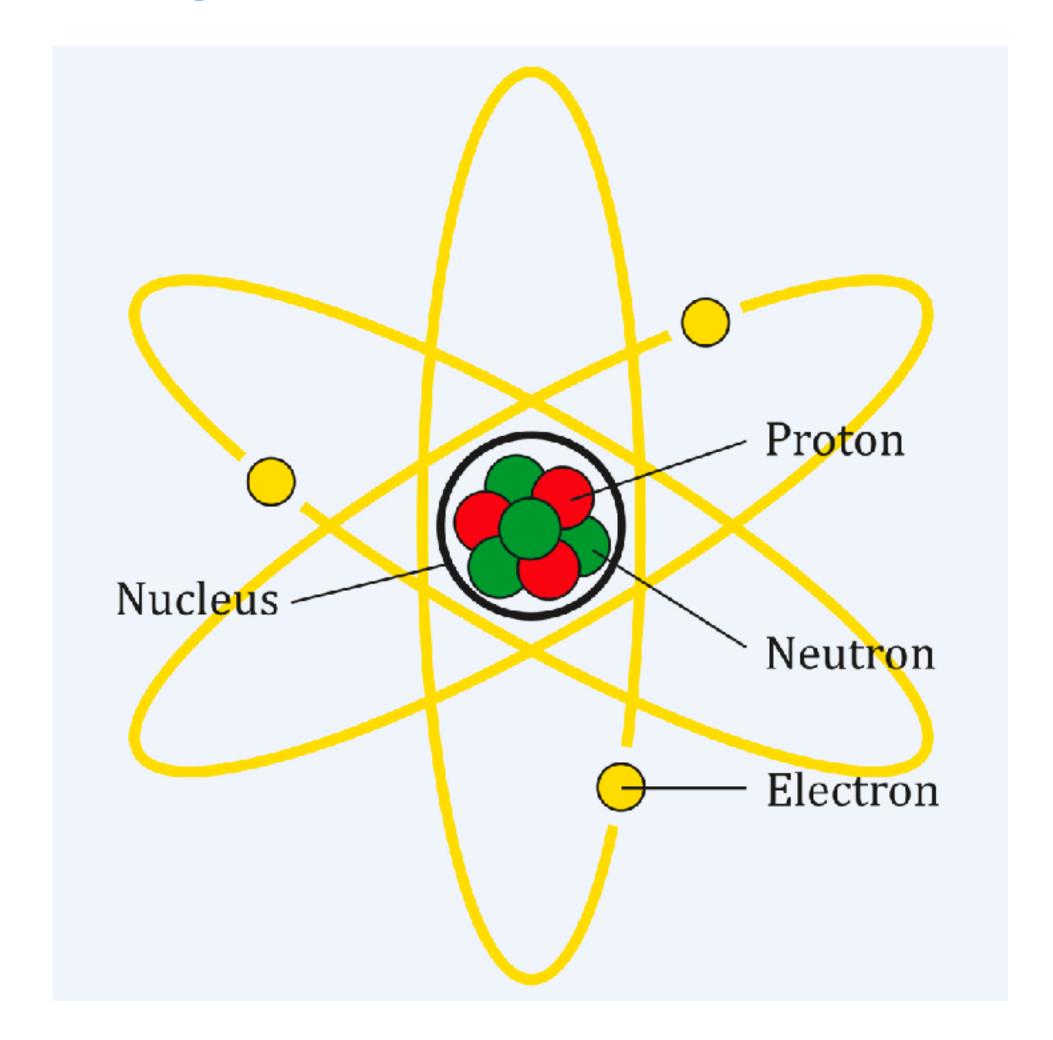






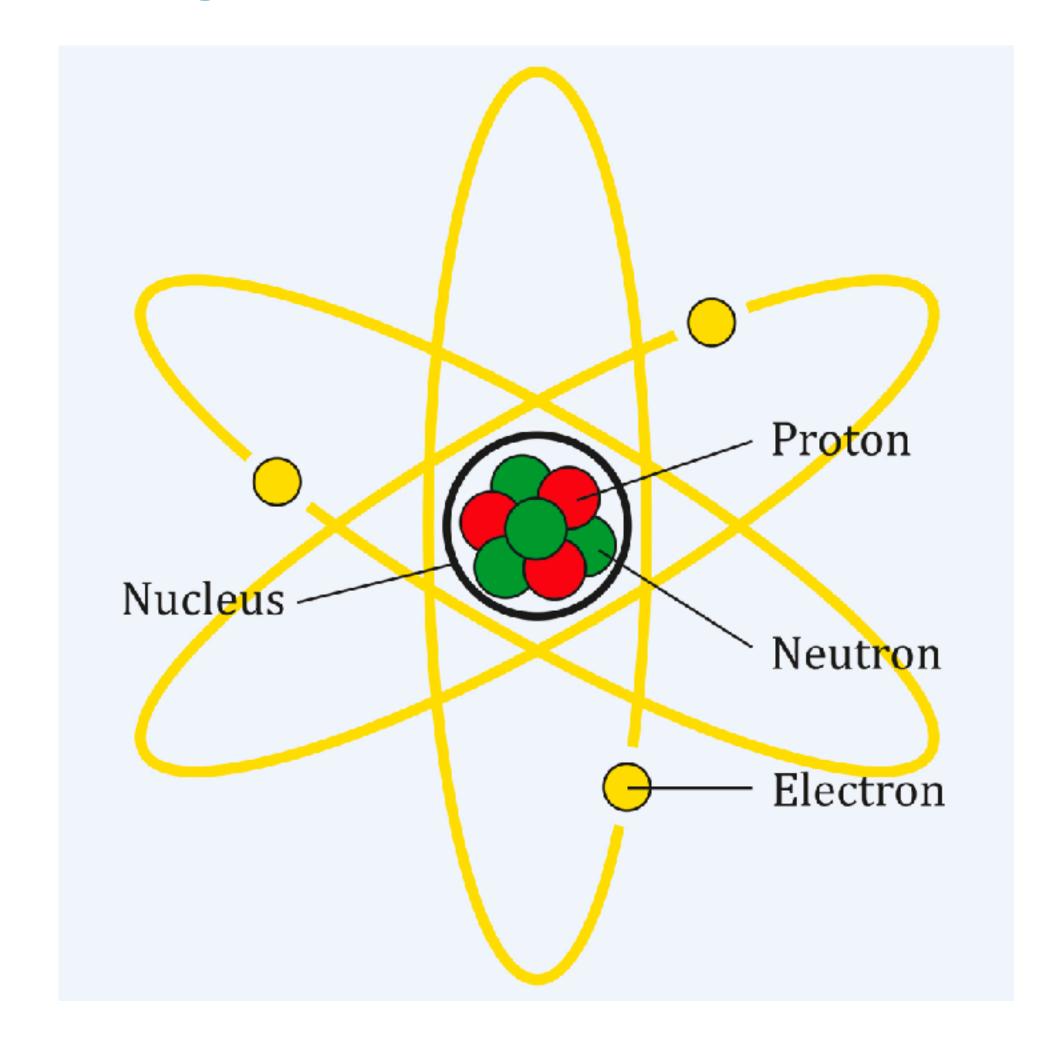
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- Protons and neutrons build up the nuclei of atoms, and electrons occupy orbitals around the nucleus.
- Different models of the atom were explored, leading to the development of quantum mechanics.



Let's put a pin in particles for now.

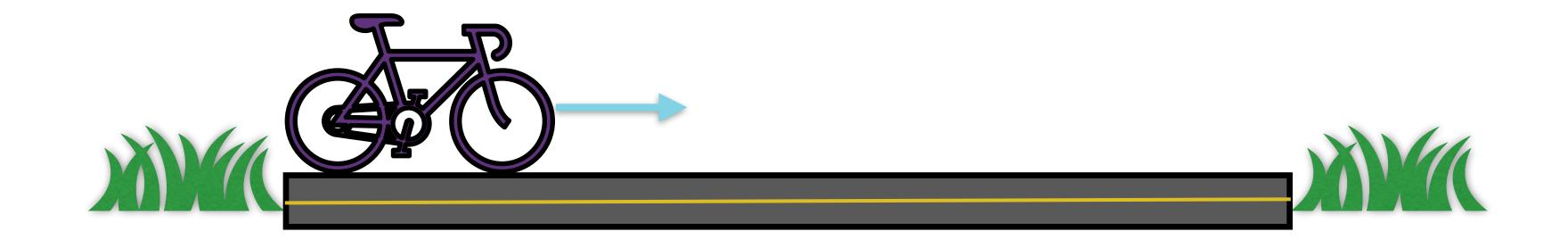


Let's put a pin in particles for now.

And pivot to every particle physicist's favorite tool: **symmetries**.

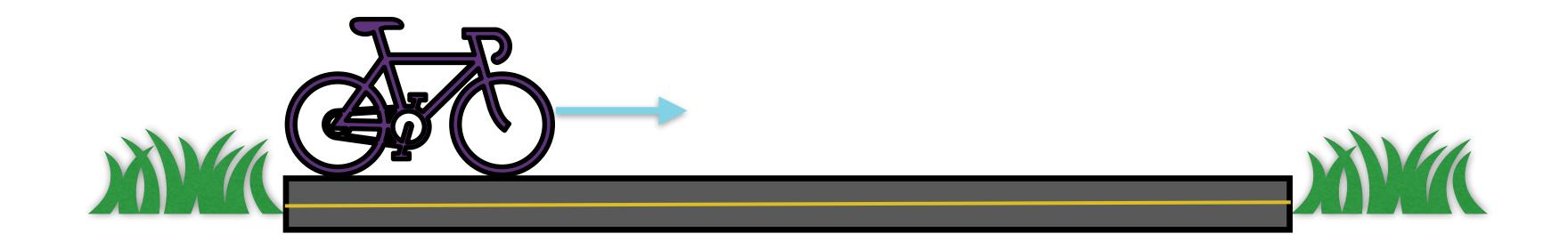


• Imagine riding a bike on a flat road...

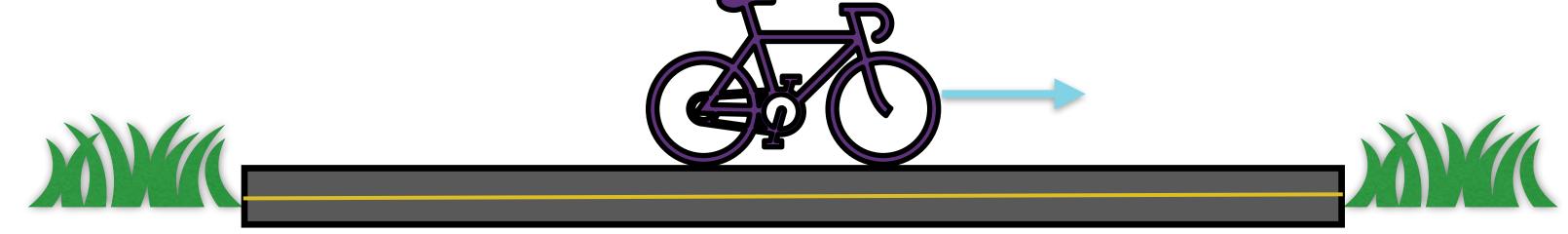




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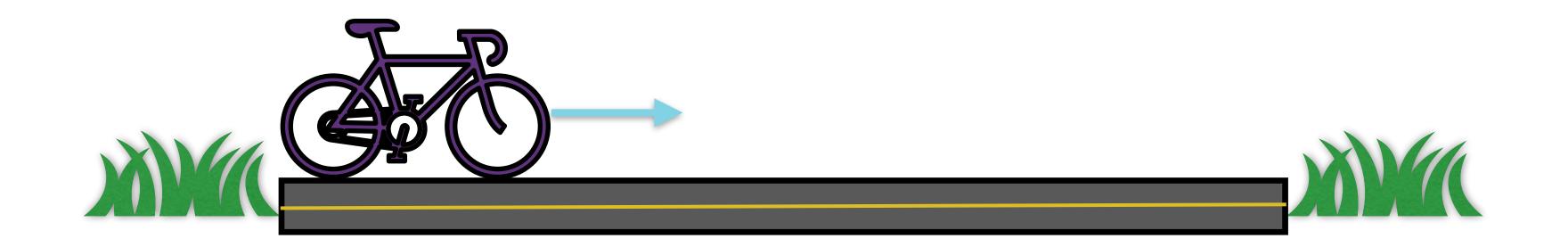


A short time later,





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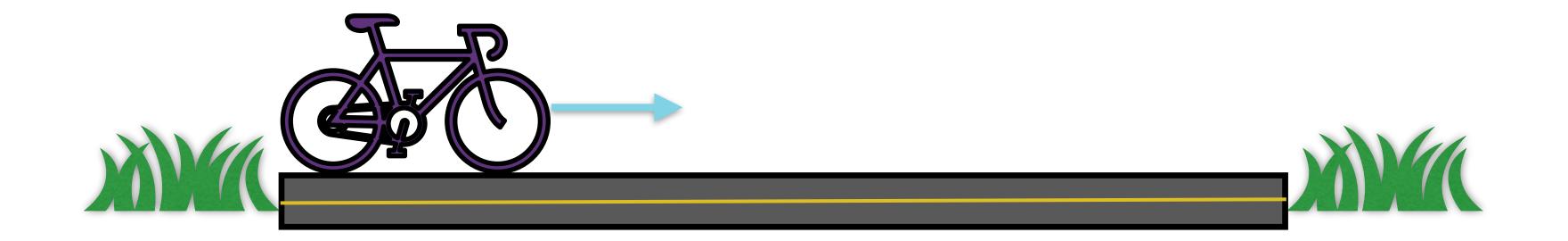


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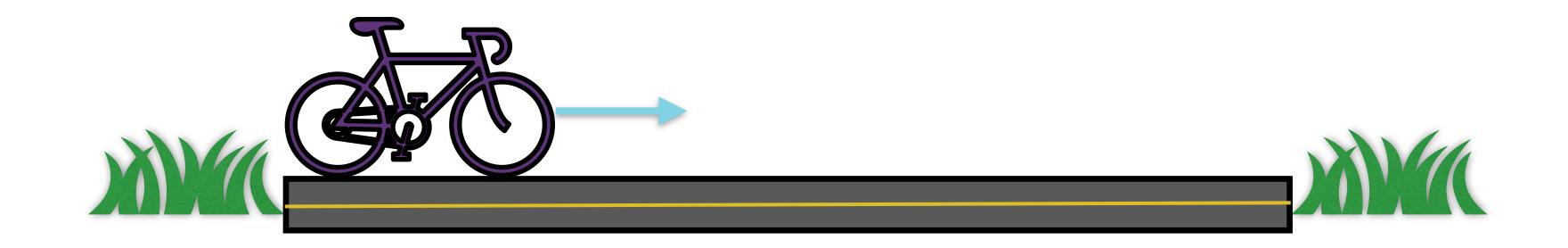


• From the bike's point of view, the surrounding road is unchanged.







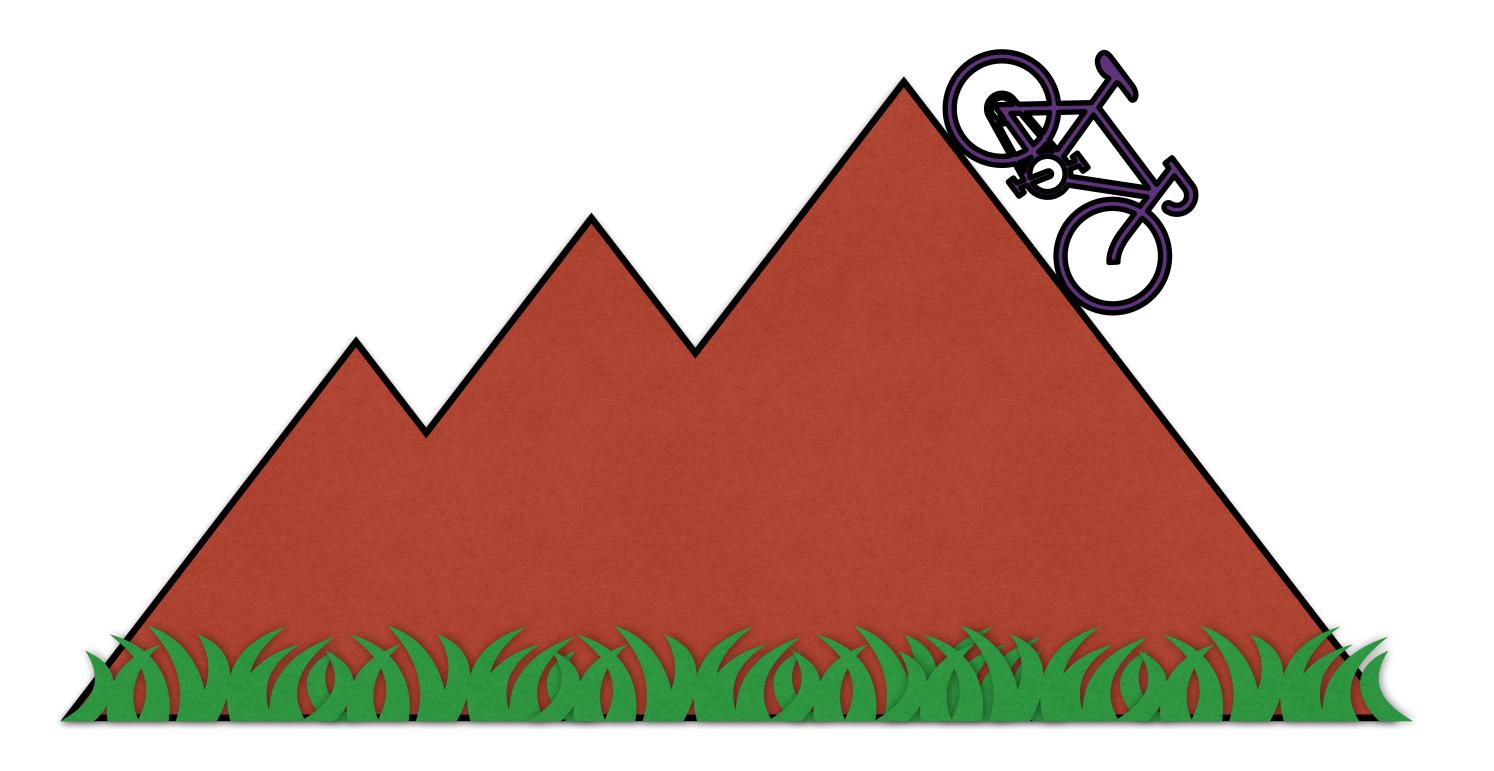


• The flat road exhibits a *translational invariance* symmetry, and, as a result, the bike conserves *linear momentum*.



Slightly more complicated: Time-translation Invariance

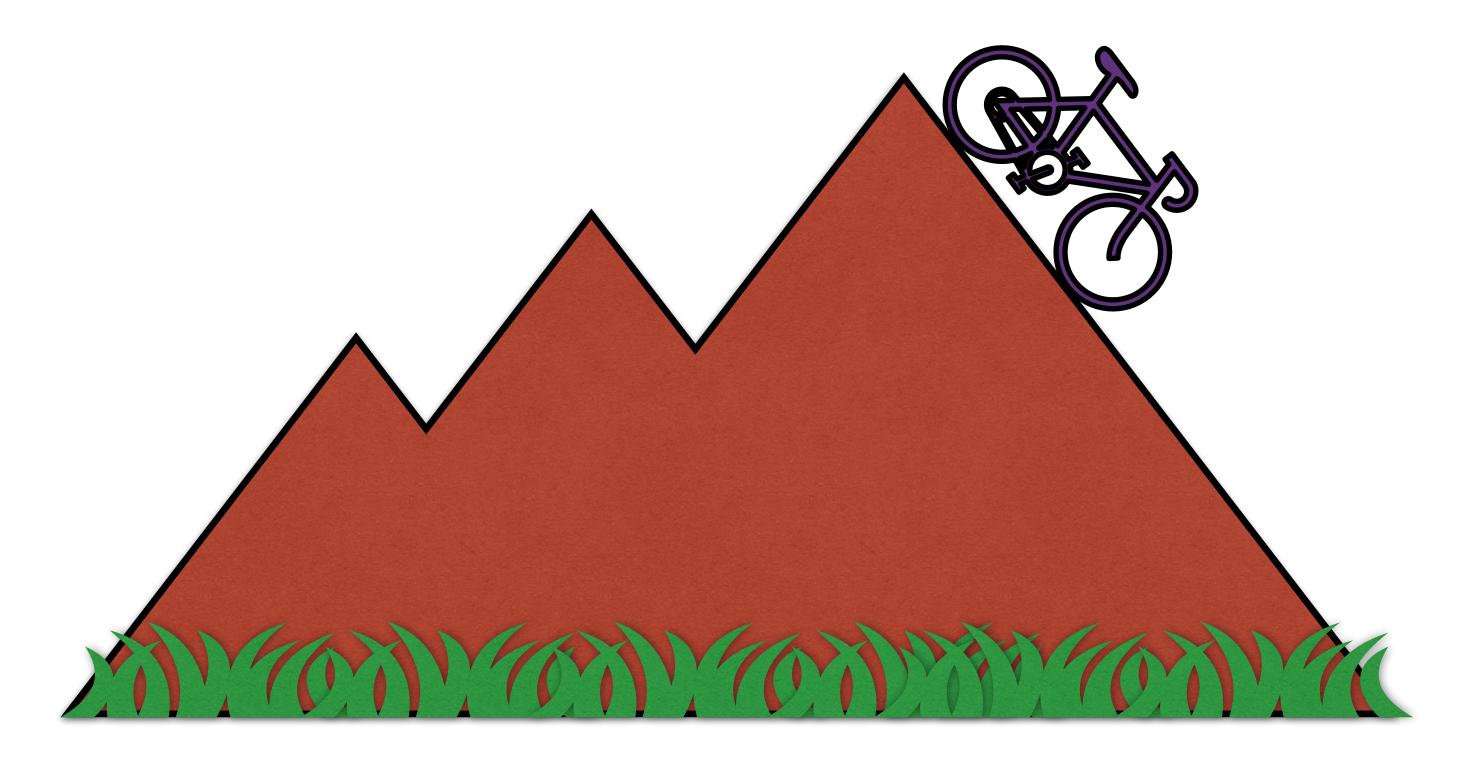
• (fancy way of saying "constant over time")





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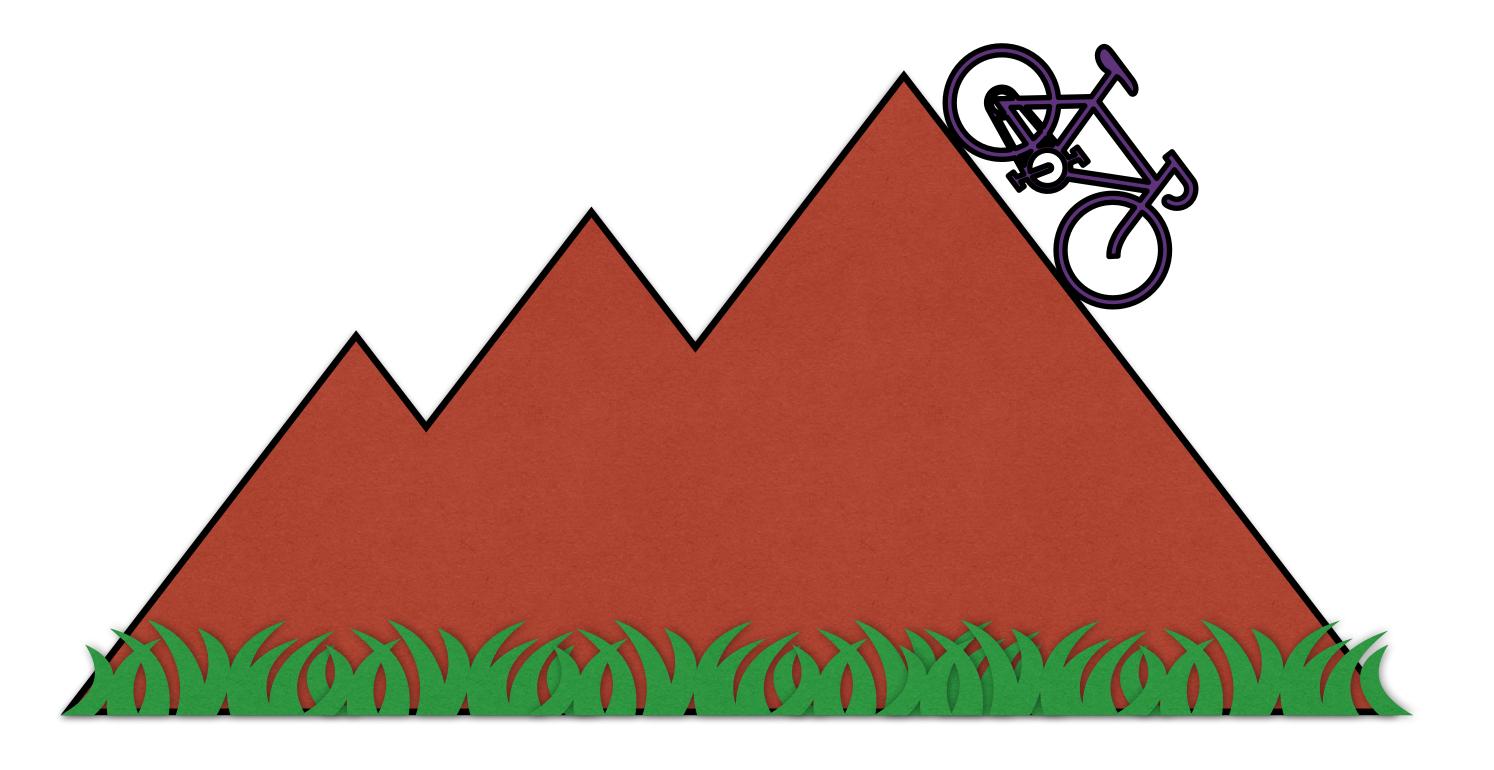


 Now with a bike riding over a hilly terrain — what conservation laws can we apply?



Slightly more complicated: Time-translation Invariance

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- Now with a bike riding over a hilly terrain — what conservation laws can we apply?
- Conservation of mechanical energy!









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Translational Invariance

Linear Momentum





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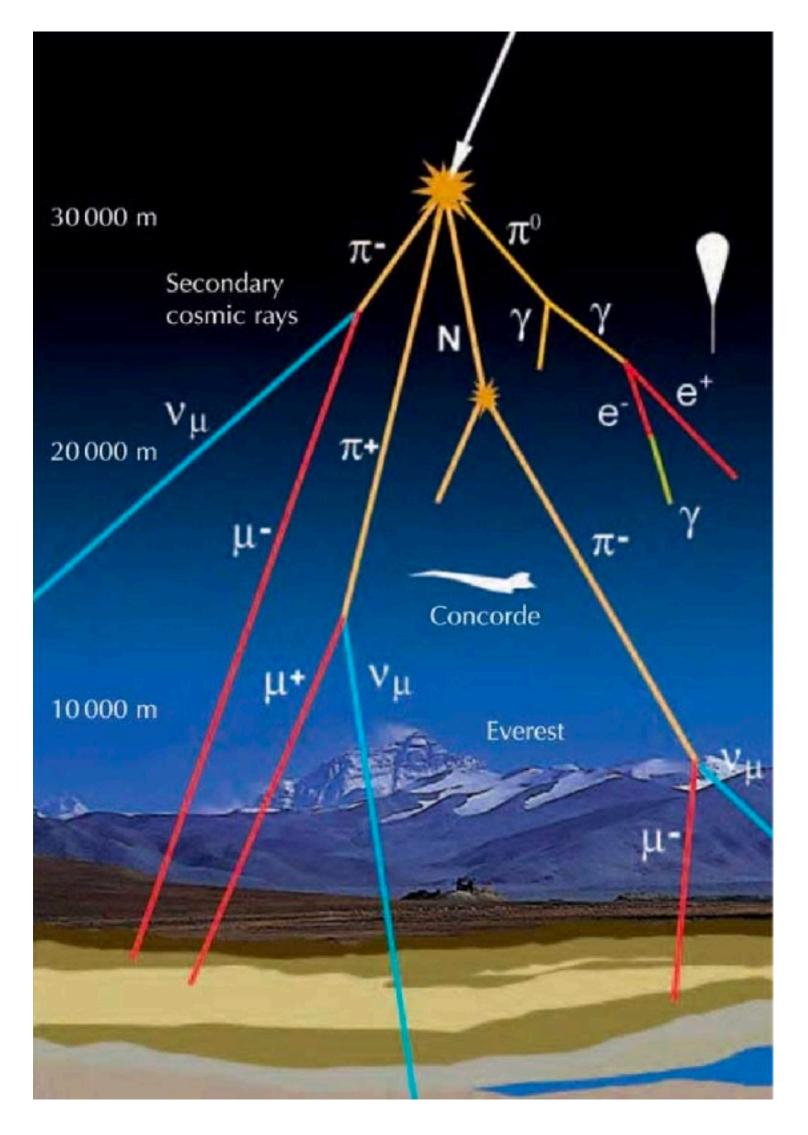
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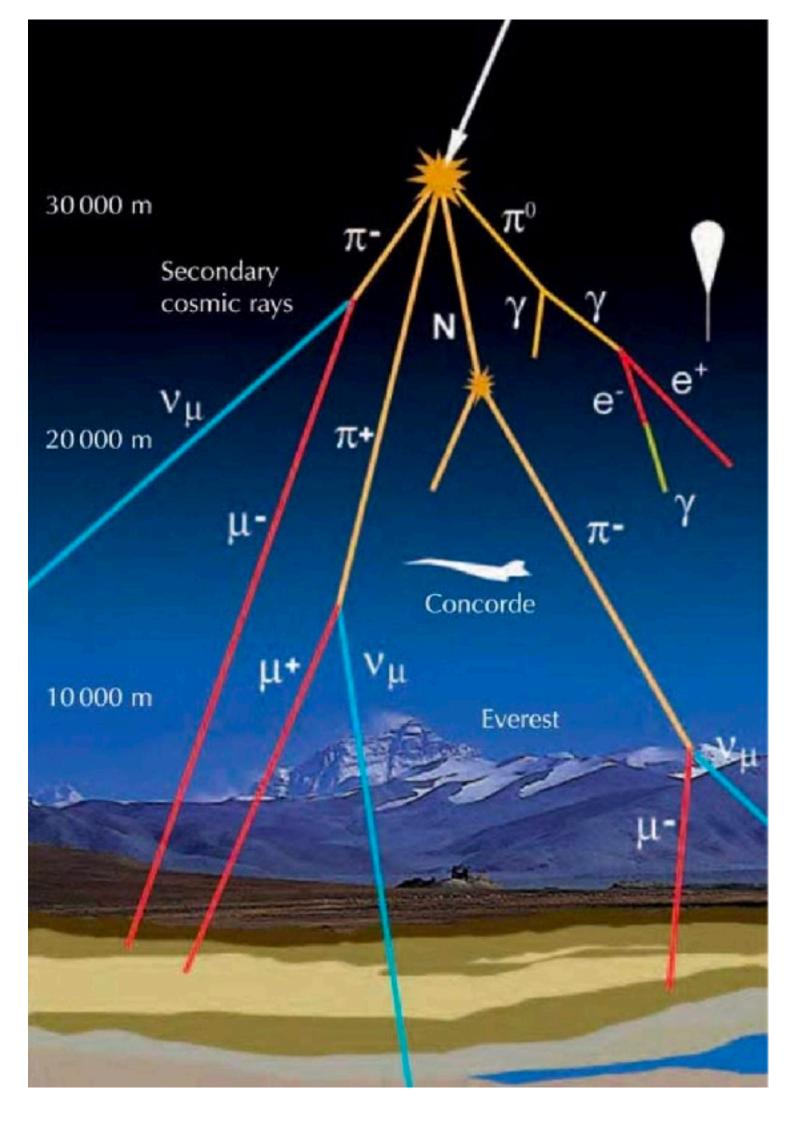
Time-translational Invariance





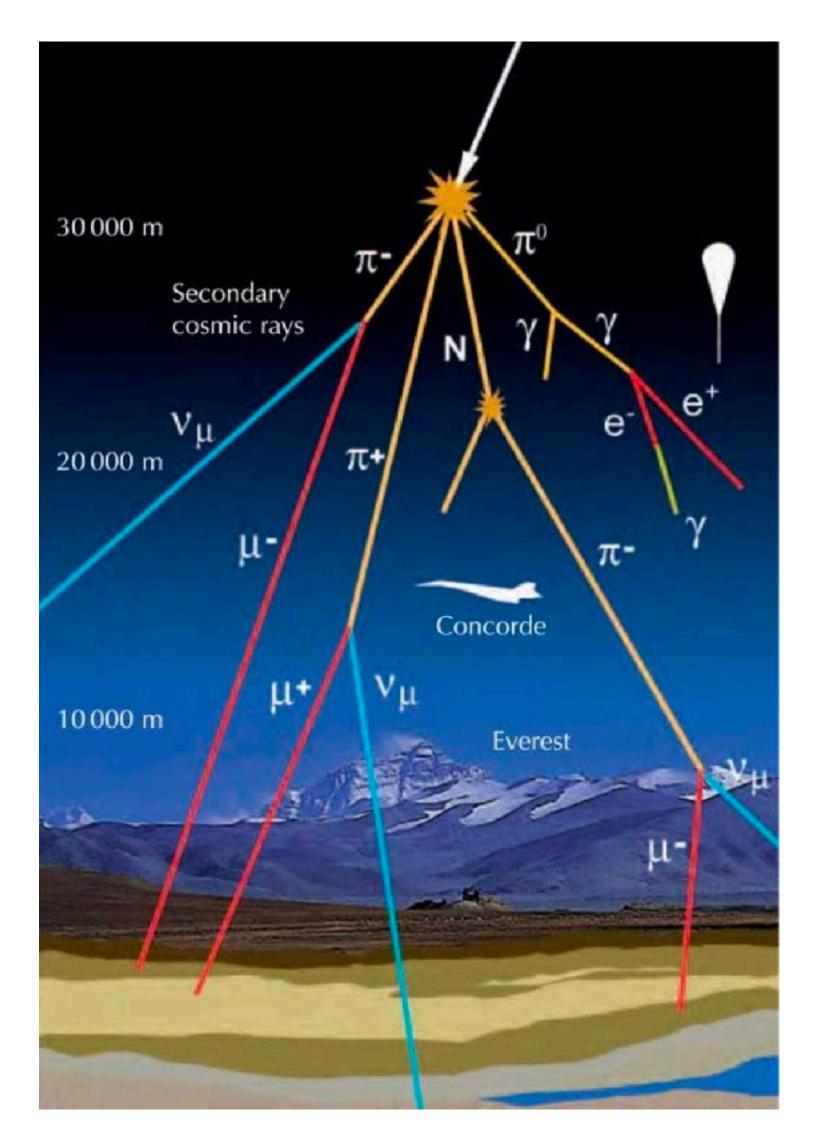


 Mid-1930s: scientists observed that protons and neutrons have very similar masses (both much larger than the electron's), and behave similarly in their interactions.



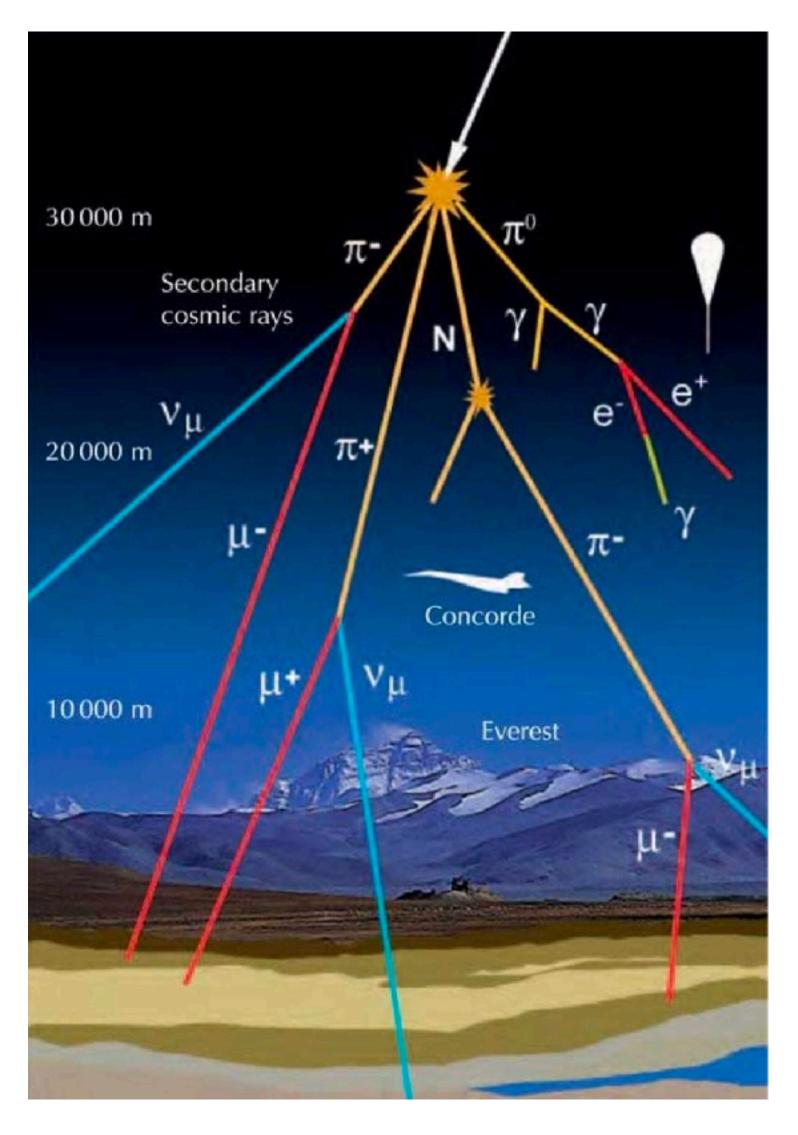


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- Introduce some spin-like symmetry where the proton has "isospin" up, and the neutron has "isospin" down.



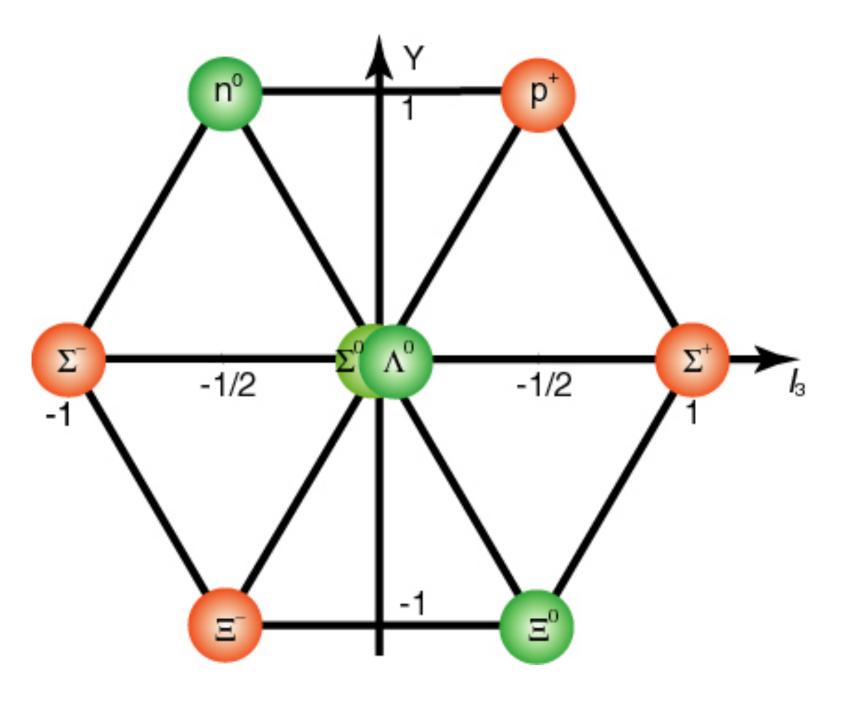


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- In the meantime, more exotic particles were being discovered in cosmic rays
 pions and muons.





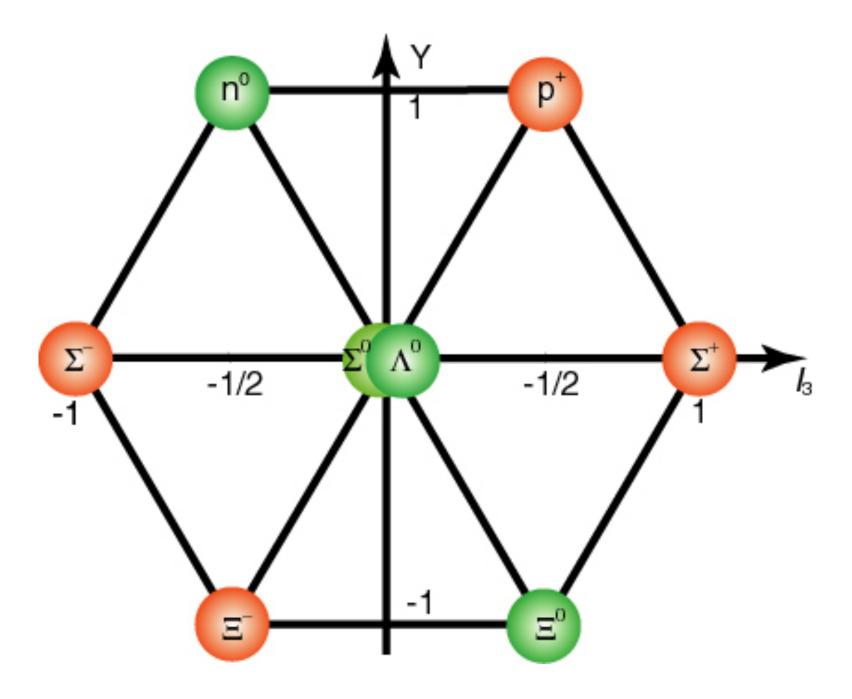
The Eightfold Way





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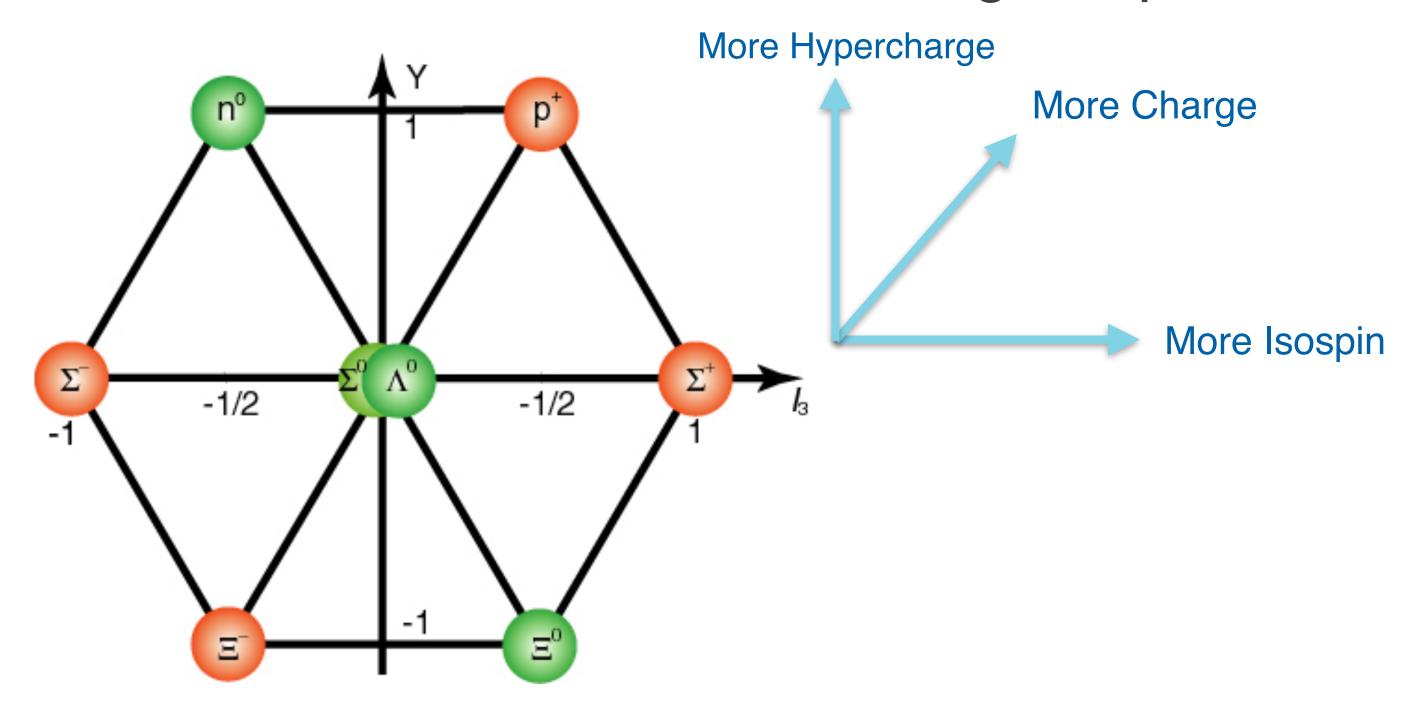
• Mesons, Baryons, and their interactions can be described using isospin and charge.





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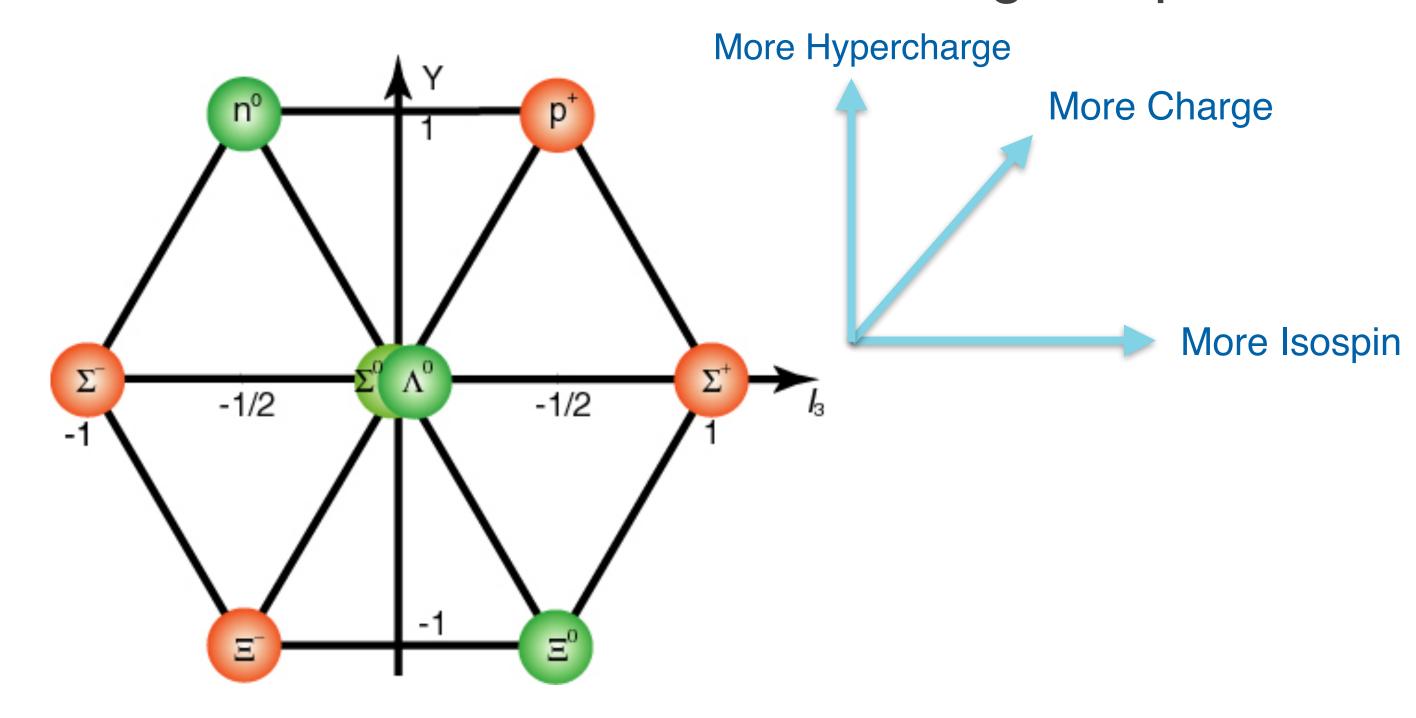
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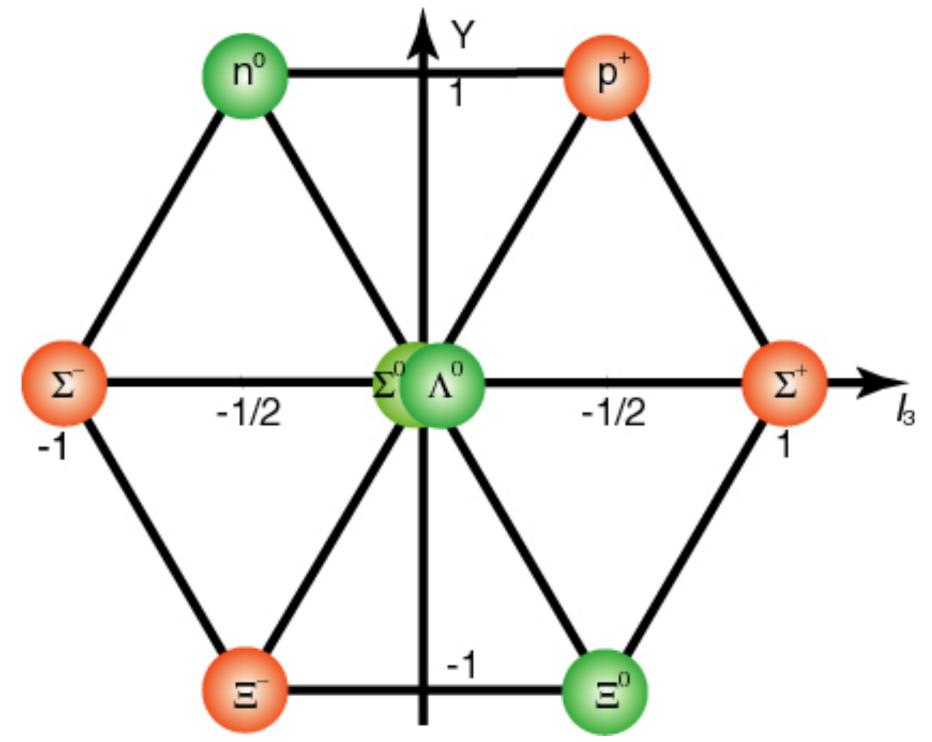


Mesons, like the charged pions, mediate interactions between these hadrons.



Gell-Mann, Zweig, and Quarks

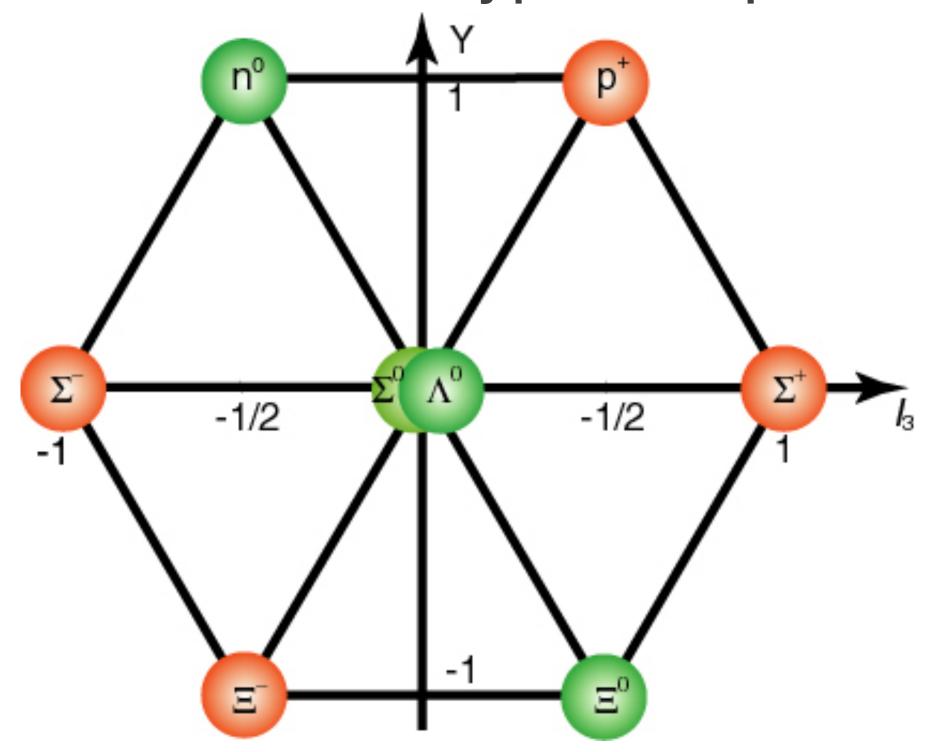
• Proposal: protons/neutrons/etc. are *not* fundamental particles, but made up of three different types of quarks: up, down, and strange.

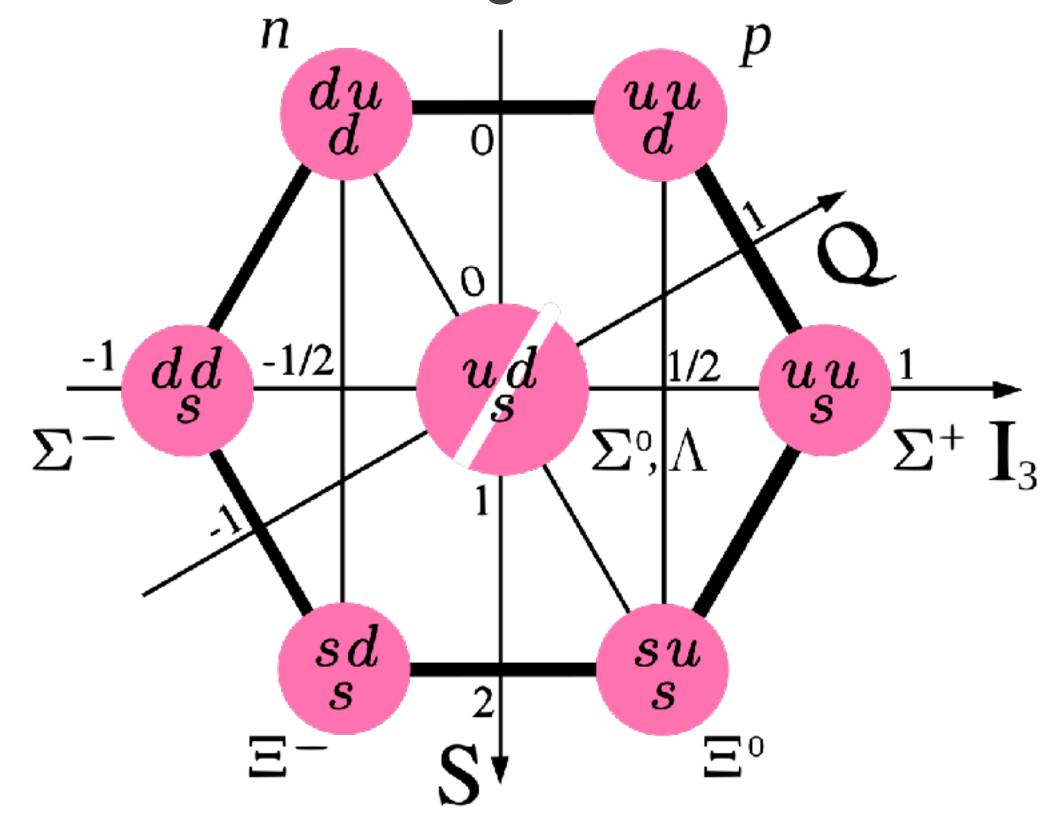




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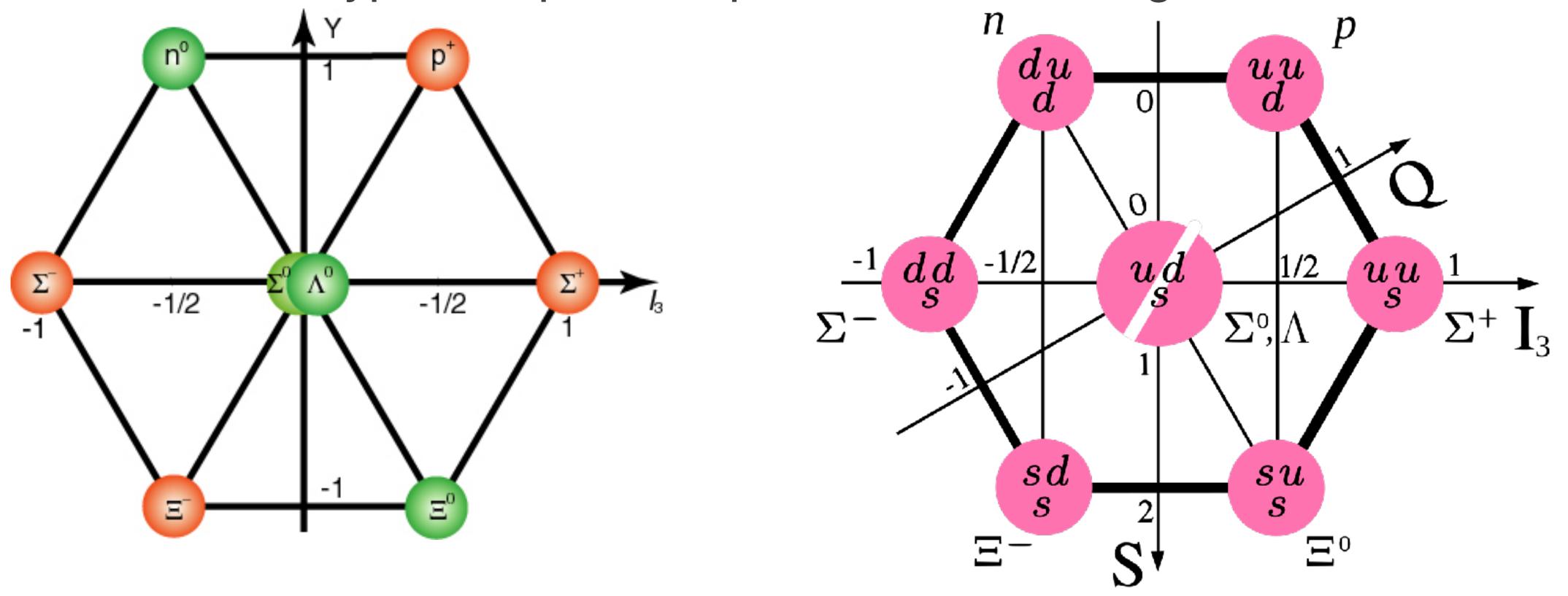






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The quarks interact under a symmetry principle called "SU(3)", a type of symmetry called a Lie group.





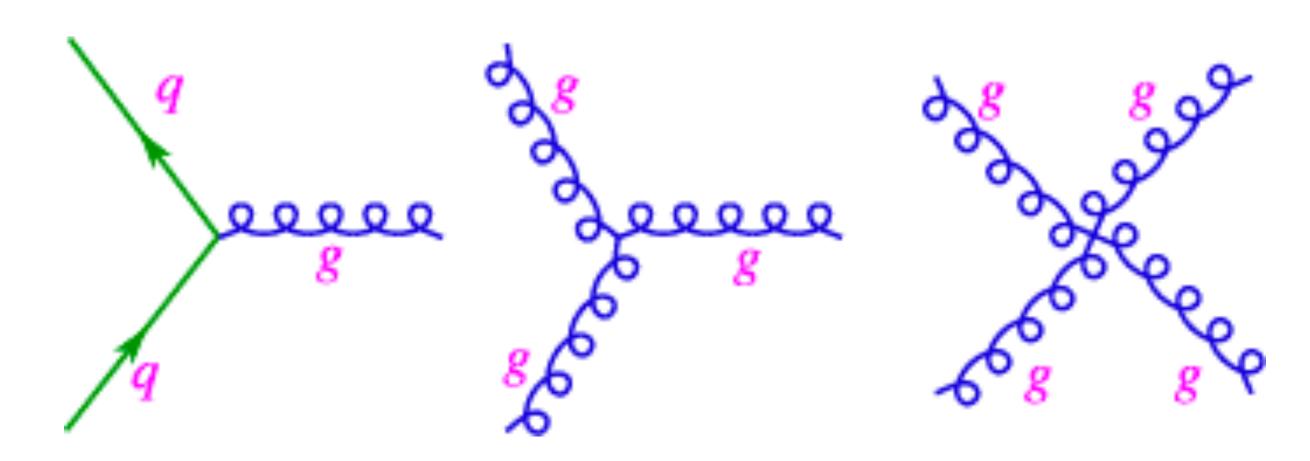
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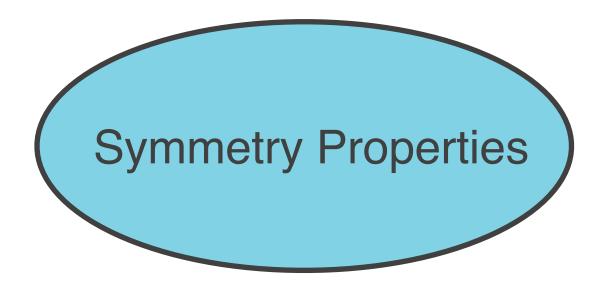


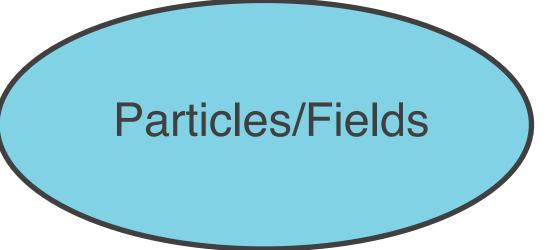
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- This led to the development of Quantum Chromodynamics (QCD), which uses a different **SU(3)** symmetry of "color" to describe all quark interactions.





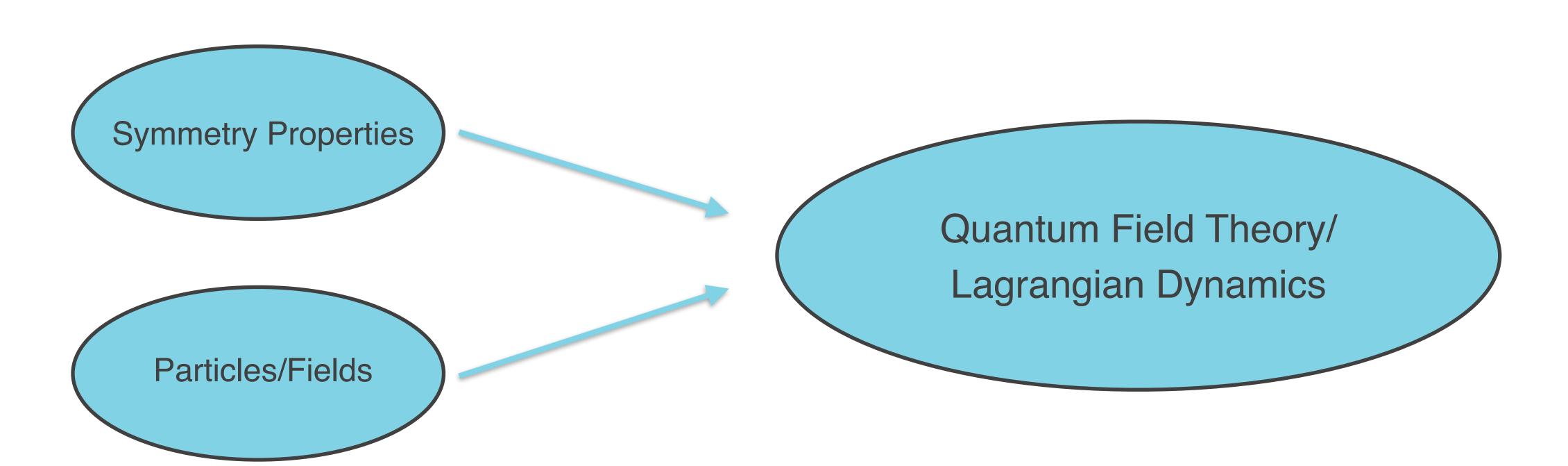
So, how do we use symmetries & particles together?







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Symmetry Properties

Particles/Fields

Quantum Field Theory/ Lagrangian Dynamics



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Particle(s) + Gauge Symmetry --- Description of Interactions



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Particle(s) + Gauge Symmetry --- Description of Interactions

Electrons + U(1) gauge symmetry ——— Quantum Electrodynamics (QED)



Side-note: Gauge Boson Masses

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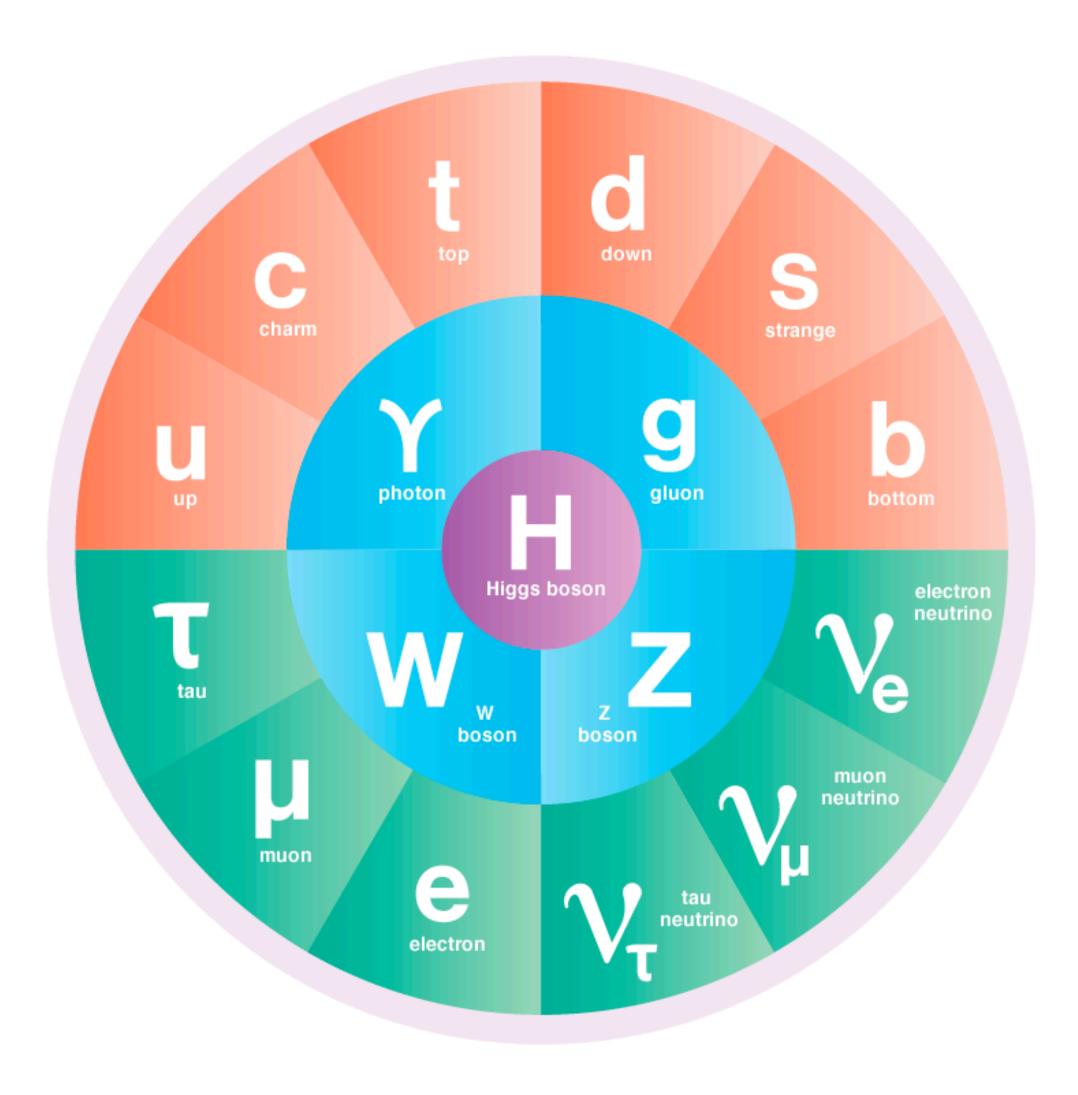
$$\mathcal{L} \supset m^2 A_\mu A^\mu$$

Under the transformation we introduced for this, though, the Lagrangian is
 not invariant — gauge bosons must be massless.

$$A^{\mu} \to A^{\mu} - \frac{1}{-} \partial^{\mu} \theta$$

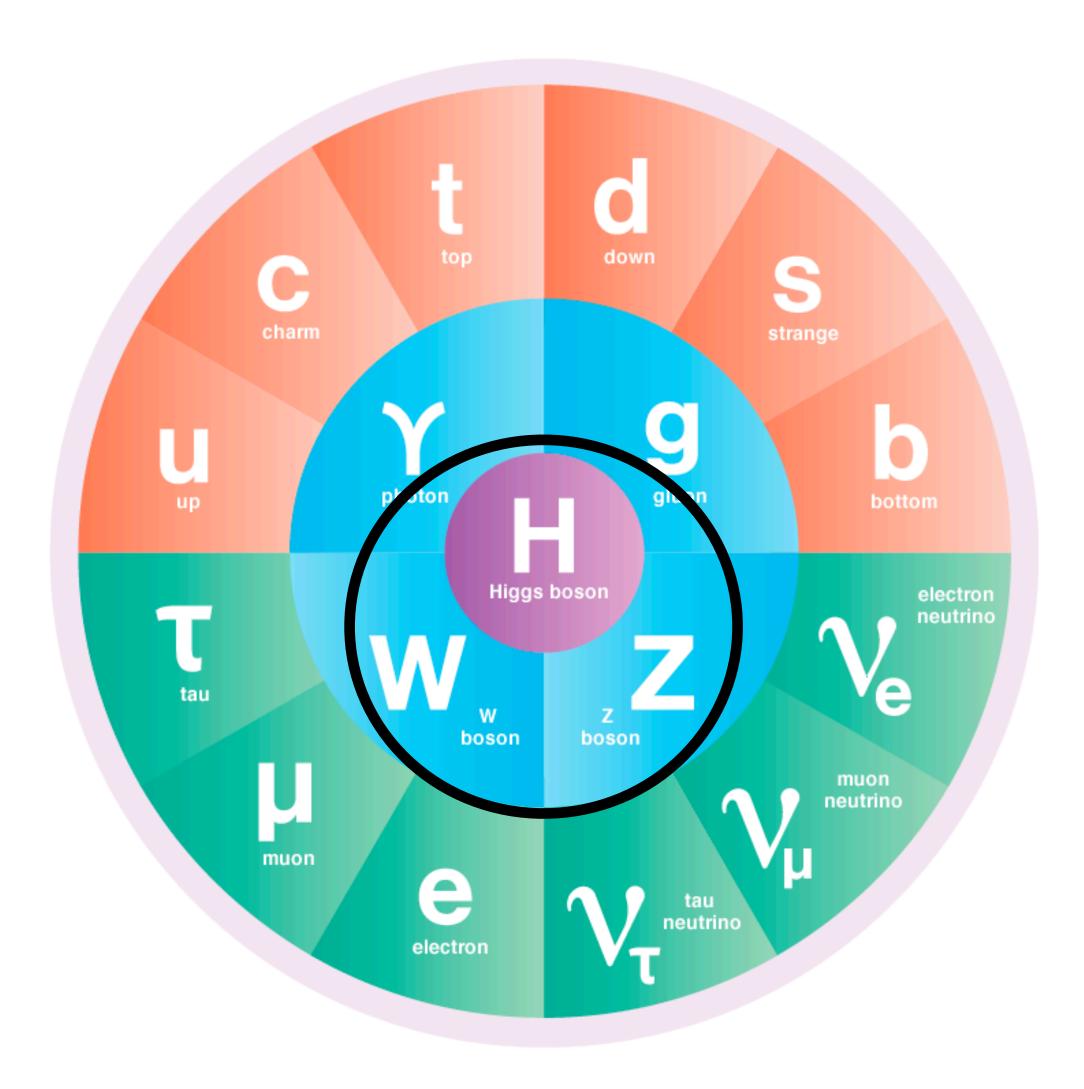


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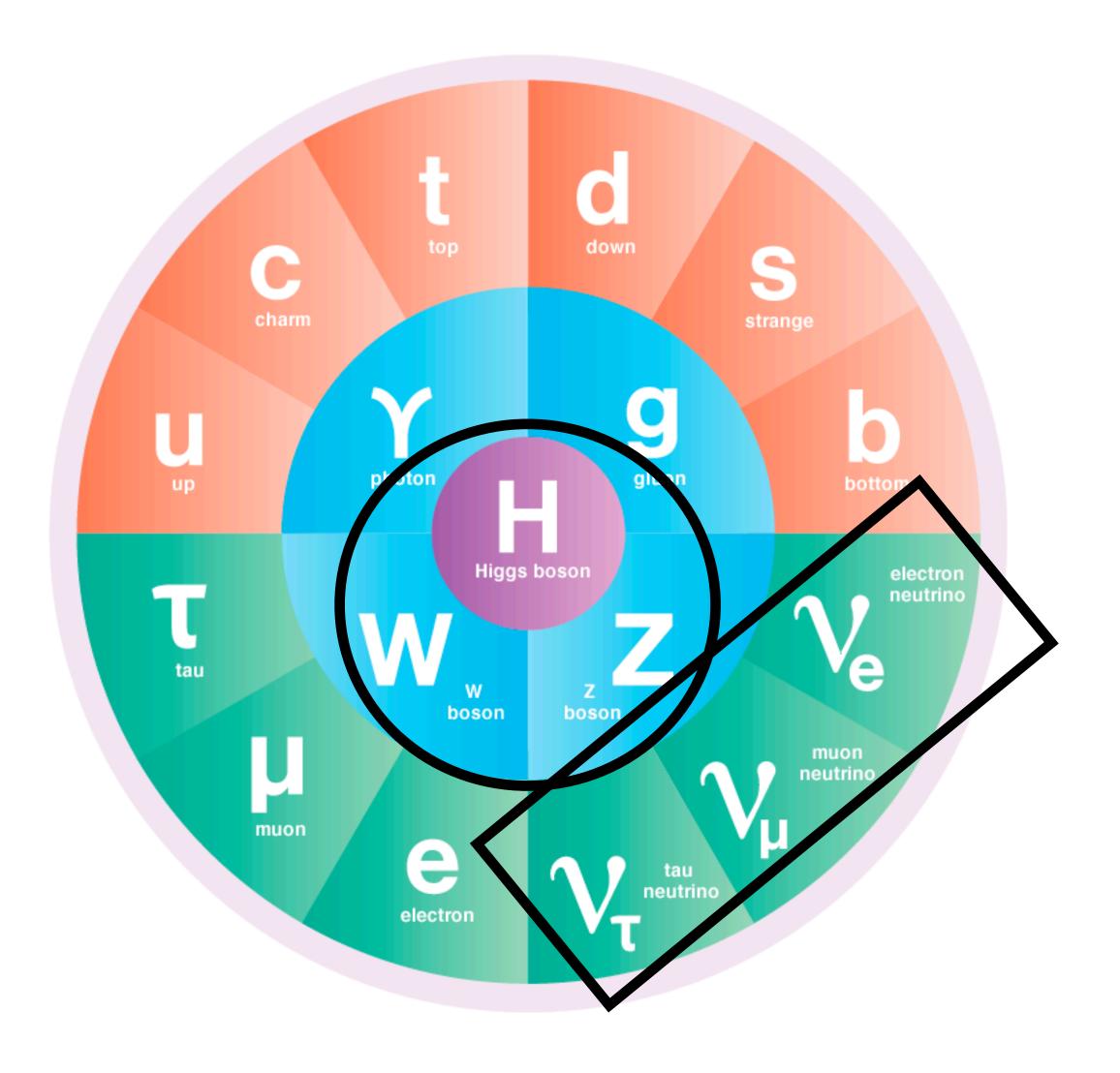


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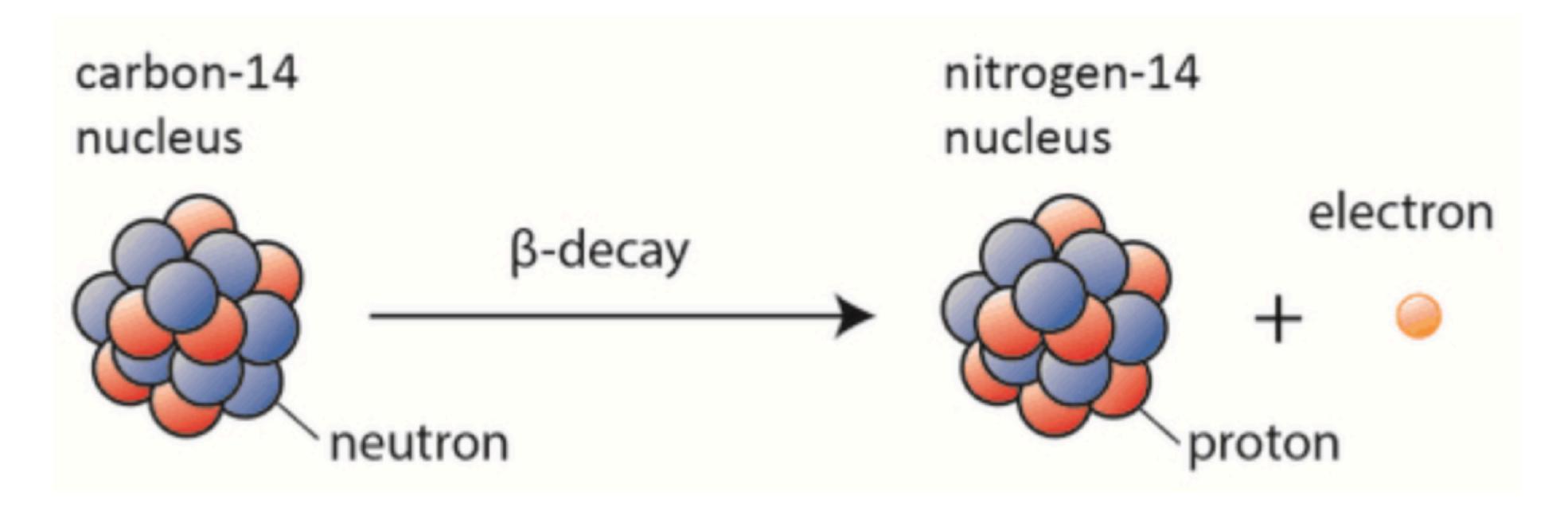


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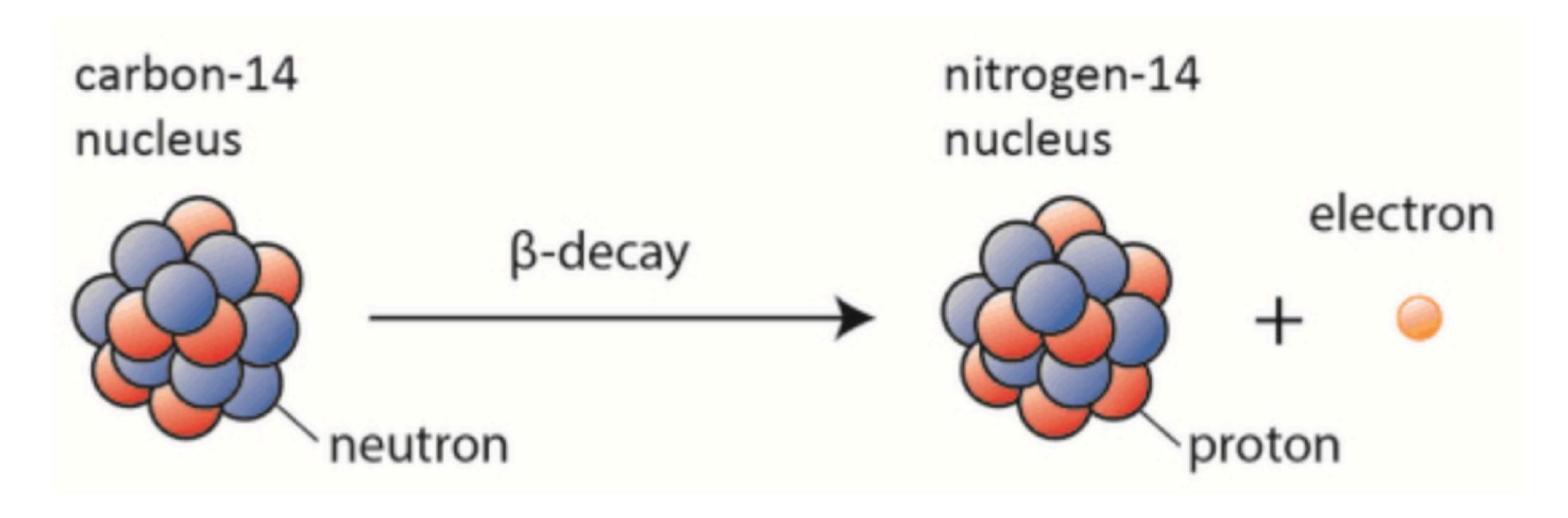


Back to the 1930s: Nuclear Beta Decays





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Example beta decay —a neutron inside a Carbon nucleus spontaneously changes to a proton (actually a down quark changing to an up quark), and an electron is emitted to conserve charge.



Conservation of Energy in Beta Decays



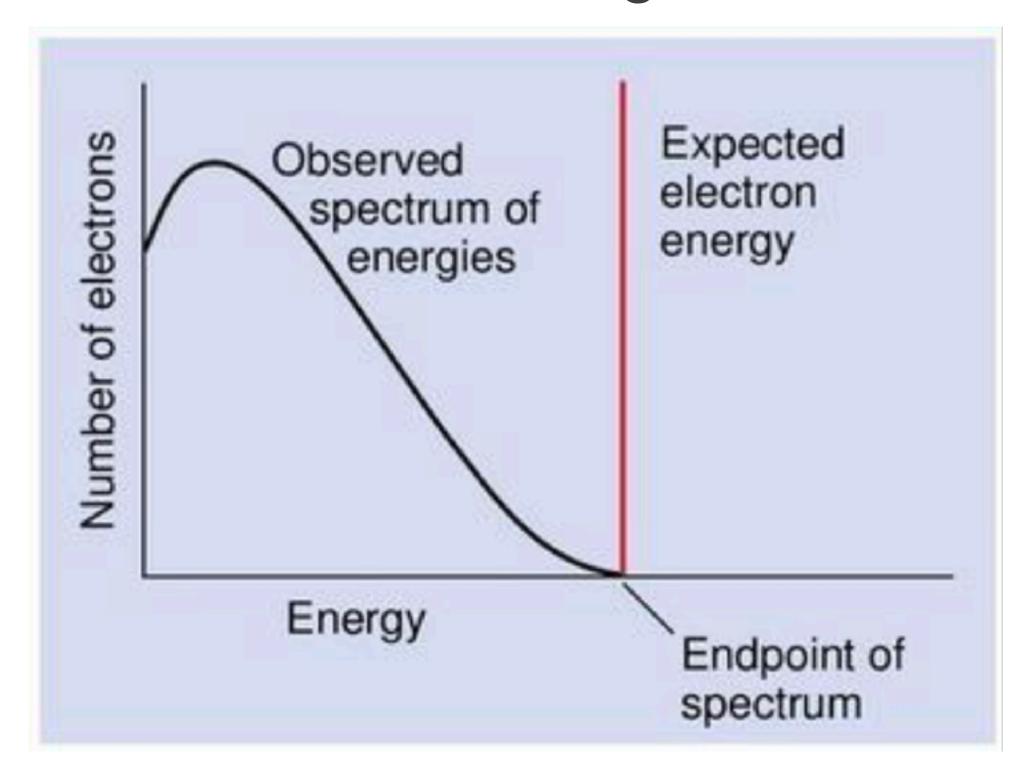
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Measurements of beta decay show that the electron has a spectrum!



Wolfgang Pauli: "Dear Radioactive Ladies and Gentlemen"

Abschrift

Physikalisches Institut der Eidg. Technischen Hochschule Zurich

Zirich, 4. Des. 1930 Cloriastrasse

Liebe Radioaktive Damen und Herren,

Wie der Ueberbringer dieser Zeilen, den ich huldvollst ansuhören bitte, Ihnen des näheren auseinandersetzen wird, bin ich angesichts der "falschen" Statistik der N- und Li-6 Kerne, sowie des kontinuierlichen beta-Spektrums auf einen versweifelten Ausweg verfallen um den "Wechselsats" (1) der Statistik und den Energiesats su retten. Mamlich die Möglichkeit, es könnten elektrisch neutrale Teilchen, die ich Neutronen nennen will, in den Kernen existieren, welche den Spin 1/2 haben und das Ausschliessungsprinzip befolgen und won lichtquanten musserden noch dadurch unterscheiden, dass sie micht mit Lichtgeschwindigkeit laufen. Die Masse der Neutronen finante von dersalben Grossenordnung wie die Elektronenmasse sein und jesenfalls nicht grosser als 0,01 Protonermasse .- Das kontimuierliche bets- Spektrum ware denn verständlich unter der Annahme, dass beim beta-Zerfall mit dem blektron jeweils noch ein Neutron emittiert Mird. derart, dass die Summe der Energien von Neutron und Elektron konstent ist.





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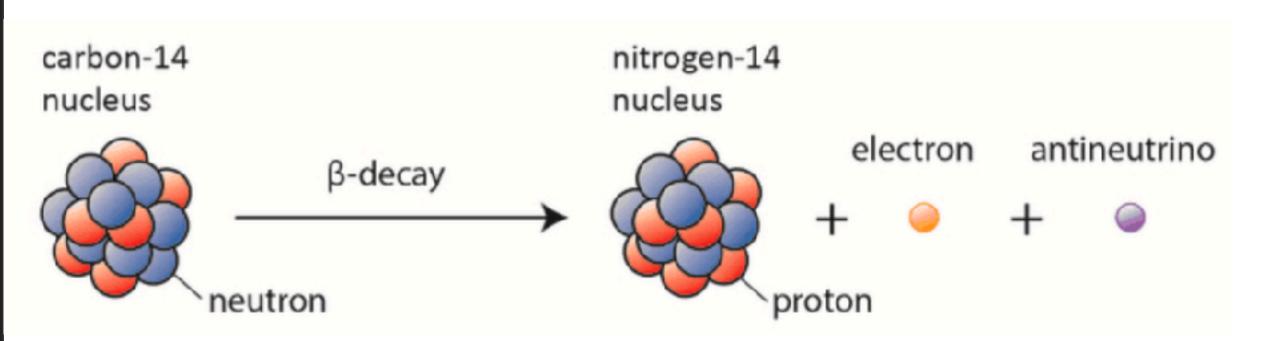
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• Pauli's solution: a new, very weakly interacting "neutrino" comes out of beta decays as well, stealing some of the energy from the outgoing electron and producing a spectrum of electron energies.



Just how "weak" is "weakly interacting"?

• Enrico Fermi, shortly after Pauli: $\mathscr{L}_{\text{Fermi}} \supset G_F\left(\overline{\psi}_p \Gamma \psi_n\right) \left(\overline{\psi}_e \Gamma' \psi_\nu\right)$



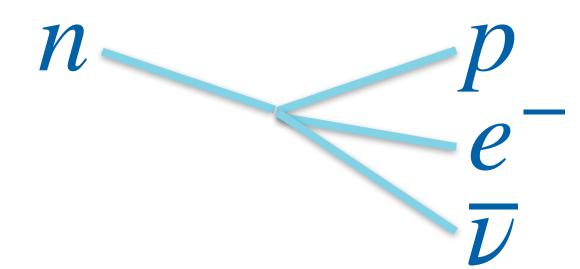
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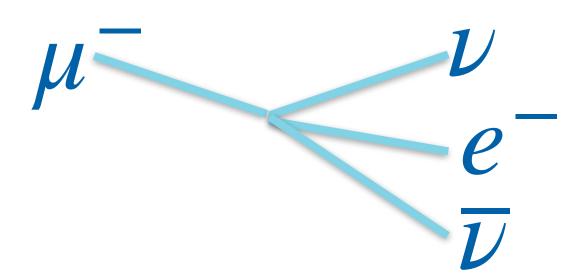


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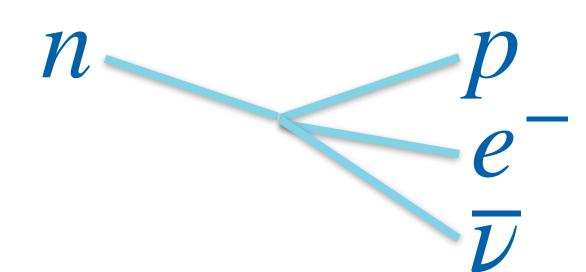


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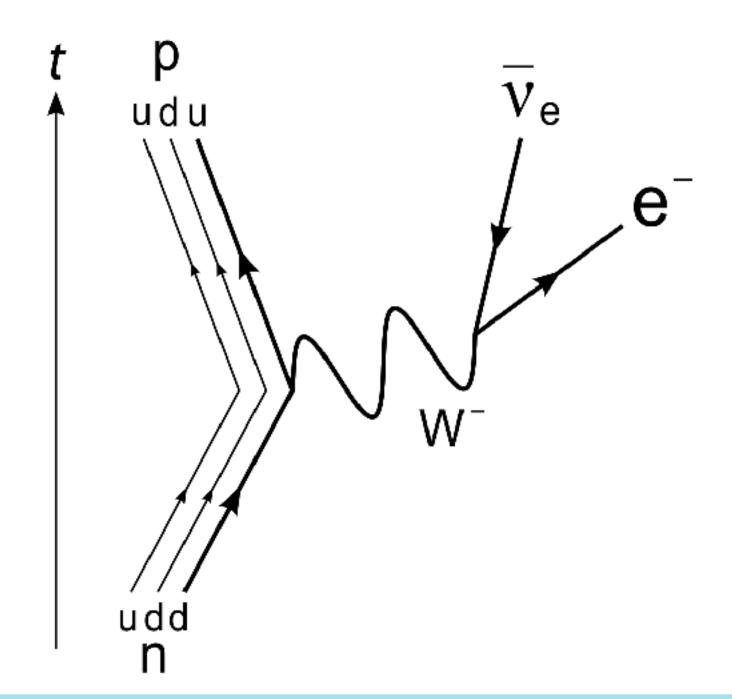
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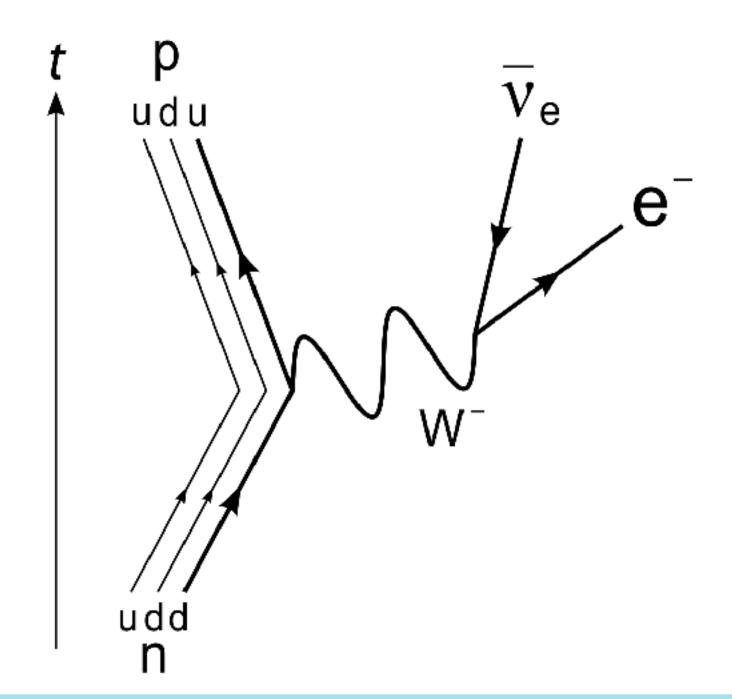
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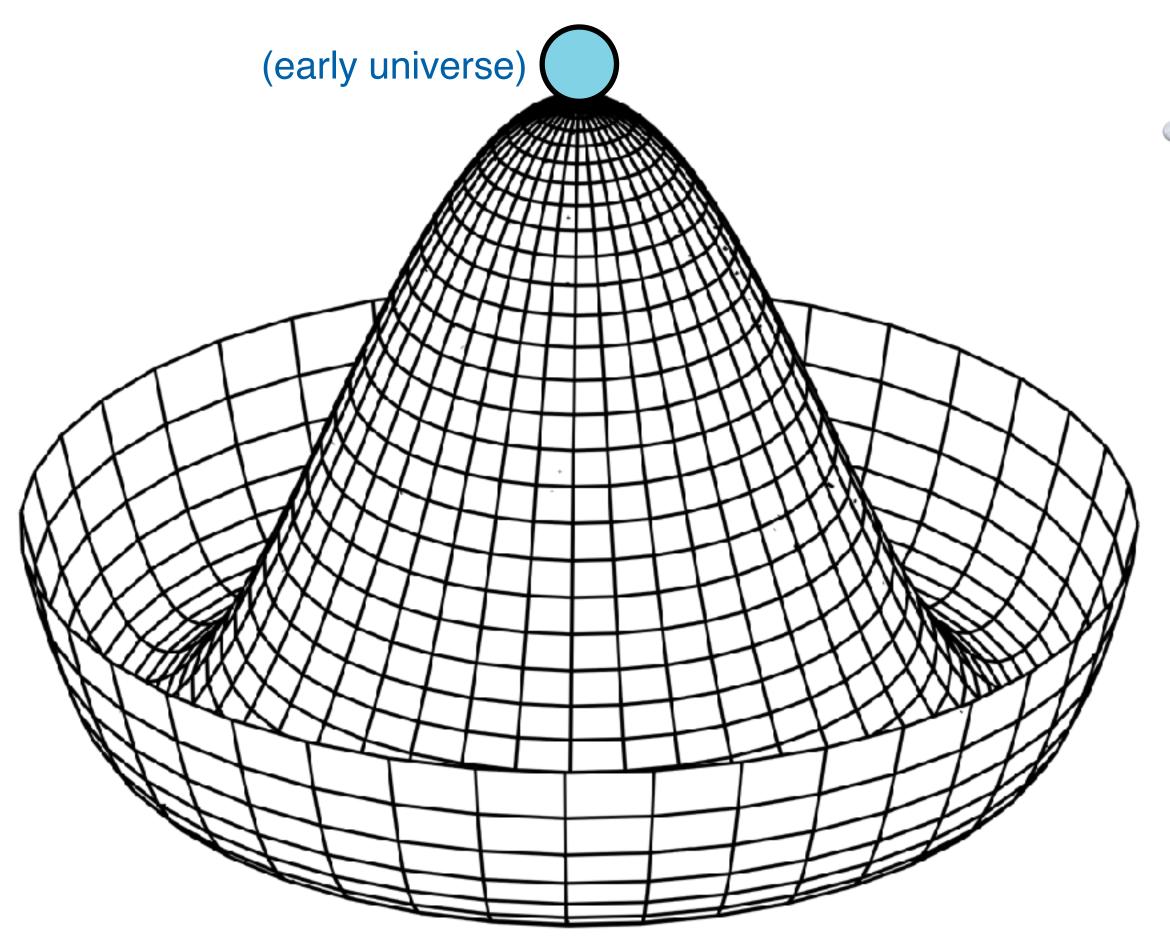
- Force carriers much heavier than the light fermions are mediating these interactions.
- 1983 the two bosons, W and Z, associated with this interaction, were discovered at CERN.
- Issue: if these are really the gauge bosons of an **SU(2)** theory, they should have zero mass!



The Higgs boson and the Higgs mechanism:



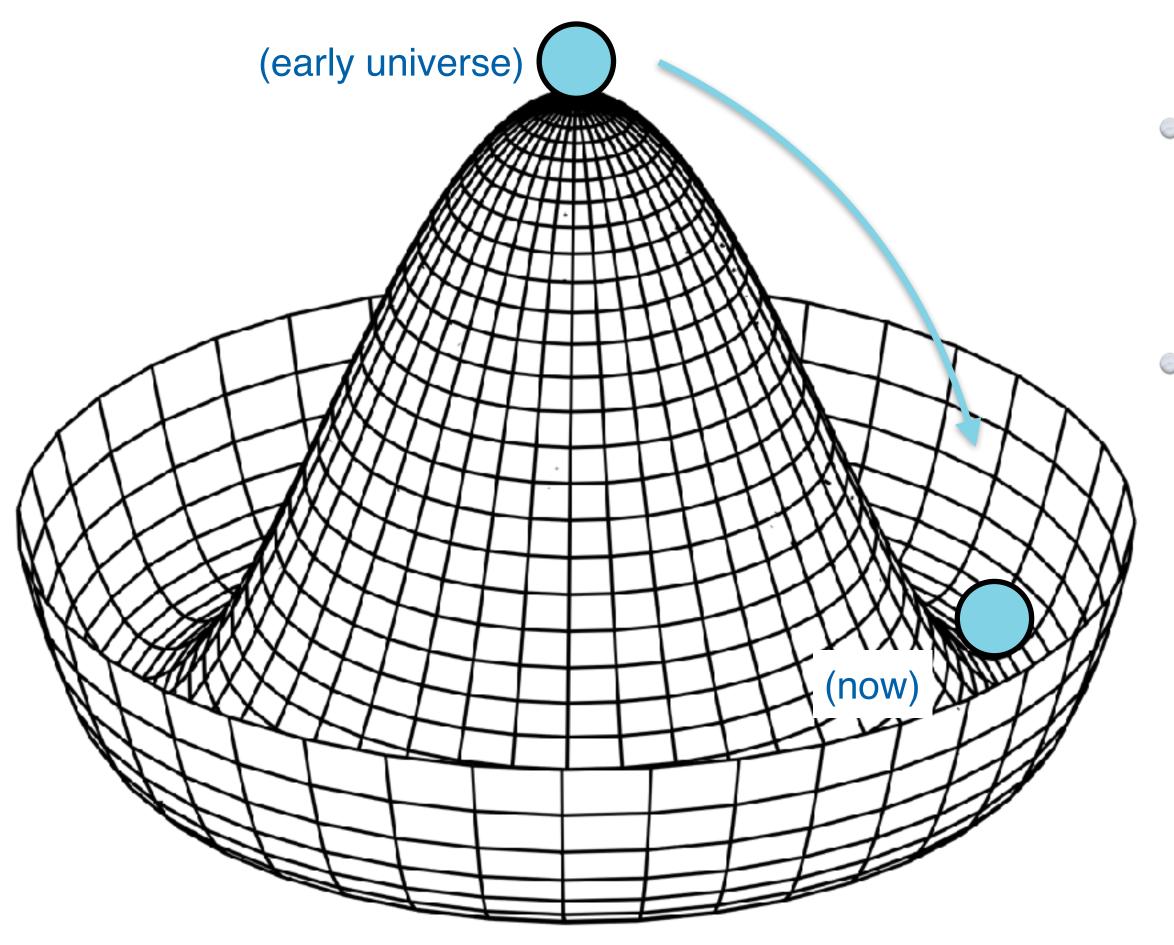
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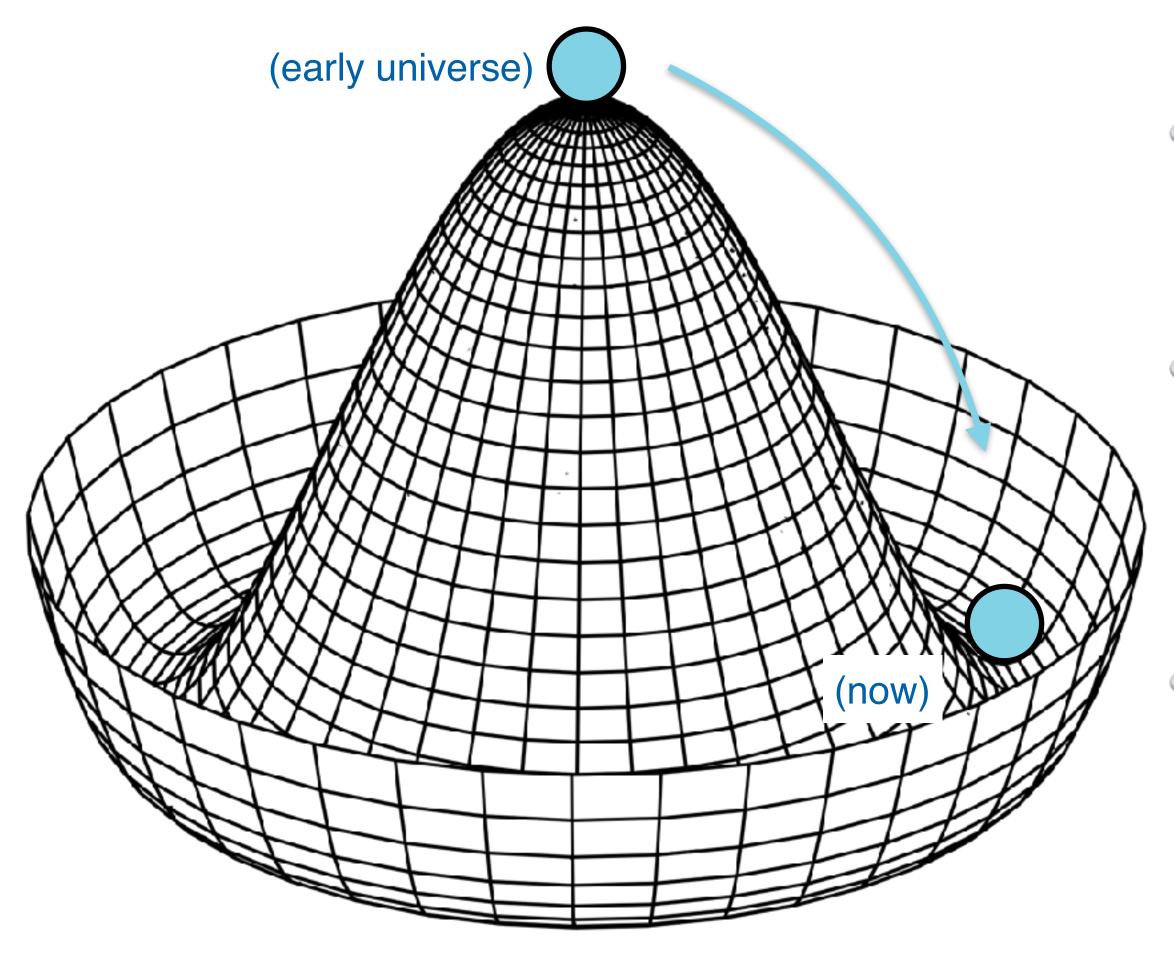
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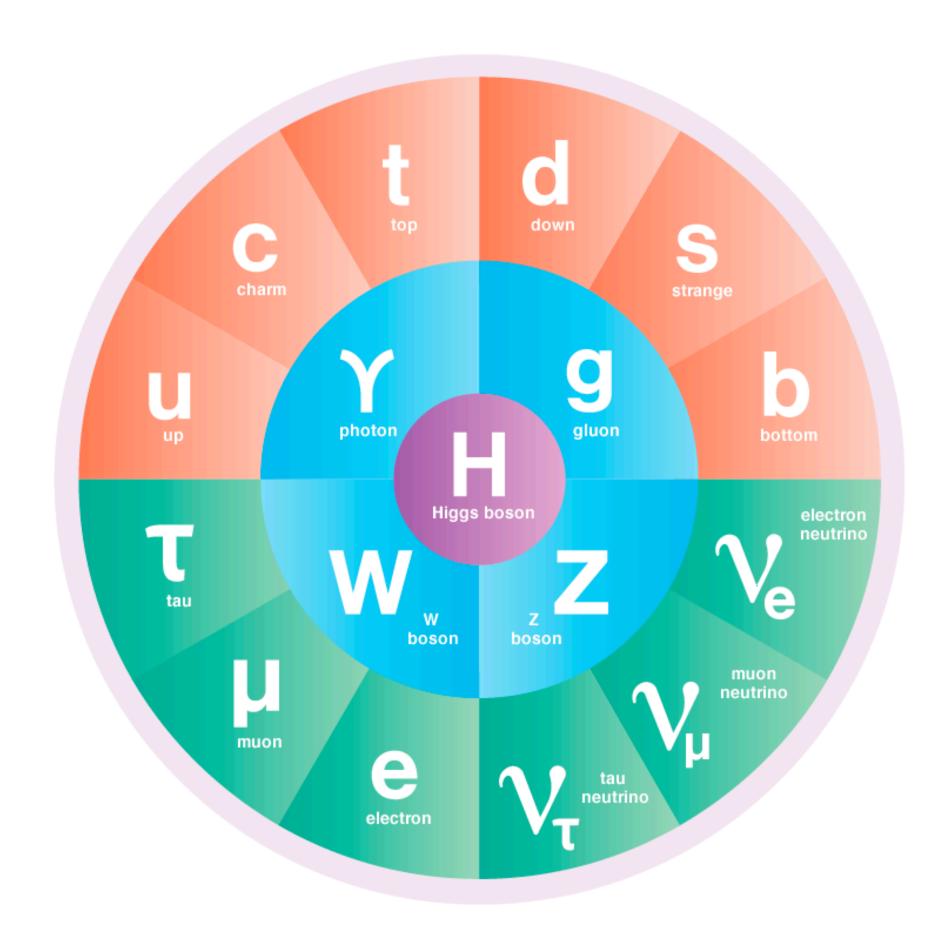
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- As the universe cools, the Higgs boson feels compelled to acquire a "vacuum expectation value" at the bottom of this "Mexican hat" potential.
- This new minimum spontaneously breaks the
 SU(2) x U(1) symmetry, resulting in things like massive particles, massive gauge bosons, and the resulting "low energy" behavior we see today.

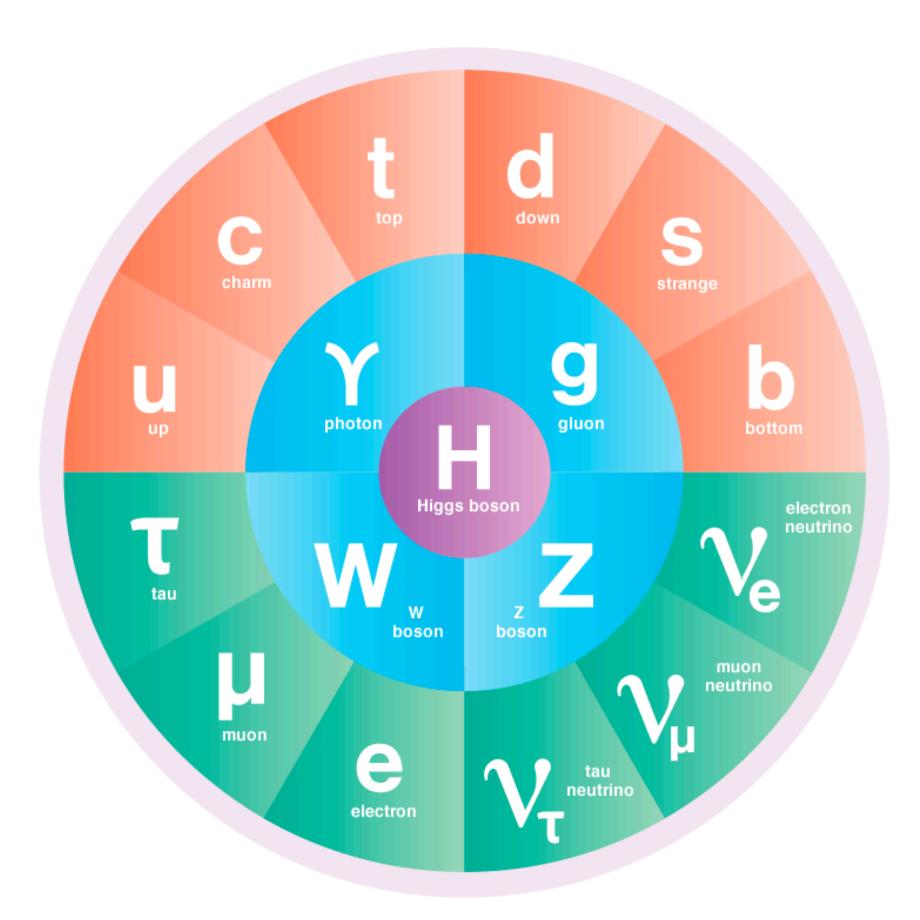


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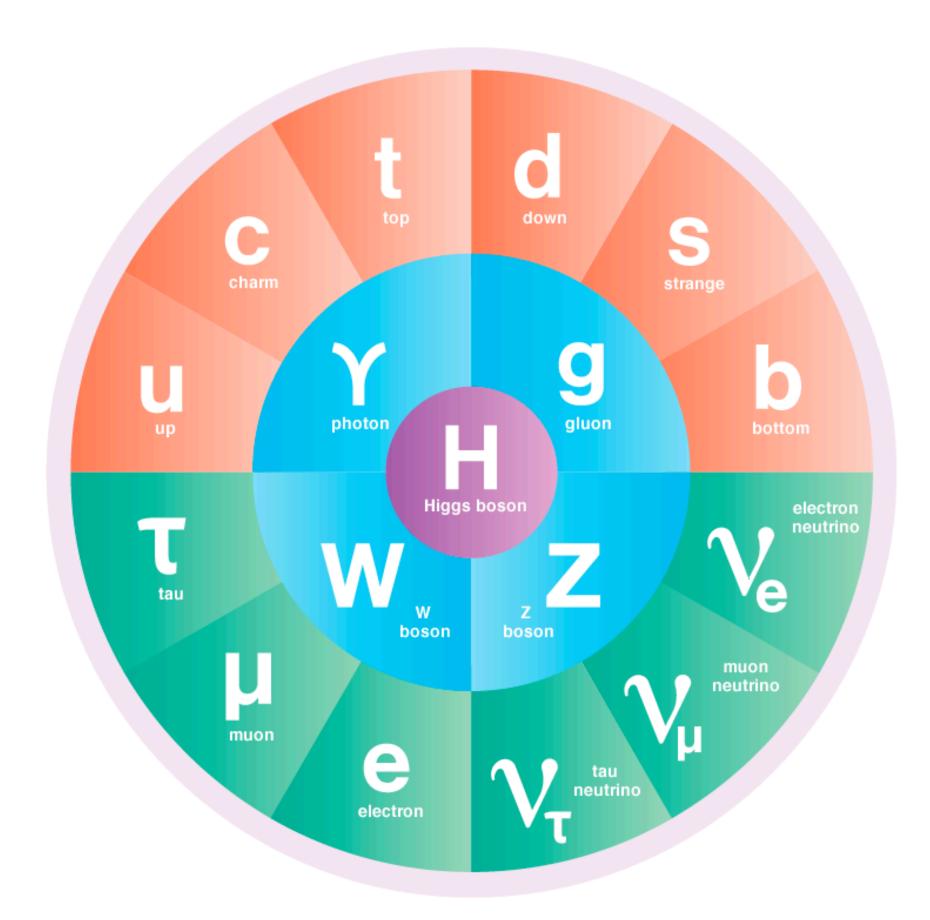
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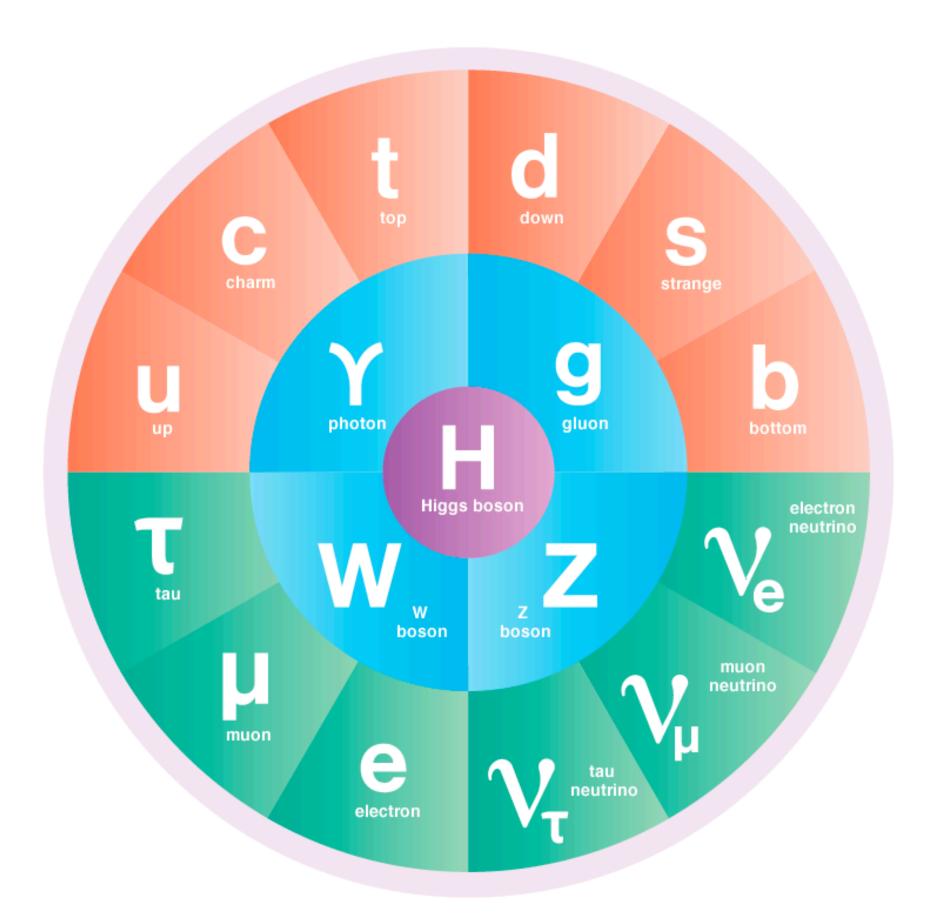
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- Force carriers are the gluons (color, interact with quarks), photons (electromagnetism, interact with charged particles), and weak bosons (weak force, interact with all fermions).



Completing the Standard Model was a HUGE triumph.



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However, there are some problems...





Dark Matter & Dark Energy



Dark Matter & Dark Energy

Neutrino Oscillations and Neutrino Masses



Dark Matter & Dark Energy

Neutrino Oscillations and Neutrino Masses

The Higgs Hierarchy Problem

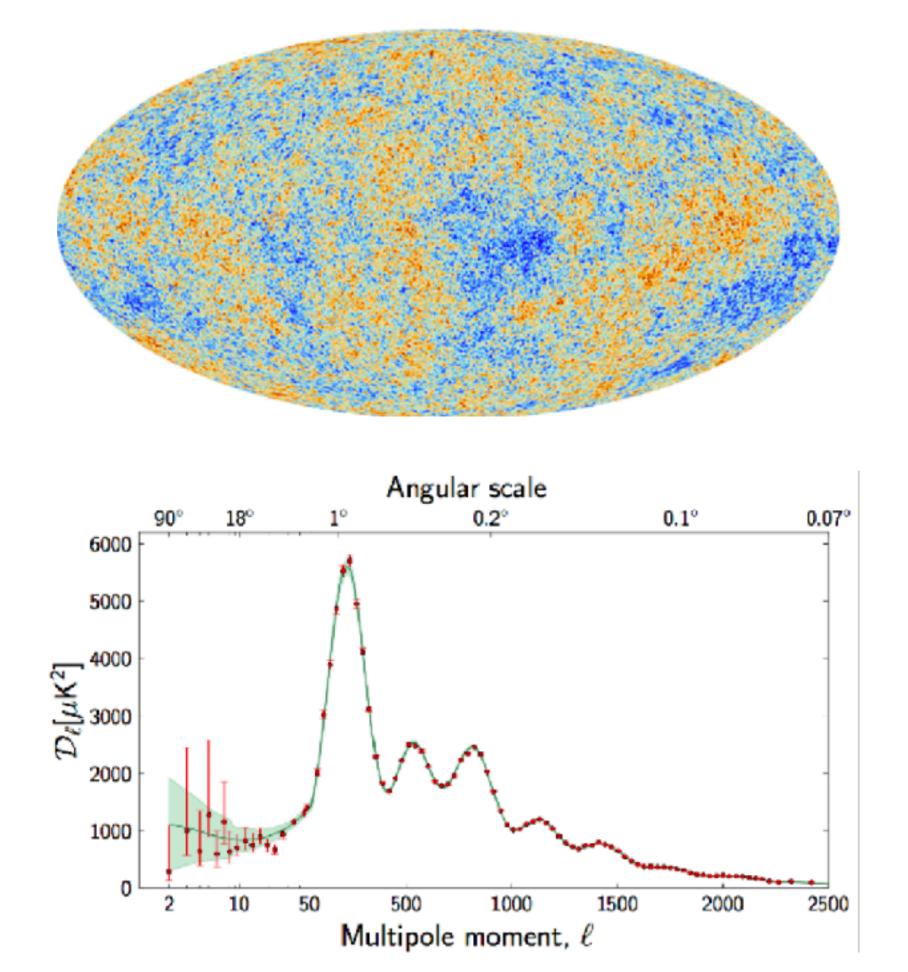


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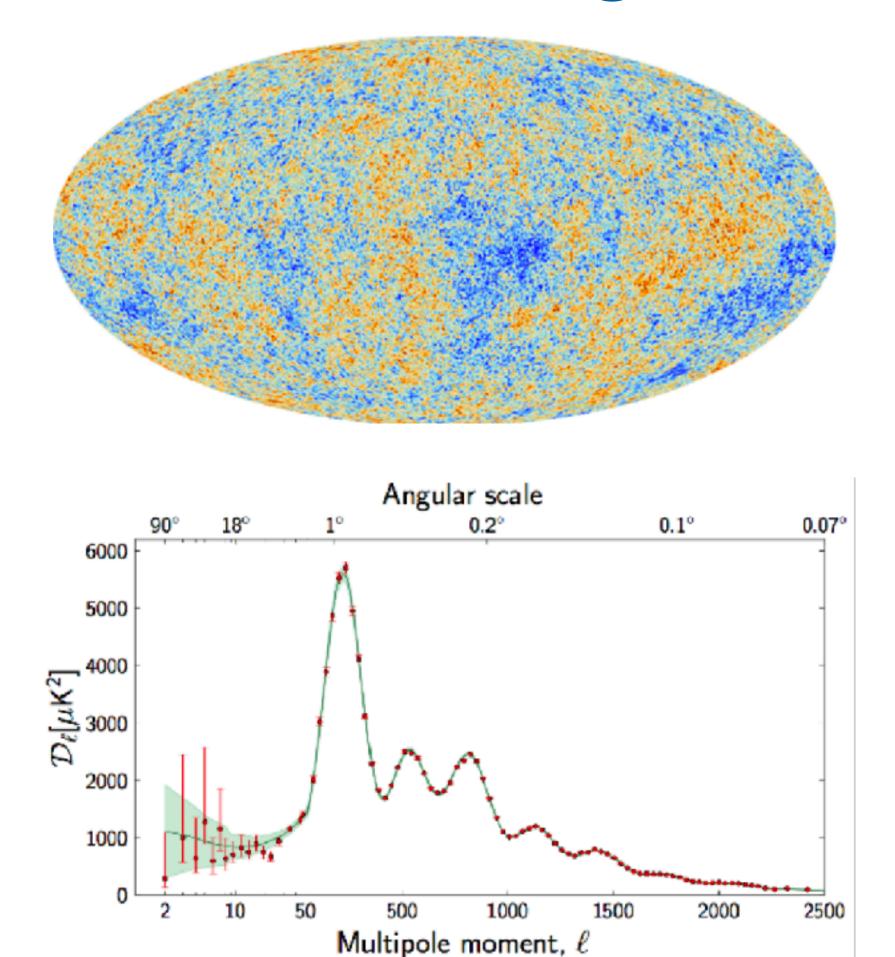






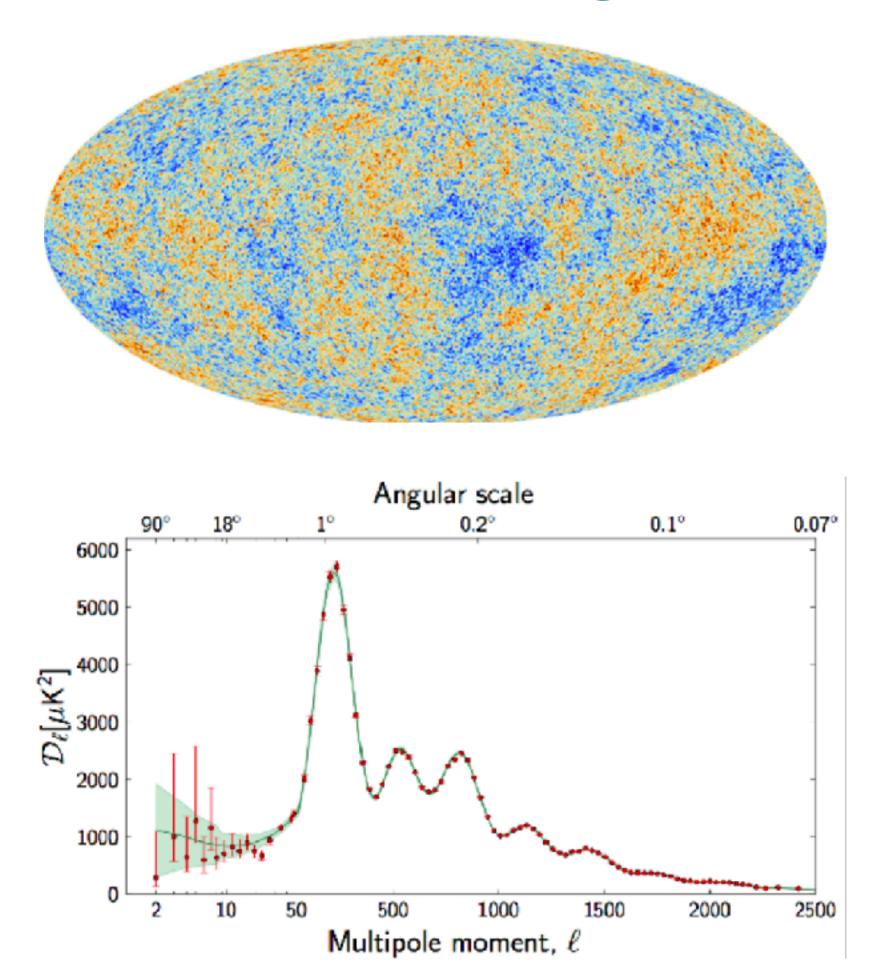


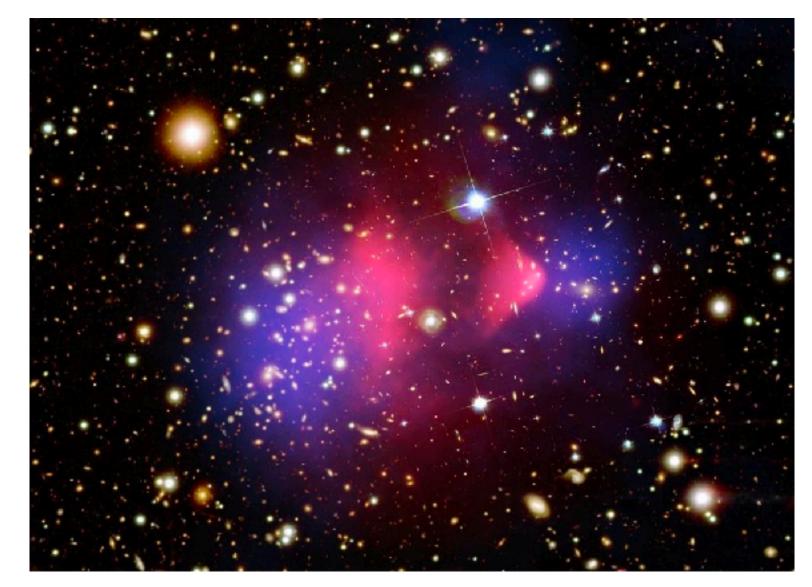


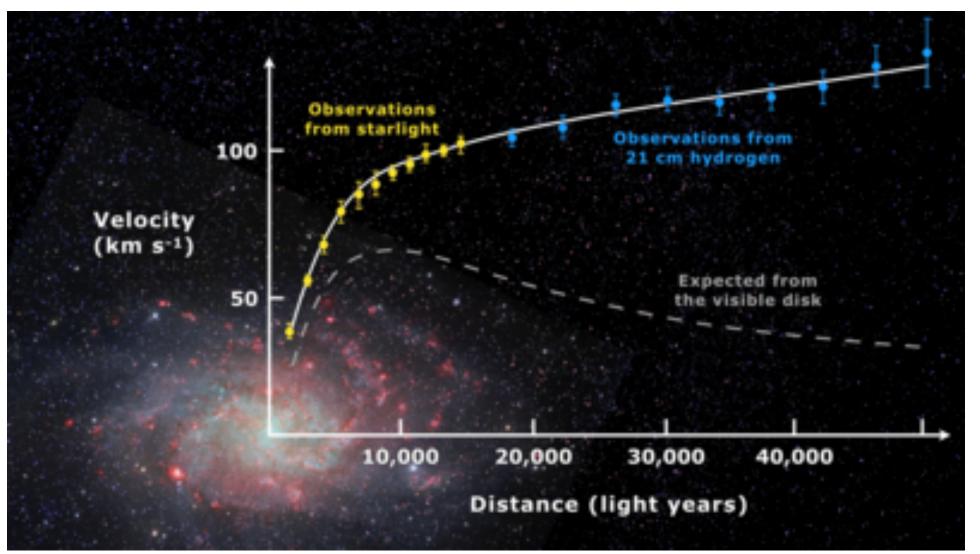




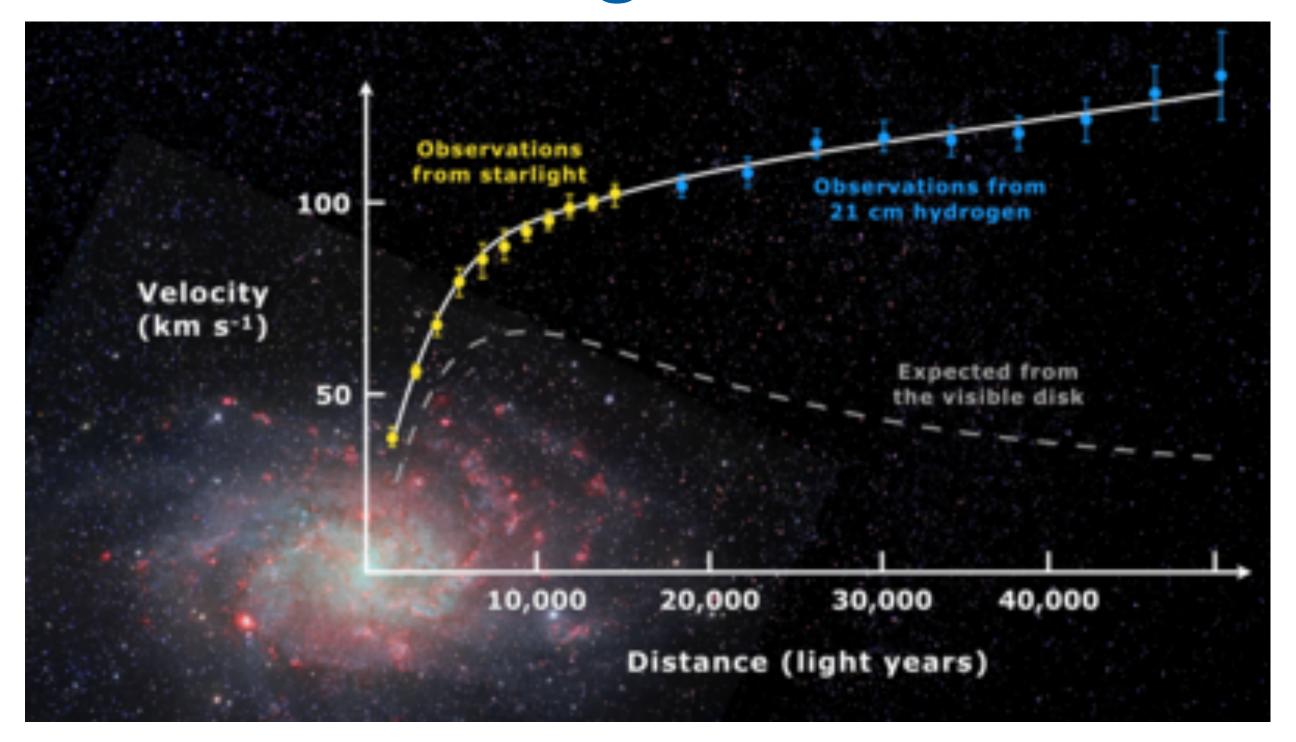




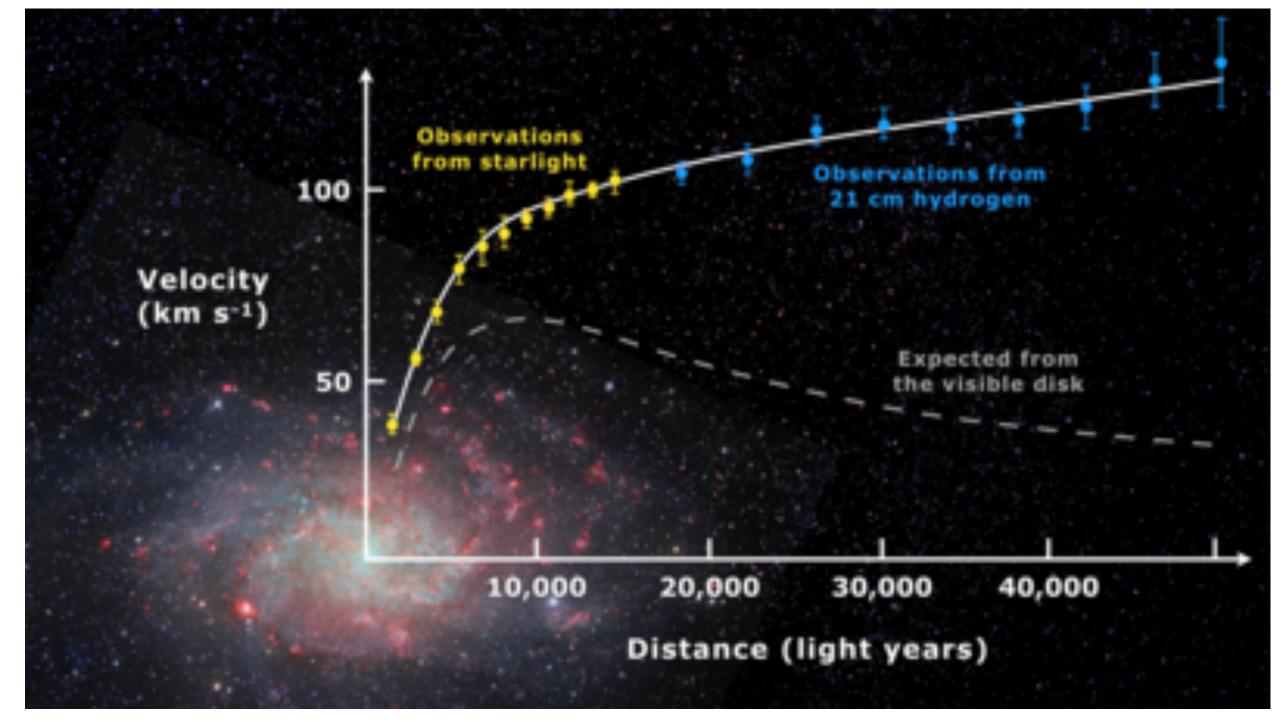






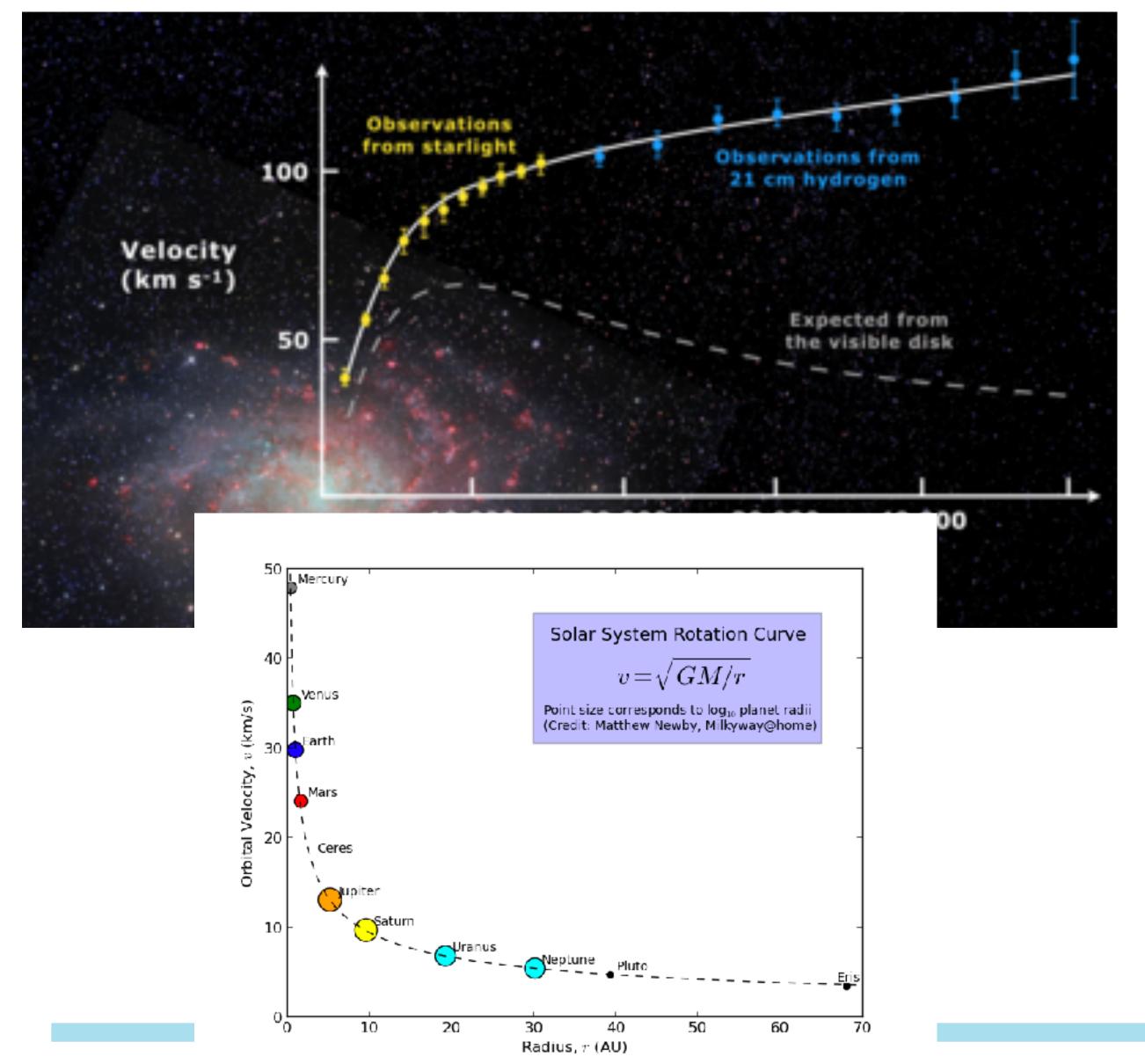






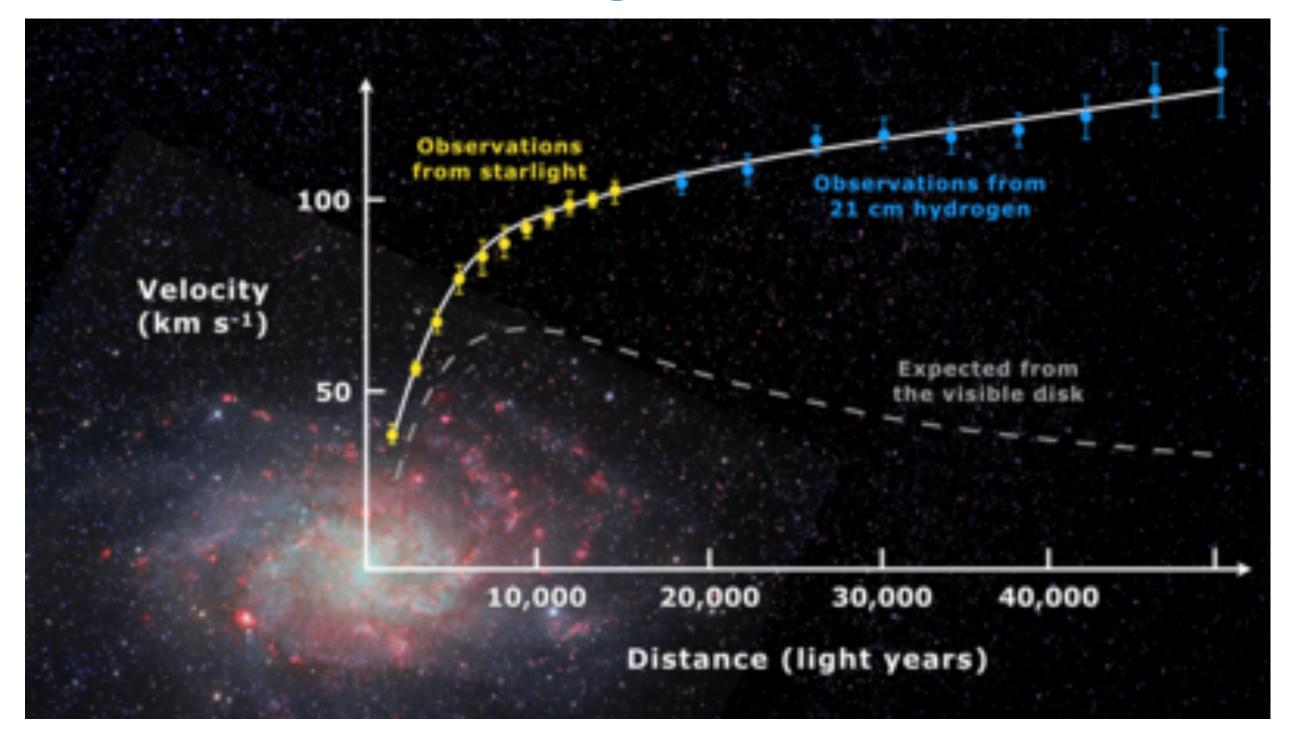
 Galactic rotation curves: one of the first pieces of evidence for dark matter.





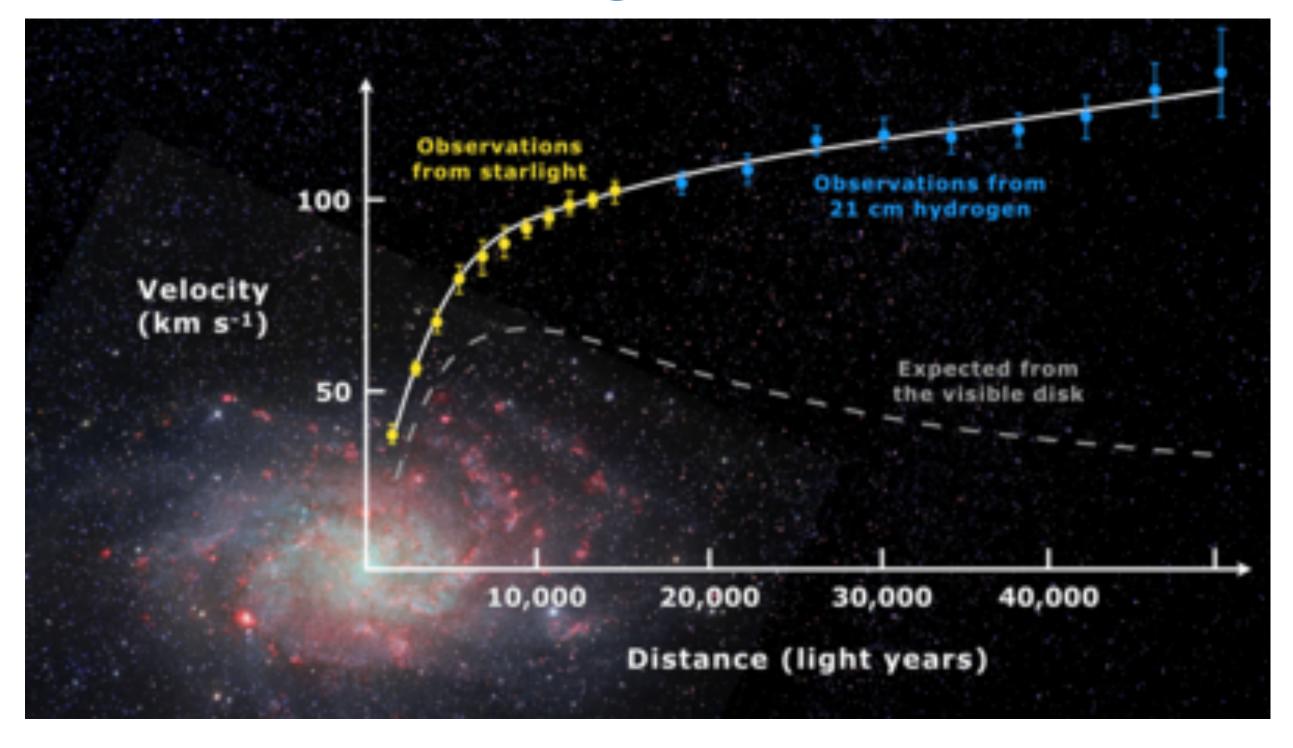
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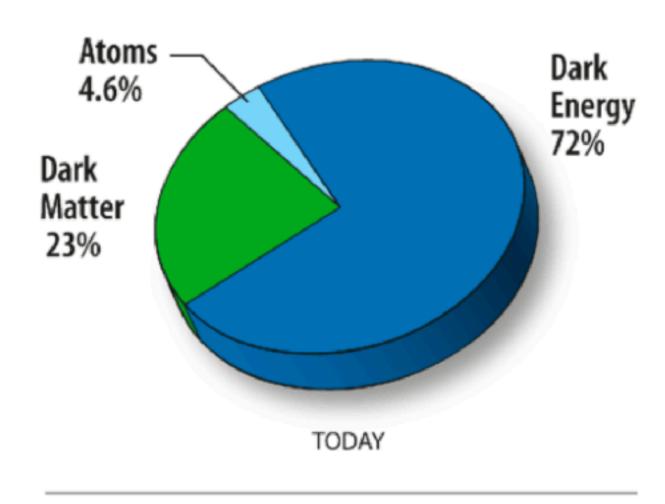
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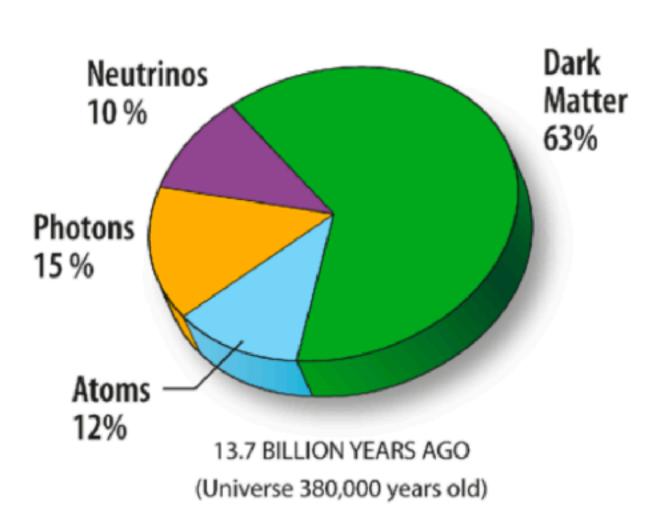




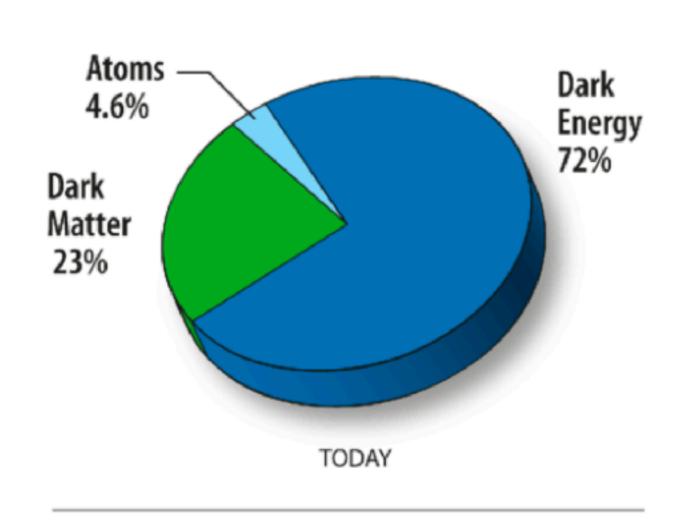
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- This implies some sort of "missing" or "dark" matter that's pulling these stars around.

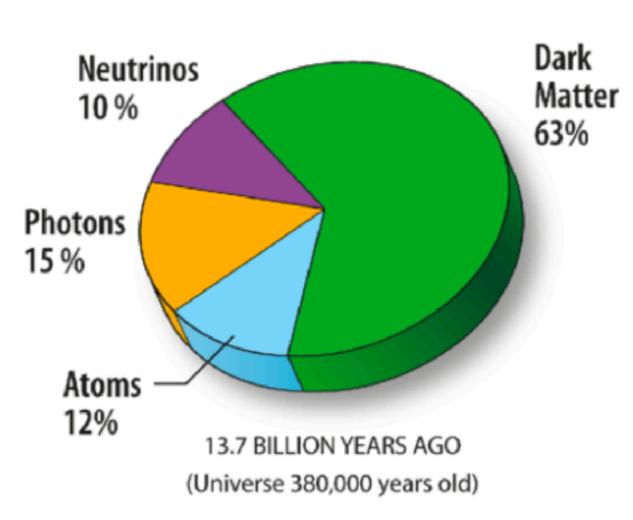






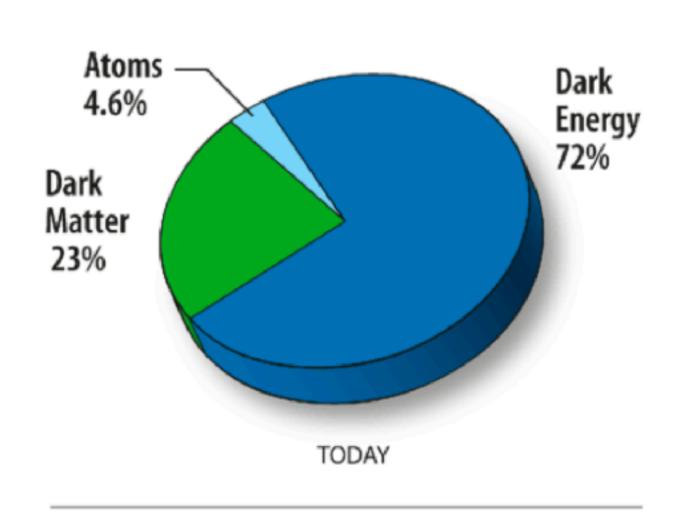


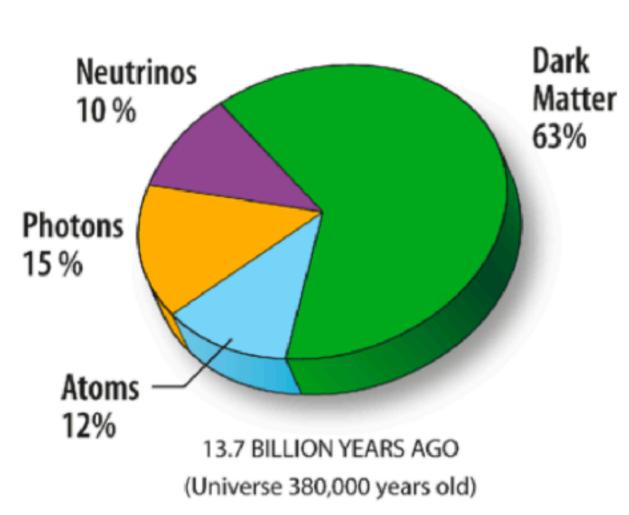




• The particles we have identified in the Standard Model make up a small fraction of all of the observed "energy density" of the universe.

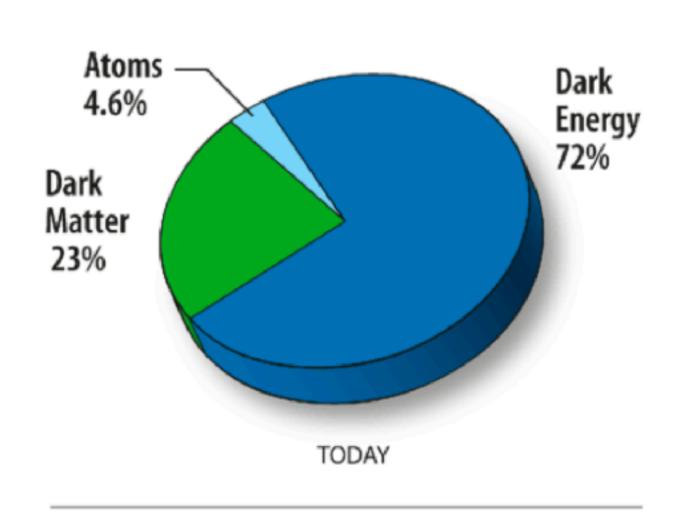


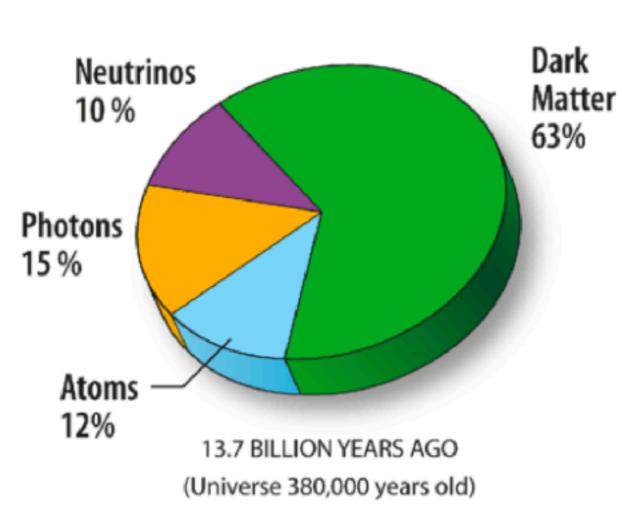




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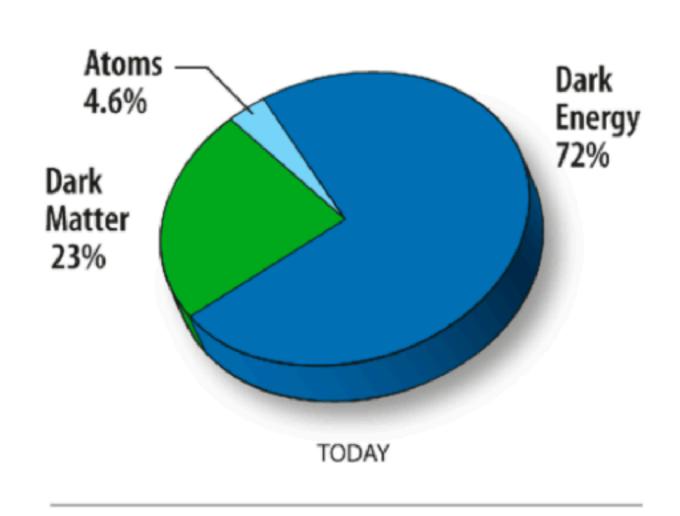


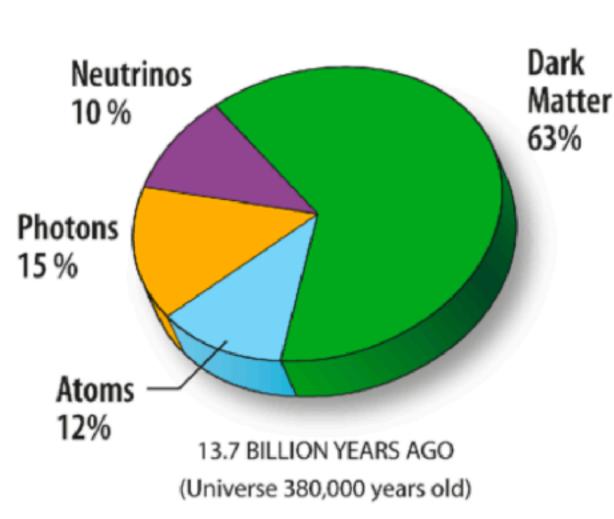




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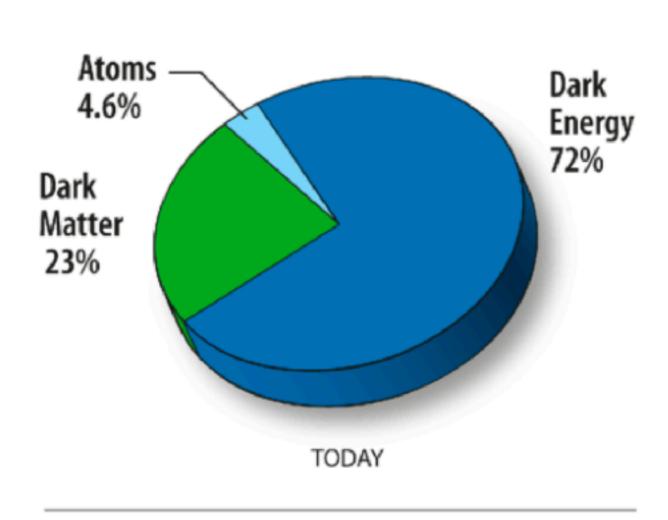


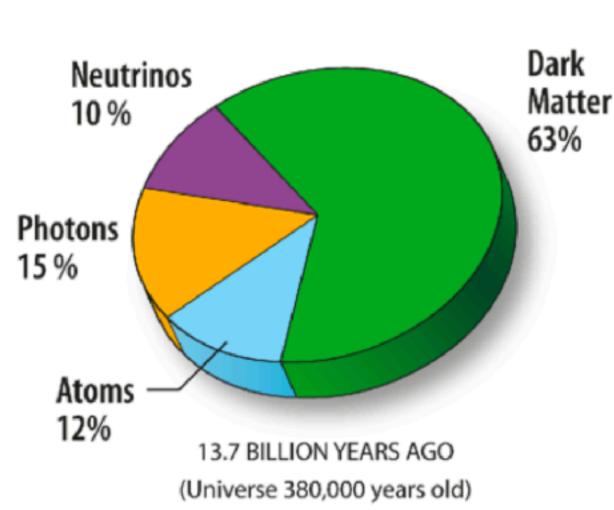


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To the best of our knowledge from these observations,





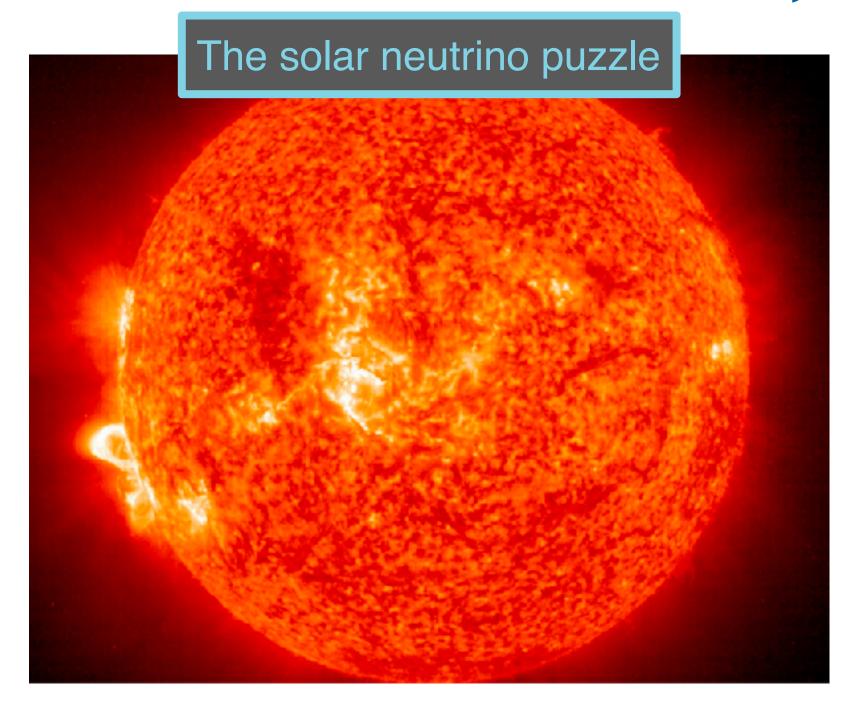
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- Worldwide efforts exist to identify and solve both of these mysteries.



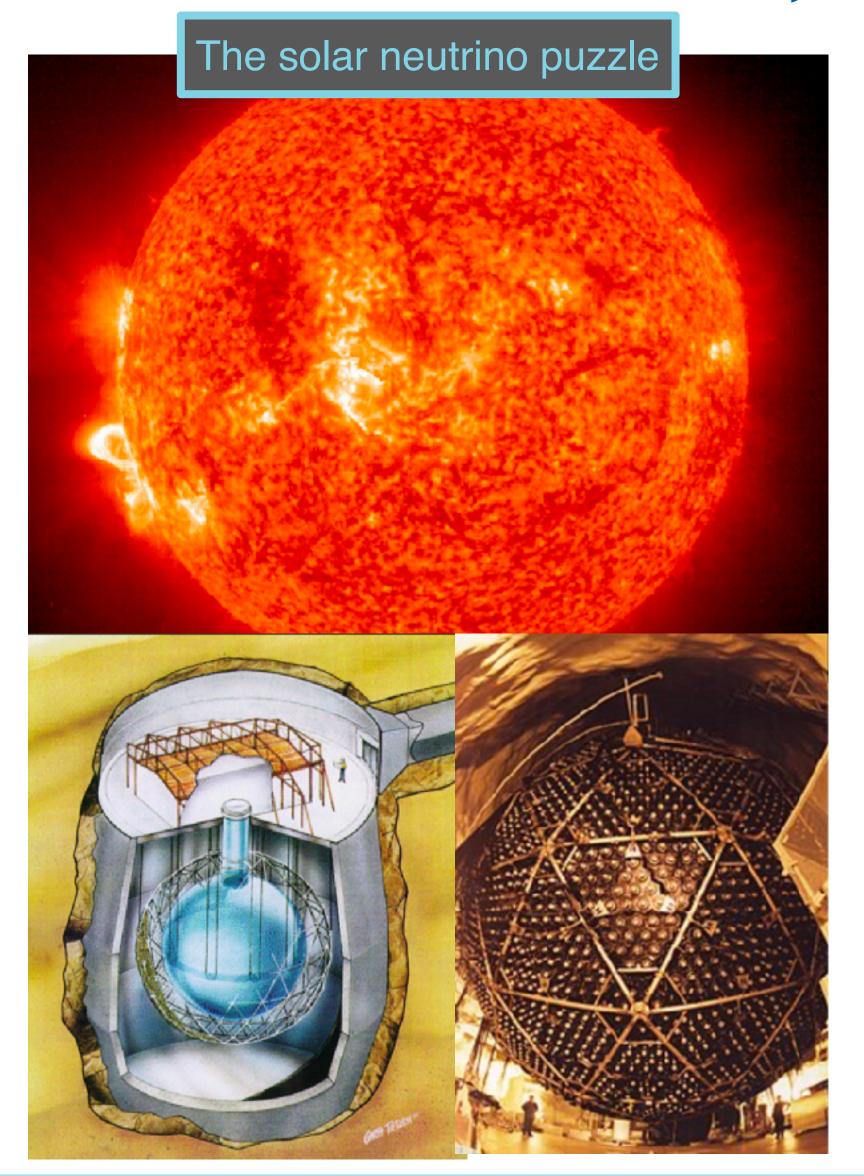
Neutrino Oscillations & Neutrino Masses







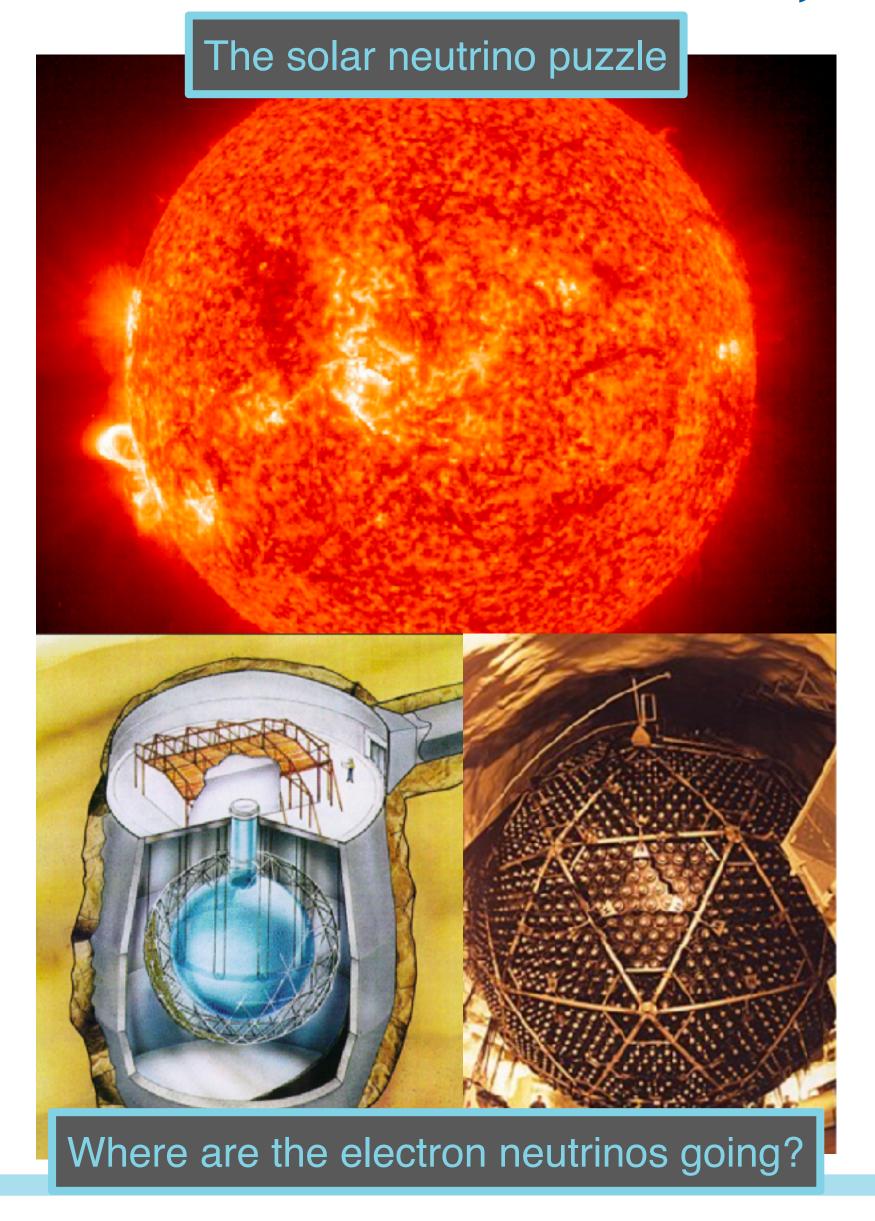


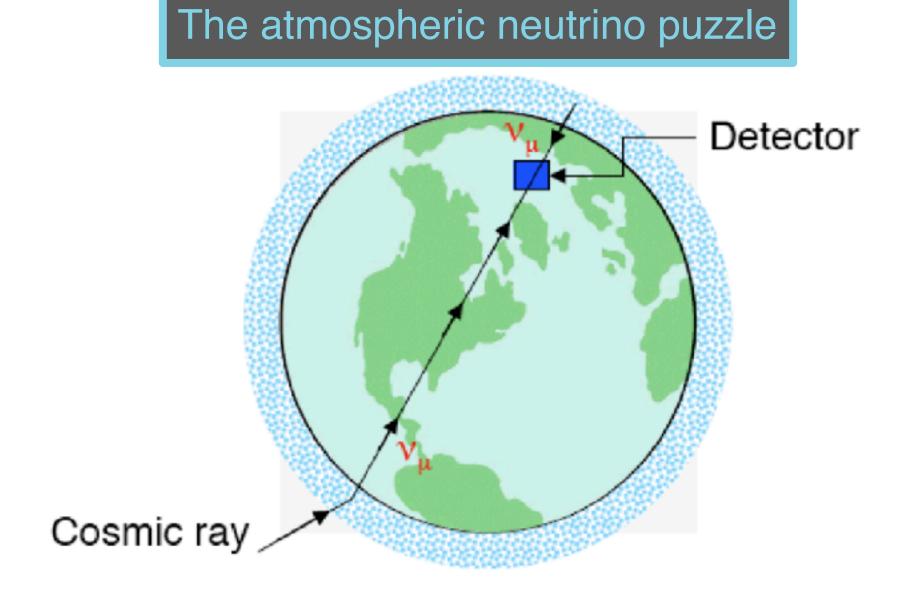




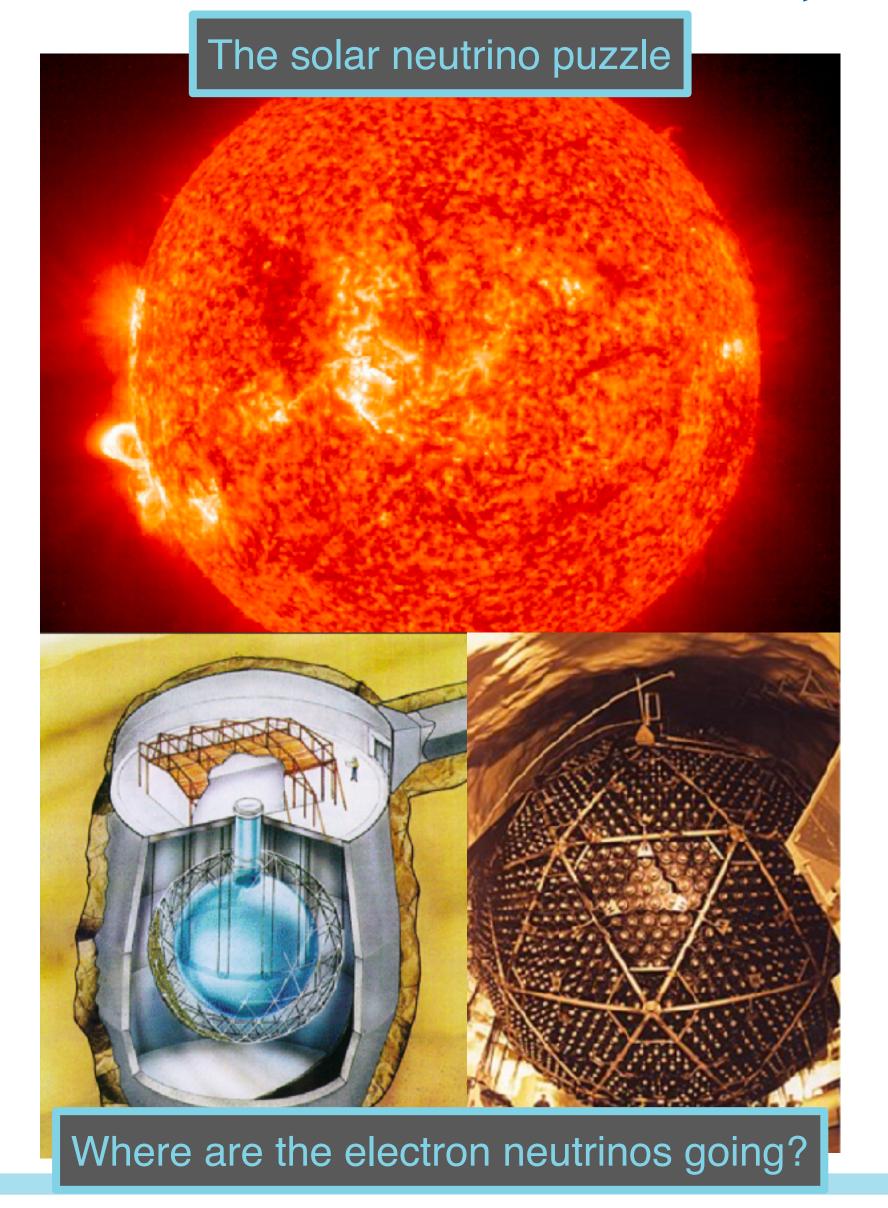


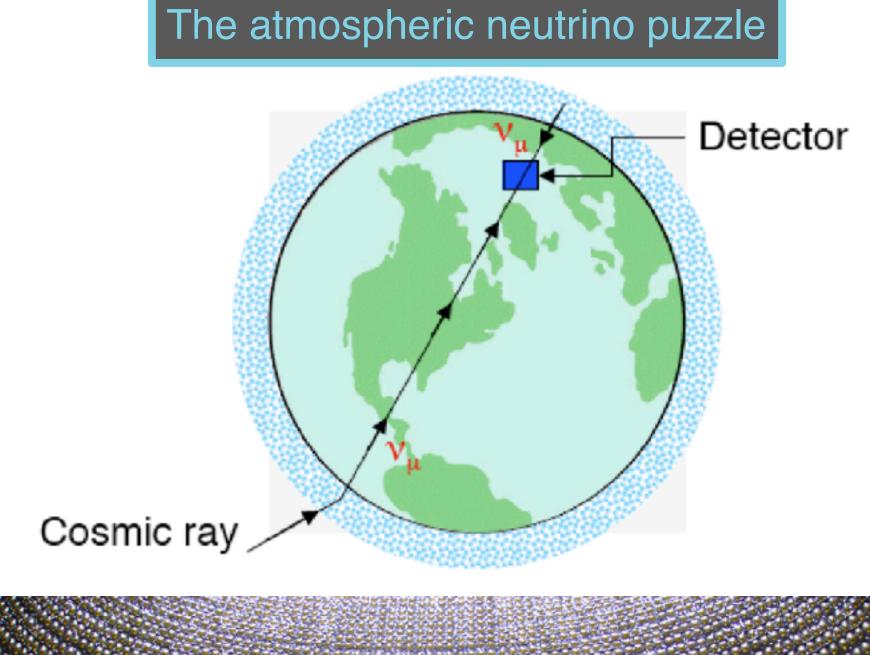


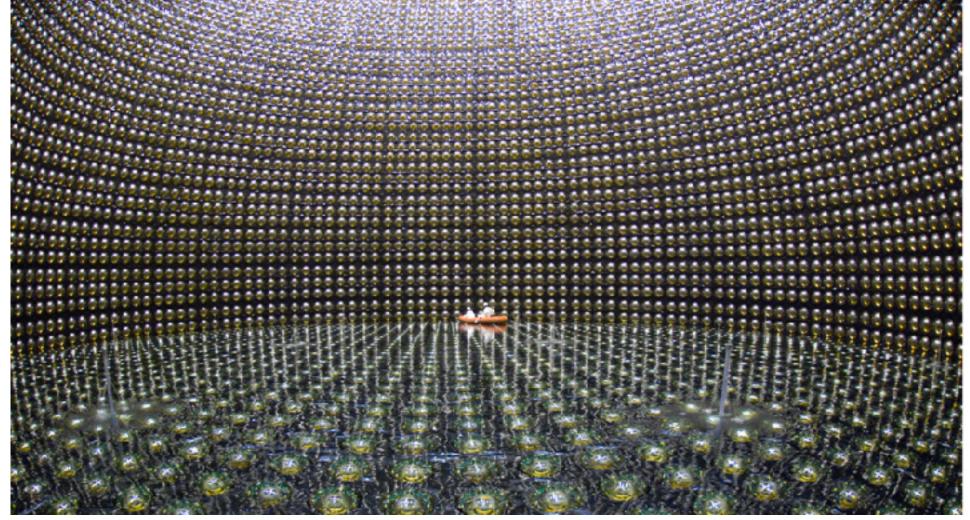






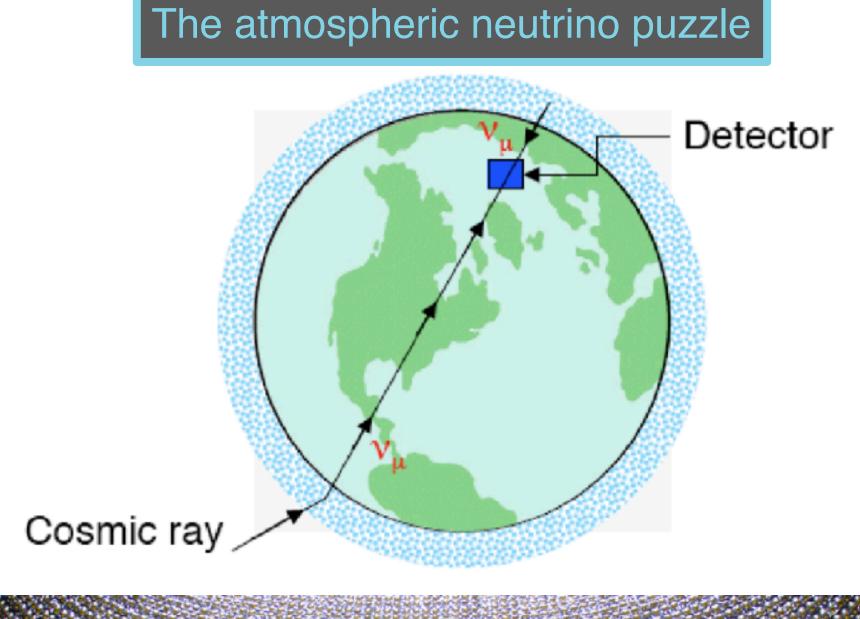


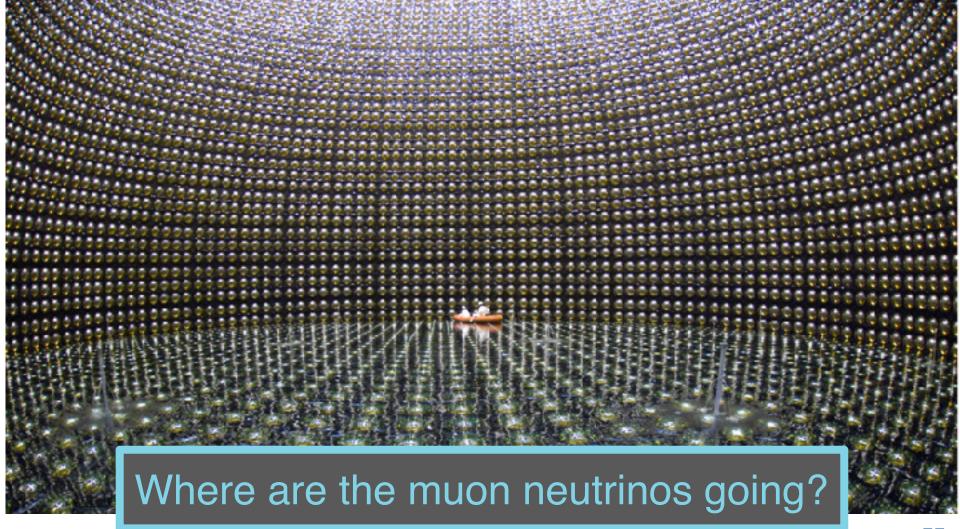




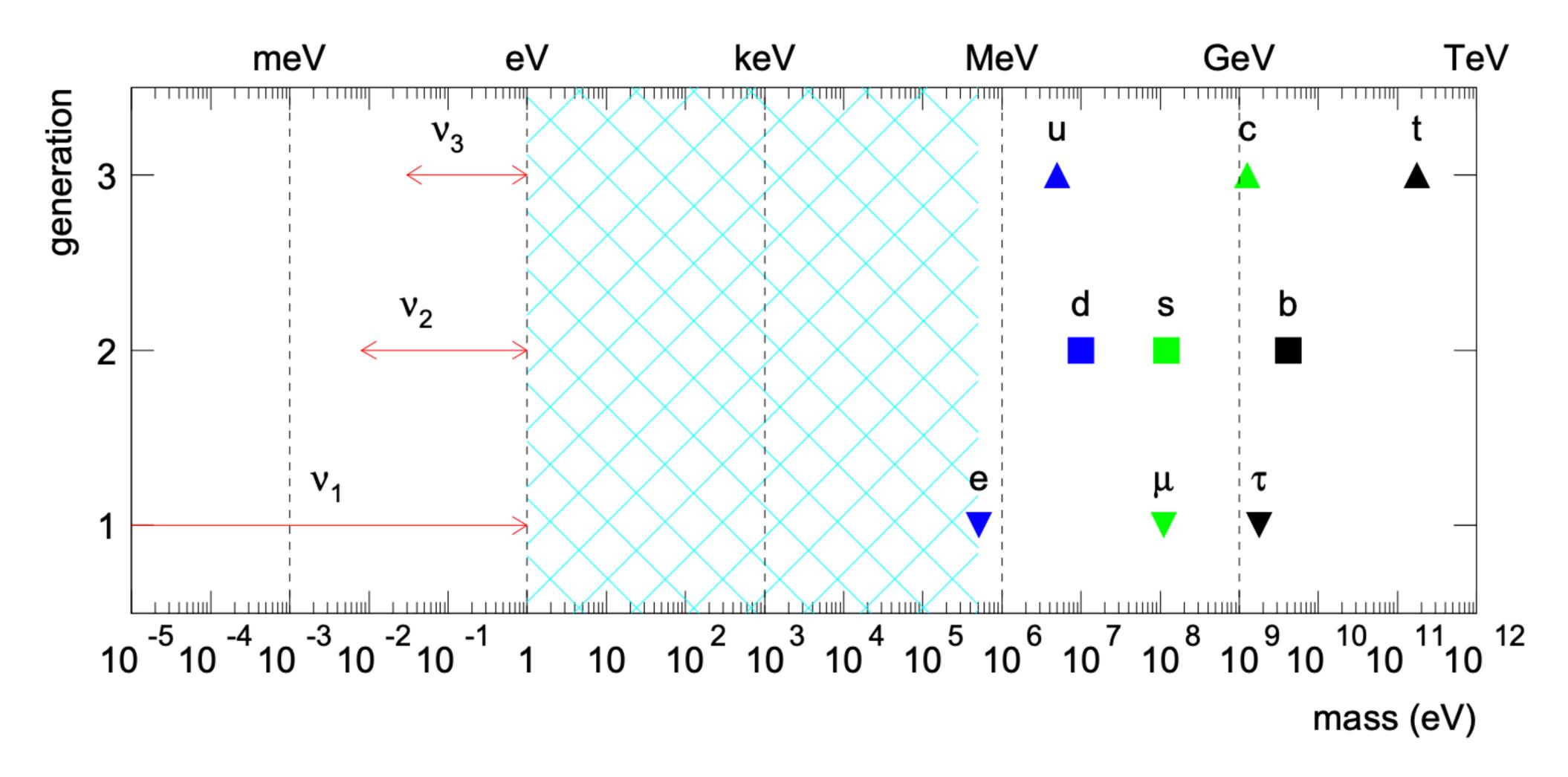




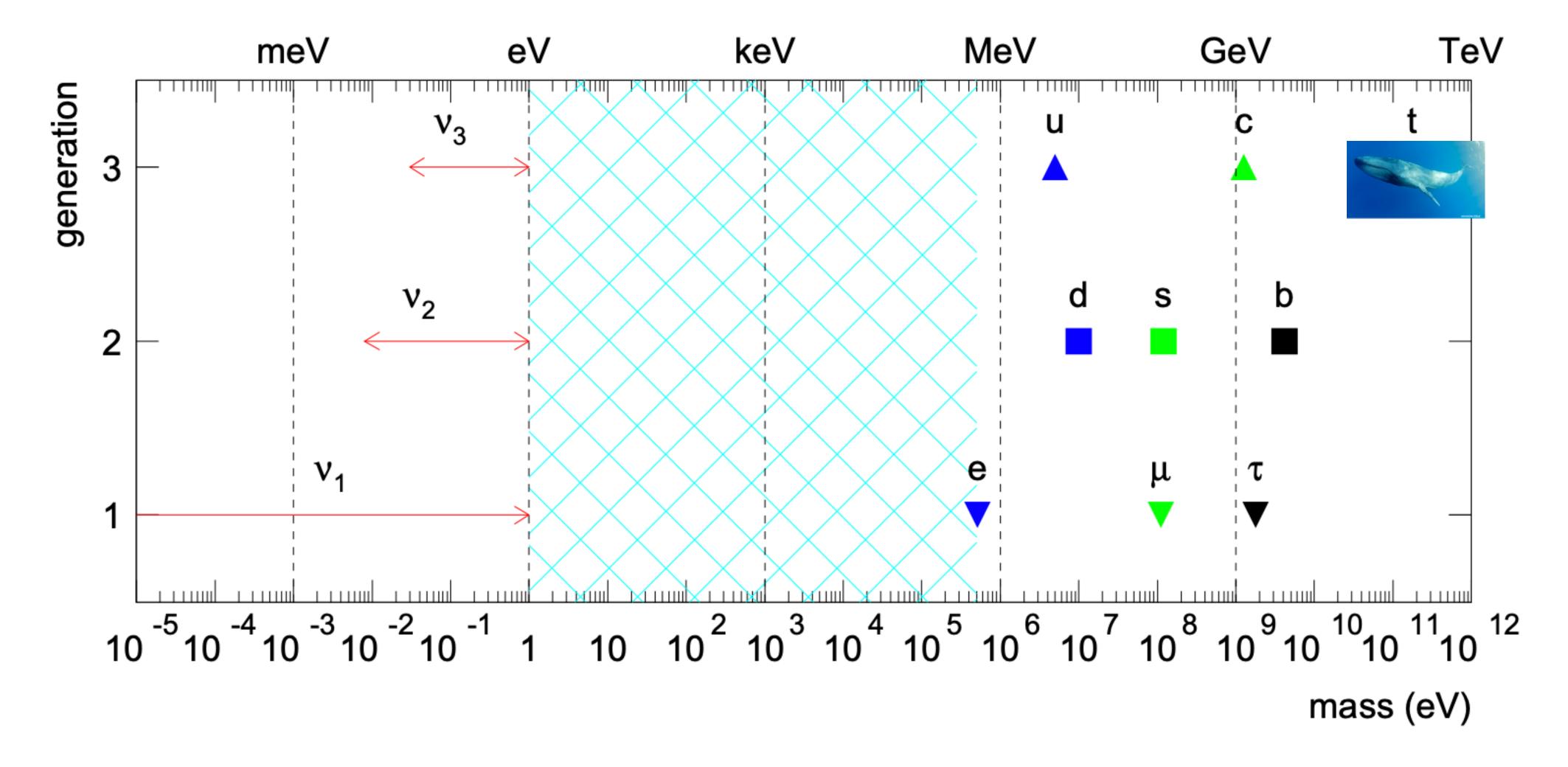




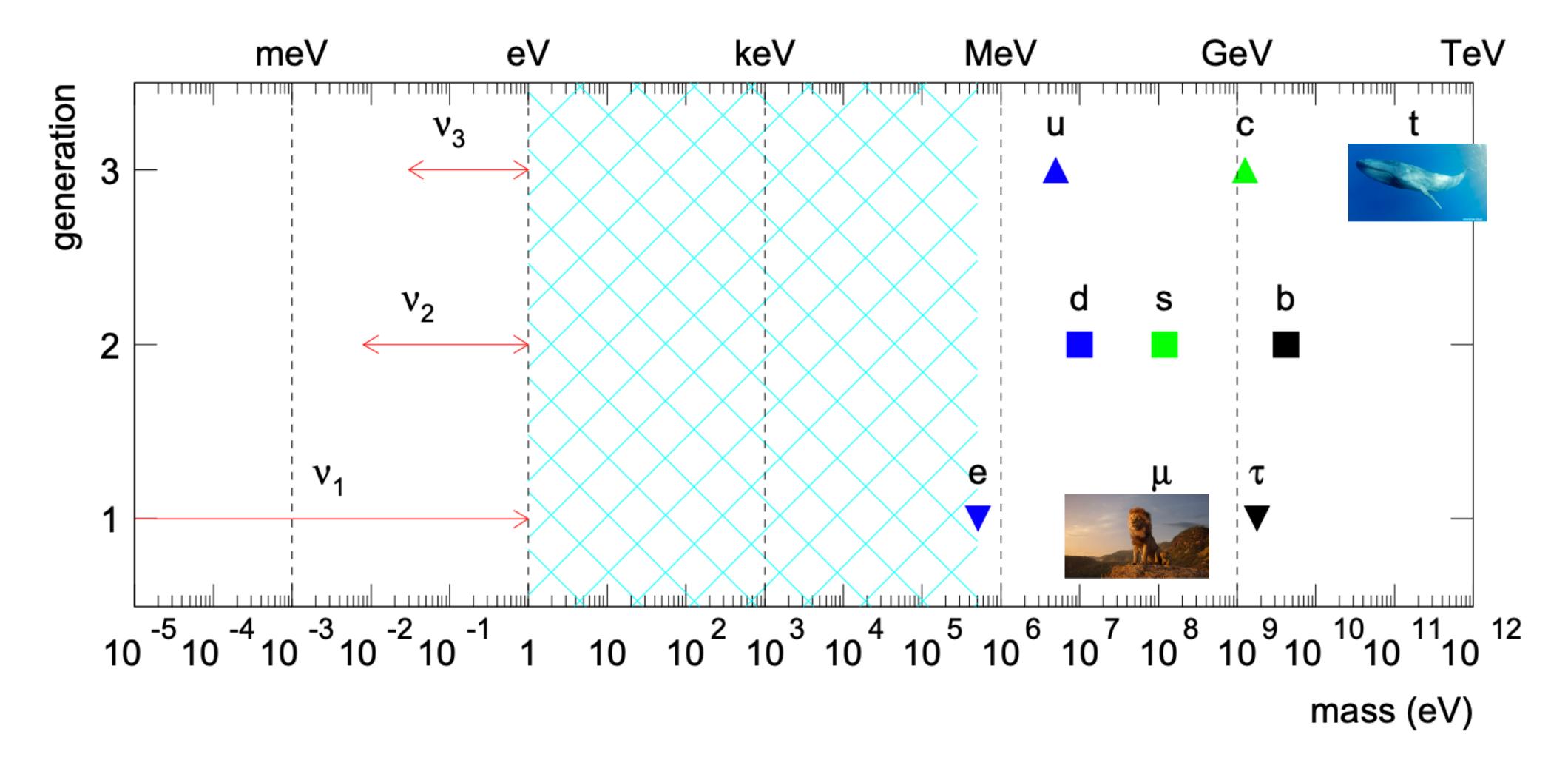




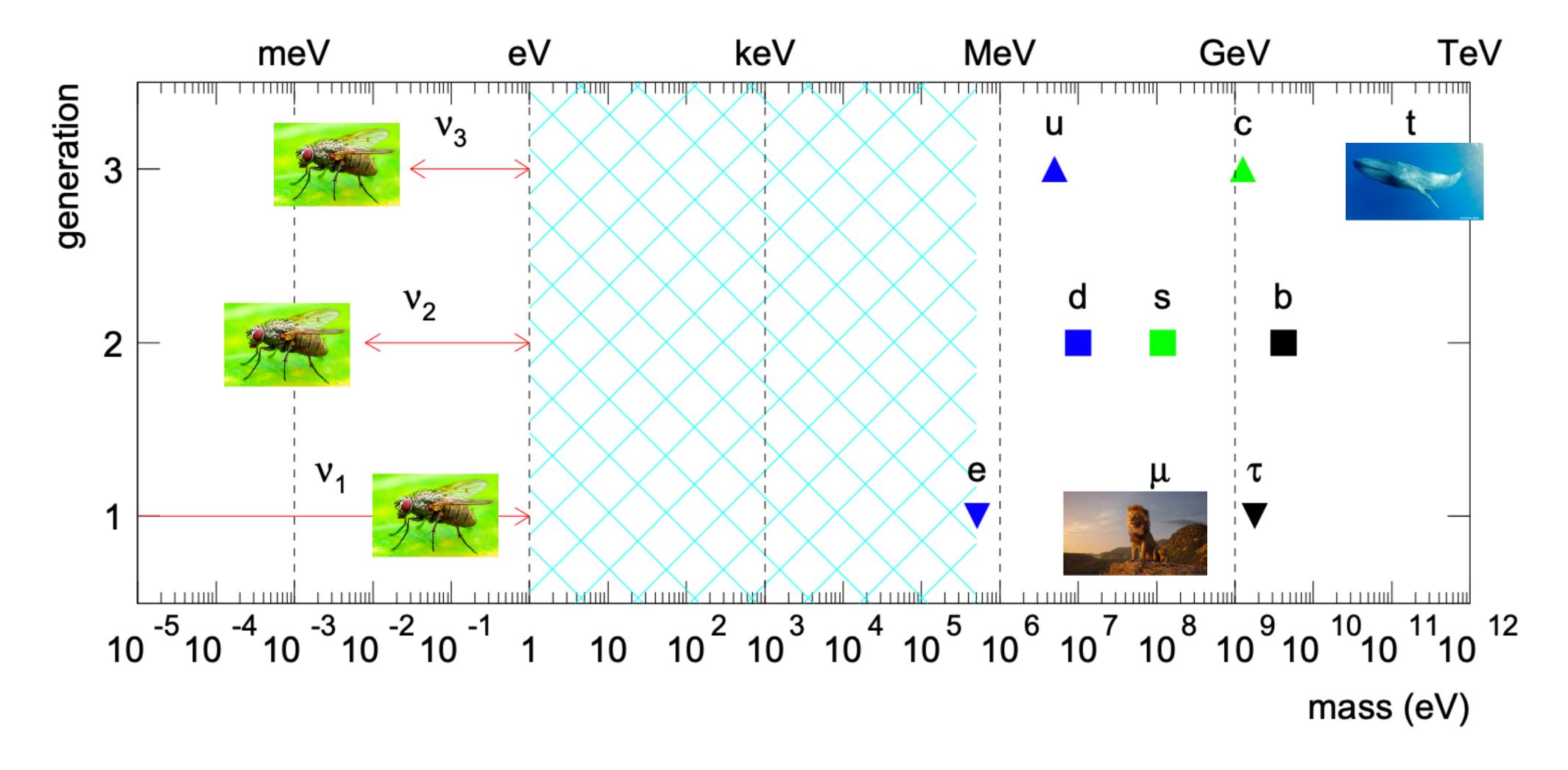




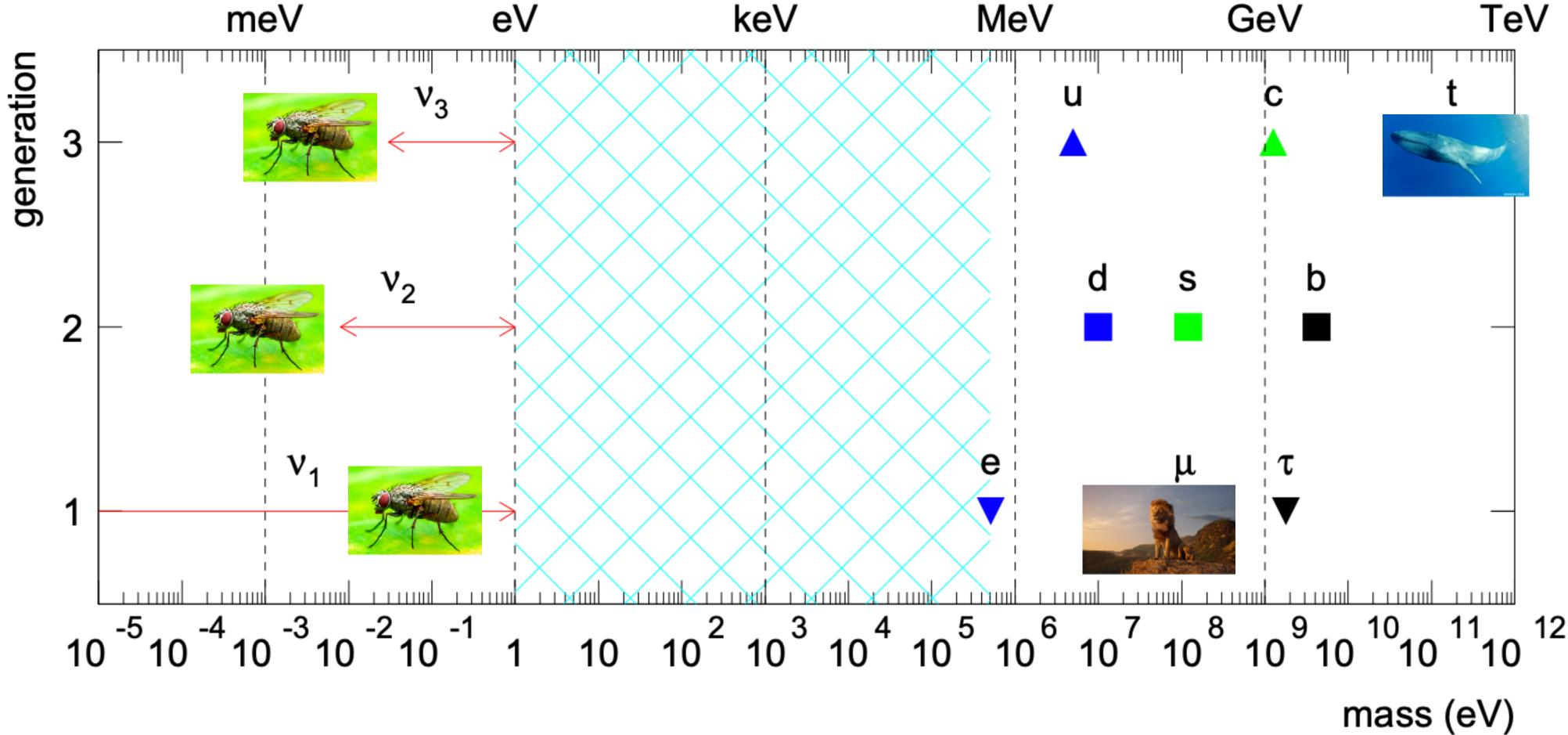








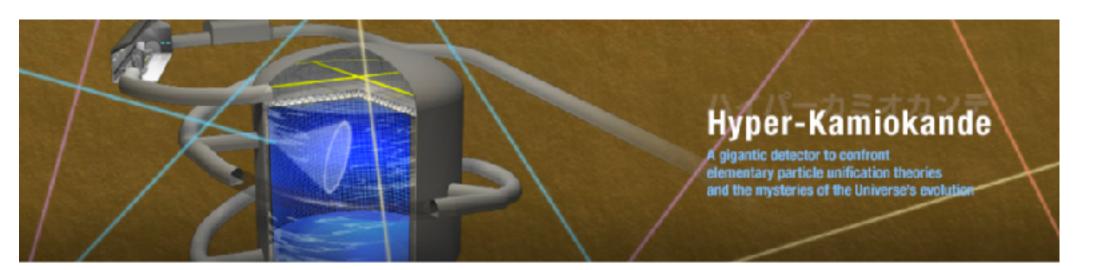




Like Dark Matter, the Standard Model does not predict neutrino masses! New physics (new particles and/or interactions) are necessary)

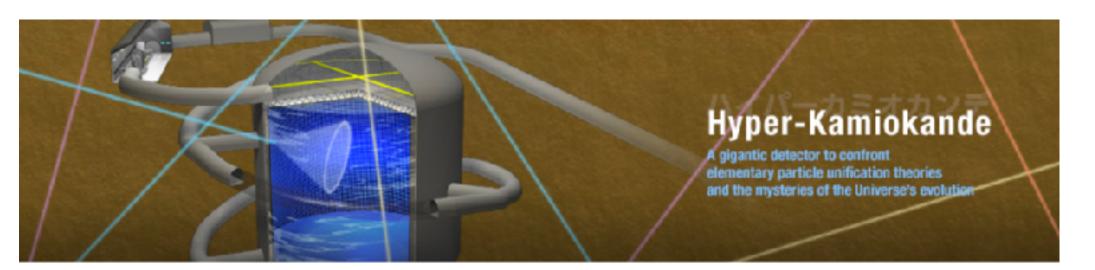








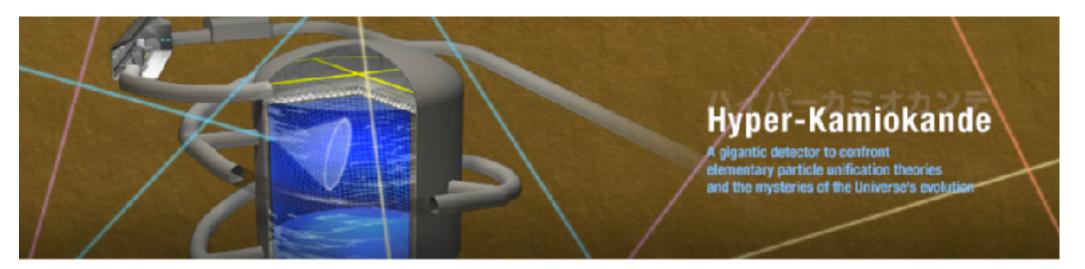




 Successor to the wildly successful Super-Kamiokande Experiment in Japan.



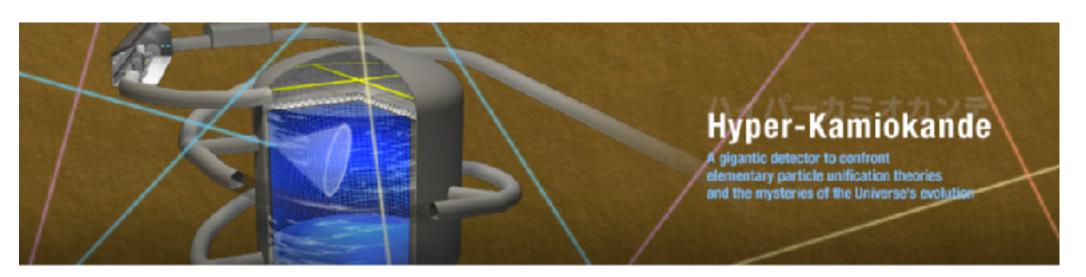
Hyper-Kamiokande



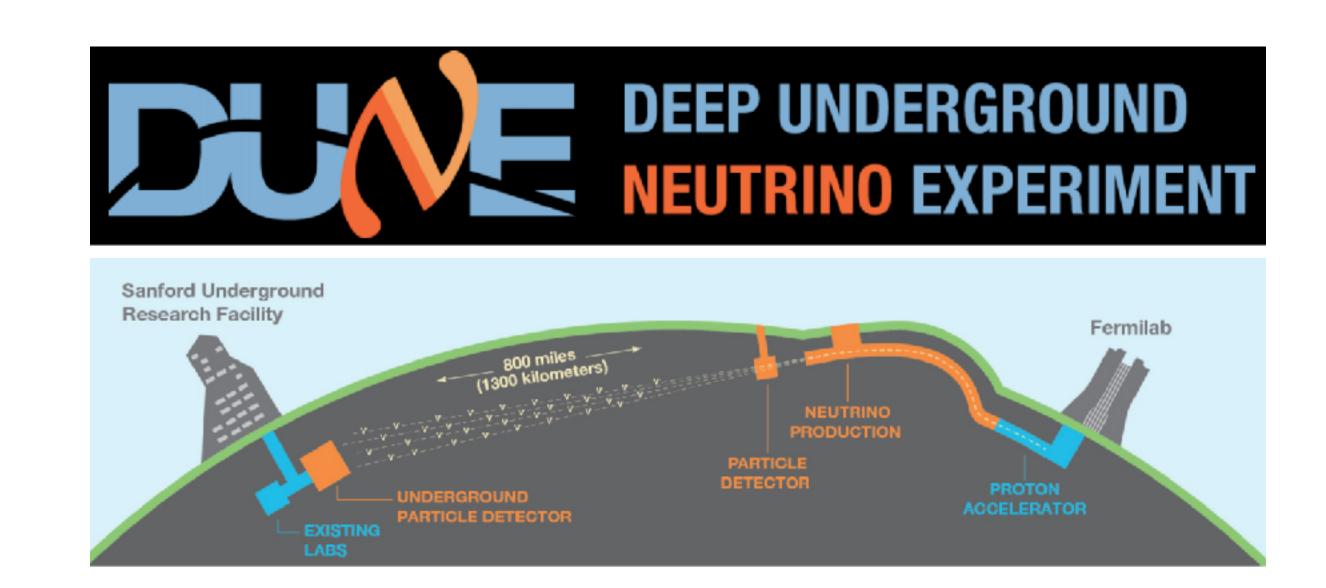
- Successor to the wildly successful Super-Kamiokande Experiment in Japan.
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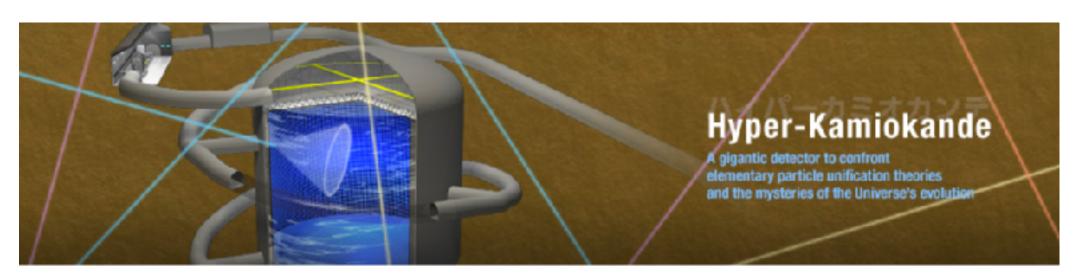


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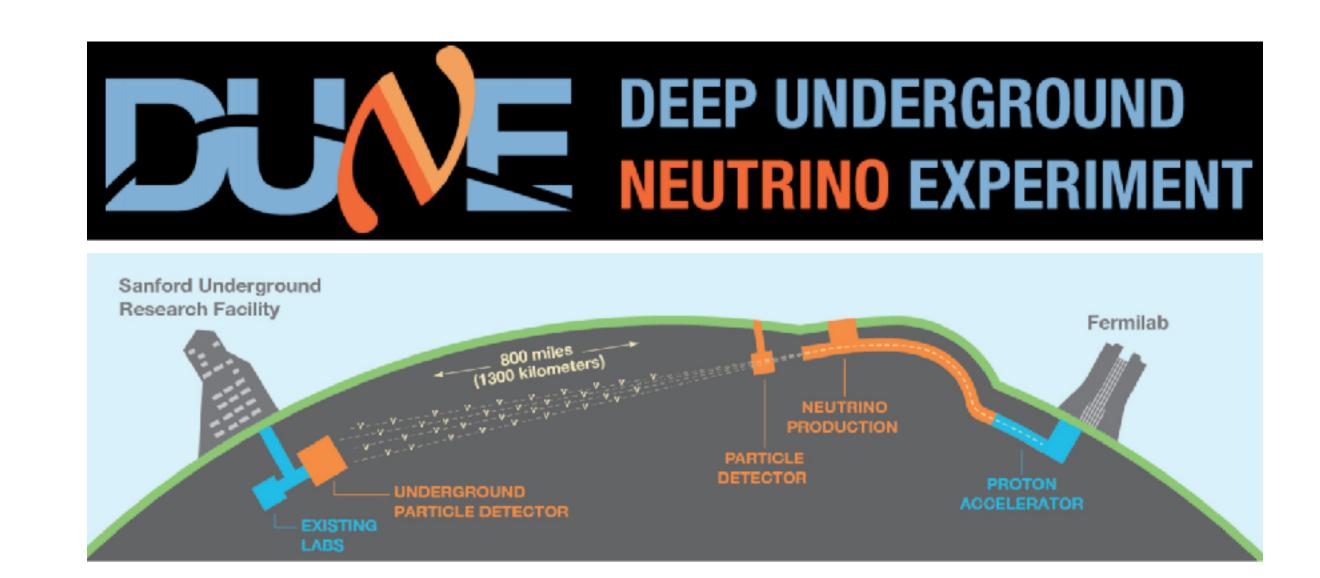








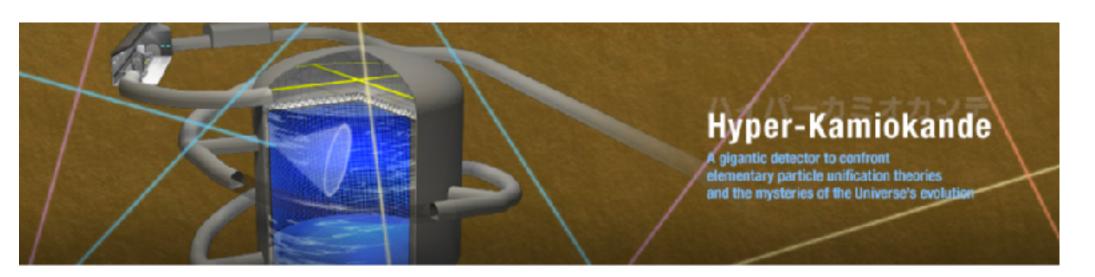
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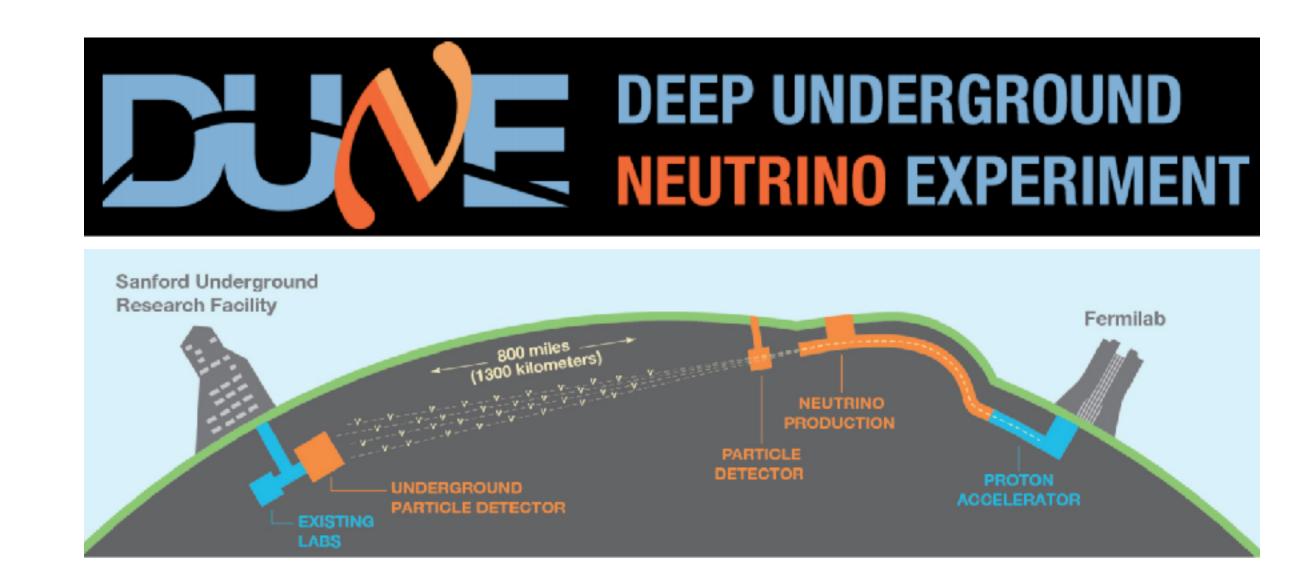
US-based precision neutrino physics project.







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- US-based precision neutrino physics project.
- Built on new, liquid-argon based technology being developed in Fermilabbased "short-baseline" neutrino program.

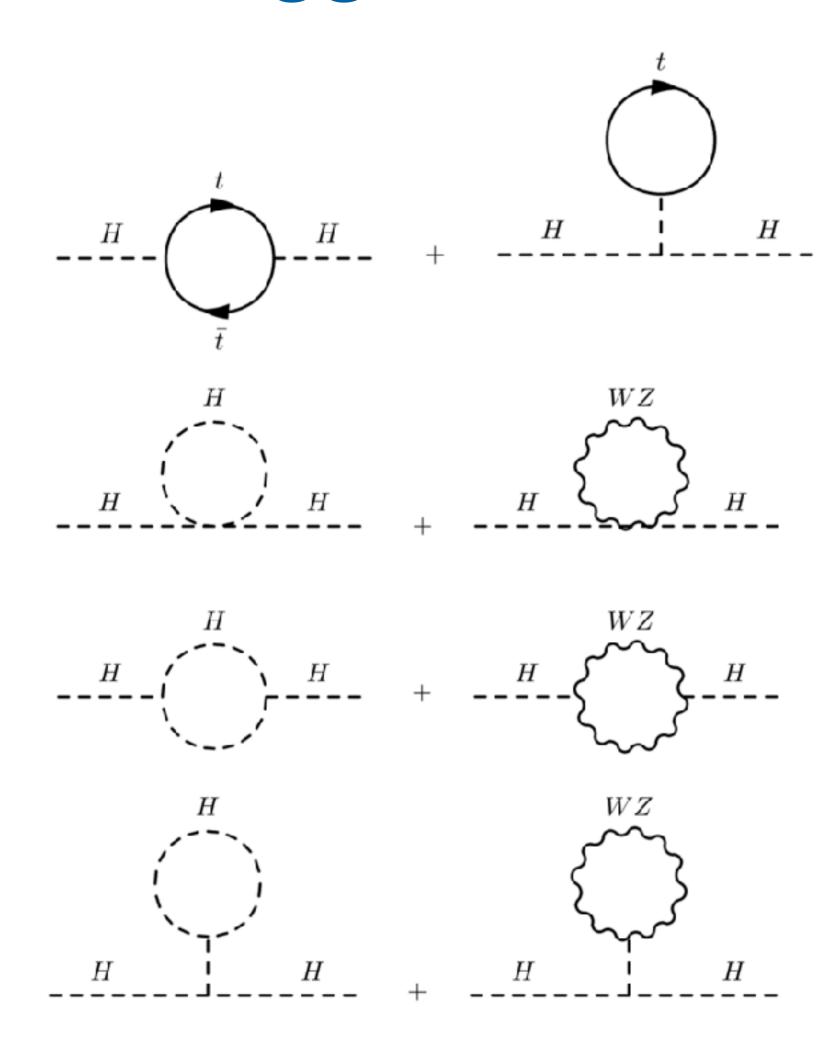


The Higgs Hierarchy Problem





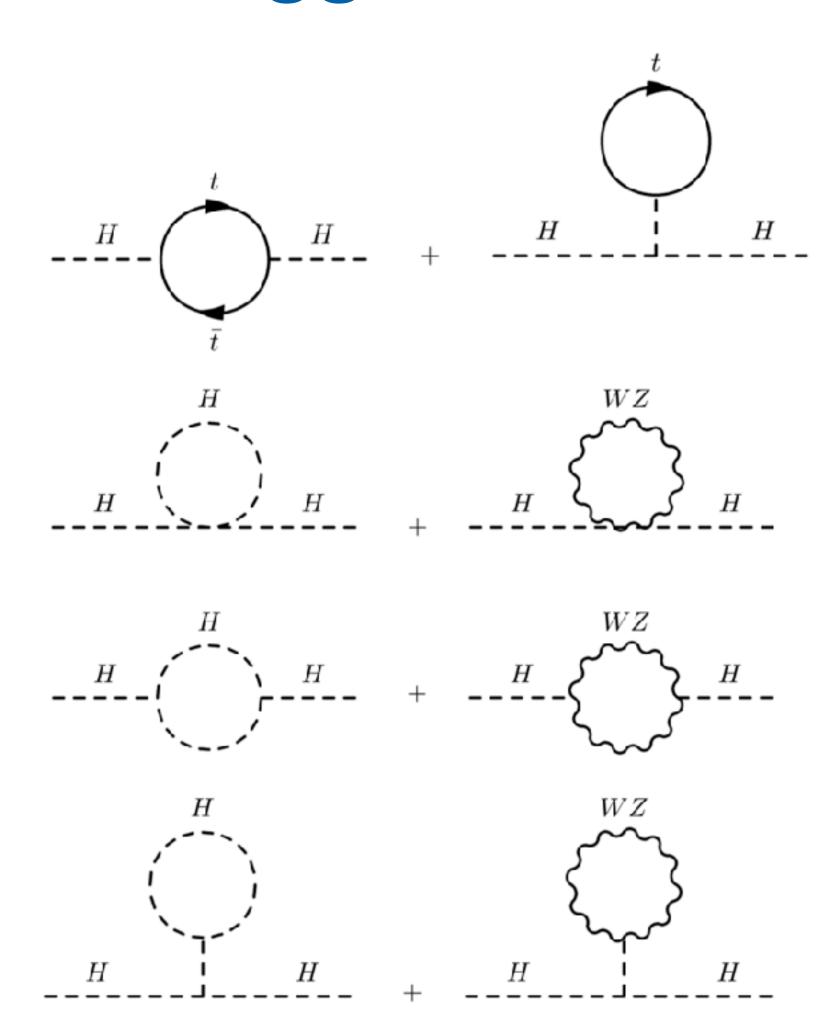




 Quantum field theory predicts that processes like these modify the Higgs boson's mass from its "Lagrangian value".

$$\Delta m_H^2 \approx -\frac{y_t^2}{8\pi^2} \Lambda^2 + \cdots$$

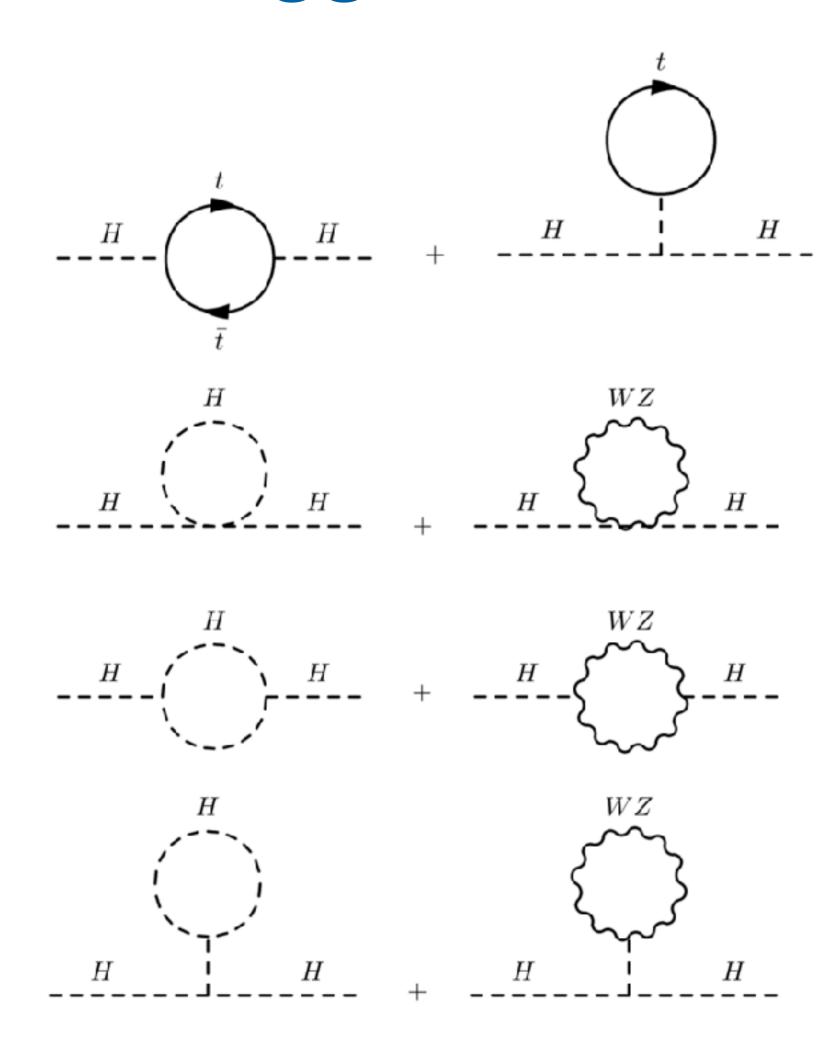




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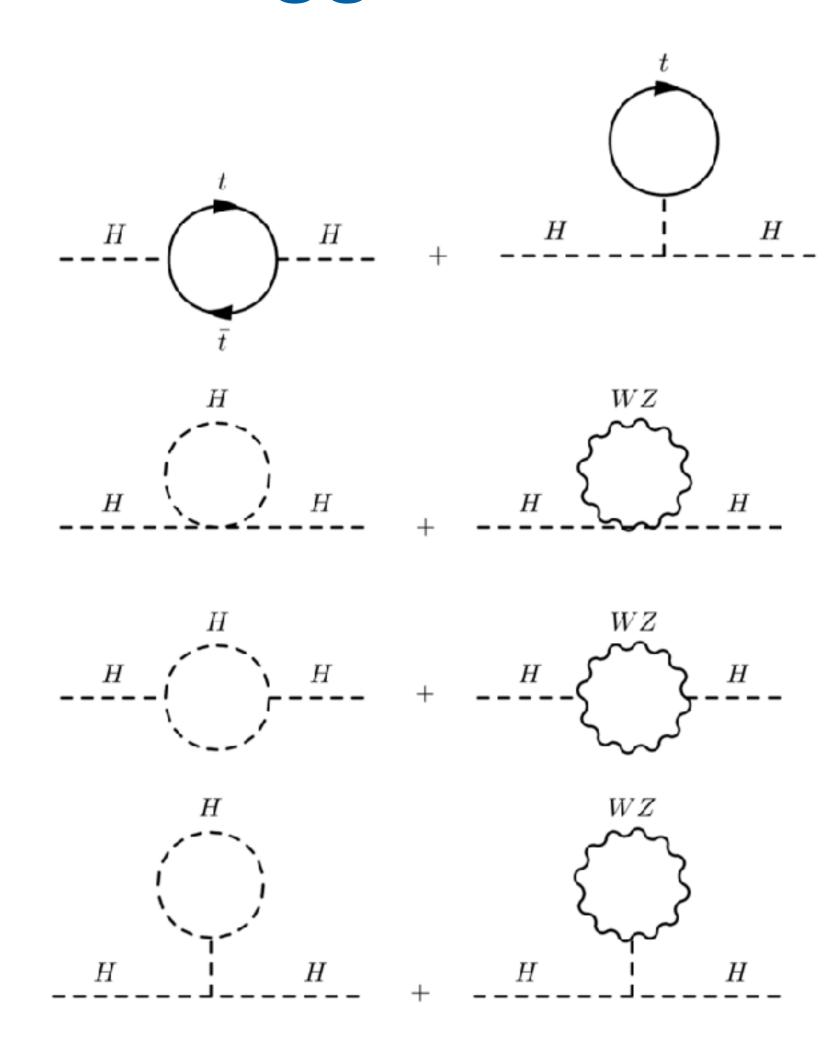




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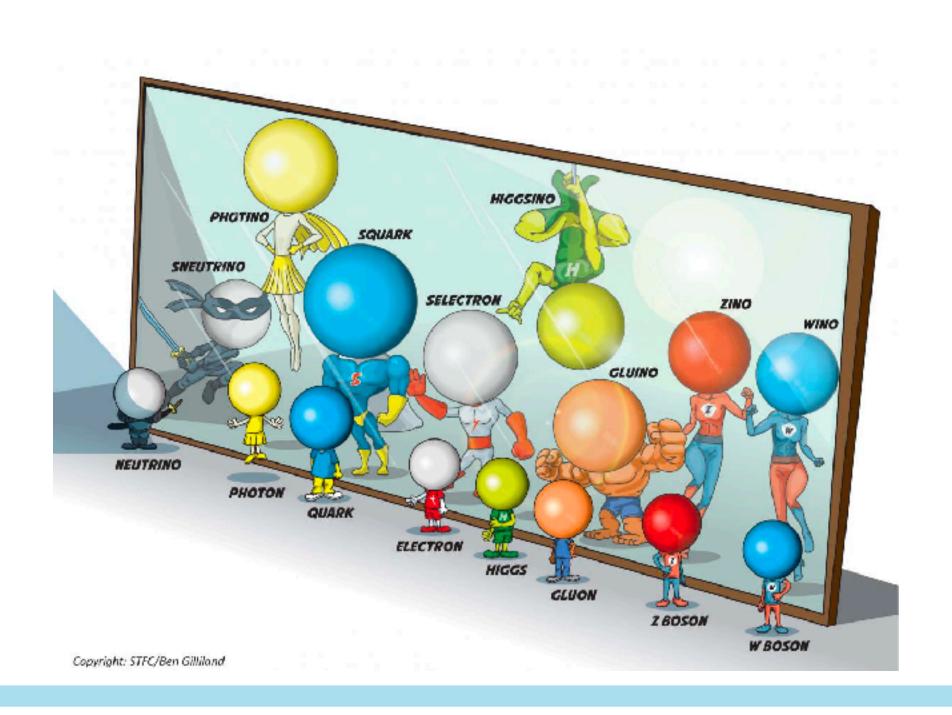
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- Why then, do we see it to be so small?

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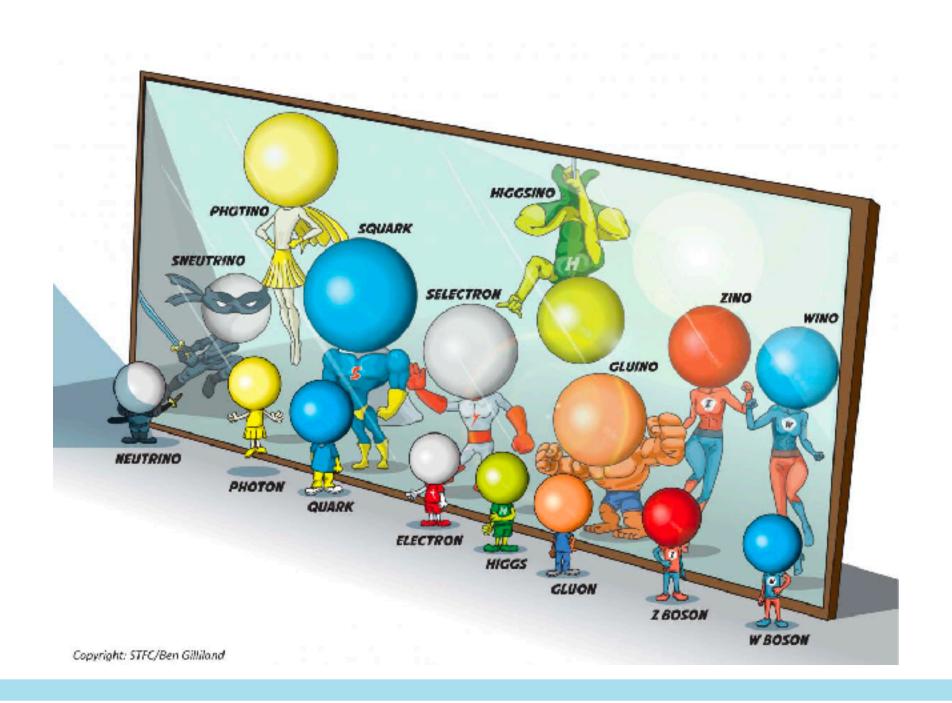


• Every standard model particle has a super-symmetric partner. These superparticles perfectly cancel out all of these quantum corrections to the Higgs mass.



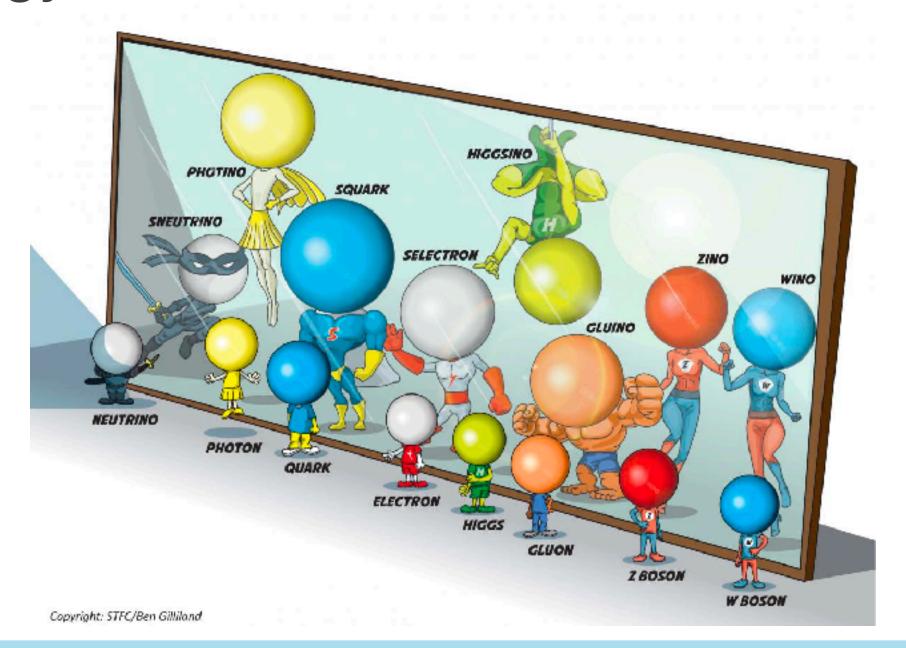


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- Many theories of supersymmetry also include a stable particle that can explain the observed Dark Matter in the universe.



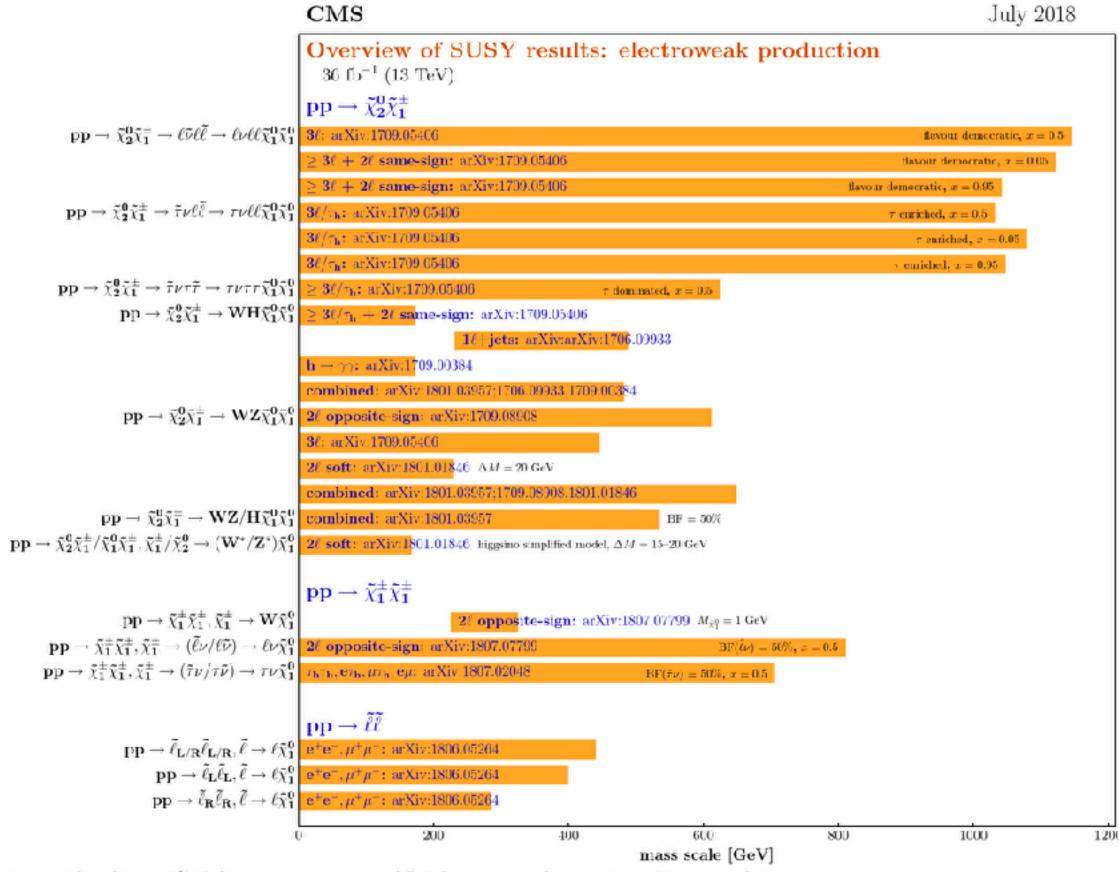


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- Supersymmetric theories also may help point to a "theory of everything", where all of the observed "low-energy" forces are unified into one at high energy.



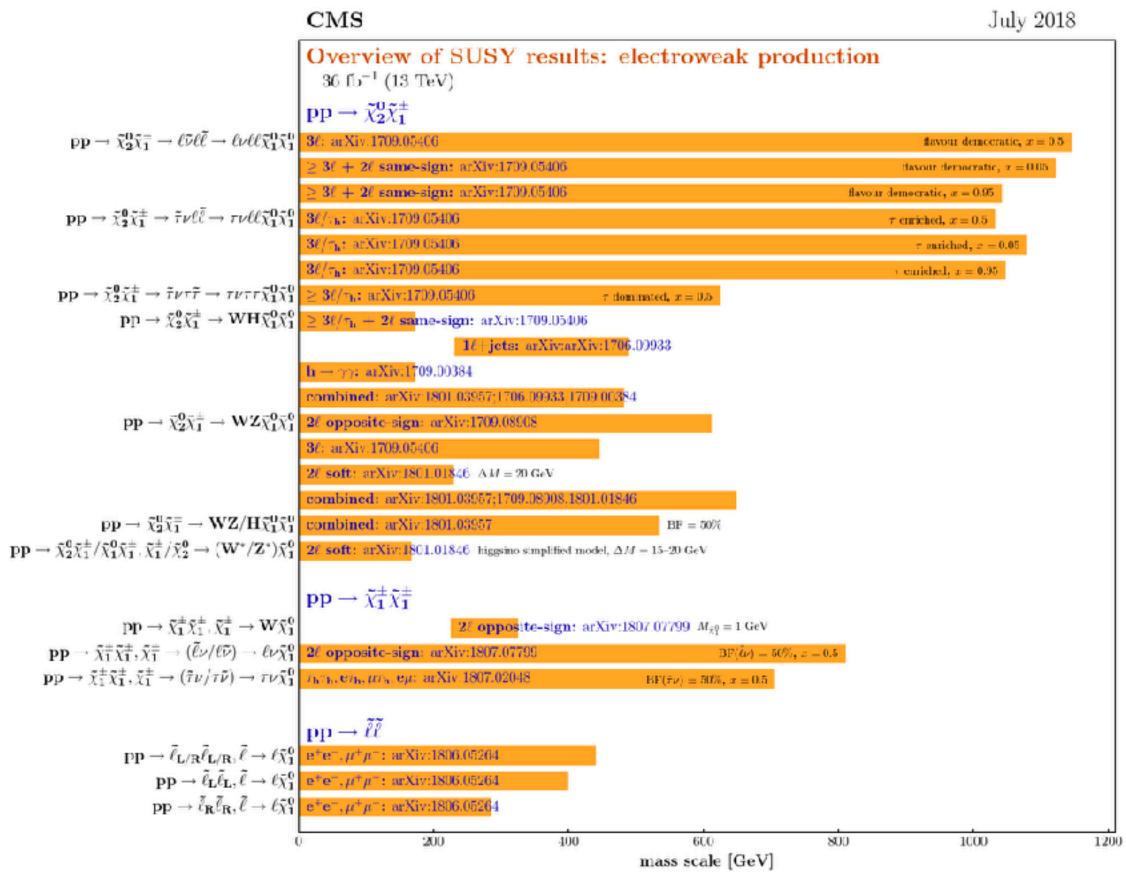






Selection of observed limits at 95% C.L. (theory uncertainties are not included). Probe up to the quoted mass limit for light LSPs unless stated otherwise. The quantities ΔM and z represent the absolute mass difference between the primary sparticle and the LSP, and the difference between the intermediate sparticle and the LSP relative to ΔM , respectively, unless indicated otherwise.





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ATLAS SUSY Searches* - 95% CL Lower Limits

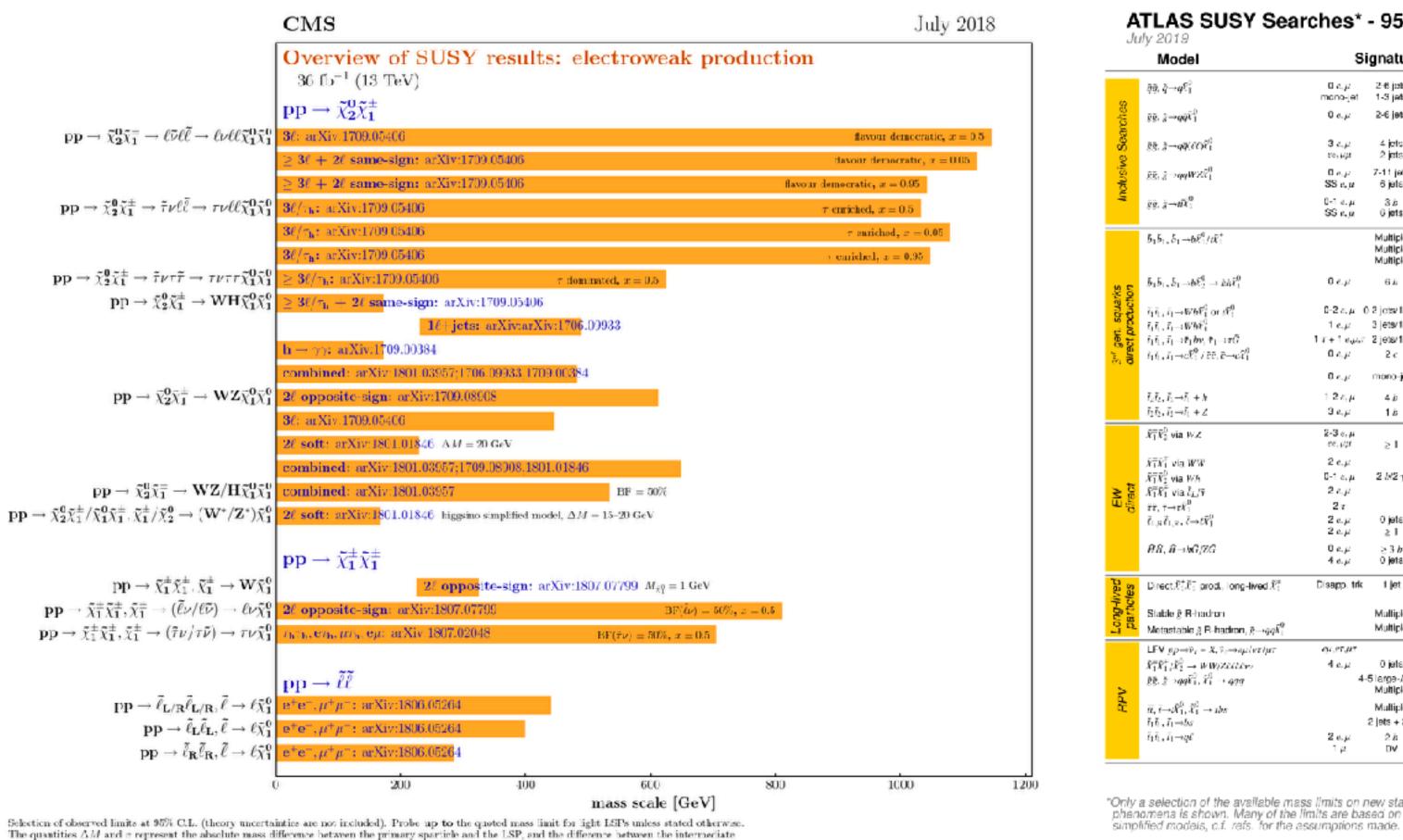
ATLAS Preliminary

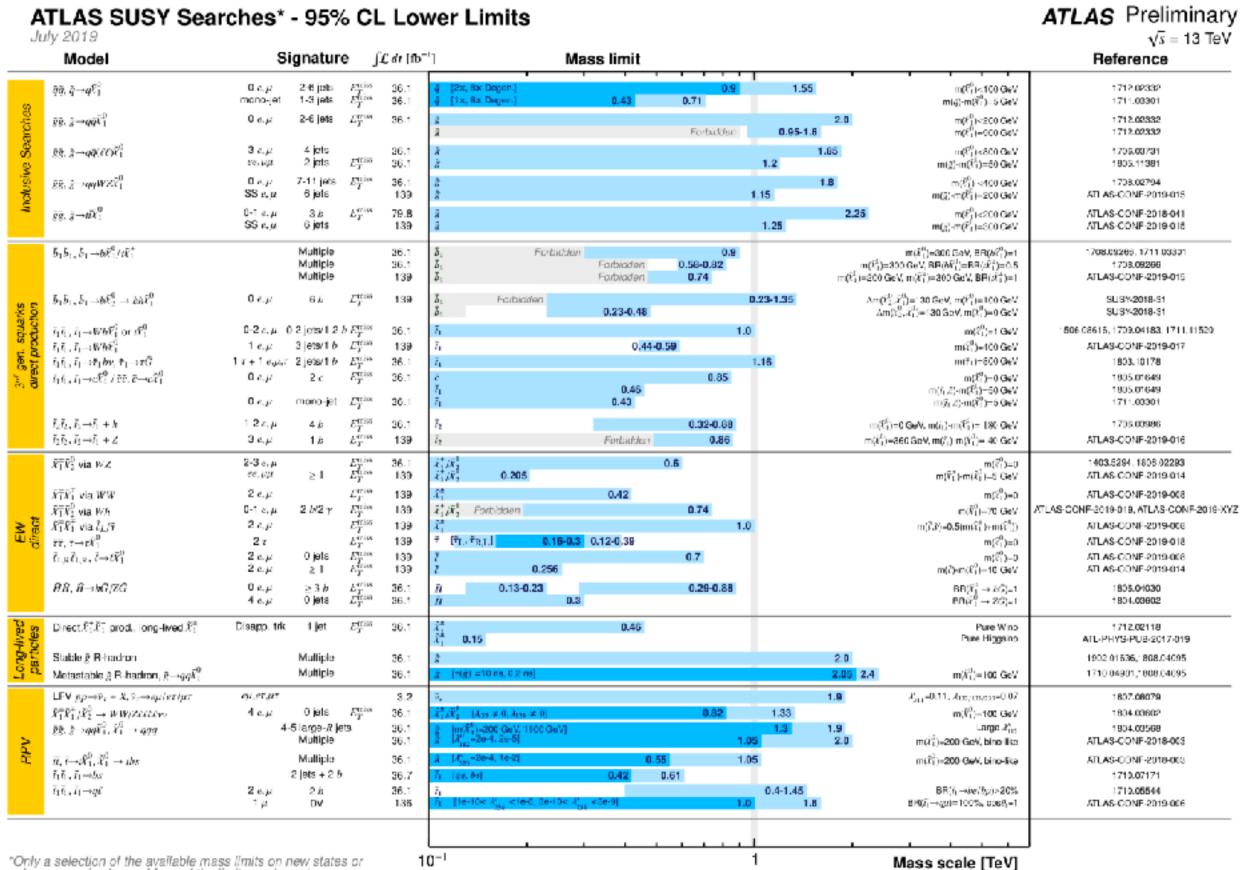
	Model	S	ignatur	e j	L de [fb	'l Ma	ss limit				$\sqrt{s} = 13 \text{ Te}^{3}$ Reference
60	$\hat{q}\hat{q}, \hat{q} \rightarrow q \hat{E}_{1}^{\circ}$	0 e,μ mono-jet	2-6 jets 1-3 jets	E_T^{miss} E_T^{miss}	36.1 36.1	\$ [2x, 8x Degen.] \$ [1x, 8x Degen.]	0.43	0.9	1.55	ກ(ເຖິງເ⊲100 GeV m(g)-m(ເຖິງ)=5 GeV	1712.02332 1711.03301
Indusive Searches	$\bar{g}\bar{g}, \bar{g} \rightarrow q\bar{q}\bar{\chi}_1^0$	$0 e, \mu$	2-6 jets	$E_T^{\rm miss}$	36.1	2		Forbidden	0.95-1.8	m(ξ ⁰ ₁)<200 GeV m(ξ ⁰ ₁)=000 GeV	1712.02332 1712.02332
	§§, 2→9\$(10)2°1	$3 e, \mu$ $ee, \mu\mu$	4 jets 2 jets	E_T^{triso}	36.1 36.1	À à			1.85	$m(\tilde{\mathcal{E}}_1^0)\! \times\! 800~{ m GeV}$ $m(\tilde{\chi})\! +\! m(\tilde{\mathcal{V}}_1^0)\! =\! 80~{ m GeV}$	1795.03791 1895.11381
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{g}_{1}^{0}$	0 v. p SS v.,u	7-11 jets 6 jets	$E_T^{\rm miss}$	36.1 139	2 2			1.15	$m(\tilde{\ell}_1^0) < 100 {\rm GeV}$ $m(\tilde{g}) < m(\tilde{\ell}_1^0) - 200 {\rm GeV}$	1708.02794 ATI AS-CONE-2019-015
	$\bar{g}g, \hat{g} \rightarrow d\bar{X}_{1}^{0}$	0-1 σ.μ SS σ.μ	3 <i>b</i> 6 jets	E_T^{miss}	79.B 139	ă ă			1.25	$m(\tilde{e}_j^0)$ <200 GeV $m(g)$ - $m(\tilde{e}_j^0)$ =500 GeV	ATLAS-CONF-2018-041 ATLAS-CONF-2019-018
	$b_1b_1, b_1 \rightarrow b\tilde{\mathcal{E}}_1^0/i\tilde{\mathcal{X}}_1^+$		Multiple Multiple Multiple		36.1 36.1 139	δ, Forbidden δ, δ,	Farbiaden Farbiaden	0.9 0.58-0.82 0.74		$m(\tilde{k}_{\perp}^{0})$ =300 GeV, BR($k\tilde{k}_{\perp}^{0}$)=1 (\tilde{k}_{\perp}^{0}) =300 GeV, BR($k\tilde{k}_{\perp}^{0}$)=BR($k\tilde{k}_{\perp}^{0}$)=0.5 200 GeV, $m(\tilde{k}_{\perp}^{0})$ =300 GeV, BR($k\tilde{k}_{\perp}^{0}$)=1	1708.09266, 1711.03301 1708.09266 ATLAS-CONF-2019-015
nks Non	$\delta_1\delta_1,\delta_1\!\rightarrow\!b\xi_2^0\to bh\xi_1^0$	$0~\epsilon.\mu$	64	$E_T^{\rm miss}$	139	Σ ₁ Farbidden Σ ₁	0.23-0.48		0.23-1.35	$\Delta m(\hat{r}_{2}^{0}, \hat{r}_{1}^{0}) = 30 \text{ GeV}, m(\hat{r}_{1}^{0}) = 100 \text{ GeV}.$ $\Delta m(\hat{r}_{2}^{0}, \hat{r}_{1}^{0}) = 130 \text{ GeV}, m(\hat{r}_{1}^{0}) = 0 \text{ GeV}.$	SUSY-2018-S1 SUSY-2018-S1
o" gen, squarks direct production	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{t}_1^0 \text{ or } s\tilde{t}_1^0$	$02\varepsilon,\mu$	0-2 jets/1-2		36.1	ħ		1.0		$m(\hat{x}_1^0)=1$ GeV	1506 08615, 1709.04183, 1711.11520
38	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wh\tilde{t}_1^0$	1 e, µ	3 jets/1 b		139	T ₁	0.44-0.			m(x ₁ ⁰)=400 GeV	ATLAS-CONF-2019-017
38	$\hat{t}_1\hat{t}_1, \hat{t}_1 \rightarrow \hat{\tau}_1 h \nu, \hat{\tau}_1 \rightarrow \tau \hat{G}$ $\hat{t}_1\hat{t}_1, \hat{t}_1 \rightarrow c\hat{\xi}_1^0 / \hat{c}\hat{c}, \hat{c} \rightarrow c\hat{\xi}_1^0$	1 T + 1 e.p.; 0 e.p.	2 jets/16/ 2 c	E_T^{miss} E_T^{miss}	36.1 36.1	iq a		0.85	1.15	m(₹₁)=800 GeV m(₹²)=0 GeV	1903.10178 1905.01649
ର ପ	116, 11→0€17 ee, e→0€1	0 e, µ	mono-jet	E_T^{miss}	36.1	h h	0.46 0.43	0.03		m(f, 2)-m(指) = 50 GeV m(f, 2)-m(指) = 5 GeV	1905.01649 1711.00301
	$\tilde{t}_1\tilde{t}_2, \tilde{t}_1 \rightarrow \tilde{t}_1 + h$	1.2 e, μ	4 b	E_T^{miss}	36.1	7,2		0.32-0.88		$m(\tilde{\mathcal{E}}_{1}^{i})$ =0 GeV, $m(\tilde{t}_{1})$ - $m(\tilde{\mathcal{E}}_{1}^{i})$ = 180 GeV	1705.00986
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	$3e,\mu$	1.6	E_T^{miss}	139	Ī ₂	Forbidden	0.86		$n(\hat{x}_{i}^{0})$ =860 GeV, $m(\hat{x}_{i}^{0})$ $m(\hat{x}_{i}^{0})$ = 40 GeV	ATLAS-CONF-2019-016
	$ar{K}_1^-ar{K}_2^0$ via WZ	2-3 ε, μ εε, μμ	≥1	E_T^{miss} E_T^{miss}	36.1 139	$\frac{\hat{x}_{1}^{+}\hat{x}_{2}^{+}}{\hat{x}_{1}^{+}\hat{x}_{2}^{+}}$ 0.205		0.6		$m(\tilde{x}_1^0)=0$ $m(\tilde{x}_1^0)=m(\tilde{x}_1^0)=5$ GeV	1403.5294, 1808.02293 ATLAS-CONF-2019-014
	$\hat{X}_{1}^{\pm}\hat{X}_{1}^{\mp}$ via WW	$2e.\mu$		E_T^{miss}	139	$\hat{\lambda}_1^{\pm}$	0.42			$m(\tilde{z}_{\perp}^{0})=0$	ATLAS-CONF-2019-008
**	$\bar{X}_1^- \bar{X}_2^0$ via Wh	0-1 ε, μ	$2 h/2 \gamma$	E_T^{miss}	139	$\tilde{\mathcal{X}}_1^{\dagger}/\tilde{\mathcal{X}}_2^{\dagger}$ Forbidden		0.74		$m(\tilde{X}_1^0)=70$ GeV	ATLAS-CONF-2019-019, ATLAS-CONF-2019
	$\hat{X}_{1}^{\pm}\hat{X}_{1}^{\mp}$ via \hat{l}_{L}/\bar{r}	$2e,\mu$		E_T^{miss}	139	x 1 ftr. fp. 1 6.18.03	040.000	1.0		$m(\vec{t}, \vec{t}) = 0.5 \text{tm}(\vec{x}_1^{(t)}) + m(\vec{x}_1^{(t)})$	ATLAS-CONF-2019-008
- 8	$\tilde{r}\tilde{r}, \tilde{\tau} \rightarrow r\tilde{X}_{1}^{0}$ $\tilde{\ell}_{1,0}\tilde{\ell}_{1,\infty}, \tilde{\ell} \rightarrow \ell\tilde{X}_{1}^{0}$	2τ 2 c. μ	0 jets	E_T^{miss} Entire	139 139	₹ [₹1. ₹8,1.] 0.16-0.3	0.12-0.38	0.7		$m(\vec{x}_1^0) = 0$ $m(\vec{x}_1^0) = 0$	ATLAS-CONF-2019-018 ATLAS-CONF-2019-008
	eli,8el,8. c→cel	2 c. µ	≥ 1	E_T^{tries} E_T^{tries}	139	7 0.256		0.1		$m(\tilde{c})$ - $m(\tilde{k}_1^0)$ =10 GeV	ATI AS-CONE-2019-014
	$\widehat{H}\widehat{H}$, $\widehat{H} \rightarrow hG/\mathbb{Z}G$	$0 e, \mu$	≥ 3 <i>b</i> 0 jets	E_T^{miss} E_T^{miss}	36.1 36.1	ju 0.13-0.23		0.29-0.88		$BB(\bar{x}_1^2 \to hG) = 1$	1805.04030
		4 α.μ		LT	26.1	Ĥ 0.3				$BD(\widetilde{x}_1^0 \to Z\widetilde{G})=1$	1804.03602
les les	Direct $\mathcal{E}_1^+\mathcal{E}_1^-$ prod., long-lived \mathcal{E}_1^a	Disapp. trk	1 jet	E_T^{miss}	36.1	$\frac{\hat{x}_{1}^{*}}{\hat{x}_{1}^{*}} = 0.15$	0.46			Pure Wino Pure Higgaino	1712.02118 ATL-PHYS-PUB-2017-019
particles	Stable & R-hadron		Multiple		36.1	2			2.0		1902.01536,1808.04095
3 15	Metastable \hat{g} B-hadron, $\hat{g} \rightarrow \hat{g} q \hat{k}_{\perp}^{0}$		Multiple		36.1	ĝ [π(ĝ) =10 ns, 0.2 ns]			2.05 2.	4 $m(\bar{X}_1^0)=100 \text{ GeV}$	1710 04901, 808,04095
	LFV $pp \rightarrow \bar{v}_T + X, \bar{v}_T \rightarrow s\mu/e\tau/\mu\tau$	еуцетит			3.2	ž,			1.9	A' ₃₁₁ =0.11, A _{122/123/233} =0.07	1807.08079
	$\bar{X}_1^-\bar{X}_1^+/\bar{X}_2^0 \rightarrow WW/ZEEEer$	4 c, µ	0 jets	E_T^{miss}	36.1	$\tilde{X}_{1}^{\pm}/\tilde{X}_{2}^{\pm} = [\lambda_{03} \neq 0, \lambda_{124} \neq 0]$		0.82	1.33	$m(\bar{V}_1^0)$ =100 GeV	1804.03602
>	$\S\S,\S o qq ilde{k}_1^0, ilde{\chi}_1^0 o qqq$	4	i-5 large- <i>R</i> ji Multiple	ets.	36.1 36.1	2 [m(\$\hat{X}_1^2) = 200 GeV, 1100 GeV] 2 \hat{X}_{10}^2 = 2e-4, 2e-5[1.0	1.3 1.9 5 2.0	Large $\mathcal{X}_{\mathrm{tot}}^{\lambda}$ m($\hat{\mathcal{X}}_{1}^{\lambda}$)=200 GeV, bino like	1904.09568 ATLAS-CONF-2018-003
RPV	$\widetilde{u}_i \in \mathcal{K}_1^0, \widetilde{\mathcal{X}}_1^0 \to ubs$		Multiple		36.1	g [JC, -2e-4, 1e-2]	0.55			m(t)=200 GeV, bino-like	ATLAS-CONF-2018-003
	$\tilde{t}_1 \tilde{t}_1 \cdot \tilde{t}_1 \rightarrow bs$		2 jets + 2 8	,	36.7	Ti Qu. Bal		1.61		minimal harmon plants and account	1710.07171
	$\tilde{t}_1\tilde{t}_1$, $\tilde{t}_1{ ightarrow} q\mathcal{E}$	$2e, \mu$	2.8		36.1	ž _l			0.4-1.45	BR($\hat{f}_t \rightarrow be/b\mu$)>20%	1710.05544
		1 μ	DV		135	T ₁ [1e-10< λ' _{1,1} < le-0, 3e-10< λ' _{1,1}	K06-9	1.0	1.8	$BH(\tilde{x}_1 \rightarrow qu) = 100\%_{s_0} \cos\theta_1 = 1$	ATLAS-CONF-2019-006

"Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made. 10^{-1}

‡ Fermilab

Mass scale [TeV]





Many, many different ways to search for SUSY at colliders. ATLAS and CMS are the two, powerful, all-purpose detectors at the Large Hadron Collider. To date, no evidence for SUSY has been found.



sparticle and the LSP relative to ΔM , respectively, unless indicated otherwise



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Thank you!