



Space Charge Simulations for LHC Injector Upgrade: Montague Resonance at PS

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In collaboration with



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- LHC injector upgrade group (C. Carli, R. Garoby, S. Gilardoni, E. Metral, F. Schmidt, et al.), CERN





- Introduction
- Computational model
- Simulation of Montague resonance experiment at PS
- Summary and future plan

Introduction



- Proton-Synchrotron (PS) is amongst the LHC injectors the oldest, and will continue to serve the LHC at least for the next 25 years.
- Space-charge effects is a dominant factor limiting the bunch intensity.
- Montague Resonance:

$$2 Qx - 2 Qy = 0$$

- can cause particle due to unequal aperture size in horizontal and vertical dimensions.
- benchmark space-charge codes



Refs: B. W. Montague, CERN-Report No. 68-38, CERN, 1968.

E. Metral et al., Proc. of EPAC 2004, p. 1894.

I. Hofmann et al., Proc. of EPAC 2004, p. 1960.

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IMPACT Code Suite

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- IMPACT-Z: parallel PIC code (z-code)
- IMPACT-T: parallel PIC code (t-code)
- Envelope code, pre- and post-processors,...
- Optimized for parallel processing
- Applied to many projects: SNS, JPARC, RIA, FRIB, PS2, future light sources, advanced streak cameras,...
- Has been used to study photoinjectors for BNL e-cooling project, Cornell ERL, FNAL/A0, LBNL/APEX, ANL, JLAB, SLAC/LCLS



IMPACT-Z

- Parallel PIC code using coordinate "z" as the independent variable
- Key Features
 - Detailed RF accelerating and focusing model
 - -Multiple 3D Poisson solvers
 - Variety of boundary conditions
 - 3D Integrated Green Function
 - -Multi-charge state
 - -Machine error studies and steering
 - —Wakes
 - —CSR (1D)
 - —Run on both serial and multiple processor computers







Particle-in-cell simulation with split-operator method

- Particle-in-cell approach:
 - Charge deposition on a grid
 - Field solution via spectral-finite difference method with transverse rectangular conducting pipe and longitudinal open
 - Field interpolation from grid to particles
- Split-operator method with $H = H_{external} + H_{space charge}$
- Thin lens kicks for nonlinear elements
- Lumped space-charge at a number locations

Poisson Solver Used in Space-Charge Calculation

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial z^2} = -\frac{\rho}{\epsilon_0}$$

with boundary conditions

$$\phi(x = 0, y, z) = 0,$$

$$\phi(x = a, y, z) = 0,$$

$$\phi(x, y = 0, z) = 0,$$

$$\phi(x, y = b, z) = 0,$$

$$\phi(x, y, z = \pm \infty) = 0,$$

$$\rho(x, y, z) = \sum_{l=1}^{N_l} \sum_{m=1}^{N_m} \rho^{lm}(z) \sin(\alpha_l x) \sin(\beta_m y),$$

$$\phi(x, y, z) = \sum_{l=1}^{N_l} \sum_{m=1}^{N_m} \phi^{lm}(z) \sin(\alpha_l x) \sin(\beta_m y),$$

where

$$\rho^{lm}(z) = \frac{4}{ab} \int_{0}^{a} \int_{0}^{b} \rho(x, y, z) \sin(\alpha_l x) \sin(\beta_m y),$$

$$\phi^{lm}(z) = \frac{4}{ab} \int_{0}^{a} \int_{0}^{b} \phi(x, y, z) \sin(\alpha_l x) \sin(\beta_m y),$$

$$\begin{aligned} \frac{\partial^2 \phi^{lm}(z)}{\partial z^2} &- \gamma_{lm}^2 \phi^{lm}(z) = -\frac{\rho^{lm}(z)}{\epsilon_0}, \\ \frac{\phi_{n+1}^{lm} - 2\phi_n^{lm} + \phi_{n-1}^{lm}}{h_z^2} - \gamma_{lm}^2 \phi_n^{lm} = -\frac{\rho_n^{lm}}{\epsilon_0}, \\ \phi_{-1}^{lm} &= \exp(-\gamma_{lm} h_z) \phi_0^{lm}, \quad n = 0, \\ \phi_{N+1}^{lm} &= \exp(-\gamma_{lm} h_z) \phi_N^{lm}, \quad n = N. \end{aligned}$$

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Physical parameters:

Vrf = ramping with f = 39.5 MHz Ek = 1.4 GeV Emit_x = 7.5 mm-mrad Emit_y = 2.5 mm-mrad Rms bunch length = 45 ns Rms dp/p = 1.7×10^{-3}

Horizontal tune: 6.15 - 6.245Vertical tune: 6.21Synchrotron period: 1.5 ms Half Aperture = 7cm x 3.5cm I = $1.0x10^{12}$

Refs: B. W. Montague, CERN-Report No. 68-38, CERN, 1968. E. Metral et al., Proc. of EPAC 2004, p. 1894. I. Hofmann et al., Proc. of EPAC 2004, p. 1960.



• Normalize 1-turn map: M=ANA⁻¹

A is the normalizing map

N is the normal form which causes only rotations in phase space

- Consider a function $g((x^2+p_x^2),(y^2+p_y^2),(t^2+p_t^2))$
- Then $f(\zeta)=g(A(x^2+p_x^2),(y^2+p_y^2),(t^2+p_t^2))$ is a matched beam.

Proof: The distribution after one turn is given him $f(M^{-1}\zeta)=g(AN A^{-1}A (x^{2}+p_{x}^{2}),(y^{2}+p_{y}^{2}),(t^{2}+p_{t}^{2})) = 0.000$ $g(AN (x^{2}+p_{x}^{2}),(y^{2}+p_{y}^{2}),(t^{2}+p_{t}^{2})) = 0.000$ $g(A (x^{2}+p_{x}^{2}),(y^{2}+p_{y}^{2}),(t^{2}+p_{t}^{2})) = 0.000$



rms emittances



rrrr

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distance (m)

Numerical Parameters: Test of Convergence (2)

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Synchrotron motion enhances the emittance exchange !

RMS Emittance Evolution with Different Horizontal Tunes

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0,003 'fort.1' u 6:1 0,002 0,001 9e-06 6.245 - X 6,245 -6.23 - X 6.23 --0,001 8e-06 6.22 -22 -0,002 6,21 6.21 7e-06 6.20 - X -0.0030,015 6,20 - Y -0,01 0,005 0,01 -0,02 -0,015 -0,005 -delE/mc^2 6.19 - X 0,0144567, -0,00303780 6.19 normalized rms emittance 6.18 - X 6e-06 6.18 - Y 6.17 - X 6.17 -6,16 -6,16 -5e-06 6,15 -6.15 -4e-06 3e-06 2e-06 2000 4000 6000 10000 12000 0 8000 turn

2865.65, 7.84388e-06

:/const

emittance exchange with constant focusing channel: square pipe vs. rectangular pipe



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RMS emittance evolution of an initial upright Gaussian beam (with different horizontal tunes)

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100 ms dynamic Crossing



Initial and the Final Phase Space Distribution of the Dynamic Resonance Crossing



X-delE phase space correlation at the beginning and the end of 44032 turns in the

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Dynamics Montague Resonance Crossing at PS (2)





- 3D self-consistent space-charge simulation reproduce the experiment data reasonably well
- Dynamic Montague resonance crossing shows no symmetry around the resonance stopband
- Longitudinal synchrotron motion helps the emittance exchange inside the resonance stopband
- Study the space-charge effects at PS with higher injection energy
- Write a paper