



QMC STA $e+^4\text{He}$ and $e+^{12}\text{C}$

NuSTEC Workshop on Electron Scattering
29 March 2022

Saori Pastore
Washington University in St Louis

<https://physics.wustl.edu/quantum-monte-carlo-group>

Quantum Monte Carlo Group @ WashU

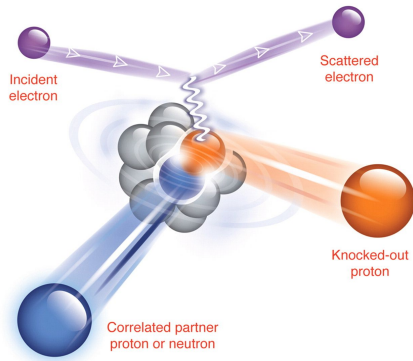
Lorenzo Andreoli (PD) Jason Bub (GS) Garrett King (GS) Maria Piarulli and Saori Pastore

Computational Resources awarded by the DOE ALCC and INCITE programs

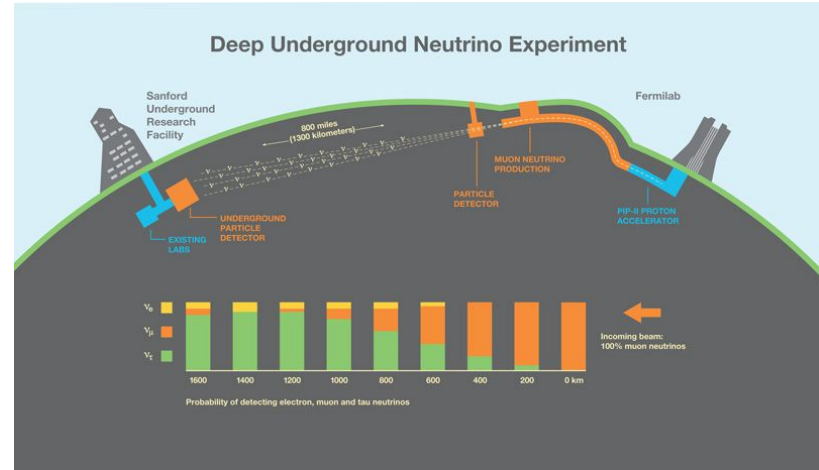
Short-Time-Approximation

Short-Time-Approximation Goals:

- Describe electroweak scattering from $A > 12$ without losing two-body physics
- Account for exclusive processes
- Incorporate relativistic effects



Subedi et al. Science320(2008)1475



[Stanford Lab article](#)

[e4u collaboration](#)

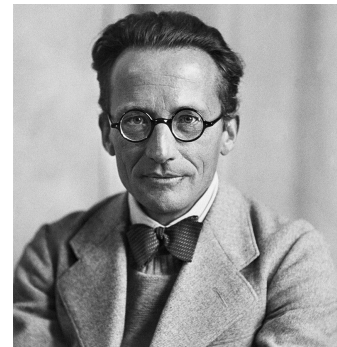


Microscopic (or *ab initio*) Description of Nuclei

Comprehensive theory that describes quantitatively and predictably nuclear structure and reactions

Requirements:

- Accurate understanding of the interactions/correlations between nucleons in **pairs, triplets, ... (two- and three-nucleon forces)**
- Accurate understanding of the electroweak interactions of external probes (electrons, neutrinos, photons) with nucleons, correlated nucleon-pairs, ... (**one- and two-body electroweak currents**)
- **Computational methods** to solve the many-body nuclear problem of strongly interacting particles

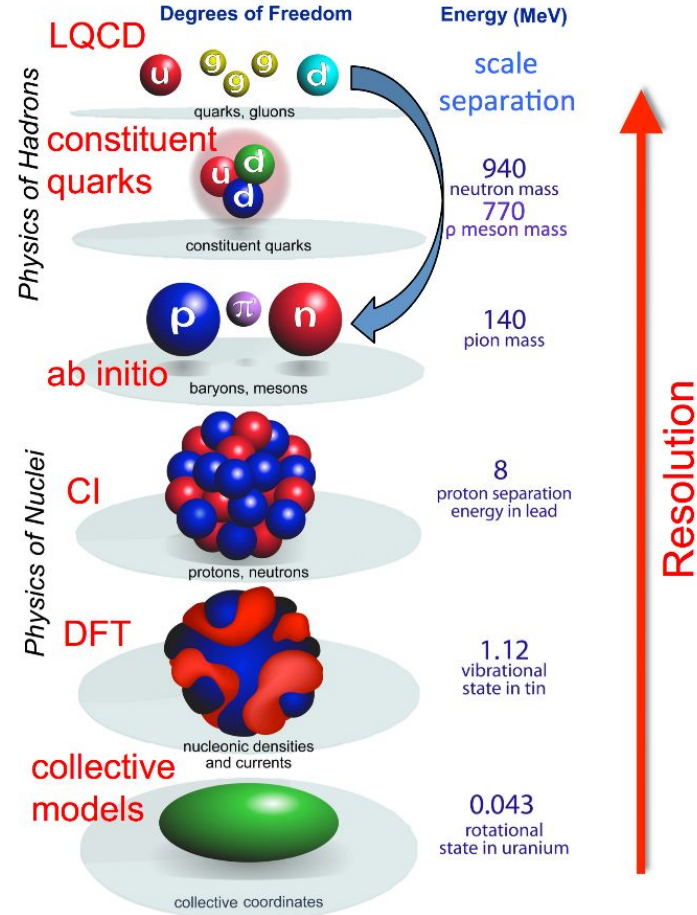


Erwin Schrödinger

$$H\Psi = E\Psi$$

From Quarks to Nuclei

- Nuclei are complex systems made of interacting **protons** and **neutrons**, which in turns are composite objects made of interacting constituent quarks.
- All fundamental forces are at play in nuclei.
- **EFTs** low-energy approximations of QCD whose d.o.f. are bound states of QCD (e.g., protons, neutrons, pions, ...)
- **EFTs** are used to construct many-nucleon interactions and currents



Many-body Nuclear Problem

Nuclear Many-body Hamiltonian

$$H = T + V = \sum_{i=1}^A t_i + \sum_{i<j} v_{ij} + \sum_{i<j<k} V_{ijk} + \dots$$

$$\Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A, s_1, s_2, \dots, s_A, t_1, t_2, \dots, t_A)$$

Ψ are **spin-isospin** vectors in **3A** dimensions with $2^A \times \frac{A!}{Z!(A-Z)!}$ components

Develop Computational Methods to solve (numerically) exactly or within approximations that are under control the many-body nuclear problem



<http://exascale.org/np/>

${}^4\text{He}$: 96
 ${}^6\text{Li}$: 1280
 ${}^8\text{Li}$: 14336
 ${}^{12}\text{C}$: 540572

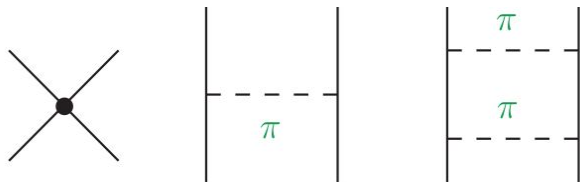
Wide availability of computing is key to critical progress in nuclear physics across all areas.

Many-body Nuclear Interactions

Many-body Nuclear Hamiltonian

$$H = T + V = \sum_{i=1}^A t_i + \sum_{i<j} v_{ij} + \sum_{i<j<k} V_{ijk} + \dots$$

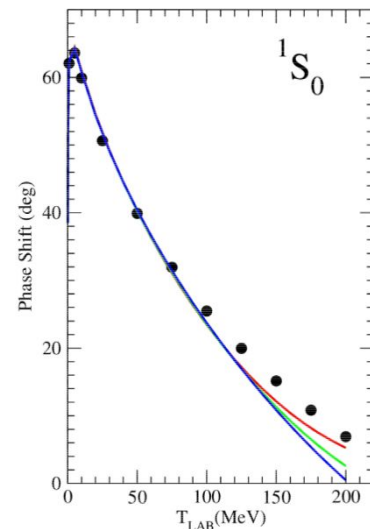
v_{ij} and V_{ijk} are two- and three-nucleon operators based on experimental data fitting; fitted parameters subsume underlying QCD dynamics



Contact term: short-range

Two-pion range: intermediate-range $r \propto (2m_\pi)^{-1}$

One-pion range: long-range $r \propto m_\pi^{-1}$



SP et al. PRC80(2009)034004

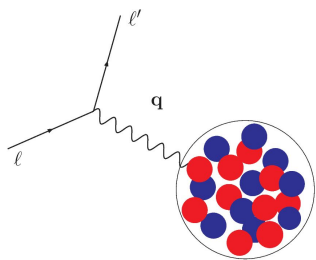
In Quantum Monte Carlo methods we use:

AV18+UIX; **AV18+IL7** Wiringa, Schiavilla, Pieper *et al.*

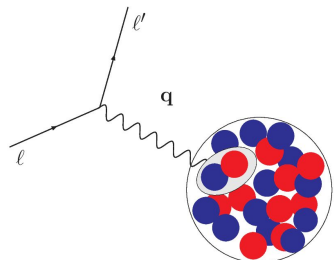
chiral π N **N2LO+N2LO** Gerzelis, Tews, Lynn *et al.*

chiral π N Δ **N3LO+N2LO** Piarulli *et al.* **Norfolk Models**

Many-body Nuclear Electroweak Currents



one-body



two-body

- One-body currents: non-relativistic reduction of covariant nucleons' currents
- Two-body currents are a manifestation of two-nucleon correlations
- Electromagnetic two-body currents are required to satisfy current conservation

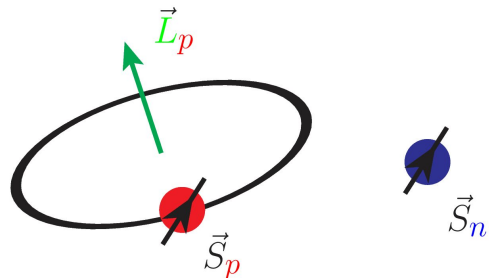
$$\mathbf{q} \cdot \mathbf{j} = [H, \rho] = [t_i + v_{ij} + V_{ijk}, \rho]$$

Nuclear Charge Operator

$$\rho = \sum_{i=1}^A \rho_i + \sum_{i<j} \rho_{ij} + \dots$$

Nuclear (Vector) Current Operator

$$\mathbf{j} = \sum_{i=1}^A \mathbf{j}_i + \sum_{i<j} \mathbf{j}_{ij} + \dots$$



Magnetic Moment: Single Particle Picture

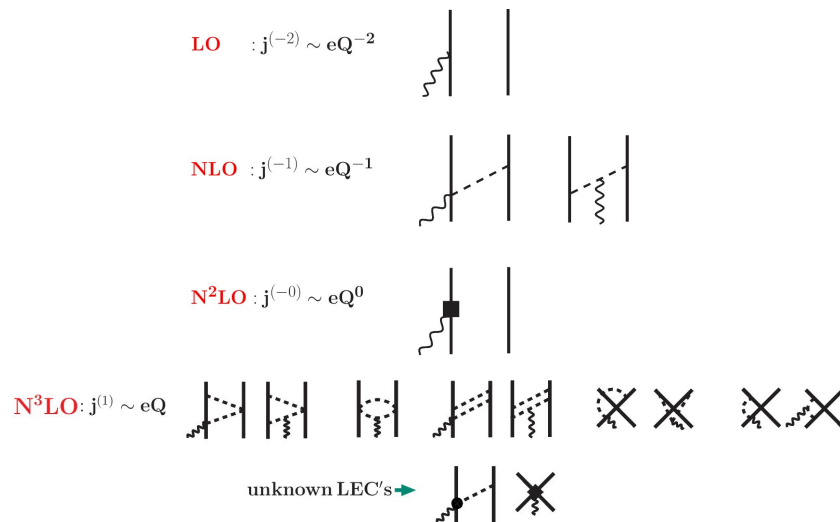
Many-body Currents: Available Models

- **Meson Exchange Currents (MEC)**

Constrain the MEC current operators by imposing that the current **conservation relation is satisfied with the given two-body potential**

- **Chiral Effective Field Theory Currents**

Are constructed consistently with the two-body chiral potential; Unknown parameters, or Low Energy Constants (LECs), need to be **determined by either fits to experimental data or by QCD calculations, as well as nucleonic form factors**



Electromagnetic Current Operator

SP *et al.* PRC78(2008)064002, PRC80(2009)034004,
 PRC84(2011)024001, PRC87(2013)014006
 Park *et al.* NPA596(1996)515, Phillips (2005)
 Kölling *et al.* PRC80(2009)045502 & PRC84(2011)054008

Lepton-Nucleus scattering: Inclusive Processes

Electromagnetic Nuclear Response Functions

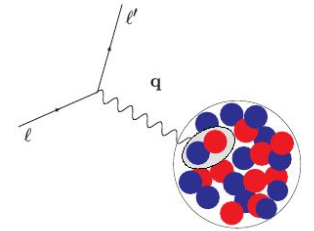
$$R_{\alpha}(q, \omega) = \sum_f \delta(\omega + E_0 - E_f) |\langle f | O_{\alpha}(\mathbf{q}) | 0 \rangle|^2$$

Longitudinal response induced by the charge operator $O_L = \rho$

Transverse response induced by the current operator $O_T = \mathbf{j}$

5 Responses in neutrino-nucleus scattering

$$\frac{d^2 \sigma}{d\omega d\Omega} = \sigma_M [v_L R_L(\mathbf{q}, \omega) + v_T R_T(\mathbf{q}, \omega)]$$



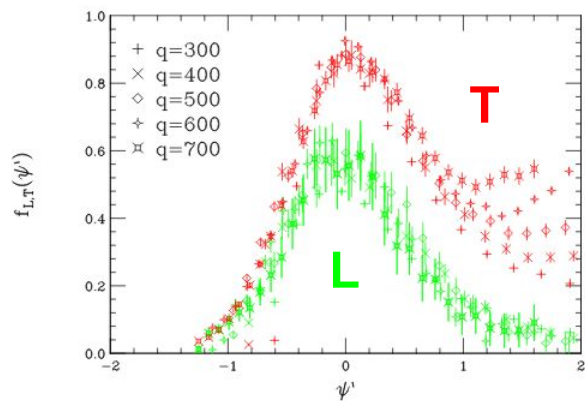
For a recent review on QMC, SF methods see

[Rocco Front. In Phys.8 \(2020\)116](#)

Lepton-Nucleus scattering: Data

Transverse Sum Rule

$$S_T(q) \propto \langle 0 | \mathbf{j}^\dagger \mathbf{j} | 0 \rangle \propto \langle 0 | \mathbf{j}_{1b}^\dagger \mathbf{j}_{1b} | 0 \rangle + \langle 0 | \mathbf{j}_{1b}^\dagger \mathbf{j}_{2b} | 0 \rangle + \dots$$



⁴He Electromagnetic Data
Carlson *et al.* PRC65(2002)024002

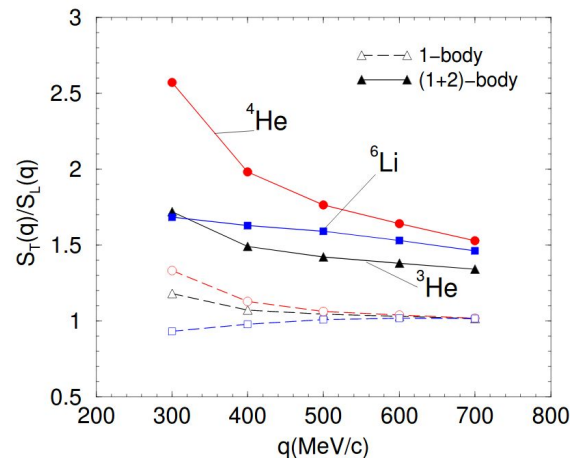
Observed transverse enhancement explained by the combined effect of two-body correlations and currents in the interference term

$$\langle \mathbf{j}_{1b}^\dagger \mathbf{j}_{1b} \rangle > 0$$

Leading one-body term

$$\langle \mathbf{j}_{1b}^\dagger \mathbf{j}_{2b} v_\pi \rangle \propto \langle v_\pi^2 \rangle > 0$$

Interference term

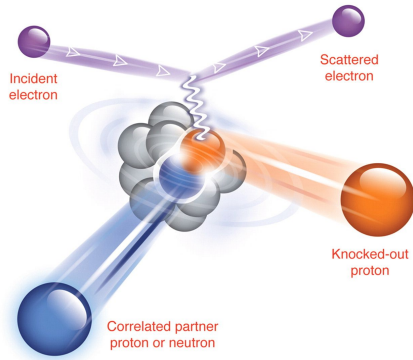


Transverse/Longitudinal Sum Rule
Carlson *et al.* PRC65(2002)024002

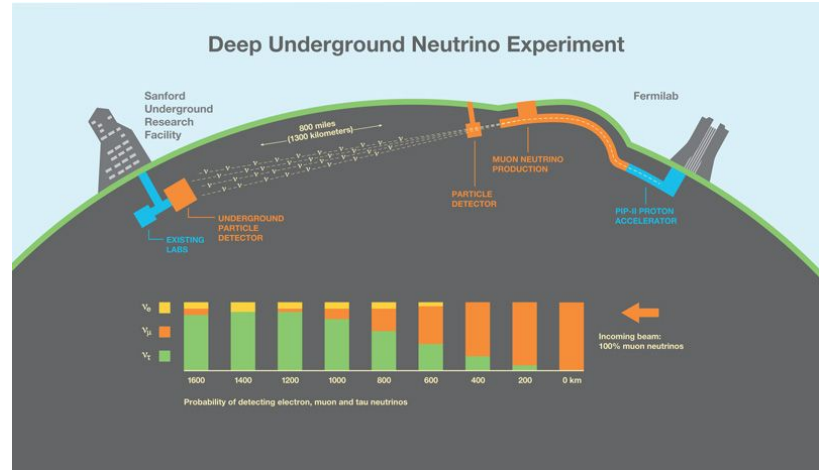
Beyond Inclusive: Short-Time-Approximation

Short-Time-Approximation Goals:

- Describe electroweak scattering from $A > 12$ without losing two-body physics
- Account for exclusive processes
- Incorporate relativistic effects



Subedi et al. Science320(2008)1475



[Stanford Lab article](#)

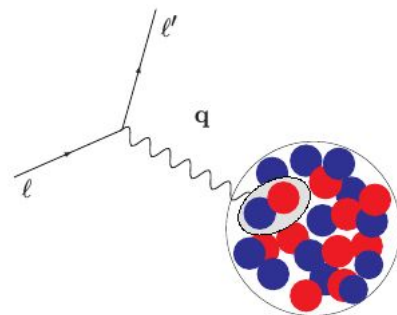
[e4u collaboration](#)



Short-Time-Approximation

Short-Time-Approximation:

- Based on Factorization
- Retain two-body physics
- Correctly accounts for interference

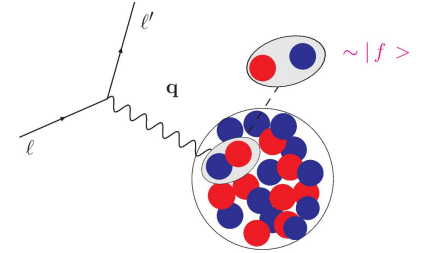


$$R(q, \omega) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{i(\omega + E_0)t} \langle 0 | O^\dagger e^{-iHt} O | 0 \rangle$$

$$O_i^\dagger e^{-iHt} O_i + O_i^\dagger e^{-iHt} O_j + O_i^\dagger e^{-iHt} O_{ij} + O_{ij}^\dagger e^{-iHt} O_{ij}$$

$$H \sim \sum_i t_i + \sum_{i < j} v_{ij}$$

Short-Time-Approximation



Short-Time-Approximation:

- Based on Factorization
- **Retains two-body physics**
- Response functions are given by the **scattering from pairs of fully interacting nucleons** that propagate into a correlated pair of nucleons
- Allows to retain both two-body correlations and currents at the vertex
- Provides **“more” exclusive information in terms of nucleon-pair kinematics via the Response Densities**

Response Functions \propto Cross Sections

$$R_{\alpha}(q, \omega) = \sum_f \delta(\omega + E_0 - E_f) |\langle f | O_{\alpha}(\mathbf{q}) | 0 \rangle|^2$$

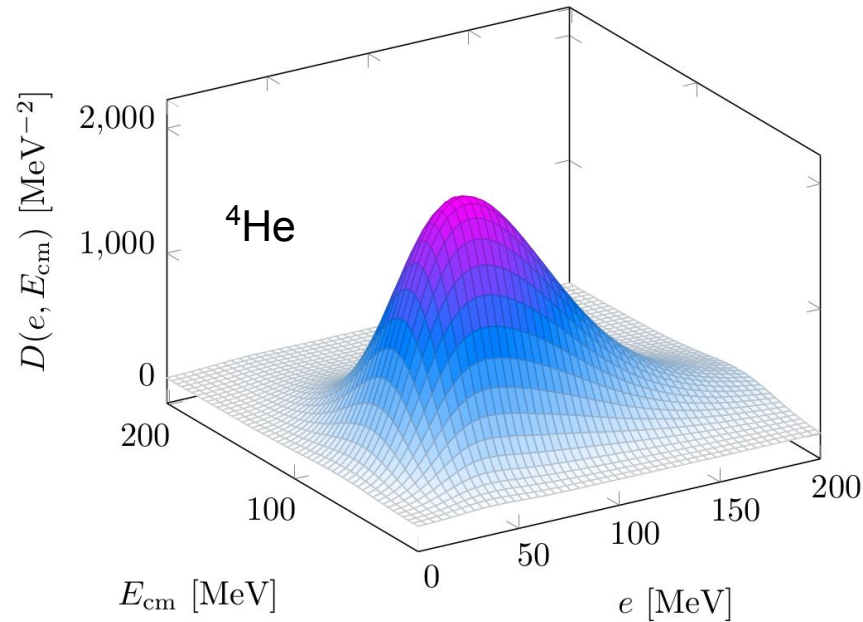
Response **Densities**

$$R(q, \omega) \sim \int \delta(\omega + E_0 - E_f) dP' dp' \mathcal{D}(p', P'; q)$$

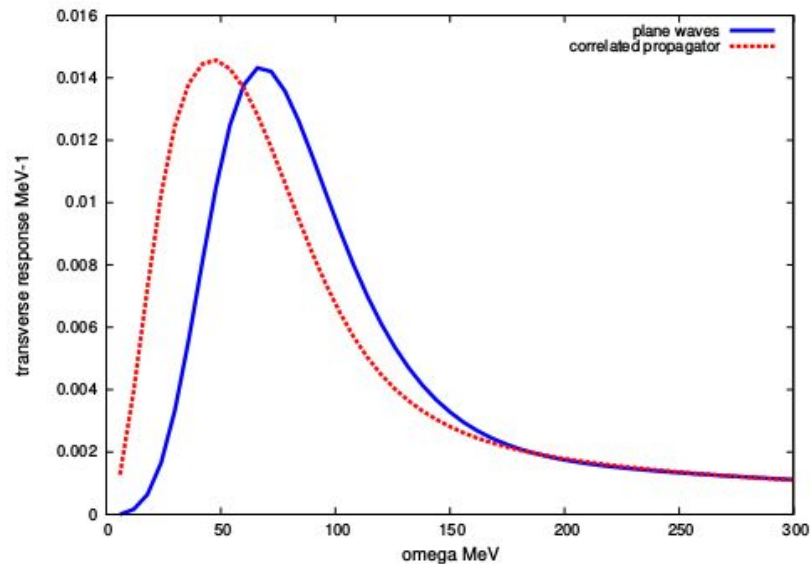
P' and p' are the CM and relative momenta of the struck nucleon pair

Transverse Response Density: e - ${}^4\text{He}$ scattering

Transverse Density $q = 500 \text{ MeV}/c$

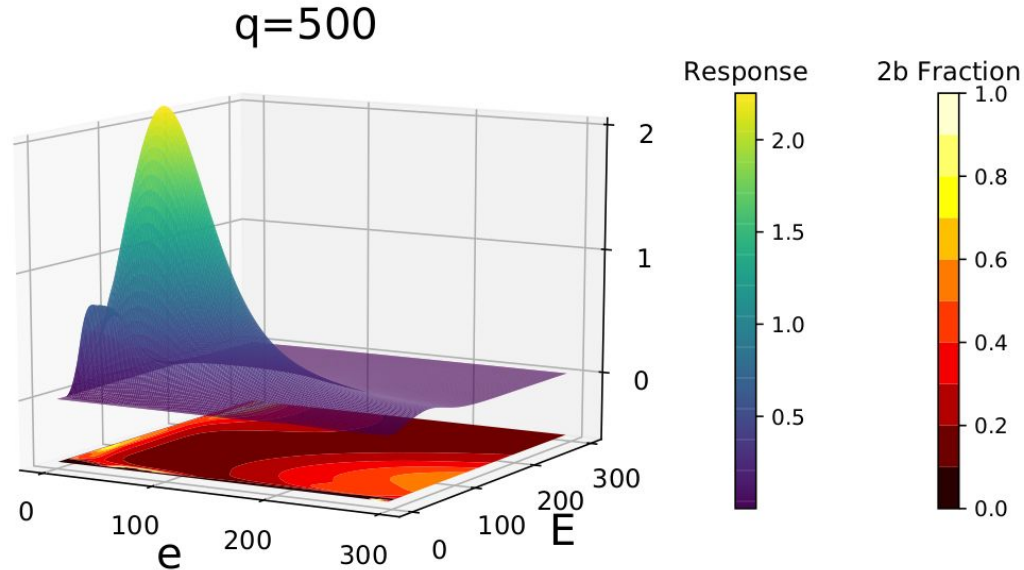


Correlated pairs vs uncorrelated pairs

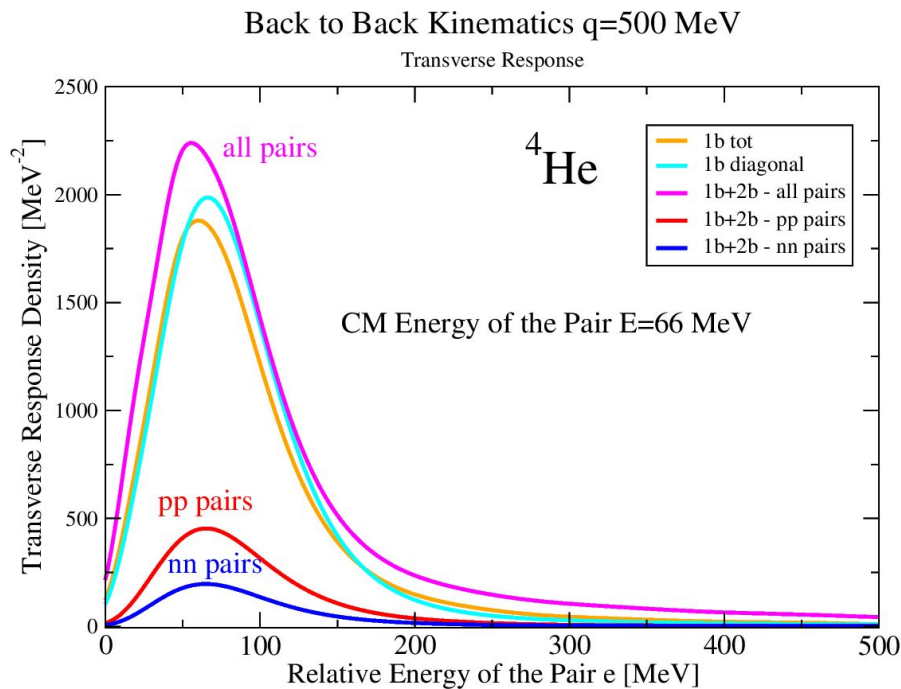


Scattering from **uncorrelated** vs **correlated** nucleon pairs

Transverse Response Density: two-body physics

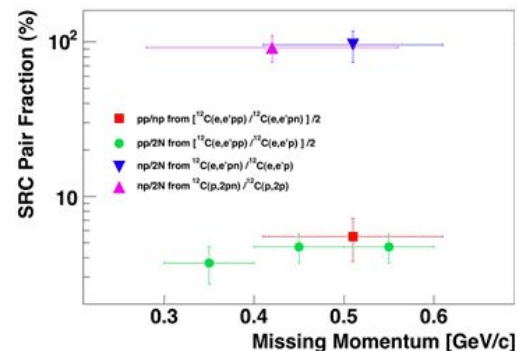


$e^{-4}\text{He}$ scattering in the back-to-back kinematic



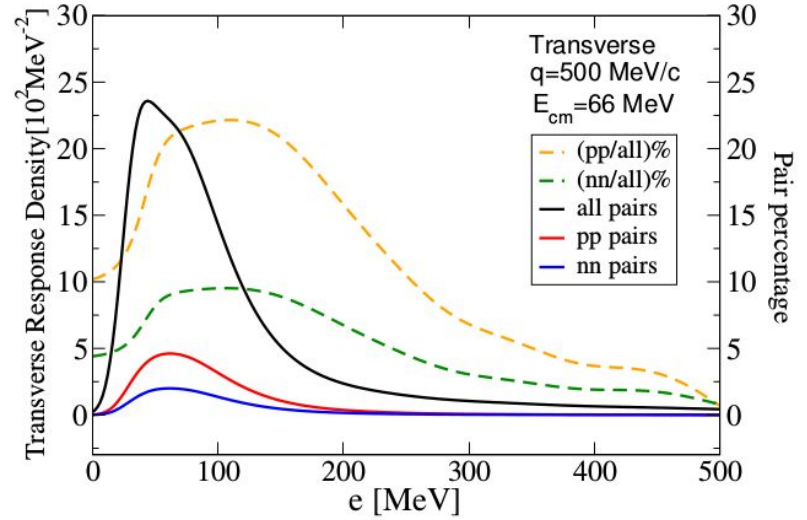
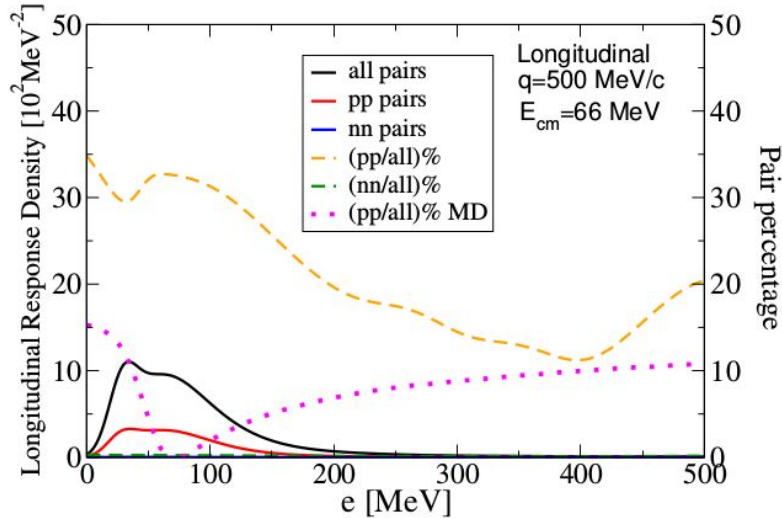
SP *et al.* PRC101(2020)044612

- pp pairs
- nn pairs
- all pairs 1body
- all pairs tot



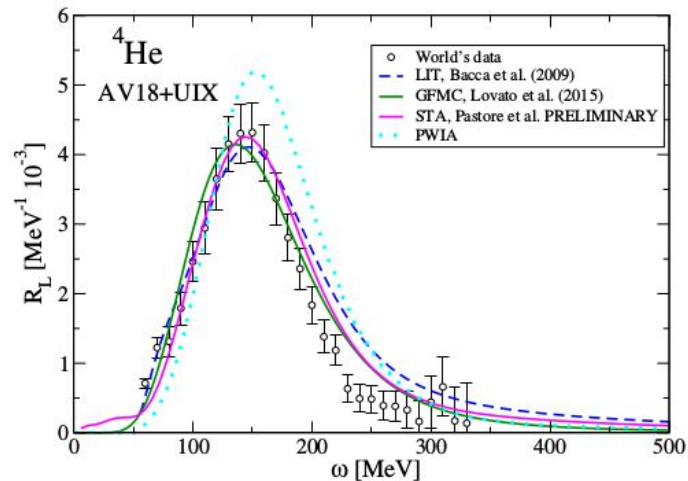
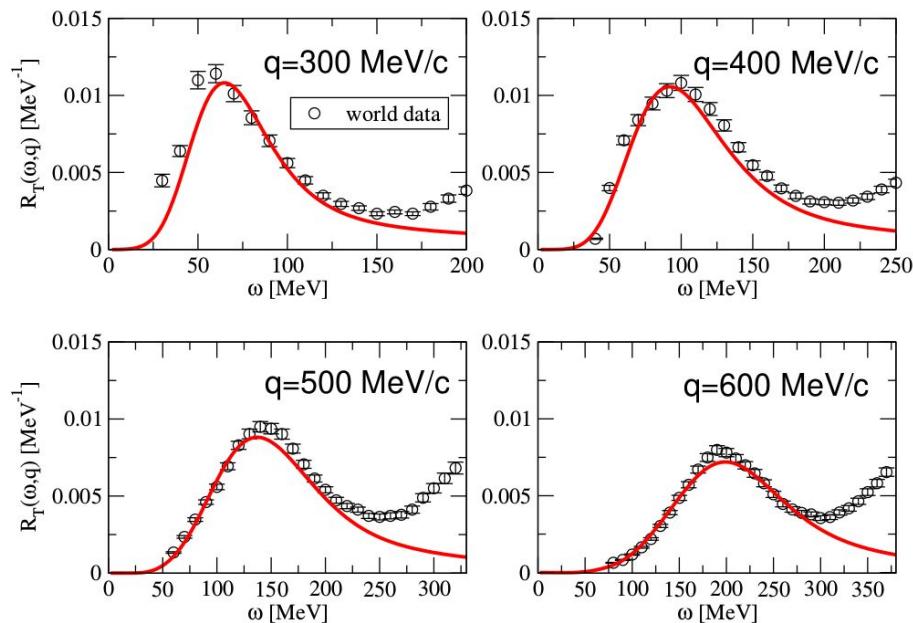
Subedi *et al.* Science320(2008)1475

Back to back scattering and particle identity



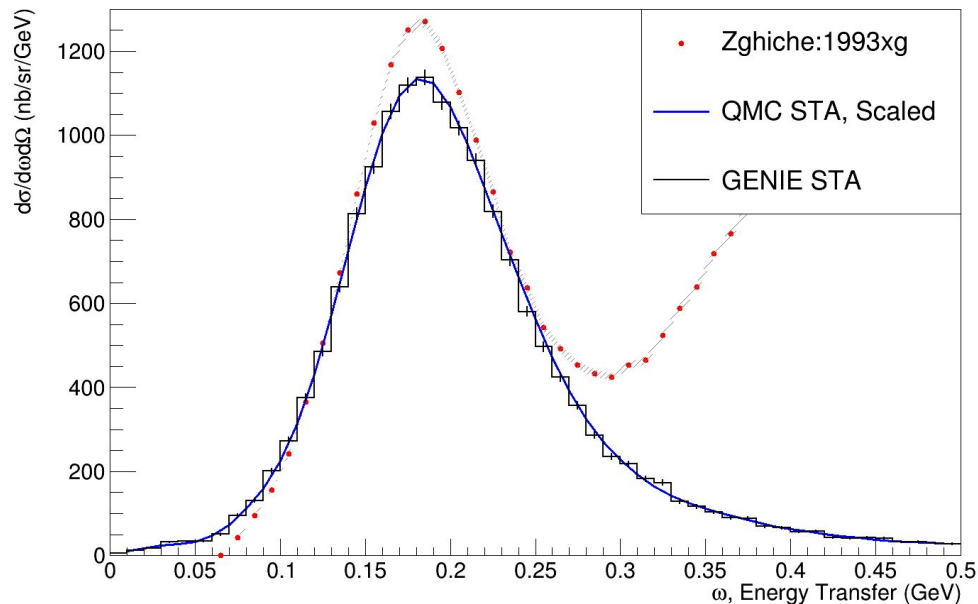
tot
 pp nn
 pp/all % pp/all % from momentum distributions
 nn/all %

Helium-4 comparison with the data



GENIE validation using e-scattering

Z = 2, A = 4, Beam Energy = 0.64 GeV, Angle = 60° ± 0.25°



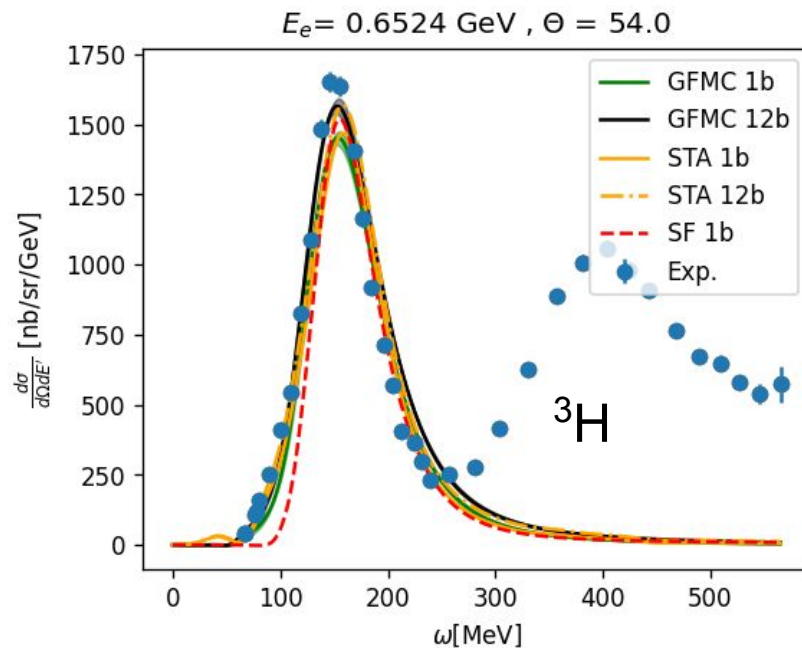
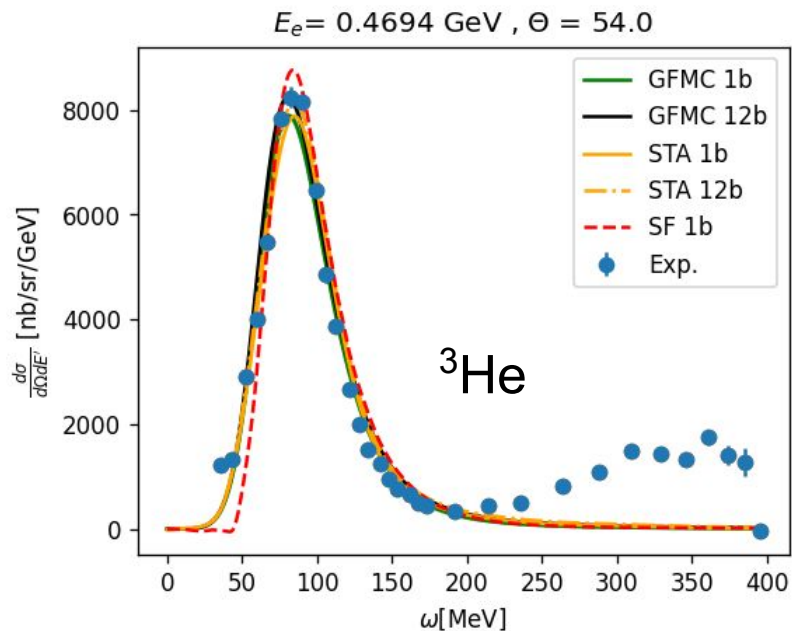
Barrow Gardiner Betancourt SP *et al.*
PRD 103 (2021) 5, 052001

Ongoing work

- Implementation of moment-morphin interpolation techniques
- Implementations of response **Densities** in GENIE
- ^{12}C response densities with [Lorenzo Andreoli](#)

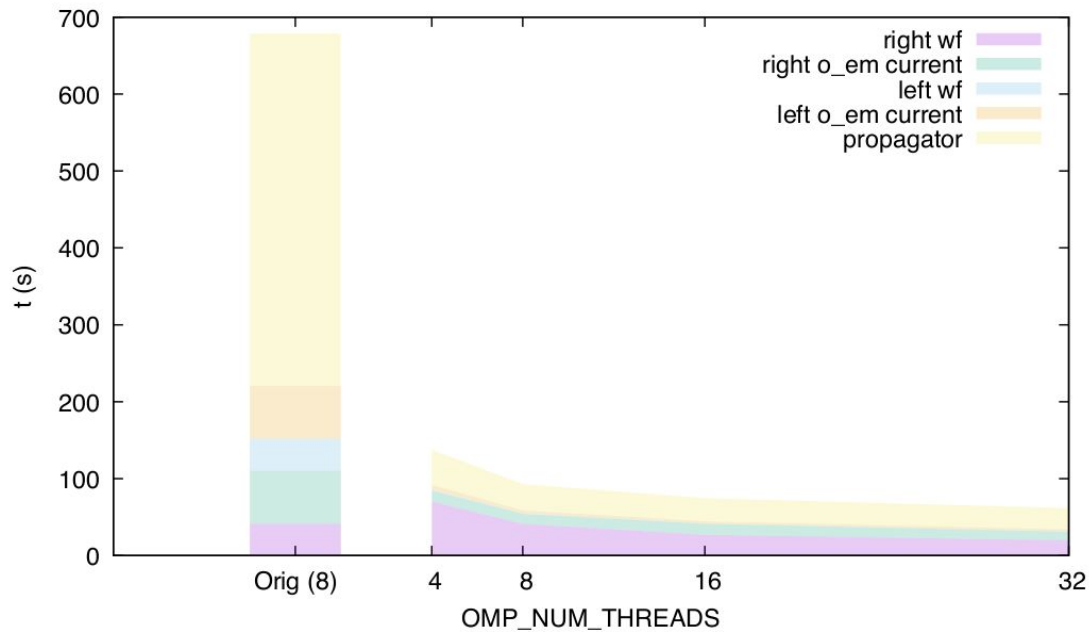
$$\frac{d^2 \sigma}{d\omega d\Omega} = \sigma_M [v_L R_L(\mathbf{q}, \omega) + v_T R_T(\mathbf{q}, \omega)]$$

GFMC SF STA: Benchmark & error estimate



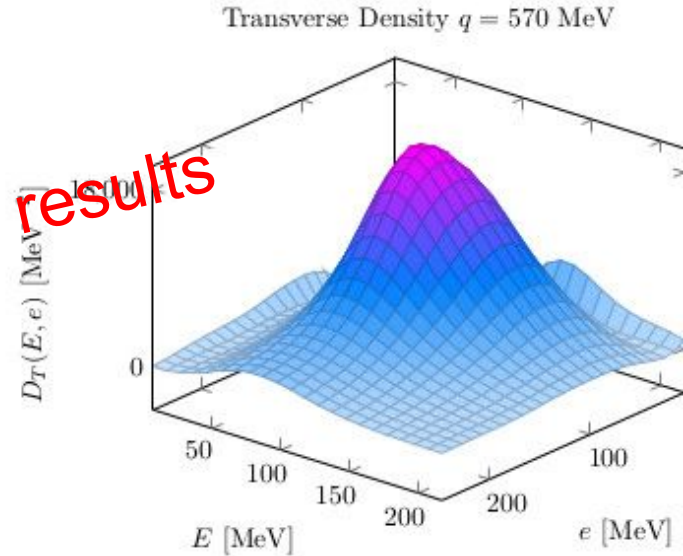
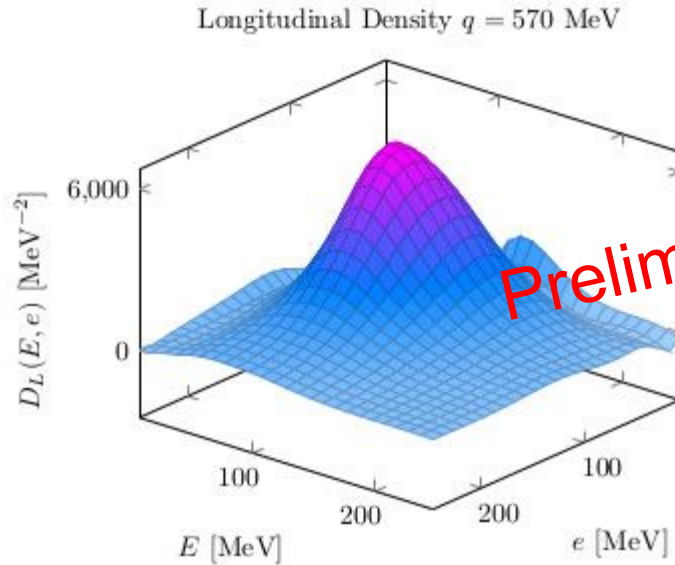
Andreoli Carlson Lovato SP Rocco
Phys.Rev.C 105 (2022) 1, 014002
PRC Editors' suggestion

Work in progress: $A=12$



Lorenzo Andreoli *et al.* in preparation

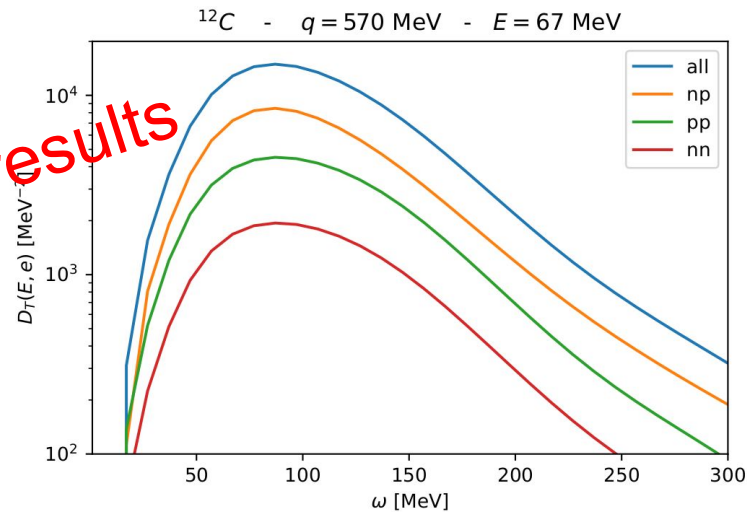
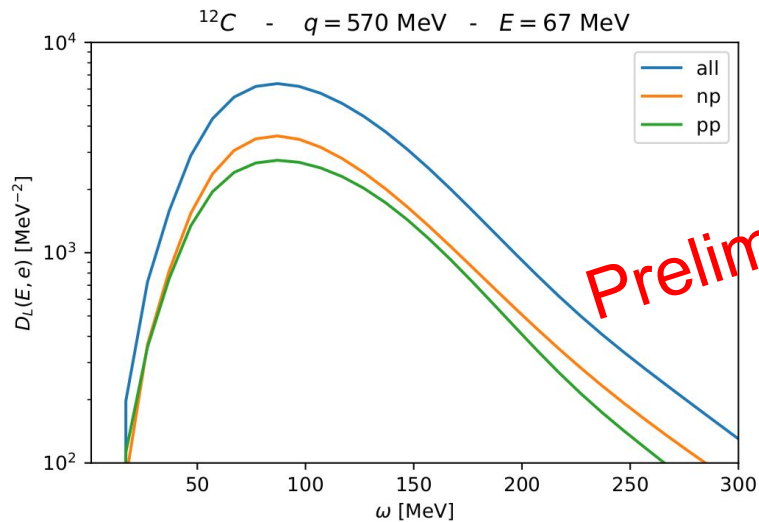
STA for Carbon 12: Response densities



Preliminary results

Andreoli *et al.* in preparation

STA for Carbon 12: Particle identity



Preliminary results

Andreoli *et al.* in preparation

Outlook

- Complete calculation on $A=12$
- Implement response densities in event generators (e.g., GENIE); leading force Barrow & Gardiner *et al.*
- Implement electroweak currents to study neutrino scattering processes
- Implement STA in other QMC methods for medium mass nuclei

Collaborators

WashU: **Andreoli Bub King** Piarulli

LANL: Carlson Cirigliano Gandolfi Hayes Mereghetti

JLab+ODU: Schiavilla

ANL: Lovato Rocco Wiringa

MIT/TAU: **Barrow**

FNAL: Betancourt **Gardiner**

UW: Dekens

Pisa U/INFN: Kievsky Marcucci Viviani

Salento U: Girlanda

Huzhou U: Dong Wang



Theory Alliance
FACILITY FOR RARE ISOTOPE BEAMS



U.S. DEPARTMENT OF
ENERGY

Office of
Science



Quantum Monte Carlo Methods

Minimize the expectation value of the nuclear Hamiltonian: $H = T + V_{ij} + V_{ijk}$

$$E_V = \frac{\langle \Psi_V | H | \Psi_V \rangle}{\langle \Psi_V | \Psi_V \rangle} \geq E_0$$

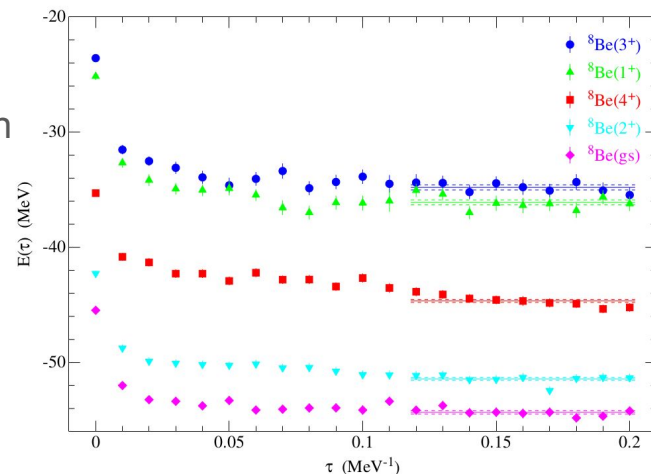
using the trial wave function:

$$|\Psi_V\rangle = \left[\mathcal{S} \prod_{i<j} (1 + U_{ij} + \sum_{k \neq i,j} U_{ijk}) \right] \left[\prod_{i<j} f_c(r_{ij}) \right] |\Phi_A(JMTT_3)\rangle$$

Further improve the trial wave function by eliminating spurious contaminations via a Green's Function Monte Carlo propagation in imaginary time

$$\Psi(\tau) = \exp[-(H - E_0)\tau] \Psi_V = \sum_n \exp[-(E_n - E_0)\tau] a_n \psi_n$$

$$\Psi(\tau \rightarrow \infty) = a_0 \psi_0$$



Carlson, Wiringa, Pieper *et al.*