



# Status of muon $g-2$ theory

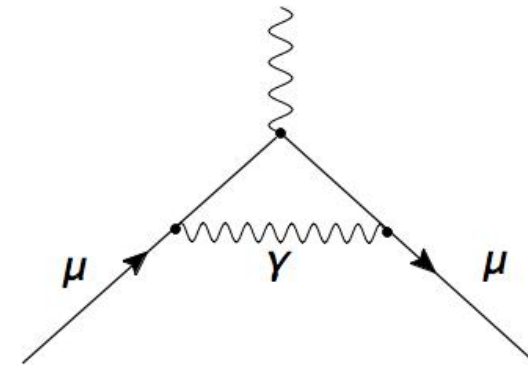
**I** Aida X. El-Khadra  
University of Illinois

Fermilab Public PAC Meeting  
16-18 November 2021

# Muon g-2: SM contributions

$$a_\mu = a_\mu(\text{QED}) + a_\mu(\text{EW}) + a_\mu(\text{hadronic})$$

QED

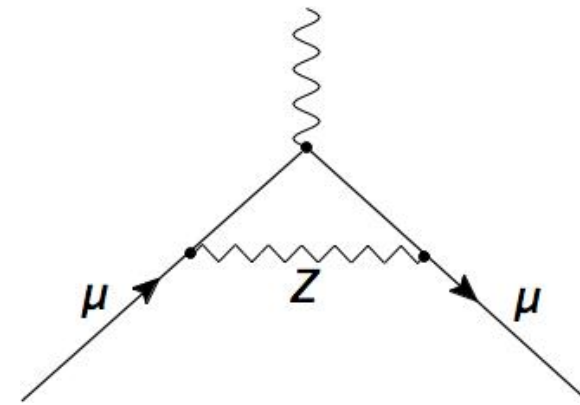


+...

$$116\,584\,718.9(1) \times 10^{-11}$$

0.001 ppm

EW

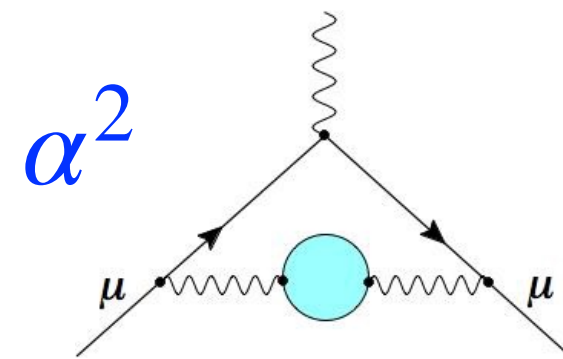


+...

$$153.6(1.0) \times 10^{-11}$$

0.01 ppm

HVP



+...

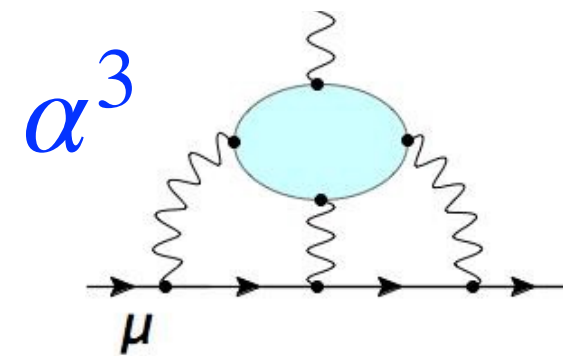
$$6845(40) \times 10^{-11}$$

[0.6%]

0.34 ppm

Hadronic corrections

HLbL

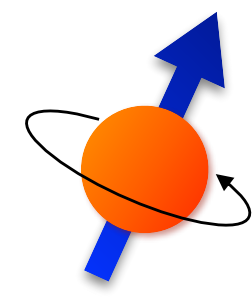


+...

$$92(18) \times 10^{-11}$$

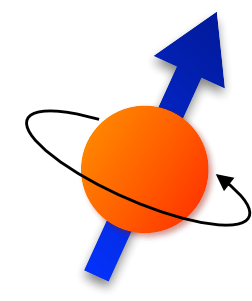
[20%]

0.15 ppm



# Muon $g-2$ Theory Initiative

- Maximize the impact of the Fermilab and J-PARC experiments
  - ▮ quantify and reduce the theoretical uncertainties on the hadronic corrections
- summarize the theory status and assess reliability of uncertainty estimates
- organize workshops to bring the different communities together:
  - [First plenary workshop @ Fermilab: 3-6 June 2017](#)
  - [HVP workshop @ KEK: 12-14 February 2018](#)
  - [HLbL workshop @ U Connecticut: 12-14 March 2018](#)
  - [Second plenary workshop @ HIM \(Mainz\): 18-22 June 2018](#)
  - [Third plenary workshop @ INT \(Seattle\): 9-13 September 2019](#)
  - [Lattice HVP at high precision workshop \(virtual\): 16-20 November 2020](#)
  - [Fourth plenary workshop @ KEK \(virtual\): 28 June - 02 July 2021](#)
  - [Fifth plenary workshop @ Higgs Centre \(Edinburgh\): 5-9 September 2022](#)
- 1<sup>st</sup> White Paper published in 2020 (132 authors, 82 institutions, 21 countries)  
[T. Aoyama et al, [arXiv:2006.04822](#), Phys. Repts. 887 (2020) 1-166.]
- 2<sup>nd</sup> White Paper (~2023): First discussions @ KEK meeting in June 2021  
expect to develop a concrete plan (outline, authors) @ Higgs Centre workshop



# Muon $g-2$ Theory Initiative

## Steering Committee

Simon Eidelman



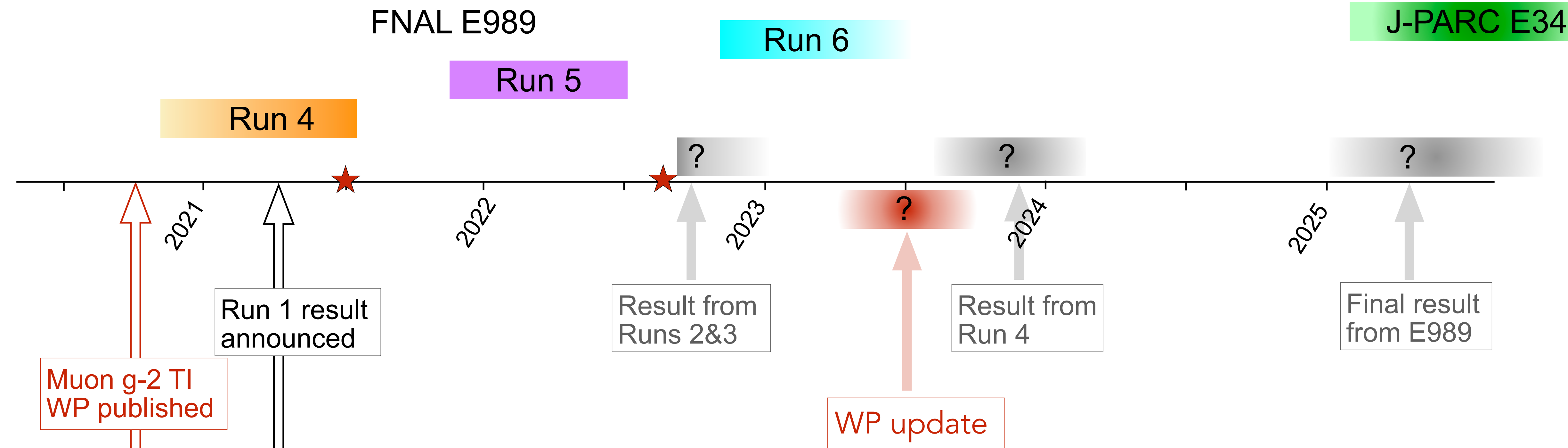
[photo by Hartmut Wittig]

(1948-2021)

- Gilberto Colangelo (Bern)
- Michel Davier (Orsay) co-chair
- Aida El-Khadra (UIUC & Fermilab) chair
- Martin Hoferichter (Bern)
- Christoph Lehner (Regensburg University & BNL) co-chair
- Laurent Lellouch (Marseille)
- Tsutomu Mibe (KEK) J-PARC Muon  $g-2$ /EDM experiment
- Lee Roberts (Boston) Fermilab Muon  $g-2$  experiment
- Thomas Teubner (Liverpool)
- Hartmut Wittig (Mainz)

website: <https://muon-gm2-theory.illinois.edu>

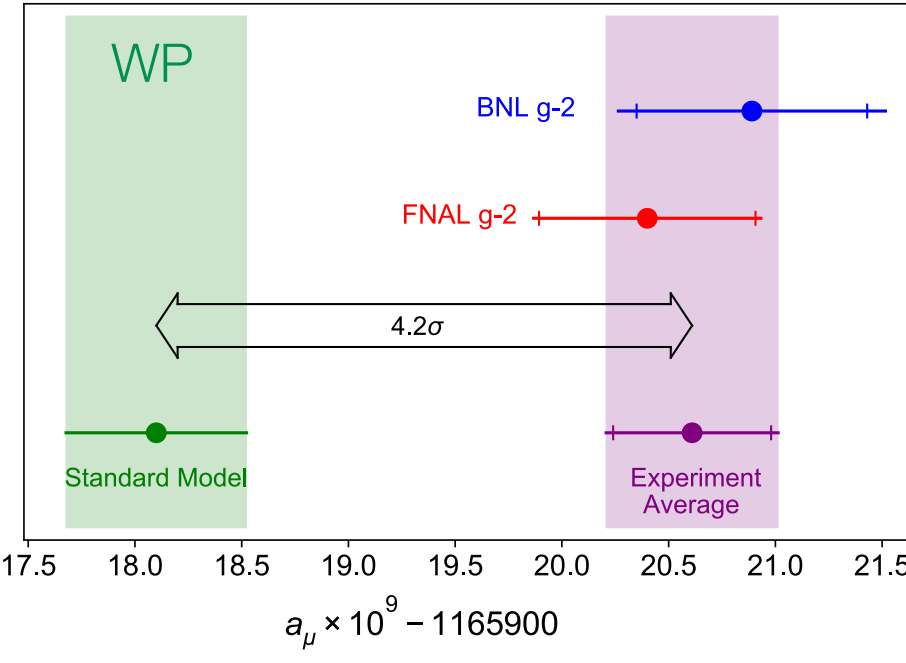
# Timeline



Physics Reports 887 (2021) 1–106  
 Contents lists available at ScienceDirect  
 Physics Reports  
 journal homepage: www.elsevier.com/locate/physrep

The anomalous magnetic moment of the muon in the Standard Model

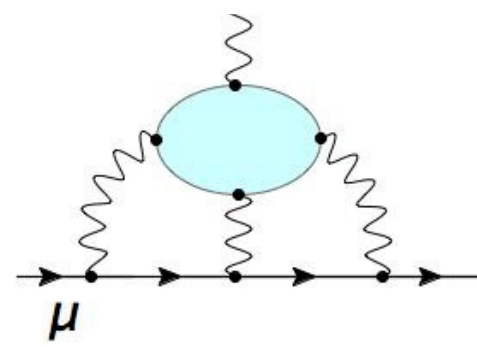
T. Aoyama<sup>1,2,3</sup>, N. Asmussen<sup>4</sup>, M. Benayoun<sup>5</sup>, J. Bijnens<sup>6</sup>, T. Blum<sup>7,8</sup>, M. Bruno<sup>9</sup>, I. Caprini<sup>10</sup>, C.M. Cariani Calame<sup>11</sup>, M. Cè<sup>12,13</sup>, G. Colangelo<sup>14</sup>, F. Crotty<sup>15</sup>, H. Czyż<sup>16</sup>, J. Danilkin<sup>17</sup>, M. Davier<sup>18</sup>, C.H. Davies<sup>19</sup>, M. Della Morte<sup>20</sup>, S.I. Eidelman<sup>21,22</sup>, A.K. Ekmacı<sup>23</sup>, A. Gérardin<sup>24</sup>, D. Giusti<sup>25</sup>, M. Golterman<sup>26</sup>, Steven Gottlieb<sup>27</sup>, V. Gubler<sup>28</sup>, F. Hagelstein<sup>29</sup>, M. Hayakawa<sup>30</sup>, G. Herdozia<sup>31</sup>, D.W. Hertzog<sup>32</sup>, A. Hoecker<sup>33</sup>, M. Hofrichter<sup>34</sup>, B.-L. Hoid<sup>35</sup>, R.J. Hudspeth<sup>36</sup>, F. Ignotov<sup>37</sup>, T. Izubuchi<sup>38</sup>, F. Jegerlehner<sup>39</sup>, L. Jin<sup>40</sup>, A. Keshavarzi<sup>41</sup>, T. Kinoshita<sup>42</sup>, B. Kubis<sup>43</sup>, A. Kupchynko<sup>44</sup>, A. Kuznetsov<sup>45</sup>, I. Laudi<sup>46</sup>, C. Lehner<sup>47</sup>, I. Lellouch<sup>48</sup>, I. Logashenko<sup>49</sup>, B. Malaescu<sup>50</sup>, K. Maltman<sup>51</sup>, M.K. Marinković<sup>52</sup>, P. Masjuan<sup>53</sup>, A.S. Meyer<sup>54</sup>, H.B. Meyer<sup>55</sup>, T. Mibe<sup>56</sup>, K. Mura<sup>57</sup>, S.E. Müller<sup>58</sup>, M. Nio<sup>59</sup>, D. Nomura<sup>60</sup>, A. Nyfeler<sup>61</sup>, V. Pascalutsa<sup>62</sup>, M. Passera<sup>63</sup>, E. Perez del Rio<sup>64</sup>, S. Peris<sup>65</sup>, A. Portelli<sup>66</sup>, M. Procura<sup>67</sup>, C.F. Redmer<sup>68</sup>, B.L. Roberts<sup>69</sup>, J. Sánchez-Puertas<sup>70</sup>, S. Seidenfaden<sup>71</sup>, B. Schwartz<sup>72</sup>, S. Simula<sup>73</sup>, D. Stöckinger<sup>74</sup>, H. Stöckinger-Kim<sup>75</sup>, P. Stoffer<sup>76</sup>, T. Teubner<sup>77</sup>, R. Van de Water<sup>78</sup>, M. Vanderhaeghe<sup>79</sup>, G. Venanzoni<sup>80</sup>, G. von Hippel<sup>81</sup>, H. Wittig<sup>82</sup>, Z. Zhang<sup>83</sup>, M.N. Acharya<sup>84</sup>, A. Bashir<sup>85</sup>, N. Cardoso<sup>86</sup>, B. Chakraborty<sup>87</sup>, E.-H. Cho<sup>88</sup>, J. Charles<sup>89</sup>, A. Crivellin<sup>90</sup>, O. Deineka<sup>91</sup>, A. Denig<sup>92</sup>, C. DeTar<sup>93</sup>, C.A. Dominguez<sup>94</sup>, A.E. Dorokhov<sup>95</sup>, V.P. Druzhinin<sup>96</sup>, G. Eichmann<sup>97</sup>, M. Fael<sup>98</sup>, C.S. Fischer<sup>99</sup>, E. Gdartz<sup>100</sup>, Z. Geiser<sup>101</sup>, J.R. Green<sup>102</sup>, S. Guellati-Khelifa<sup>103</sup>, D. Hatton<sup>104</sup>, R. Herrmannsson-Truesdell<sup>105</sup>, S. Holz<sup>106</sup>, B. Hörz<sup>107</sup>, M. Knecht<sup>108</sup>, J. Koponen<sup>109</sup>, A.S. Kronfeld<sup>110</sup>, I. Laiso<sup>111</sup>, S. Leupold<sup>112</sup>, P.B. Mackenzie<sup>113</sup>, W.J. Marciano<sup>114</sup>, C. McNeile<sup>115</sup>, D. Mohler<sup>116</sup>, J. Monnard<sup>117</sup>, E.T. Neil<sup>118</sup>, A.V. Nesterenko<sup>119</sup>, K. Ottnaad<sup>120</sup>, V. Pauk<sup>121</sup>, A.E. Radhabov<sup>122</sup>, E. de Rafael<sup>123</sup>, K. Raya<sup>124</sup>, A. Rich<sup>125</sup>, A. Rodríguez-Sánchez<sup>126</sup>, P. Roig<sup>127</sup>, T. San José<sup>128</sup>, E.P. Solodov<sup>129</sup>, R. Sugar<sup>130</sup>, K. Yu. Todoshyn<sup>131</sup>, A. Vainshtein<sup>132</sup>, A. Vagstad Avilés-Casco<sup>133</sup>, E. Weil<sup>134</sup>, J. Wilhelm<sup>135</sup>, R. Williams<sup>136</sup>, A.S. Zhelezovskiy<sup>137</sup>



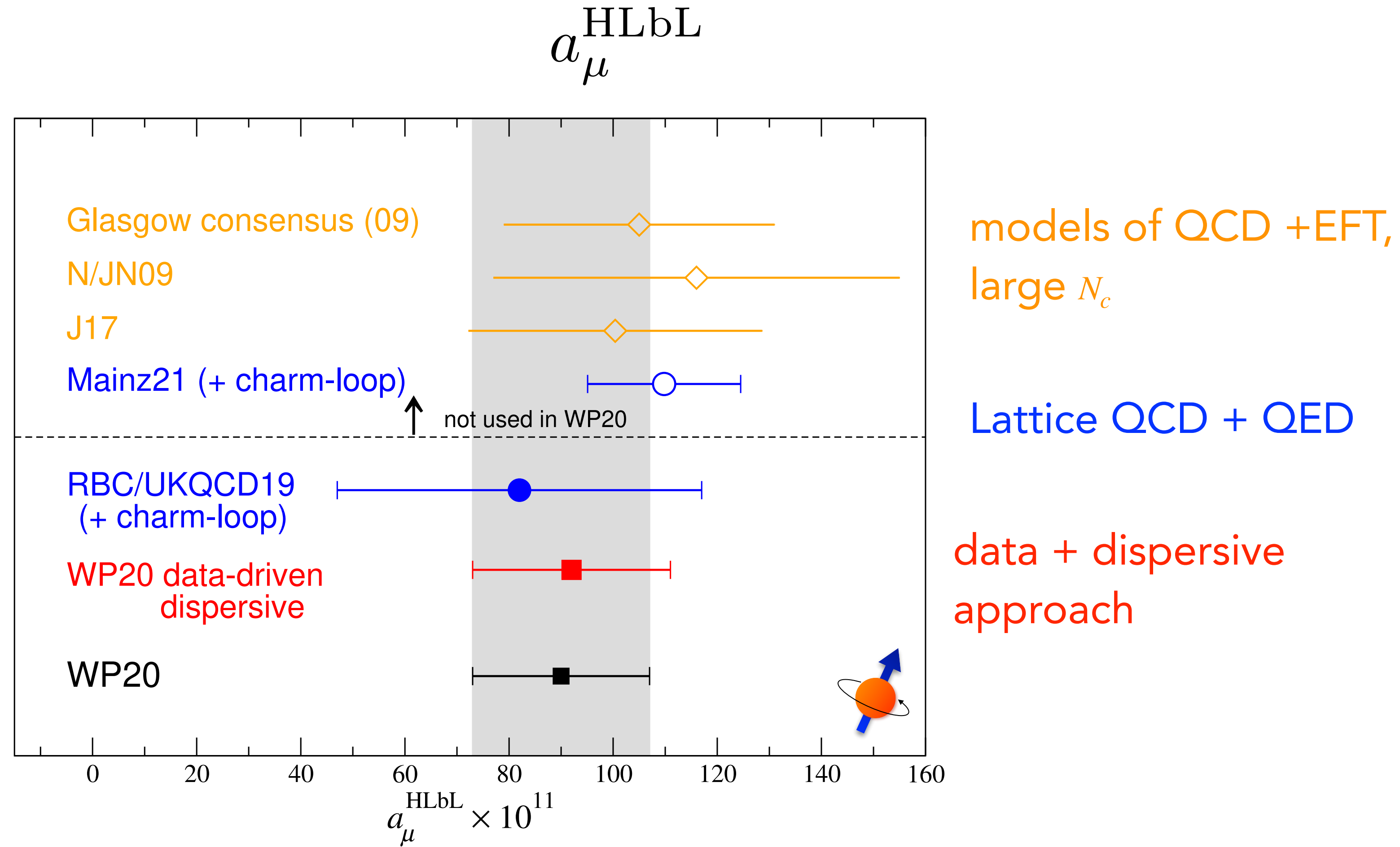
## Theory Initiative:

- ★ ongoing activities: develop method average for Lattice HVP
- ★ plan to update WP with new SM predictions (~ 2023)

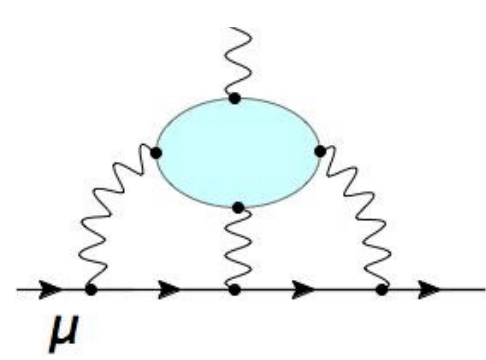
★ TI workshops: Jun 2021 @ KEK (virtual)  
 Sep 2022 @ Higgscentre



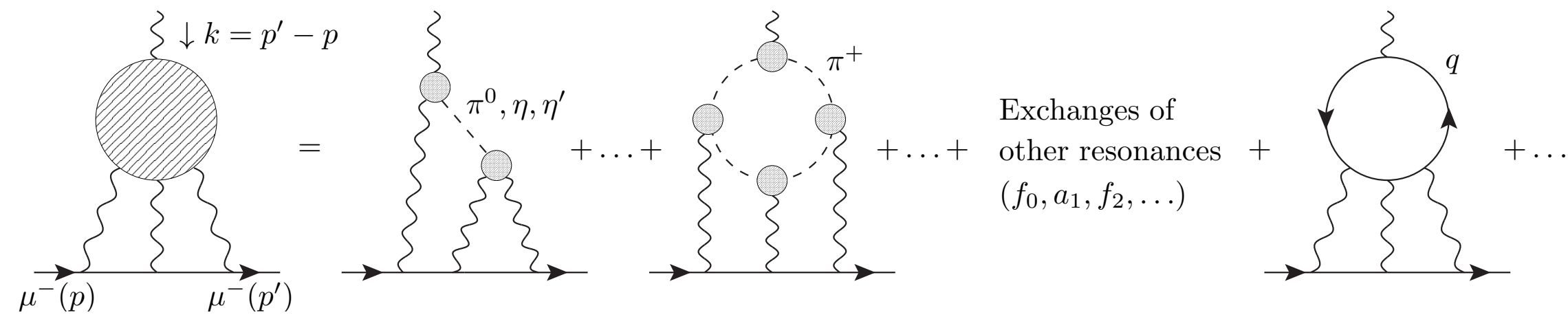
# HLbL: Comparison



Now well-determined in two independent approaches, systematically improvable



# Hadronic Light-by-light

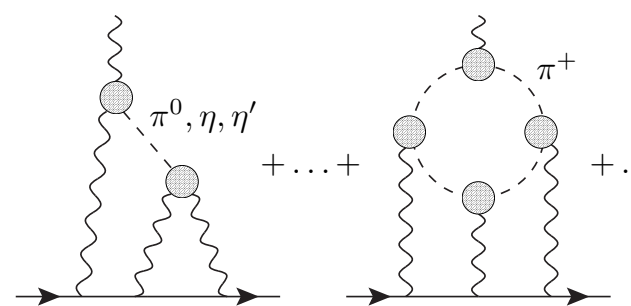


## Dispersive approach:

[Colangelo et al, 2014; Pauk & Vanderhaegen 2014; ...]

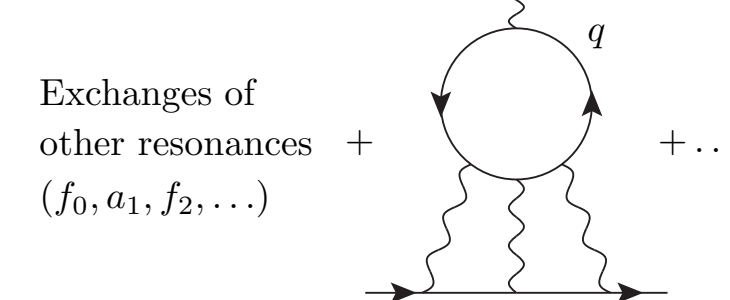
- ◆ model independent
- ◆ significantly more complicated than for HVP
- ◆ provides a framework for data-driven evaluations
- ➡◆ can also use lattice results as inputs

## Dominant contributions (≈ 75 % of total):



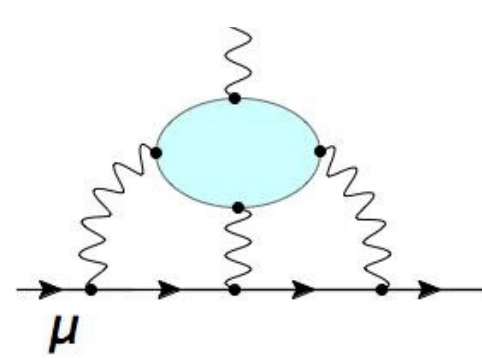
- ◆ Well quantified with ≈ 6 % uncertainty
- ◆  $\eta, \eta'$  pole contributions: Canterbury approximants only
- ◆ Ongoing work: consolidation of  $\eta, \eta'$  pole contributions using disp. relations and LQCD

## Subleading contributions (≈ 15 % of total):



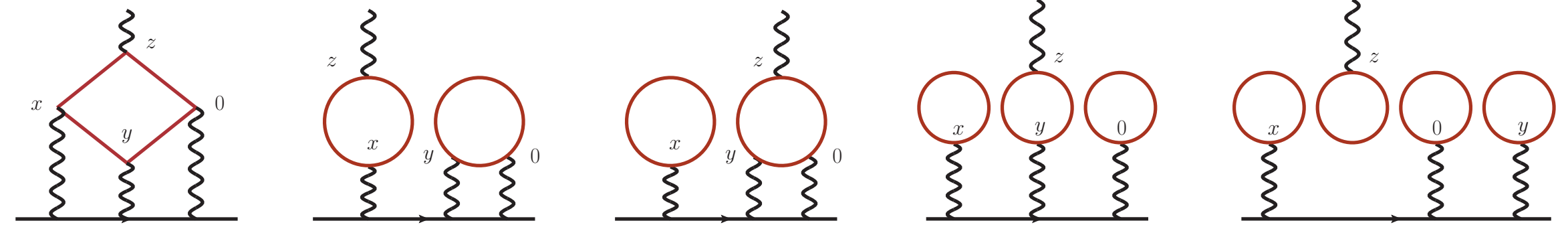
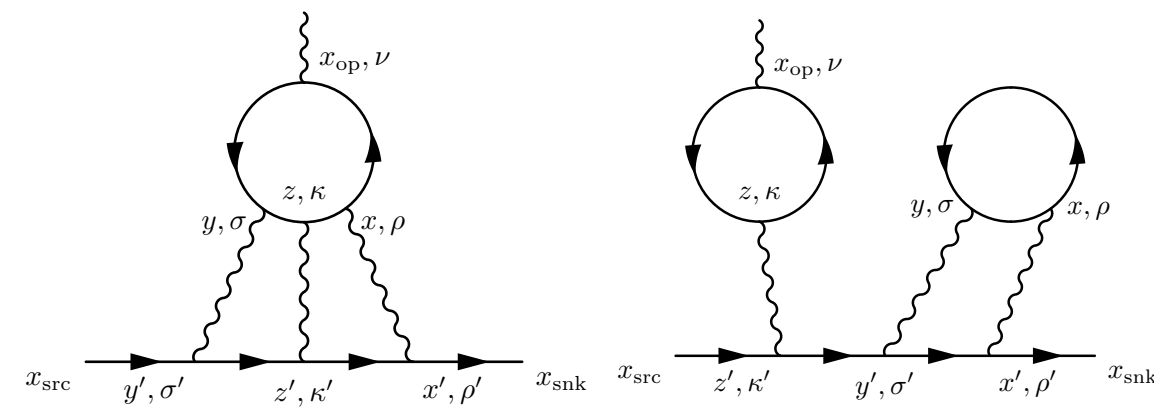
- ◆ Not yet well known
  - ➡◆ dominant contribution to total uncertainty
- ◆ Ongoing work:
  - Implementation of short-distance constraints (now at 2-loop)
  - DR implementation for axial vector contributions
  - BESIII ramping up  $\gamma^{(*)}\gamma^*$  program

Dispersive, data-driven evaluation of HLbL with  $\leq 10\%$  total uncertainty feasible by ~2025.



# Hadronic Light-by-light

Lattice QCD+QED: Two independent and complete direct calculations of  $a_\mu^{\text{HLbL}}$



## ◆ RBC/UKQCD

[T. Blum et al, arXiv:1610.04603, 2016 PRL; arXiv:1911.08123, 2020 PRL]

## ◆ QCD + QED<sub>L</sub> (finite volume)

DWF ensembles at/near phys mass,  
 $a \approx 0.08 - 0.2 \text{ fm}$ ,  $L \sim 4.5 - 9.3 \text{ fm}$

## ◆ Mainz group

[E. Chao et al, arXiv:2104.02632]

## ◆ QCD + QED (infinite volume & continuum)

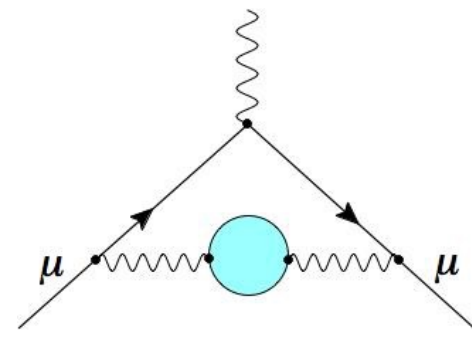
CLS (2+1 Wilson-clover) ensembles

$m_\pi \sim 200 - 430 \text{ MeV}$ ,  $a \approx 0.05 - 0.1 \text{ fm}$ ,  $m_\pi L > 4$

- ◆ Cross checks between RBC/UKQCD & Mainz approaches in White Paper at unphysical pion mass
- ◆ Both groups will continue to improve their calculations, adding more statistics, lattice spacings, physical mass ensemble (Mainz)

Lattice HLbL results with 10% total uncertainty feasible by ~2025





# HVP: Comparison

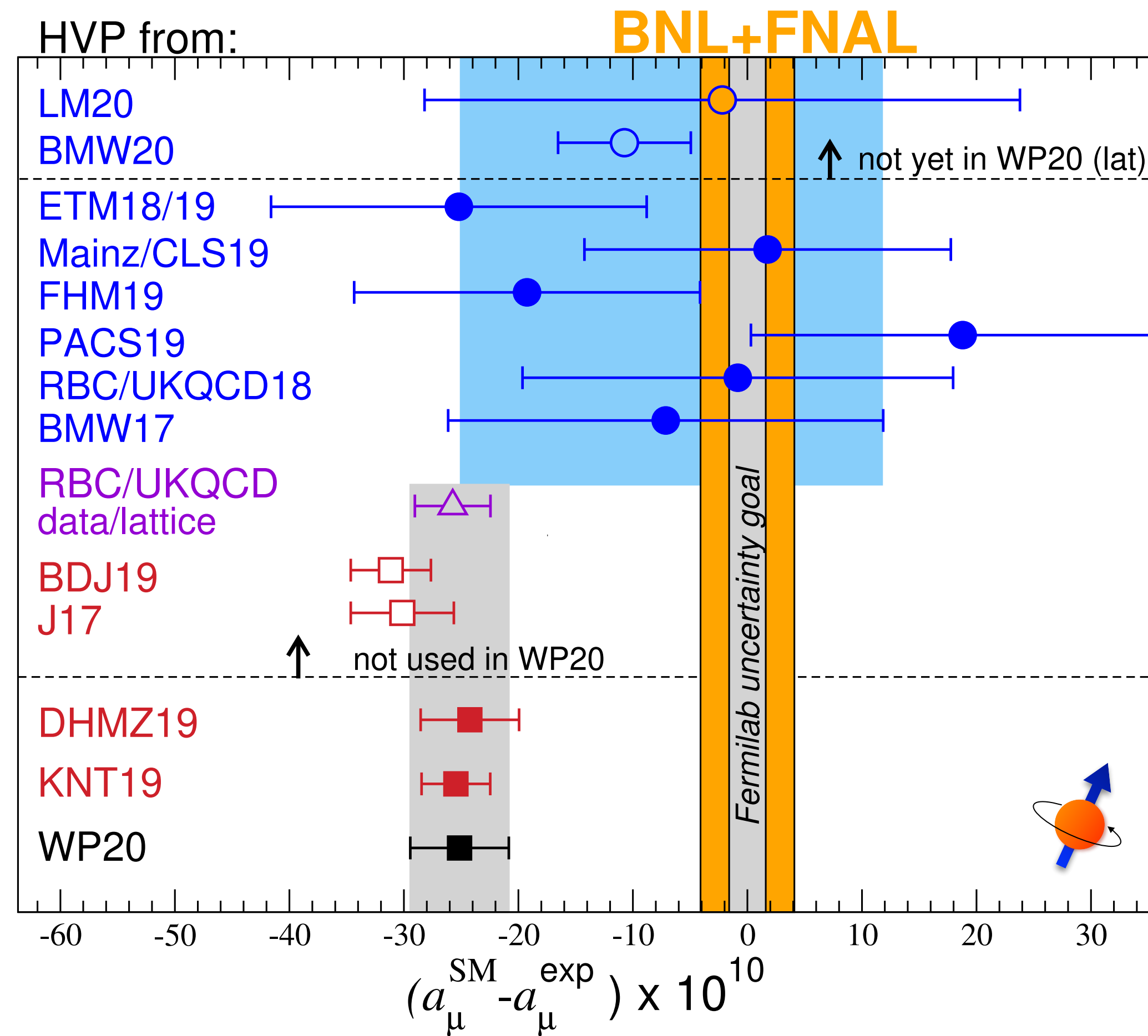
$$a_{\mu}^{\text{SM}} = a_{\mu}^{\text{HVP}} + [a_{\mu}^{\text{QED}} + a_{\mu}^{\text{Weak}} + a_{\mu}^{\text{HLbL}}]$$

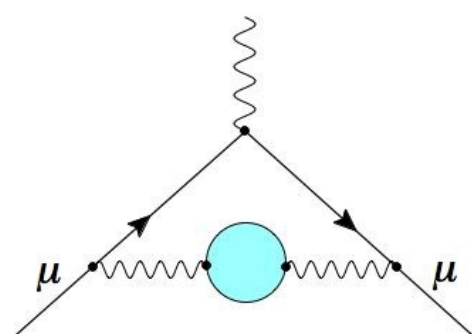
Lattice QCD + QED

hybrid: combine data & lattice

data driven

+ unitarity/analyticity constraints





# HVP: data-driven

## In 2020 WP:

Conservative merging procedure to obtain a realistic assessment of the underlying uncertainties:

- account for tensions between data sets
- account for differences in methodologies for compilation of experimental inputs
- include correlations between systematic errors
- cross checks from unitarity & analyticity constraints

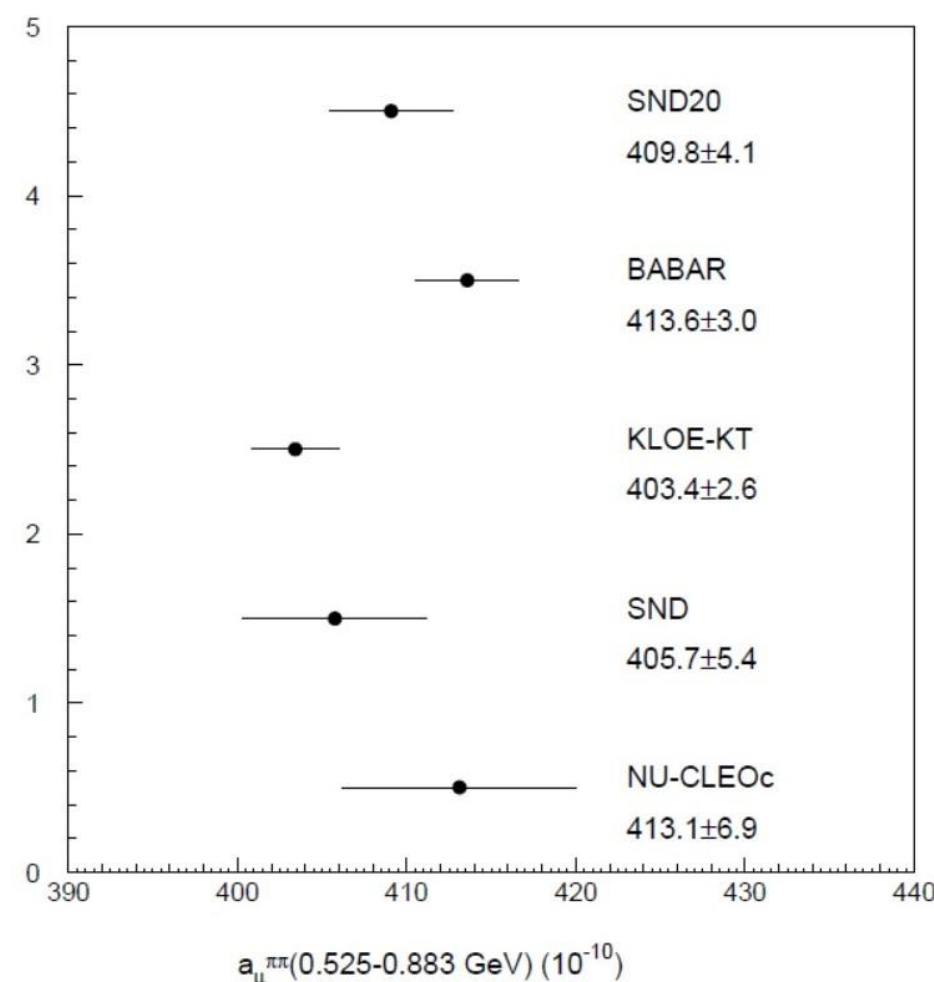
[Colangelo et al, 2018; Anantharayan et al, 2018; Davier et al, 2019; Hoferichter et al, 2019]

- Full NLO radiative corrections [Campanario et al, 2019]

$$a_{\mu}^{\text{HVP,LO}} = 693.1 (2.8)_{\text{exp}} (0.7)_{\text{DV+pQCD}} (2.8)_{\text{BaBar-KLOE}} \times 10^{-10}$$

$$= 693.1 (4.0) \times 10^{-10}$$

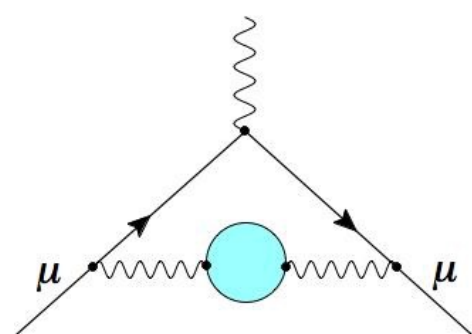
[M. Davier @ KEK workshop]



## Ongoing work:

- BaBar: new analysis of large (7x) data set in  $\pi\pi$  channel (1-2 years), also  $\pi\pi\pi$ , other channels
- SND: new results for  $\pi\pi$  channel, other channels in progress
- CMD-3: ongoing analyses for  $\pi\pi$  and other channels
- BESIII: new results in 2021 for  $\pi\pi$  channel, continued analysis also for  $\pi\pi\pi$ , other channels
- Belle II: will have high-statistics data for low-energy cross sections.
- Need blind analyses to resolve the tensions (esp. for  $\pi\pi$  channel)
- Developing NNLO Monte Carlo generators (STRONG 2020 workshop next week <https://agenda.infn.it/event/28089/>)

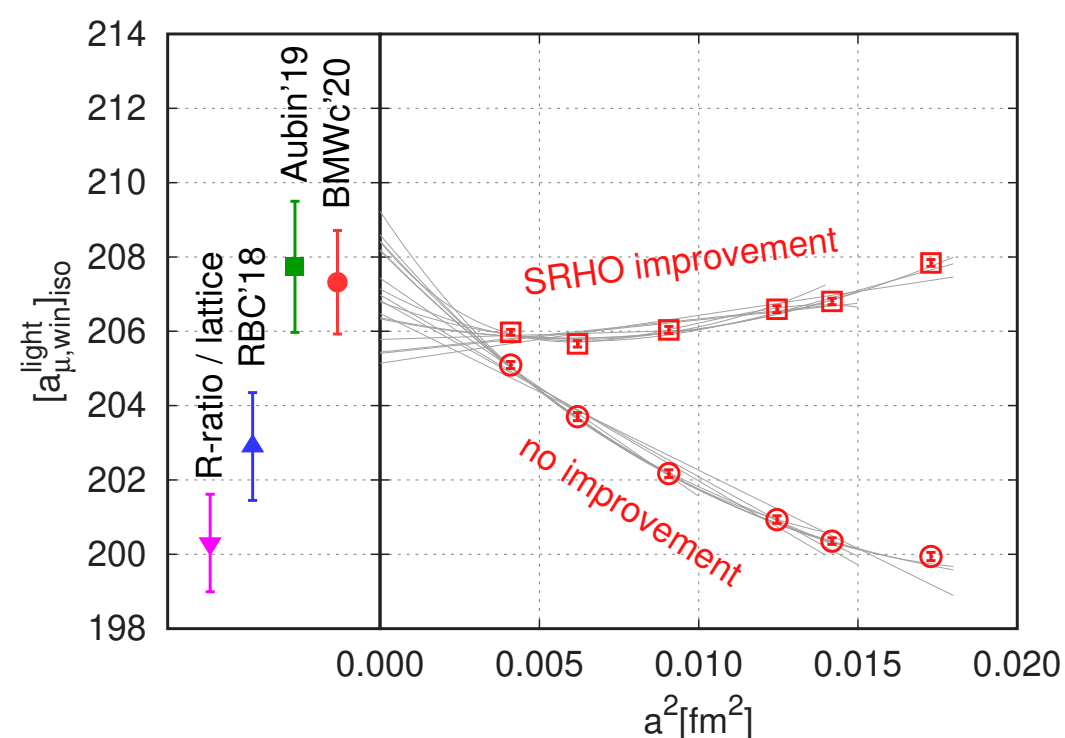
Data-driven evaluations of HVP with  $\sim 0.3\%$  feasible by  $\sim 2025$ , if tensions between experiments are resolved.



# HVP: lattice

## In 2020 WP:

- Lattice HVP average at 2.6 % total uncertainty:  
 $a_\mu^{\text{HVP,LO}} = 711.6 (18.4) \times 10^{10}$
- BMW 20 (published in 2021)  
 first LQCD calculation with sub-percent (0.8 %) error  
 but in tension with data-driven HVP ( $2.1\sigma$ )
- Further tensions for intermediate window:



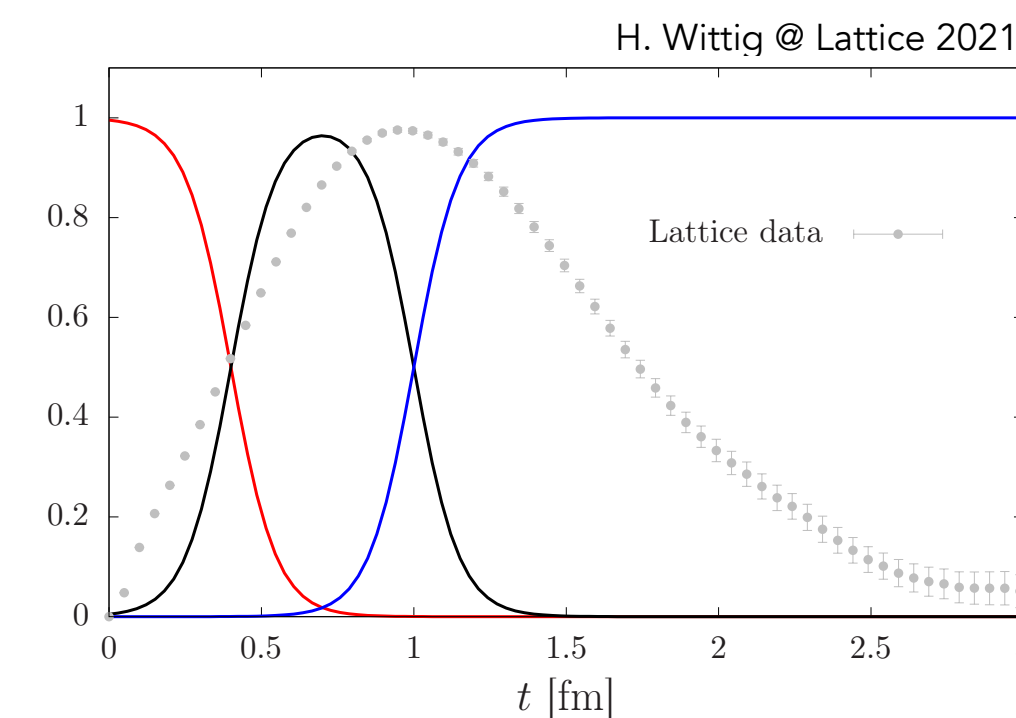
-3.7 $\sigma$  tension with data-driven evaluation  
 -2.2 $\sigma$  tension with RBC/UKQCD18

$$a_\mu^{\text{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dt \tilde{w}(t) C(t)$$

- Use windows in Euclidean time to consider the different time regions separately. [T. Blum et al, arXiv:1801.07224, 2018 PRL]

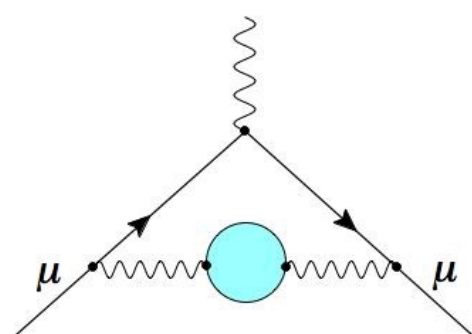
Short Distance (SD)  $t : 0 \rightarrow t_0$   
 Intermediate (W)  $t : t_0 \rightarrow t_1$   
 Long Distance (LD)  $t : t_1 \rightarrow \infty$

$$t_0 = 0.4 \text{ fm}, t_1 = 1.0 \text{ fm}$$



- disentangle systematics/statistics from long distance/FV and discretization effects
- intermediate window: easy to compute in lattice QCD & using disperse approach:
- Internal cross check:  
 Compute each window separately (in continuum, infinite volume limits,...) and combine:

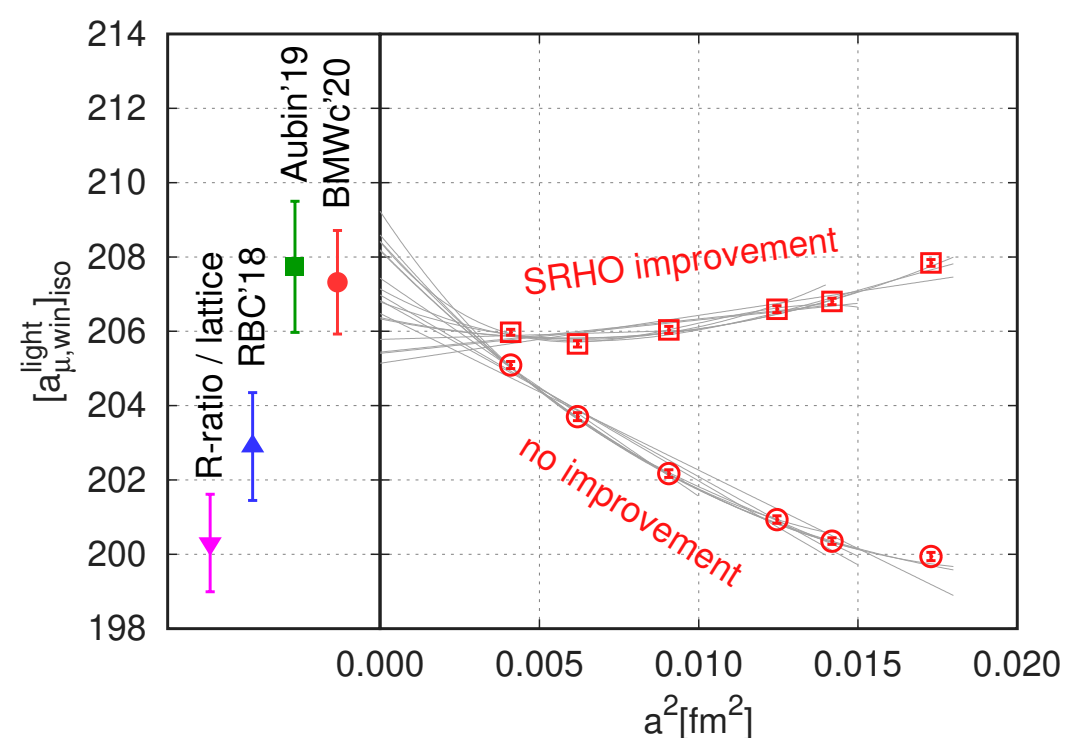
$$a_\mu = a_\mu^{\text{SD}} + a_\mu^{\text{W}} + a_\mu^{\text{LD}}$$



# HVP: lattice

## In 2020 WP:

- Lattice HVP average at 2.6 % total uncertainty:  
 $a_{\mu}^{\text{HVP,LO}} = 711.6 (18.4) \times 10^{10}$
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 first LQCD calculation with sub-percent (0.8 %) error  
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- Further tensions for intermediate window:

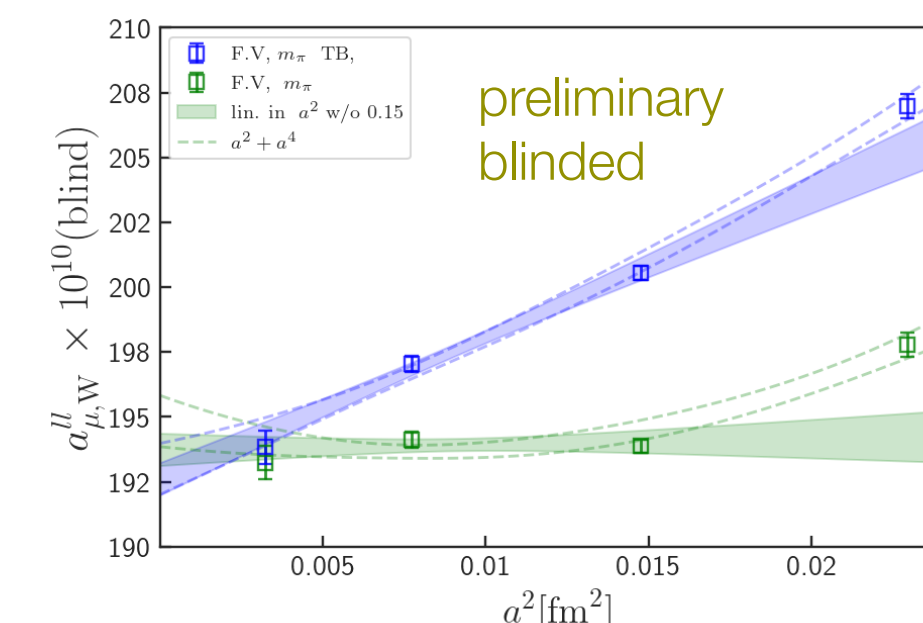


-3.7 $\sigma$  tension with data-driven evaluation  
 -2.2 $\sigma$  tension with RBC/UKQCD18

## Ongoing work:

- Expect new results from RBC/UKQCD and FNAL/MILC (and likely other lattice groups) in the coming months:

S. Lahert  
(FNAL/MILC)



- need blind analyses (already being done in FNAL/MILC and RBC/UKQCD)
- Including  $\pi\pi$  states for refined long-distance computation (Mainz, RBC/UKQCD, FNAL/MILC)
- Developing method average for lattice HVP — started at KEK workshop (June 2021), based on detailed comparisons
  - list of sub quantities (and their definitions)
  - common prescription for separating QCD & QED
  - quality criteria for inclusion
- Most groups plan to include smaller lattice spacings to test continuum extrapolations (needs adequate computational resources)

If results are consistent, Lattice HVP (average) with  $\lesssim 0.5\%$  errors feasible by 2025

# Summary and Outlook

## ★ Theory Initiative:

- WP update ~2023 will include any new available results and a method average for lattice HVP, HLbL
- Concrete plans for writing WP update (outline, authors,...) @ next workshop

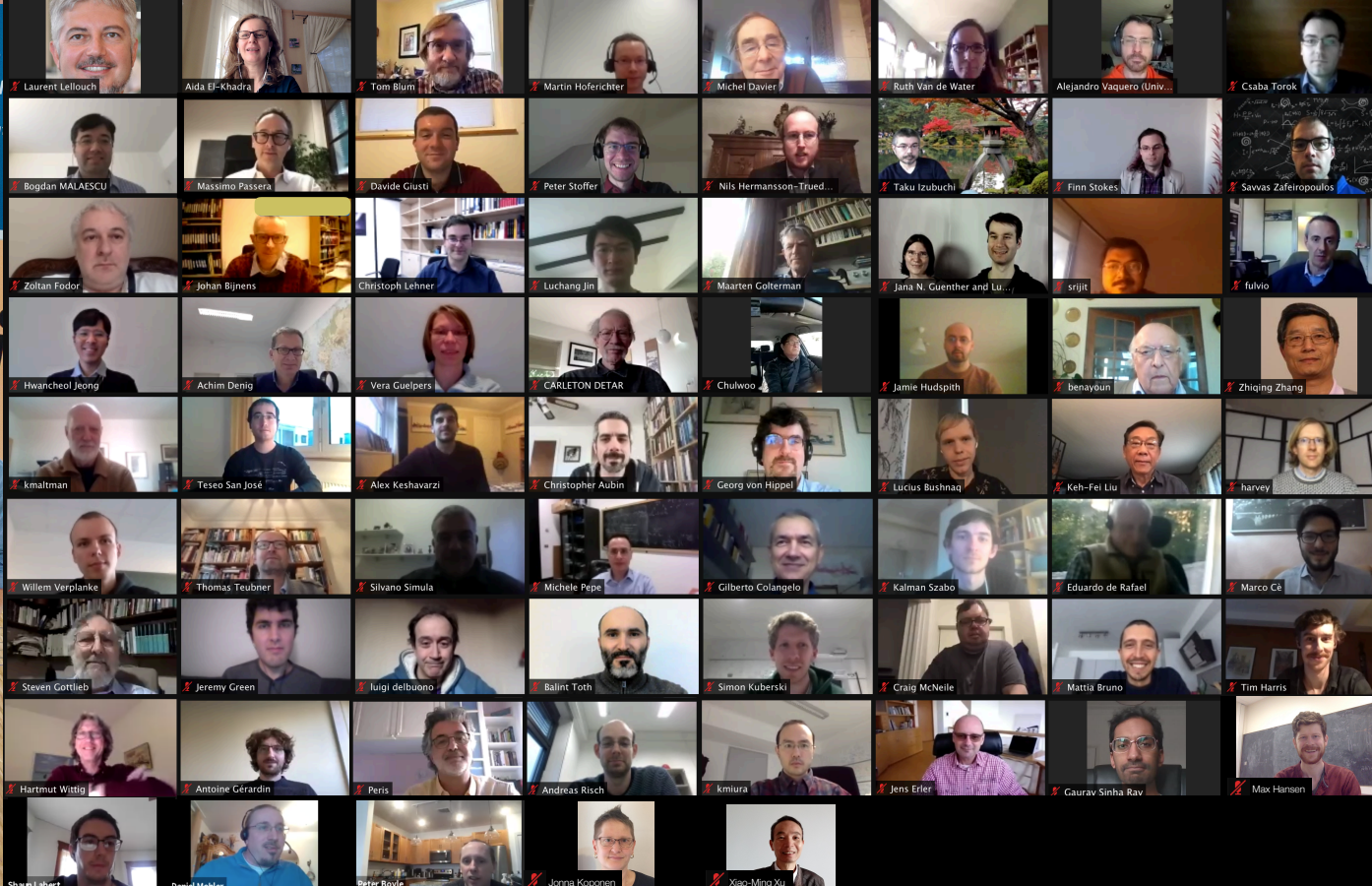
## ★ Programs and plans in place to improve:

- data-driven HVP  $\sim 0.3\%$
- lattice HVP  $\lesssim 0.5\%$
- dispersive HLbL and lattice HLbL:  $\sim 10\%$

... assuming tensions are resolved.

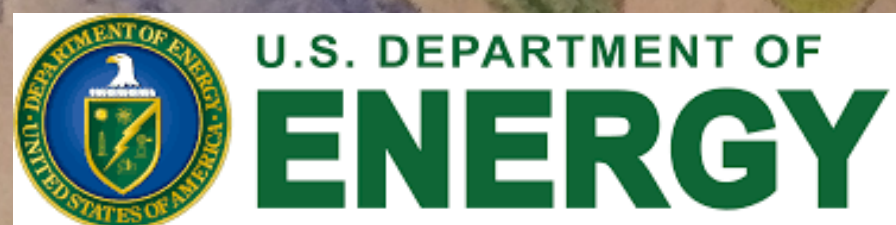
- ★ If tensions between data-driven HVP and lattice HVP are resolved, SM predictions will likely reach desired precision
- ★ Beyond 2025: MUonE (space-like momentum measurement of  $\Delta\alpha$ ) will provide more information/cross checks.

▣▣▣▣ Next workshop of the Muon g-2 Theory Initiative: 5-9 Sep 2022





UNIVERSITY of WASHINGTON



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Thank you!



# Appendix



# Updated WP Summary Table

Contribution	Value $\times 10^{11}$	References
Experimental average (E989+E821)	116592061(41)	<u>Phys.Rev.Lett. 124, 141801</u>
HVP LO ( $e^+e^-$ )	6931(40)	Refs. [2–7]
HVP NLO ( $e^+e^-$ )	−98.3(7)	Ref. [7]
HVP NNLO ( $e^+e^-$ )	12.4(1)	Ref. [8]
HVP LO (lattice, $udsc$ )	7116(184)	Refs. [9–17]
HLbL (phenomenology)	92(19)	Refs. [18–30]
HLbL NLO (phenomenology)	2(1)	Ref. [31]
HLbL (lattice, $uds$ )	79(35)	Ref. [32]
HLbL (phenomenology + lattice)	90(17)	Refs. [18–30, 32]
QED	116 584 718.931(104)	Refs. [33, 34]
Electroweak	153.6(1.0)	Refs. [35, 36]
HVP ( $e^+e^-$ , LO + NLO + NNLO)	6845(40)	Refs. [2–8]
HLbL (phenomenology + lattice + NLO)	92(18)	Refs. [18–32]
Total SM Value	116 591 810(43)	Refs. [2–8, 18–24, 31–36]
Difference: $\Delta a_\mu := a_\mu^{\text{exp}} - a_\mu^{\text{SM}}$	251(59)	

website: <https://muon-gm2-theory.illinois.edu>

# Muon g-2: SM contributions

$$a_\mu = a_\mu(\text{QED}) + a_\mu(\text{EW}) + a_\mu(\text{hadronic})$$

## QED

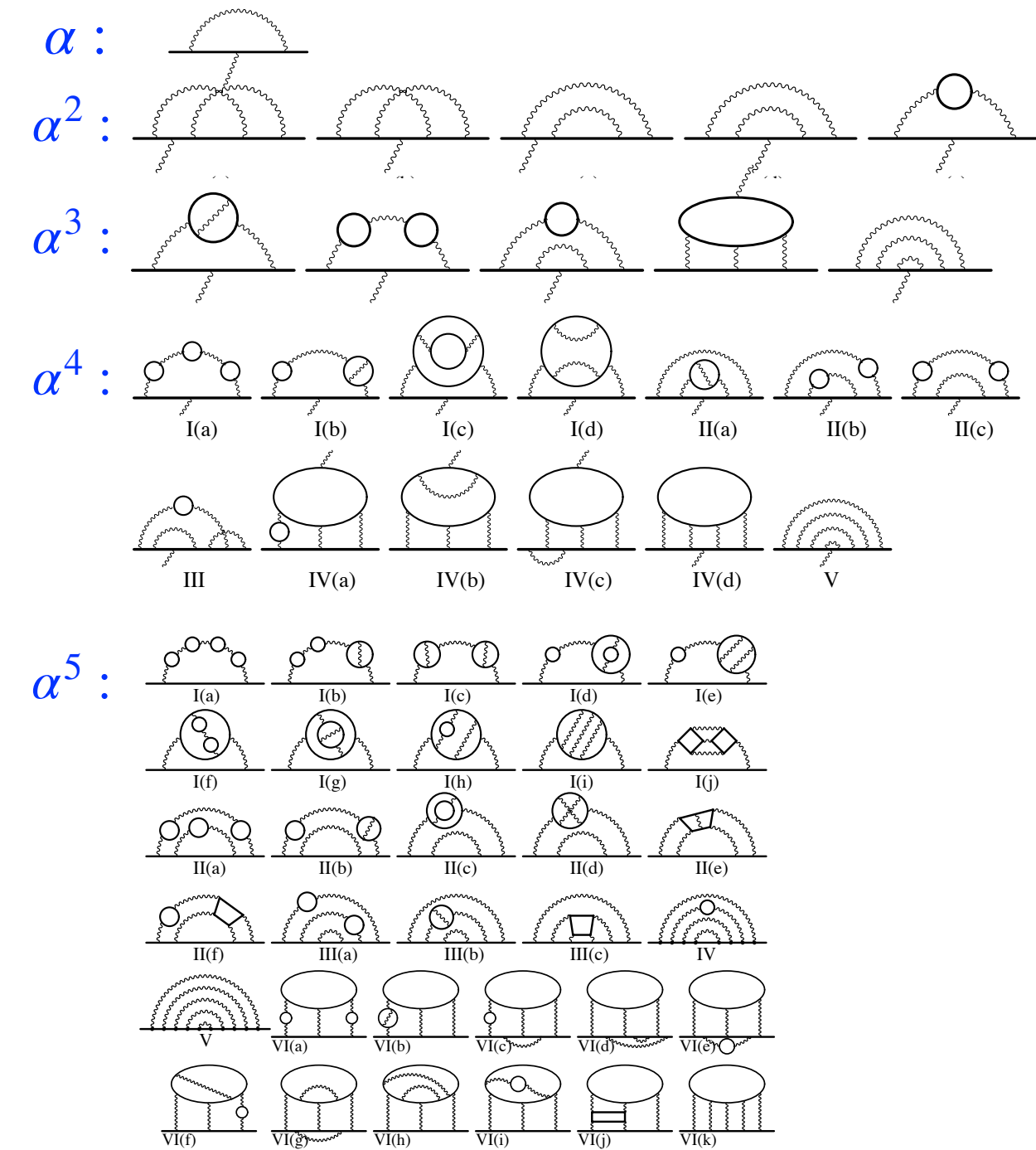
$$a_\mu(\text{QED}) = A_1 + A_2 \left( \frac{m_\mu}{m_e} \right) + A_2 \left( \frac{m_\mu}{m_\tau} \right) + A_3 \left( \frac{m_\mu}{m_e}, \frac{m_\mu}{m_\tau} \right)$$

$$A_i = \sum_{n=0} \left( \frac{\alpha}{\pi} \right)^n A_i^{2n}$$

$n$	# of diagrams	Contribution x $10^{11}$
1	1	116140973.32
2	7	413 217.63
3	71	30141.90
4	891	381.00
5	12672	5.08

$$a_\mu(\text{QED}) = 116\,584\,718.9(1) \times 10^{-11}$$

[T. Aoyama et al, arXiv:1205.5370, PRL;  
T. Aoyama, T. Kinoshita, M. Nio, Atoms 7 (1) (2019) 28]



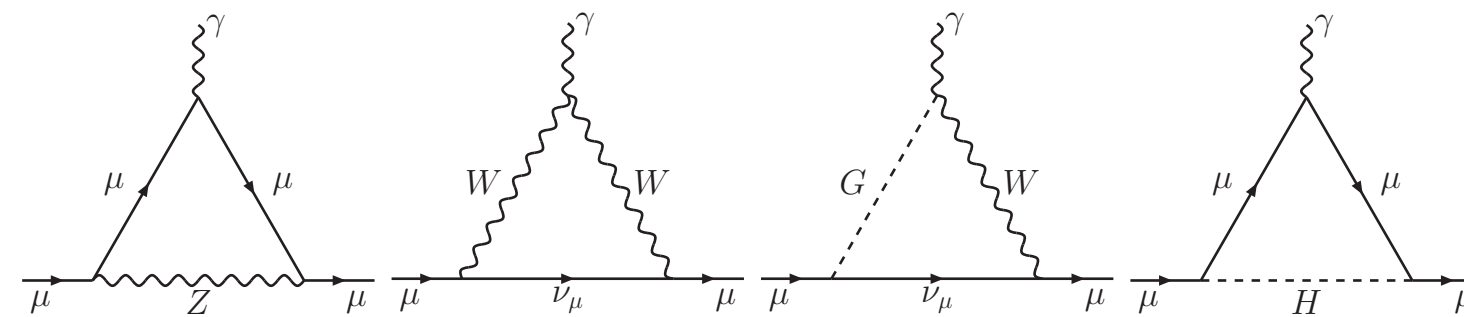
# Muon g-2: SM contributions

$$a_\mu = a_\mu(\text{QED}) + a_\mu(\text{EW}) + a_\mu(\text{hadronic})$$

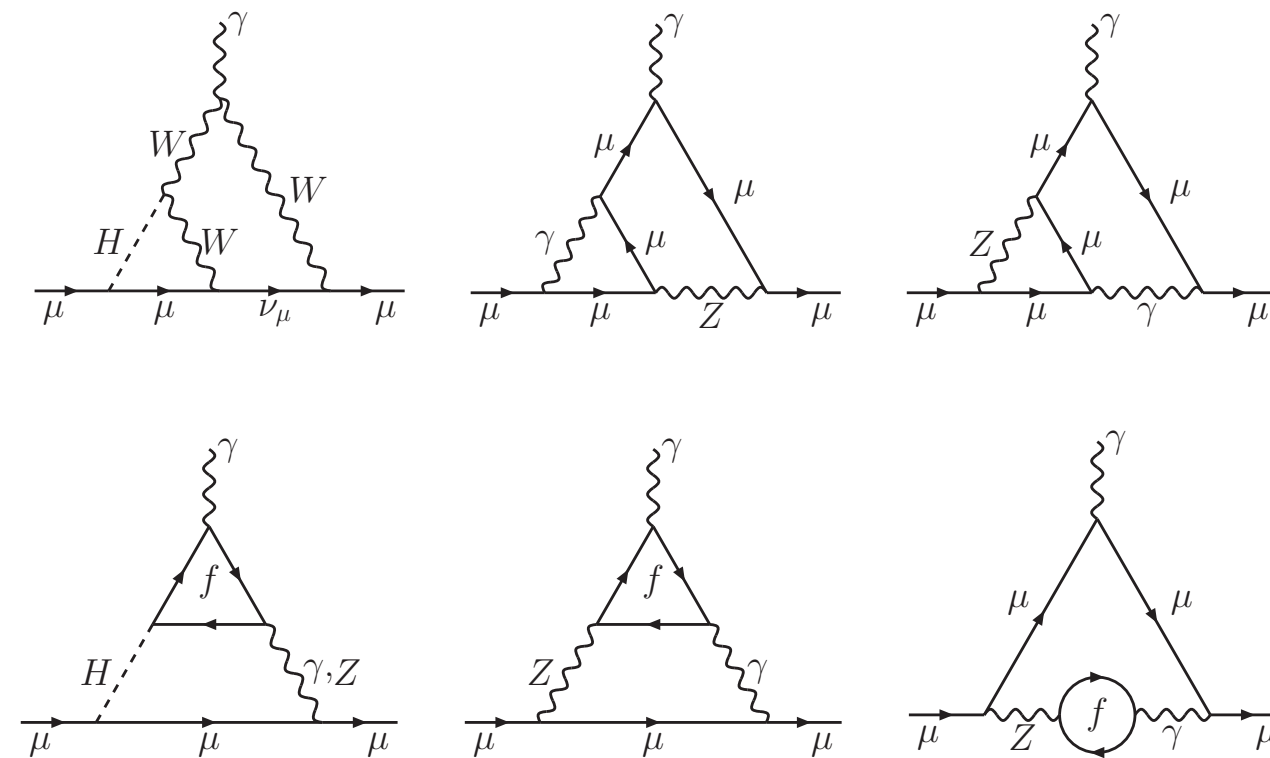
## Electroweak

(contributions from W,Z,H,.. bosons)

1-loop



2-loop



$$a_\mu(\text{EW}) = 153.6 (1.0) \times 10^{-11}$$

[A. Czarnecki et al, hep-ph/0212229, PRD;  
C. Gnendinger et al, arXiv:1306.5546, PRD]

# Muon g-2: SM contributions

$$a_\mu = a_\mu(\text{QED}) + a_\mu(\text{EW}) + a_\mu(\text{hadronic})$$

leading hadronic



◆ The hadronic contributions are written as:

$$a_\ell(\text{hadronic}) = a_\ell^{\text{HVP, LO}} + a_\ell^{\text{HVP, NLO}} + a_\ell^{\text{HVP, NNLO}} + \dots$$

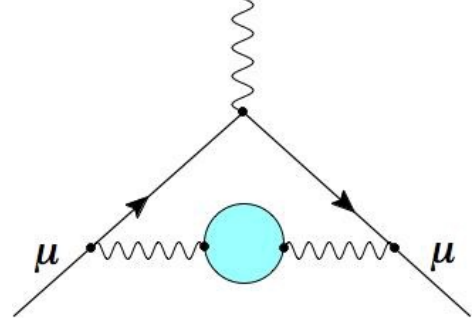
$$+ a_\ell^{\text{HLbL}} + a_\ell^{\text{HLbL, NLO}} + \dots$$

$\alpha^2$

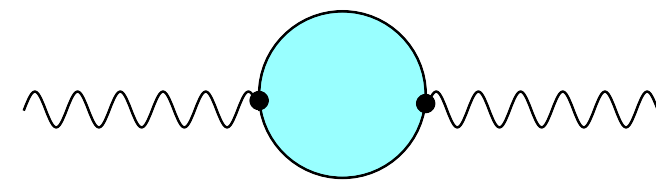
$\alpha^3$

$\alpha^4$

$\sim 10^{-7}$



# Hadronic vacuum polarization



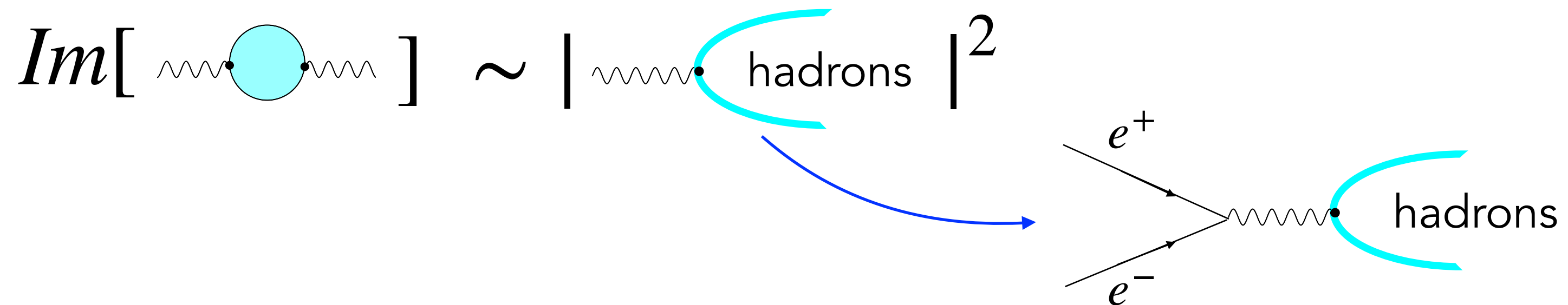
$$\hat{\Pi}(q^2) = \Pi(q^2) - \Pi(0)$$

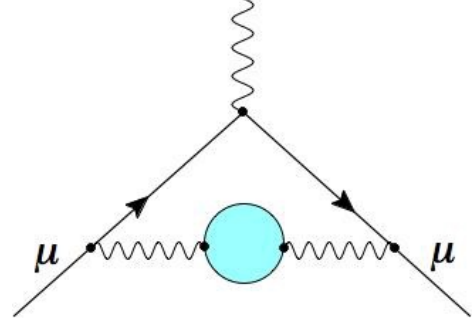
$$\Pi_{\mu\nu} = \int d^4x e^{iqx} \langle j_\mu(x) j_\nu(0) \rangle = (q_\mu q_\nu - q^2 g_{\mu\nu}) \Pi(q^2)$$

Leading order HVP correction:

$$a_\mu^{\text{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int dq^2 \omega(q^2) \hat{\Pi}(q^2)$$

- Use optical theorem and dispersion relation to rewrite the integral in terms of the hadronic  $e^+e^-$  cross section:

$$\text{Im} \left[ \text{Feynman diagram} \right] \sim \left| \text{Feynman diagram} \right|^2$$




# Hadronic vacuum polarization

$$\hat{\Pi}(q^2) = \Pi(q^2) - \Pi(0)$$

$$\Pi_{\mu\nu} = \int d^4x e^{iqx} \langle j_\mu(x) j_\nu(0) \rangle = (q_\mu q_\nu - q^2 g_{\mu\nu}) \Pi(q^2)$$

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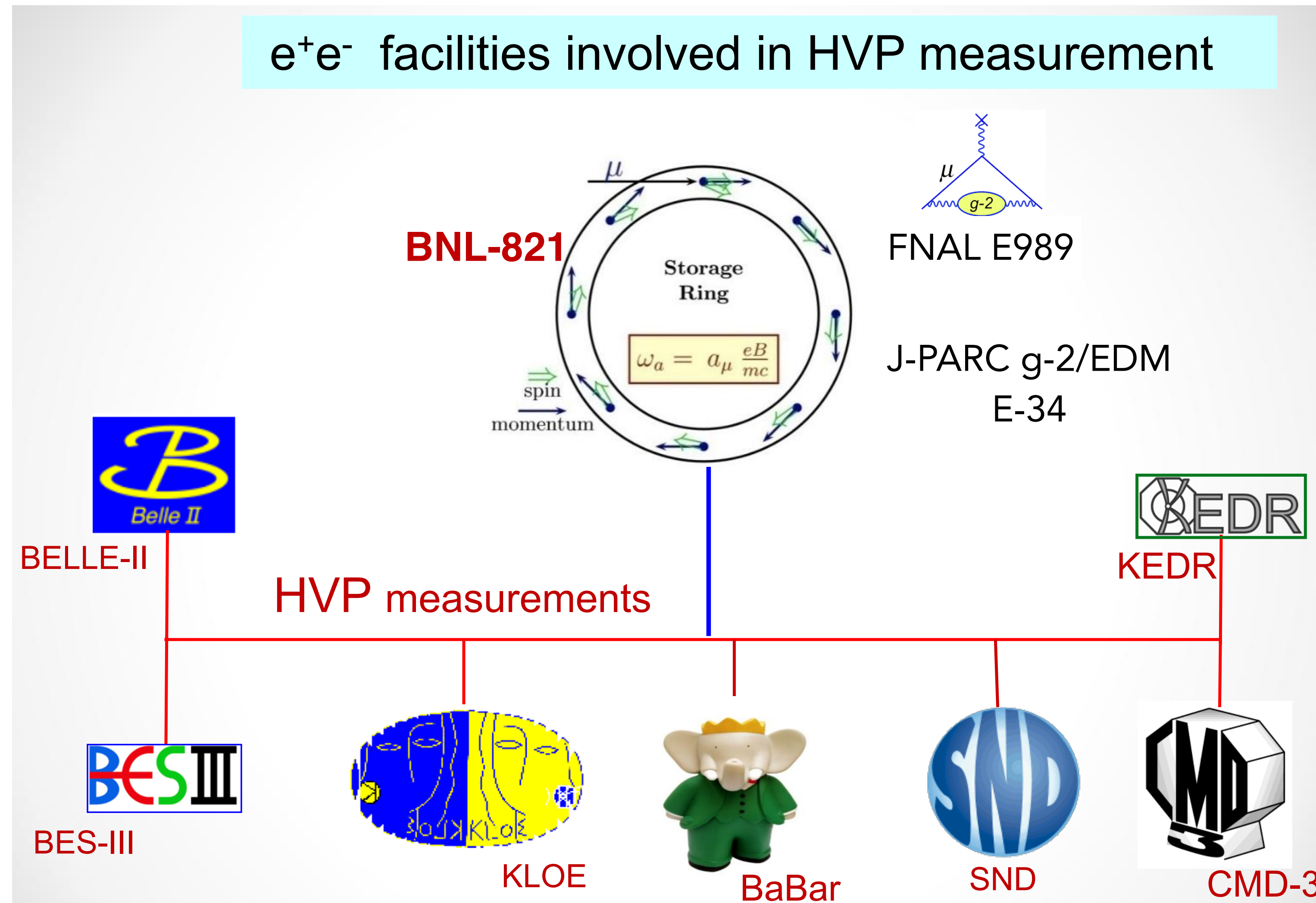
$$a_\mu^{\text{HVP,LO}} = \frac{m_\mu^2}{12\pi^3} \int ds \frac{\hat{K}(s)}{s} \sigma_{\text{exp}}(s)$$

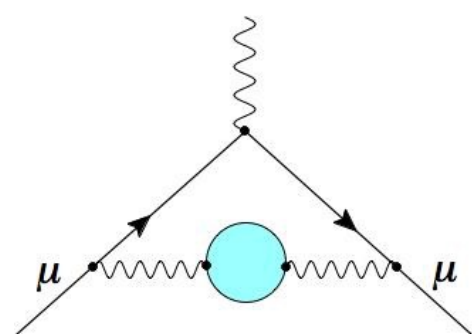
Dominant contributions from low energies  
 $\pi^+\pi^-$  channel: 73% of total

- Use direct integration method, summing up cross sections for all possible hadronic channels up to  $\sim 2$  GeV

# Experimental Inputs to HVP

S. Serednyakov (for SND) @ HVP KEK workshop





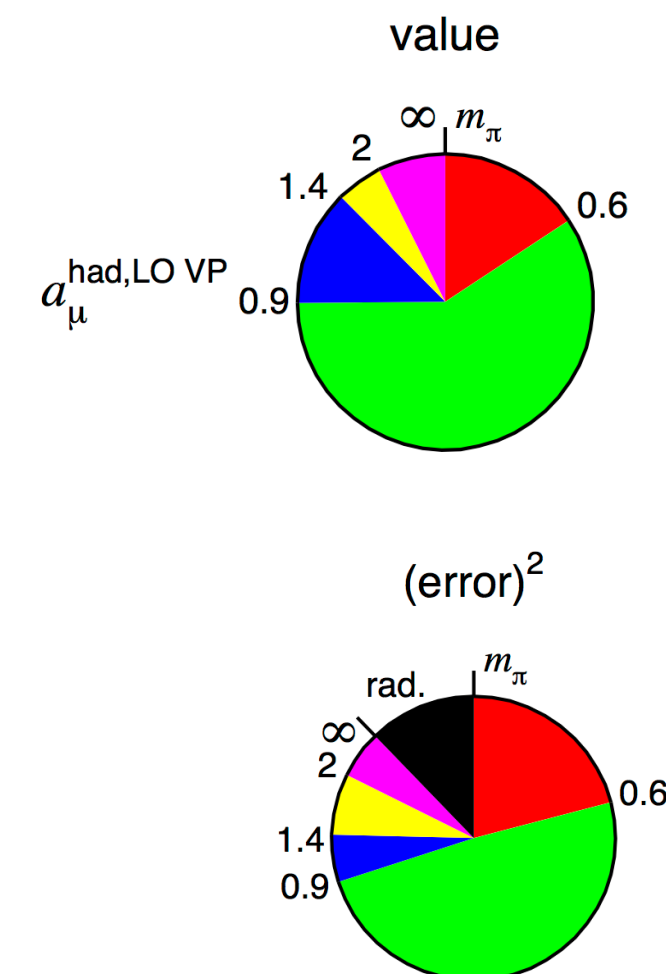
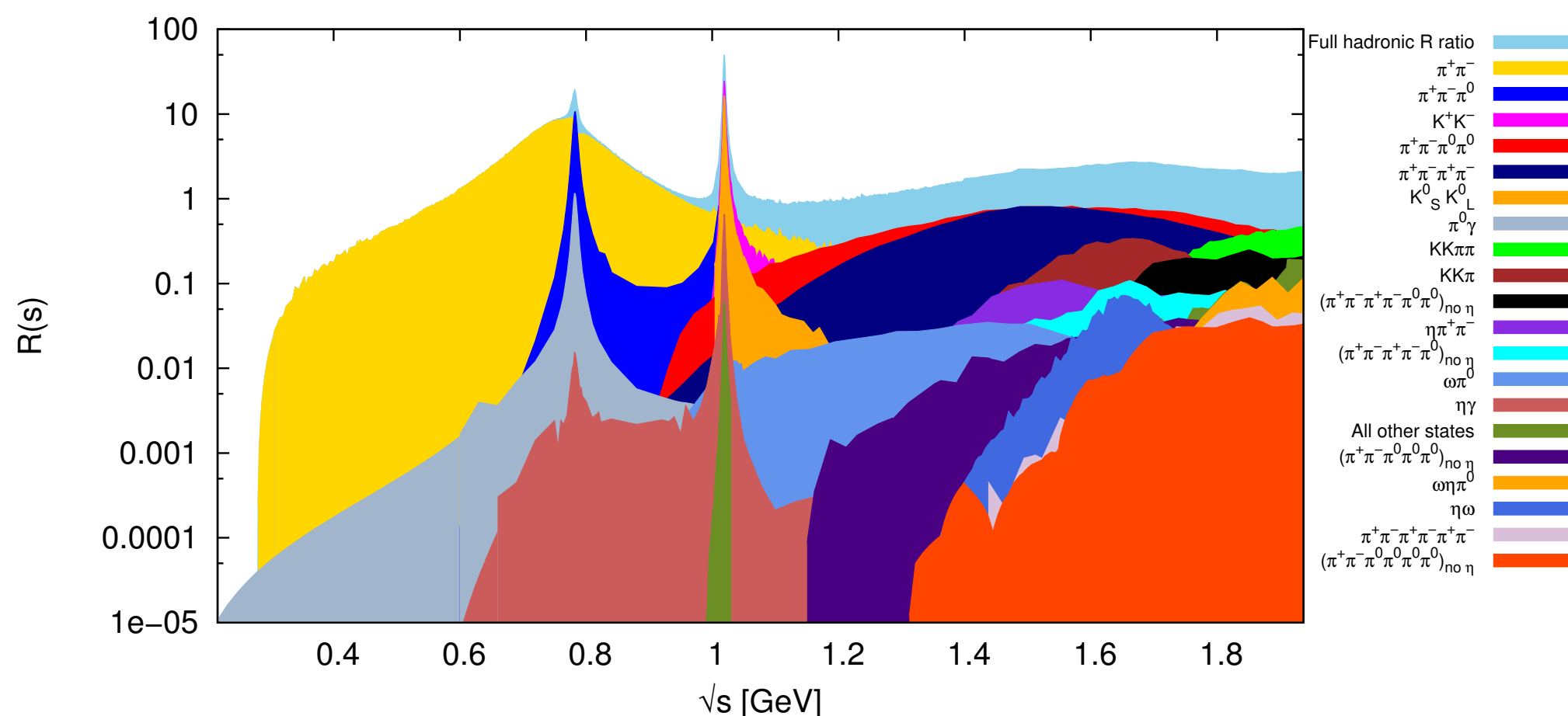
# HVP: data-driven

Z. Zhang for DHMZ @ INT g-2 workshop

[M. Davier et al, arXiv:1908.00921]

Channel	$a_\mu^{\text{had, LO}} [10^{-10}]$
$\pi^0\gamma$	$4.29 \pm 0.06 \pm 0.04 \pm 0.07$
$\eta\gamma$	$0.65 \pm 0.02 \pm 0.01 \pm 0.01$
$\pi^+\pi^-$	$507.80 \pm 0.83 \pm 3.19 \pm 0.60$
$\pi^+\pi^-\pi^0$	$46.20 \pm 0.40 \pm 1.10 \pm 0.86$
$2\pi^+2\pi^-$	$13.68 \pm 0.03 \pm 0.27 \pm 0.14$
$\pi^+\pi^-2\pi^0$	$18.03 \pm 0.06 \pm 0.48 \pm 0.26$
$2\pi^+2\pi^-\pi^0$ ( $\eta$ excl.)	$0.69 \pm 0.04 \pm 0.06 \pm 0.03$
$\pi^+\pi^-3\pi^0$ ( $\eta$ excl.)	$0.49 \pm 0.03 \pm 0.09 \pm 0.00$
$3\pi^+3\pi^-$	$0.11 \pm 0.00 \pm 0.01 \pm 0.00$
$2\pi^+2\pi^-2\pi^0$ ( $\eta$ excl.)	$0.71 \pm 0.06 \pm 0.07 \pm 0.14$
$\pi^+\pi^-4\pi^0$ ( $\eta$ excl., isospin)	$0.08 \pm 0.01 \pm 0.08 \pm 0.00$
$\eta\pi^+\pi^-$	$1.19 \pm 0.02 \pm 0.04 \pm 0.02$
$\eta\omega$	$0.35 \pm 0.01 \pm 0.02 \pm 0.01$
$\eta\pi^+\pi^-\pi^0$ (non- $\omega, \phi$ )	$0.34 \pm 0.03 \pm 0.03 \pm 0.04$
$\eta 2\pi^+2\pi^-$	$0.02 \pm 0.01 \pm 0.00 \pm 0.00$
$\omega\eta\pi^0$	$0.06 \pm 0.01 \pm 0.01 \pm 0.00$
$\omega\pi^0$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.94 \pm 0.01 \pm 0.03 \pm 0.00$
$\omega(\pi\pi)^0$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.07 \pm 0.00 \pm 0.00 \pm 0.00$
$\omega$ (non- $3\pi, \pi\gamma, \eta\gamma$ )	$0.04 \pm 0.00 \pm 0.00 \pm 0.00$
$K^+K^-$	$23.08 \pm 0.20 \pm 0.33 \pm 0.21$
$K_S K_L$	$12.82 \pm 0.06 \pm 0.18 \pm 0.15$
$\phi$ (non- $K\bar{K}, 3\pi, \pi\gamma, \eta\gamma$ )	$0.05 \pm 0.00 \pm 0.00 \pm 0.00$
$K\bar{K}\pi$	$2.45 \pm 0.05 \pm 0.10 \pm 0.06$
$K\bar{K}2\pi$	$0.85 \pm 0.02 \pm 0.05 \pm 0.01$
$K\bar{K}3\pi$ (estimate)	$-0.02 \pm 0.01 \pm 0.01 \pm 0.00$
$\eta\phi$	$0.33 \pm 0.01 \pm 0.01 \pm 0.00$
$\eta K\bar{K}$ (non- $\phi$ )	$0.01 \pm 0.01 \pm 0.01 \pm 0.00$
$\omega K\bar{K}$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.01 \pm 0.00 \pm 0.00 \pm 0.00$
$\omega 3\pi$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.06 \pm 0.01 \pm 0.01 \pm 0.01$
$7\pi$ ( $3\pi^+3\pi^-\pi^0$ + estimate)	$0.02 \pm 0.00 \pm 0.01 \pm 0.00$
<hr/>	
$J/\psi$ (BW integral)	$6.28 \pm 0.07$
$\psi(2S)$ (BW integral)	$1.57 \pm 0.03$
<hr/>	
$R$ data [3.7 – 5.0] GeV	$7.29 \pm 0.05 \pm 0.30 \pm 0.00$
$R_{\text{QCD}} [1.8 - 3.7 \text{ GeV}]_{uds}$	$33.45 \pm 0.28 \pm 0.65_{\text{dual}}$
$R_{\text{QCD}} [5.0 - 9.3 \text{ GeV}]_{udsc}$	$6.86 \pm 0.04$
$R_{\text{QCD}} [9.3 - 12.0 \text{ GeV}]_{udscb}$	$1.21 \pm 0.01$
$R_{\text{QCD}} [12.0 - 40.0 \text{ GeV}]_{udscb}$	$1.64 \pm 0.00$
$R_{\text{QCD}} [> 40.0 \text{ GeV}]_{udscb}$	$0.16 \pm 0.00$
$R_{\text{QCD}} [> 40.0 \text{ GeV}]_t$	$0.00 \pm 0.00$
<b>Sum</b>	$693.9 \pm 1.0 \pm 3.4 \pm 1.6 \pm 0.1_\psi \pm 0.7_{\text{QCD}}$

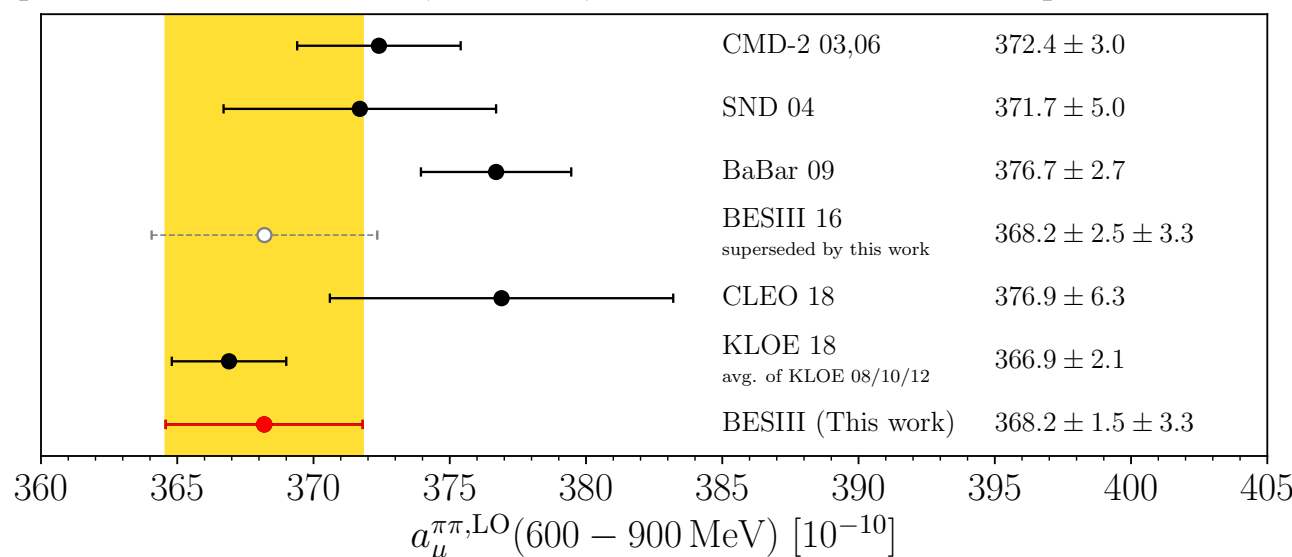
[A. Keshavarzi et al, arXiv:1802.02995]



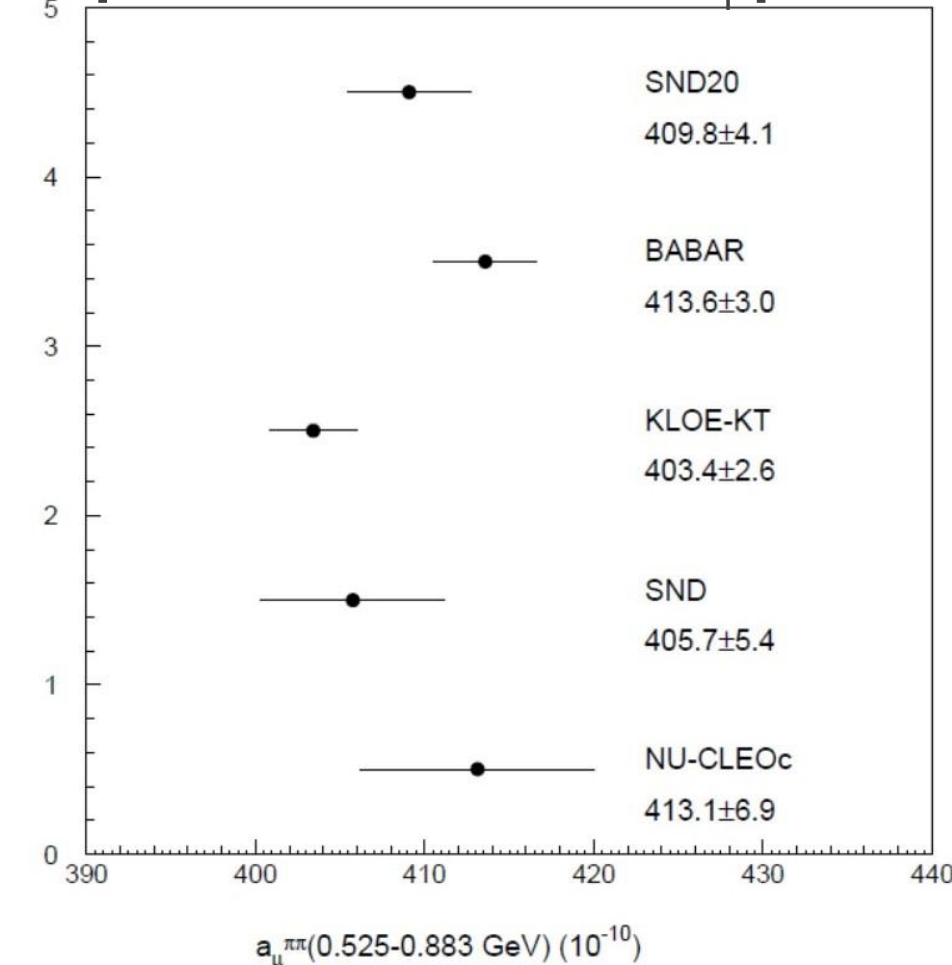
## Tensions between BaBar and KLOE data sets:

- Cross checks using analyticity and unitarity relating pion form factor to  $\pi\pi$  scattering
- Combinations of data sets affected by tensions
- conservative merging procedure

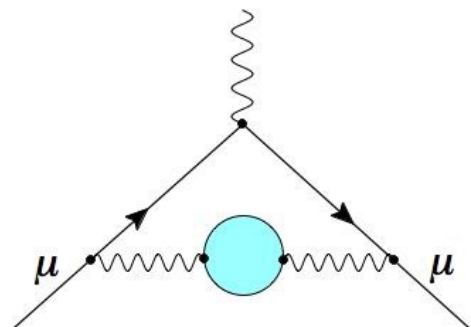
[M. Ablikim et al (BES III), arXiv:2009.05011]



[M. Davier @ KEK workshop]

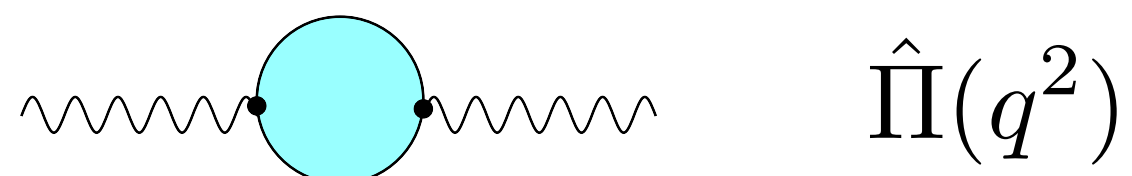






# Lattice HVP: Introduction

[B. Lautrup, A. Peterman, E. de Rafael, Phys. Rep 1972;  
E. de Rafael, Phys. Let. B 1994; T. Blum, PRL 2002]



$$\hat{\Pi}(q^2)$$

Leading order HVP correction:

$$a_{\mu}^{\text{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int dq^2 \omega(q^2) \hat{\Pi}(q^2)$$

- Calculate  $a_{\mu}^{\text{HVP,LO}}$  in Lattice QCD

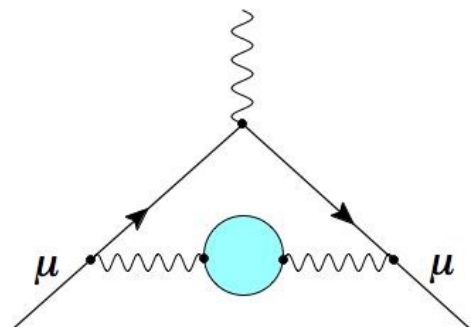
Compute correlation function: 
$$C(t) = \frac{1}{3} \sum_{i,x} \langle j_i(x,t) j_i(0,0) \rangle$$

and 
$$\hat{\Pi}(Q^2) = 4\pi^2 \int_0^{\infty} dt C(t) \left[ t^2 - \frac{4}{Q^2} \sin^2 \left( \frac{Qt}{2} \right) \right]$$

[D. Bernecker and H. Meyer, arXiv:1107.4388,  
EPJA 2011]

Obtain  $a_{\mu}^{\text{HVP,LO}}$  from an integral over Euclidean time:

$$a_{\mu}^{\text{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^{\infty} dt \tilde{w}(t) C(t)$$



# Lattice HVP: Introduction

Calculate  $a_\mu^{\text{HVP}}$  in Lattice QCD:

$$a_\mu^{\text{HVP,LO}} = \sum_f a_{\mu,f}^{\text{HVP,LO}} + a_{\mu,\text{disc}}^{\text{HVP,LO}}$$

- Separate into connected for each quark flavor + disconnected contributions (gluon and sea-quark background not shown in diagrams)

Note: almost always  $m_u = m_d$

$$\sum_f \left[ \text{quark loop with photon} \right] + \left[ \text{quark self-energy} \right] + \left[ \text{ghost loop} \right] \quad f = ud, s, c, b$$

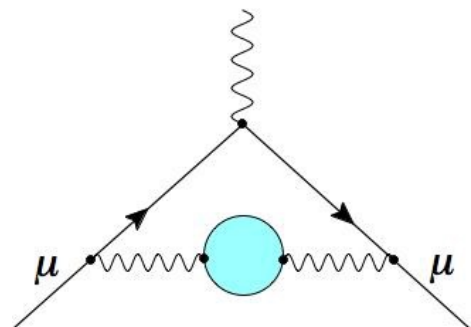
The diagram shows a sum over quark flavors  $f$ . The first term is a quark loop with a photon line attached to the top vertex. The second term is a quark self-energy loop. The third term is a ghost loop. The quark lines are labeled with  $f$  and  $\bar{f}$ .

- need to add QED and strong isospin breaking ( $\sim m_u - m_d$ ) corrections:

$$\left[ \text{quark loop with gluon} \right] + \dots$$

The diagram shows a quark loop with a gluon line attached to the top vertex, followed by an ellipsis indicating further terms.

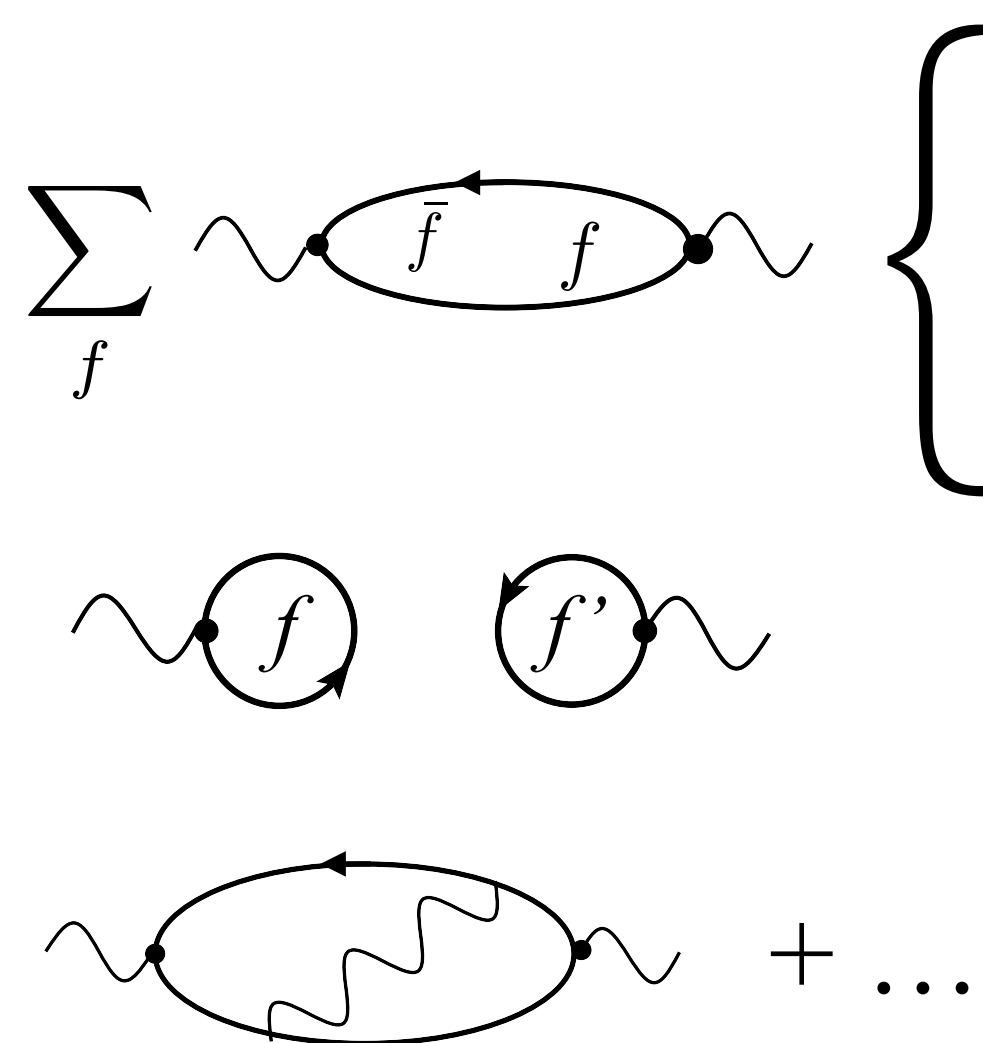
- either perturbatively on isospin symmetric QCD background
- or by using QCD + QED ensembles with  $m_u \neq m_d$



# Lattice HVP: Introduction

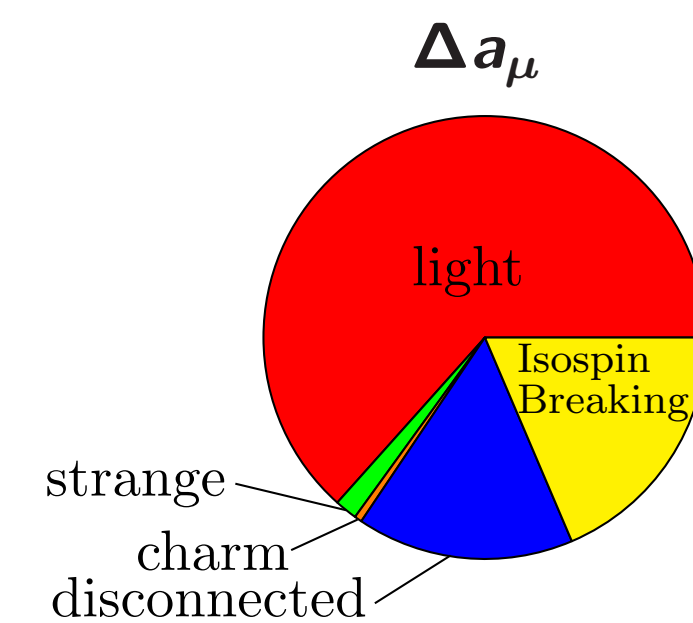
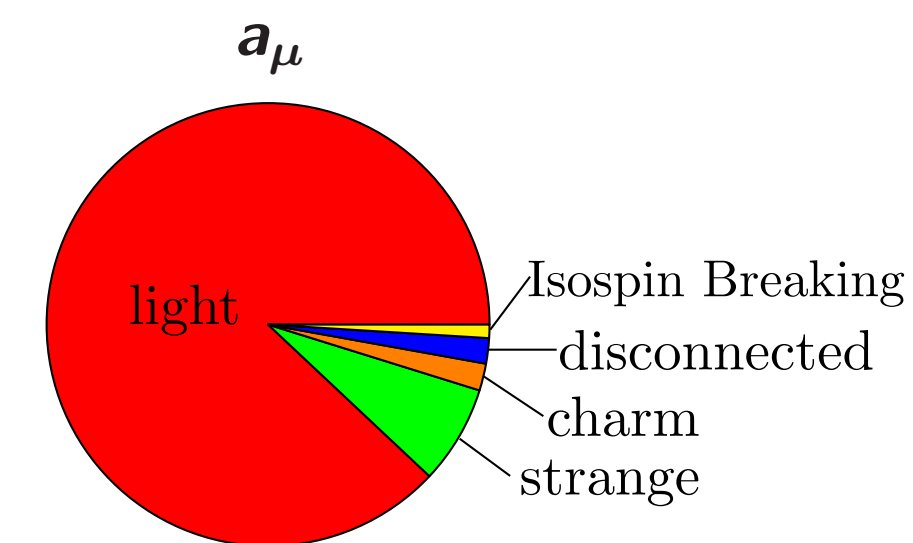
V. Gülpers @ Lattice HVP workshop

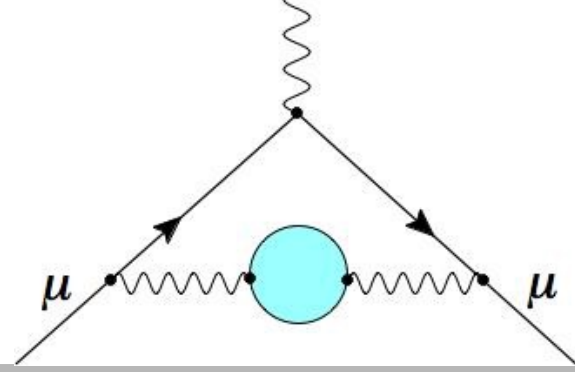
Target: ~ 0.2% total error



- light-quark connected contribution:  
 $a_\mu^{\text{HVP,LO}}(ud) \sim 90\%$  of total
- s,c,b-quark contributions  
 $a_\mu^{\text{HVP,LO}}(s, c, b) \sim 8\%, 2\%, 0.05\%$  of total
- disconnected contribution:  
 $a_{\mu, \text{disc}}^{\text{HVP,LO}} \sim 2\%$  of total
- Isospinbreaking (QED +  $m_u \neq m_d$ ) corrections:  
 $\delta a_\mu^{\text{HVP,LO}} \sim 1\%$  of total

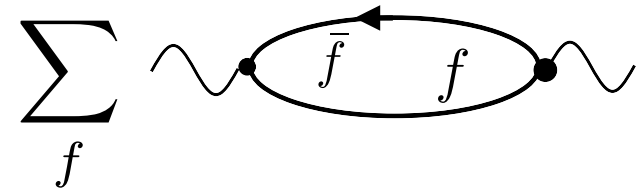
$$a_\mu^{\text{HVP,LO}} = a_\mu^{\text{HVP,LO}}(ud) + a_\mu^{\text{HVP,LO}}(s) + a_\mu^{\text{HVP,LO}}(c) + a_{\mu, \text{disc}}^{\text{HVP,LO}} + \delta a_\mu^{\text{HVP,LO}}$$



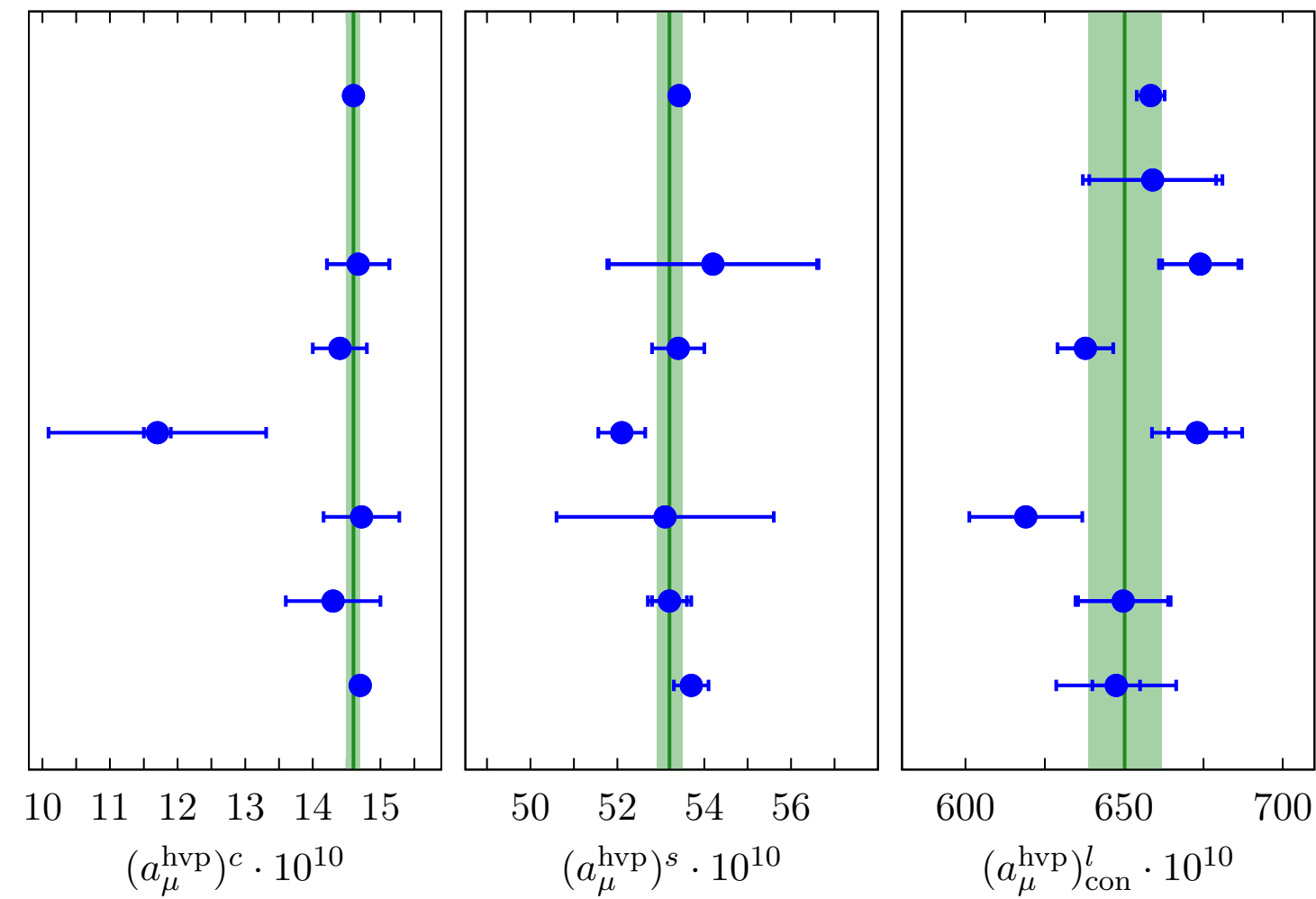


# Lattice HVP: Isospin corrections

H. Wittig @ Lattice HVP workshop

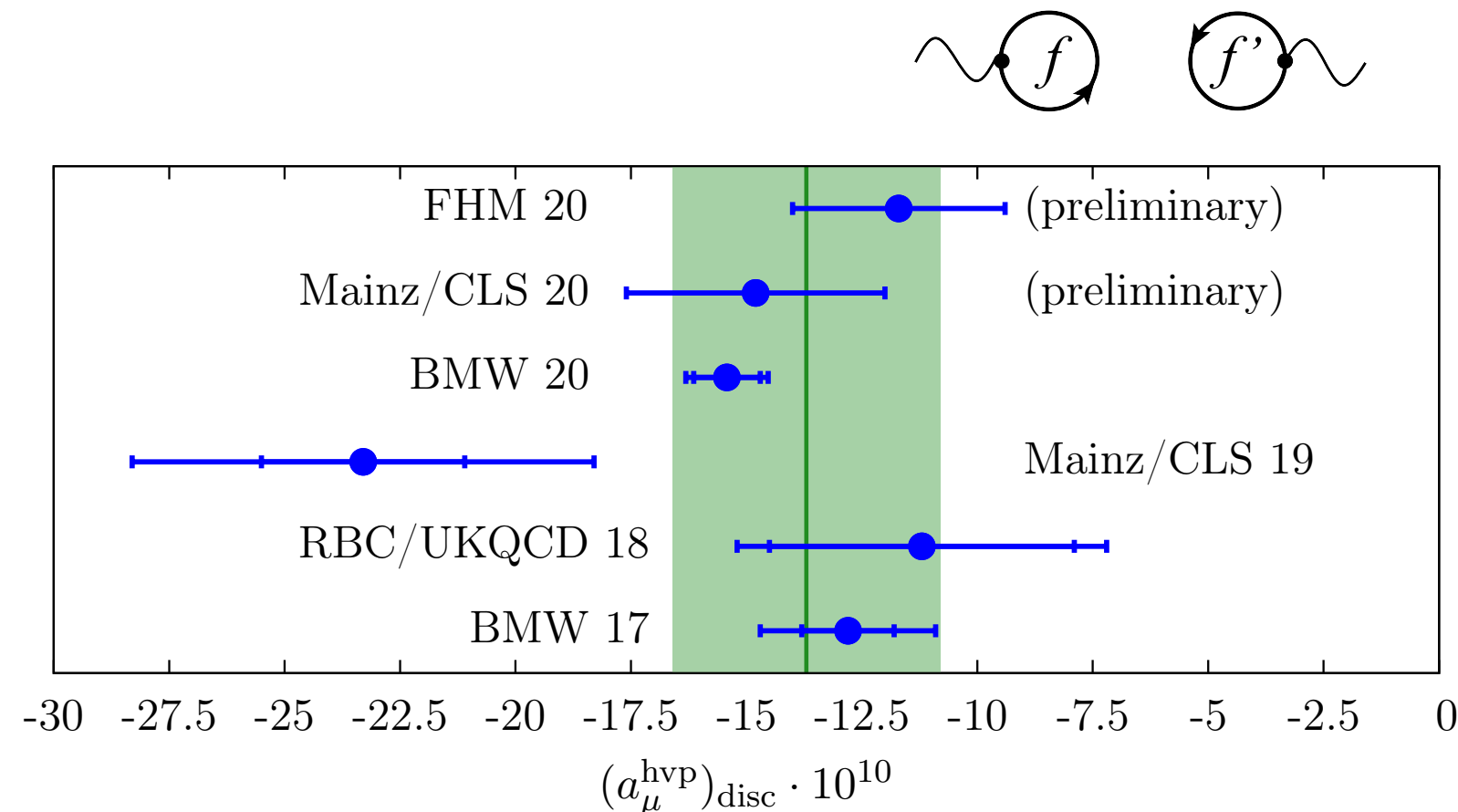


- Charm, strange contributions already well determined.
- Mild tensions for light contribution



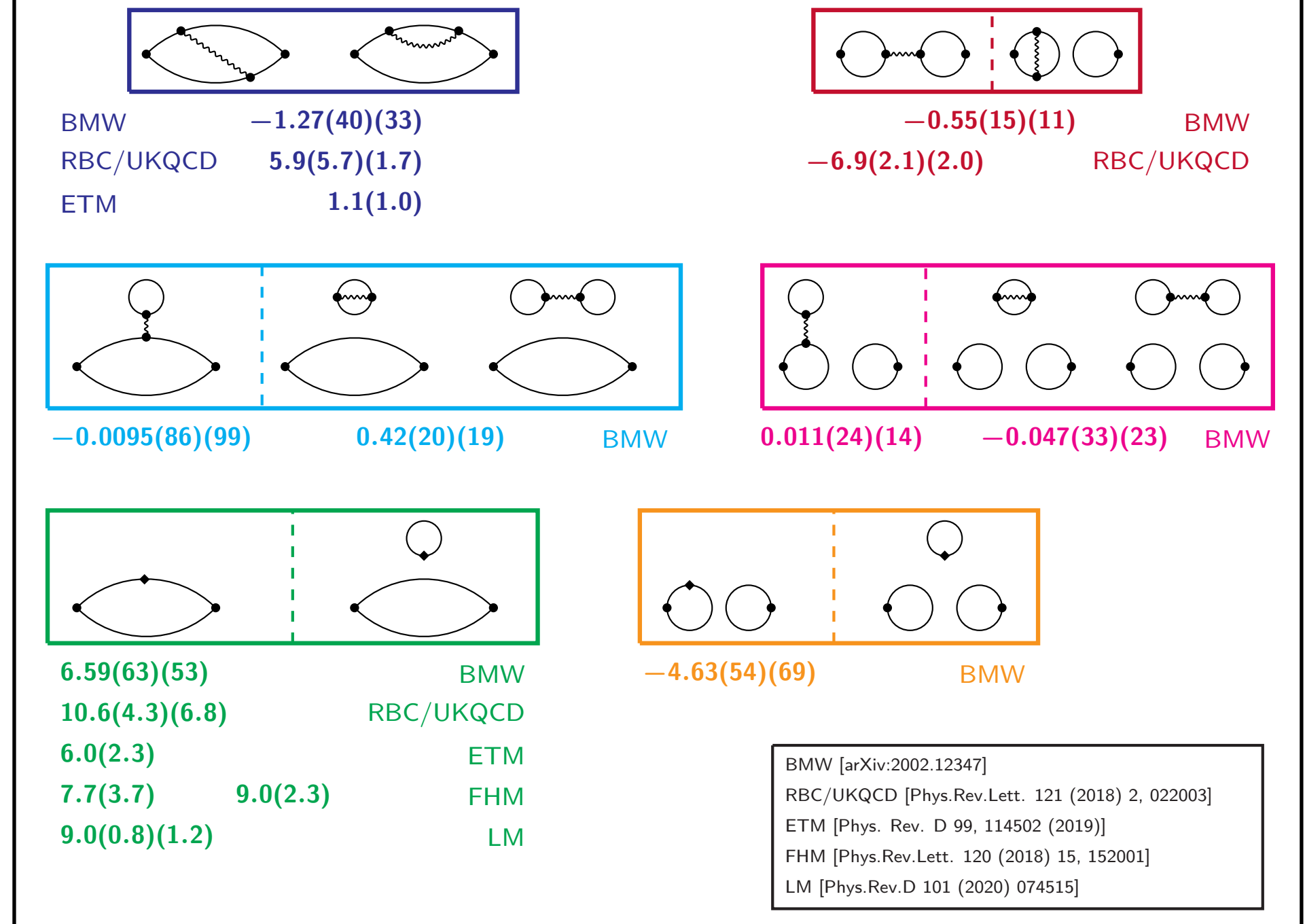
Ongoing efforts by  
FNAL-HPQCD-MILC  
RBC/UKQCD, Mainz

Consistent results with  
increasing precision

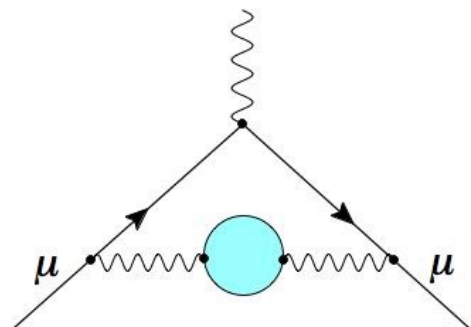


V. Gülpers @ Lattice HVP workshop

Overview of published results - contributions to  $a_\mu \times 10^{10}$



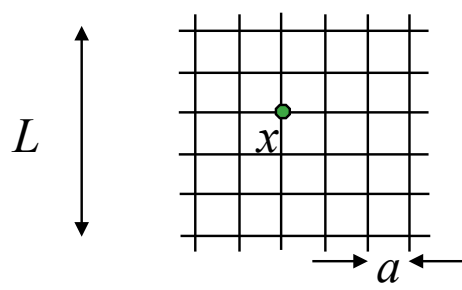
- Some tensions between lattice results for individual contributions.
- Large cancellations between individual contributions:  
 $\delta a_\mu^{\text{IB}} \lesssim 1\%$



# Lattice HVP: Introduction

- Target:  $\sim 0.2\%$  total error
- Challenges:
  - ✓ needs ensembles with (light sea) quark masses at their physical values
  - ✓ finite volume corrections
  - growth of statistical errors at long-distances
  - Continuum extrapolation
  - scale setting
  - disconnected contribution
  - QED and strong isospin breaking corrections ( $m_u \neq m_d$ )
- Focus on windows in Euclidean times [T. Blum et al, arXiv:1801.07224, 2018 PRL]
  - disentangle systematics/statistics from long distance/FV and discretization effects
    - ▮▮▮▮▮ valuable cross checks
  - intermediate window easy to compute & compare with dispersive methods

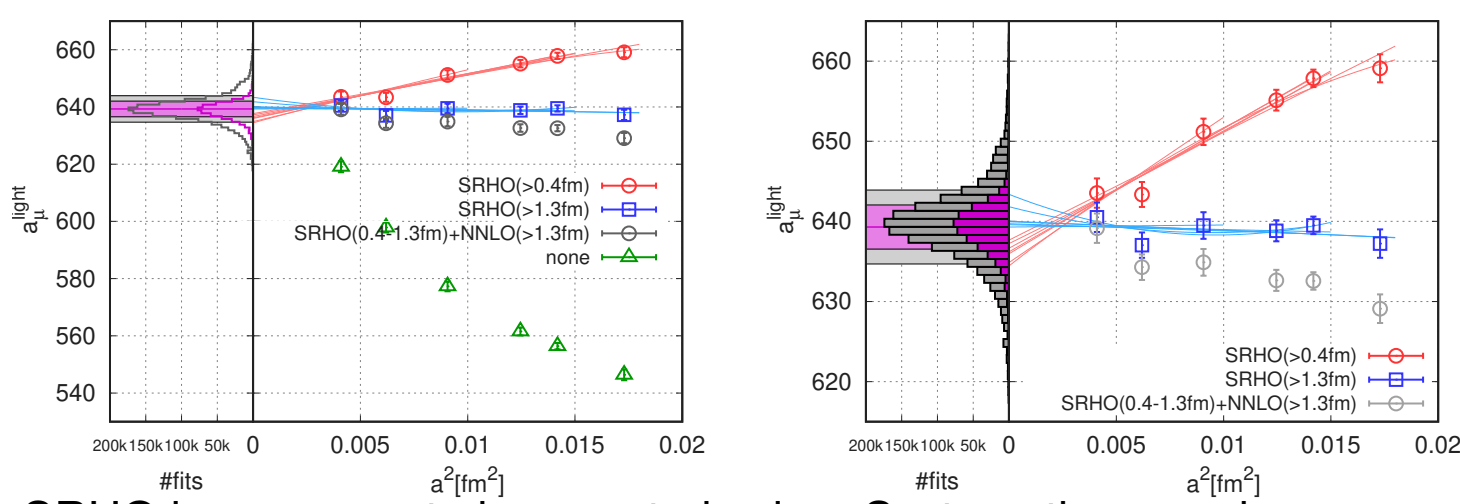
# Continuum extrapolation



Kalman Szabo (BMWc) @ Lattice 2021

## Taste improvement II

- $a_\mu(a) \rightarrow a_\mu(a) - a_\mu^{\text{SRHO}}(a) + a_\mu^{\text{RHO}}$
- reduces lattice artefact, also makes  $a^2$  dependence linear



SRHO improvement gives central value. Systematic errors by:

- change starting point of improvement  $t = 0.4 \rightarrow 1.3$  fm
- skip coarse lattices
- change  $\Gamma = 0$  and  $\Gamma = 3$
- replace SRHO by NNLO SXPT above 1.3 fm

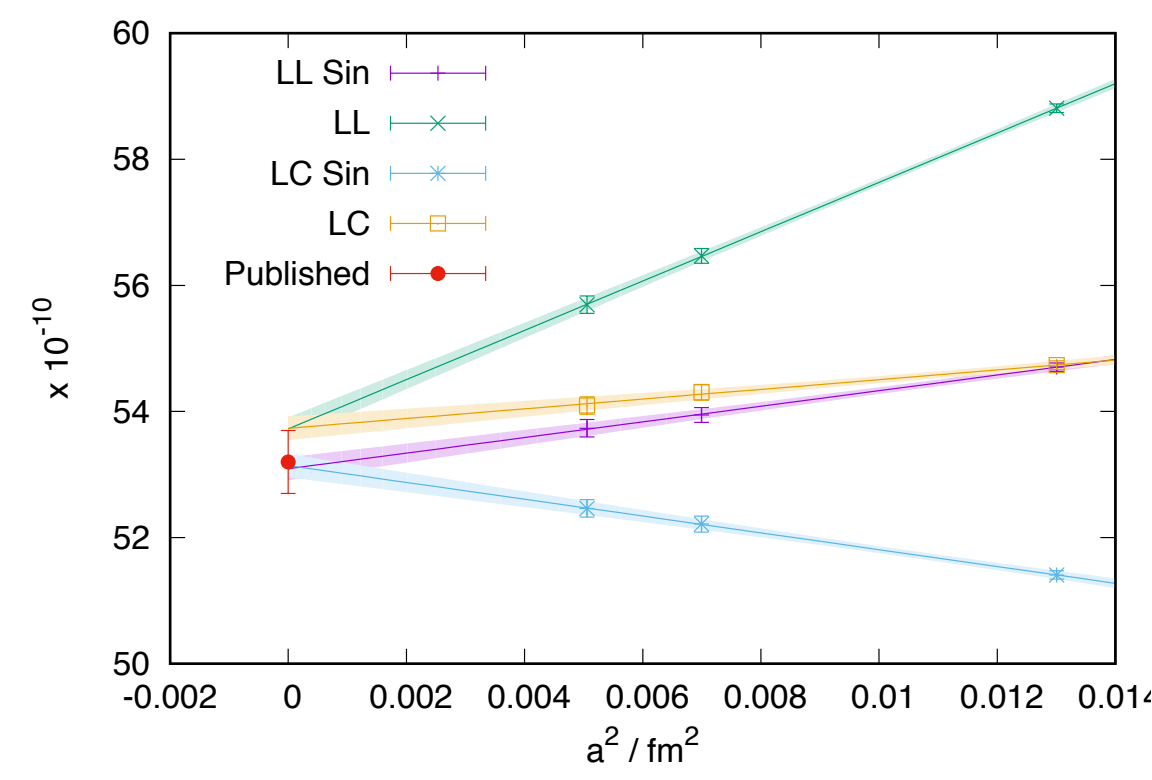
- Large taste-breaking effects with BMW set-up
  - uncorrected data not easily fit to power series, i.e.

$$1 \quad A_0 + A_1 [a^2] + A_2 [a^2]^2$$

$$2 \quad A_0 + A_1 [a^2 \alpha_s^3(\frac{1}{a})] + A_2 [a^2 \alpha_s^3(\frac{1}{a})]^2$$

Christoph Lehner (RBC/UKQCD) @ Lattice 2021

- Third lattice spacing for strange data ( $a^{-1} = 2.77$  GeV with  $m_\pi = 234$  MeV with sea light-quark mass corrected from global fit):

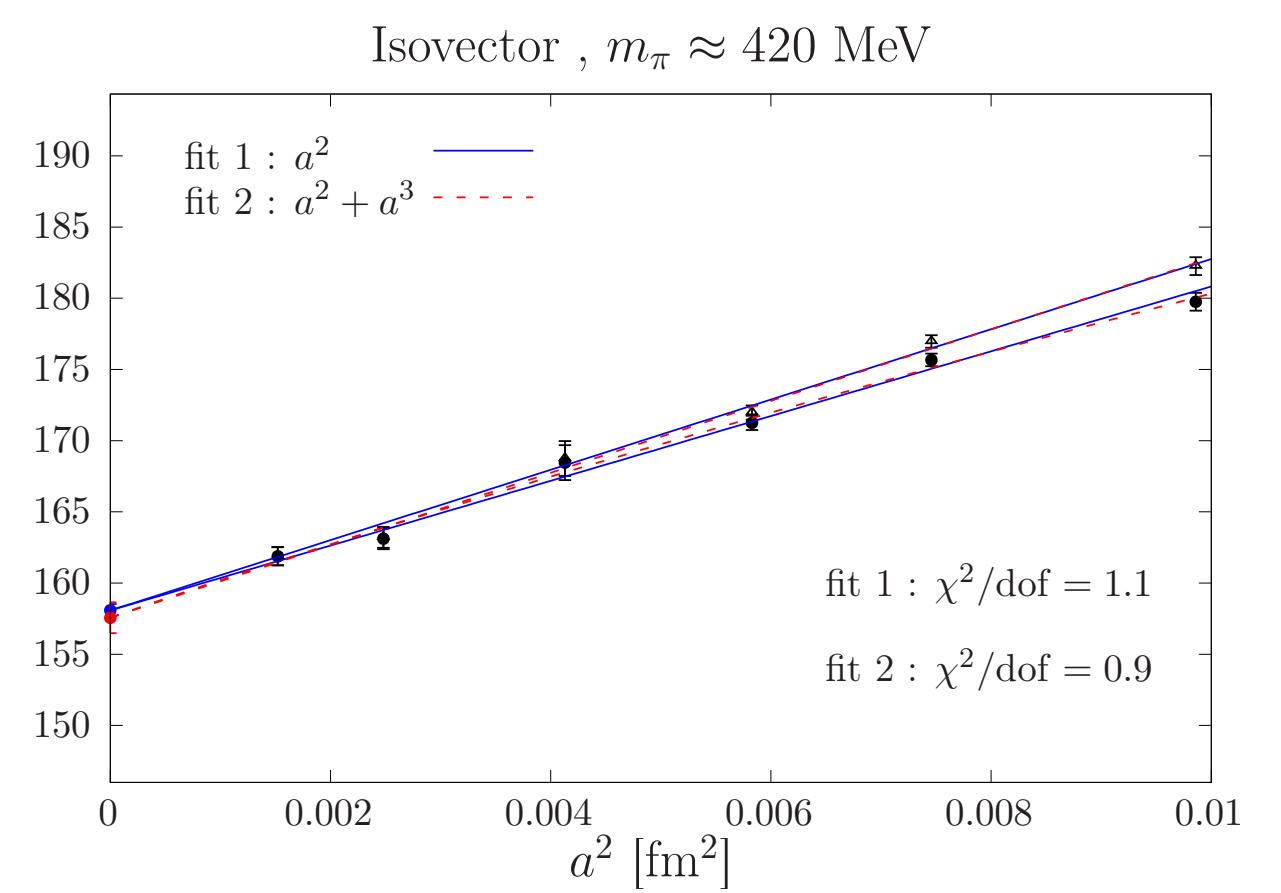


- For light quark use new 96l ensemble at physical pion mass. Data still being generated on Summit in USA and Booster in Germany ( $a^{-1} = 2.77$  GeV with  $m_\pi = 139$  MeV)

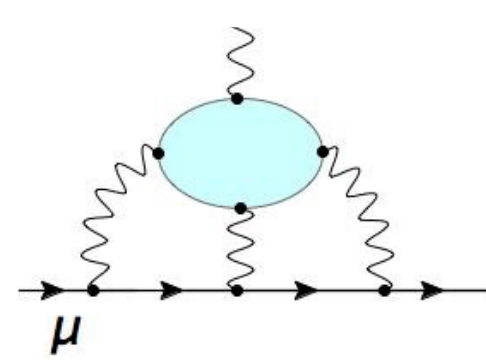
- RBC: Currently adding add a third lattice spacing

- Fermilab-HPQCD-MILC: planning to add a 5th lattice spacing (0.042 fm).

Hartmut Wittig (Mainz) @ Lattice 2021



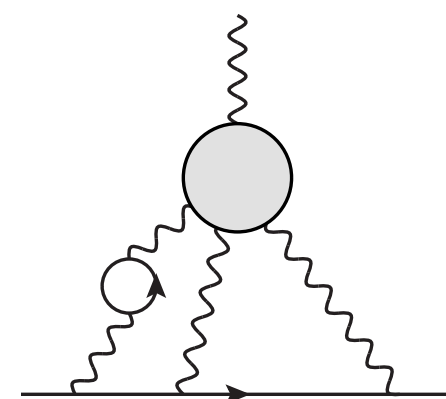
- Mainz and ETMc perform combined chair and continuum extrapolation



# HLbL: dispersive

## Comparison:

Contribution	PdRV(09) [471]	N/JN(09) [472, 573]	J(17) [27]	Our estimate
$\pi^0, \eta, \eta'$ -poles	114(13)	99(16)	95.45(12.40)	93.8(4.0)
$\pi, K$ -loops/boxes	-19(19)	-19(13)	-20(5)	-16.4(2)
$S$ -wave $\pi\pi$ rescattering	-7(7)	-7(2)	-5.98(1.20)	-8(1)
subtotal	88(24)	73(21)	69.5(13.4)	69.4(4.1)
scalars	-	-	-	} -1(3)
tensors	-	-	1.1(1)	
axial vectors	15(10)	22(5)	7.55(2.71)	6(6)
$u, d, s$ -loops / short-distance	-	21(3)	20(4)	15(10)
$c$ -loop	2.3	-	2.3(2)	3(1)
total	105(26)	116(39)	100.4(28.2)	92(19)



NLO HLbL contribution:

$$a_{\mu}^{\text{HLbL,NLO}} = 2(1) \times 10^{-11}$$

# Connections

$$\sigma(e^+e^- \rightarrow \text{hadrons}) \Leftrightarrow a_\mu^{\text{HVP}} \Leftrightarrow \Delta\alpha_{\text{had}}(M_Z^2)$$

- $\Delta\alpha_{\text{had}}(M_Z^2)$  also depends on the hadronic vacuum polarization function, and can be written as an integral over  $\sigma(e^+e^- \rightarrow \text{hadrons})$ , but weighted towards higher energies.
- a shift in  $a_\mu^{\text{HVP}}$  also changes  $\Delta\alpha_{\text{had}}(M_Z^2)$ :  $\Rightarrow$  EW fits  
[Crivellin et al 2020, Keshavarsi et al 2020, Malaescu & Scott 2020]  
If the shift is due to differences in the low ( $\lesssim 2 \text{ GeV}$ ) energy region, the impact on  $\Delta\alpha_{\text{had}}(M_Z^2)$  and EW fits is small.
- A shift in  $a_\mu^{\text{HVP}}$  from low ( $\lesssim 2 \text{ GeV}$ ) energies  $\Rightarrow \sigma(e^+e^- \rightarrow \pi\pi)$   
must satisfy unitarity & analyticity constraints  $\Rightarrow F_\pi^V(s)$   
can be tested with lattice calculations  
[Colangelo, Hoferichter, Stoffer 2021]



# Connections

Martin Hoferichter @ Lattice HVP workshop

## Hadronic running of $\alpha$ and global EW fit

	$e^+e^-$ KNT, DHMZ	EW fit HEPFit	EW fit GFitter	guess based on BMWc
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) \times 10^4$	276.1(1.1)	270.2(3.0)	271.6(3.9)	277.8(1.3)
difference to $e^+e^-$		$-1.8\sigma$	$-1.1\sigma$	$+1.0\sigma$

- Time-like formulation:

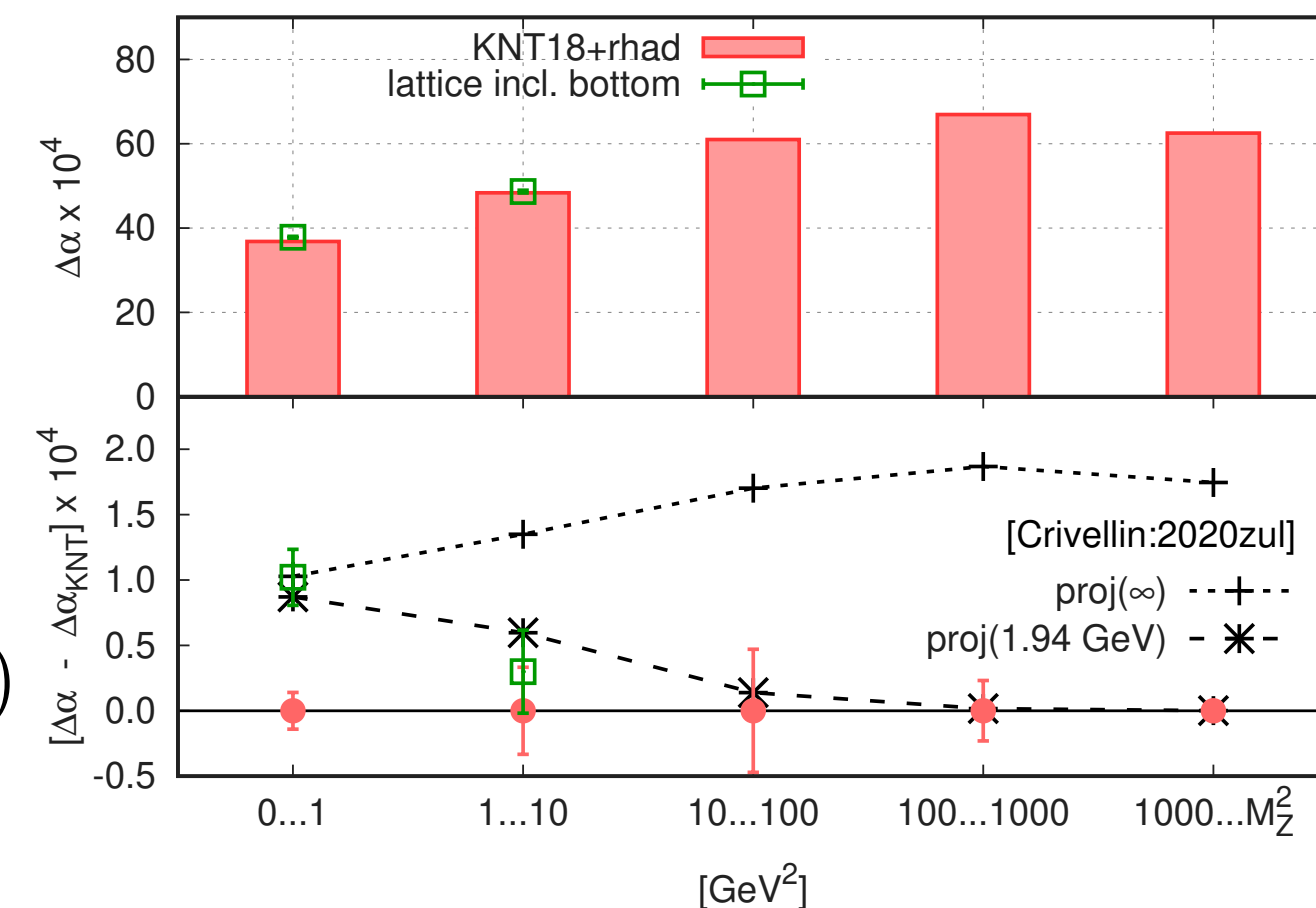
$$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) = \frac{\alpha M_Z^2}{3\pi} P \int_{s_{\text{thr}}}^{\infty} ds \frac{R_{\text{had}}(s)}{s(M_Z^2 - s)}$$

- Space-like formulation:

$$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) = \frac{\alpha}{\pi} \hat{\Pi}(-M_Z^2) + \frac{\alpha}{\pi} (\hat{\Pi}(M_Z^2) - \hat{\Pi}(-M_Z^2))$$

- Global EW fit

- Difference between HEPFit and GFitter implementation mainly treatment of  $M_W$
- Pull goes into **opposite direction**



BMWc 2020

More in talks by M. Passera, B. Malaescu (phenomenology) and K. Miura, T. San José (lattice)



# Connections

Peter Stoffer @ Lattice HVP workshop

Constraints on the two-pion contribution to HVP

arXiv:2010.07943 [hep-ph]

Modifying  $a_{\mu}^{\pi\pi} |_{\leq 1 \text{ GeV}}$

- “low-energy” scenario: local changes in cross section of  $\sim 8\%$  around  $\rho$
- “high-energy” scenario: impact on **pion charge radius** and space-like VFF  $\Rightarrow$  chance for **independent lattice-QCD checks**

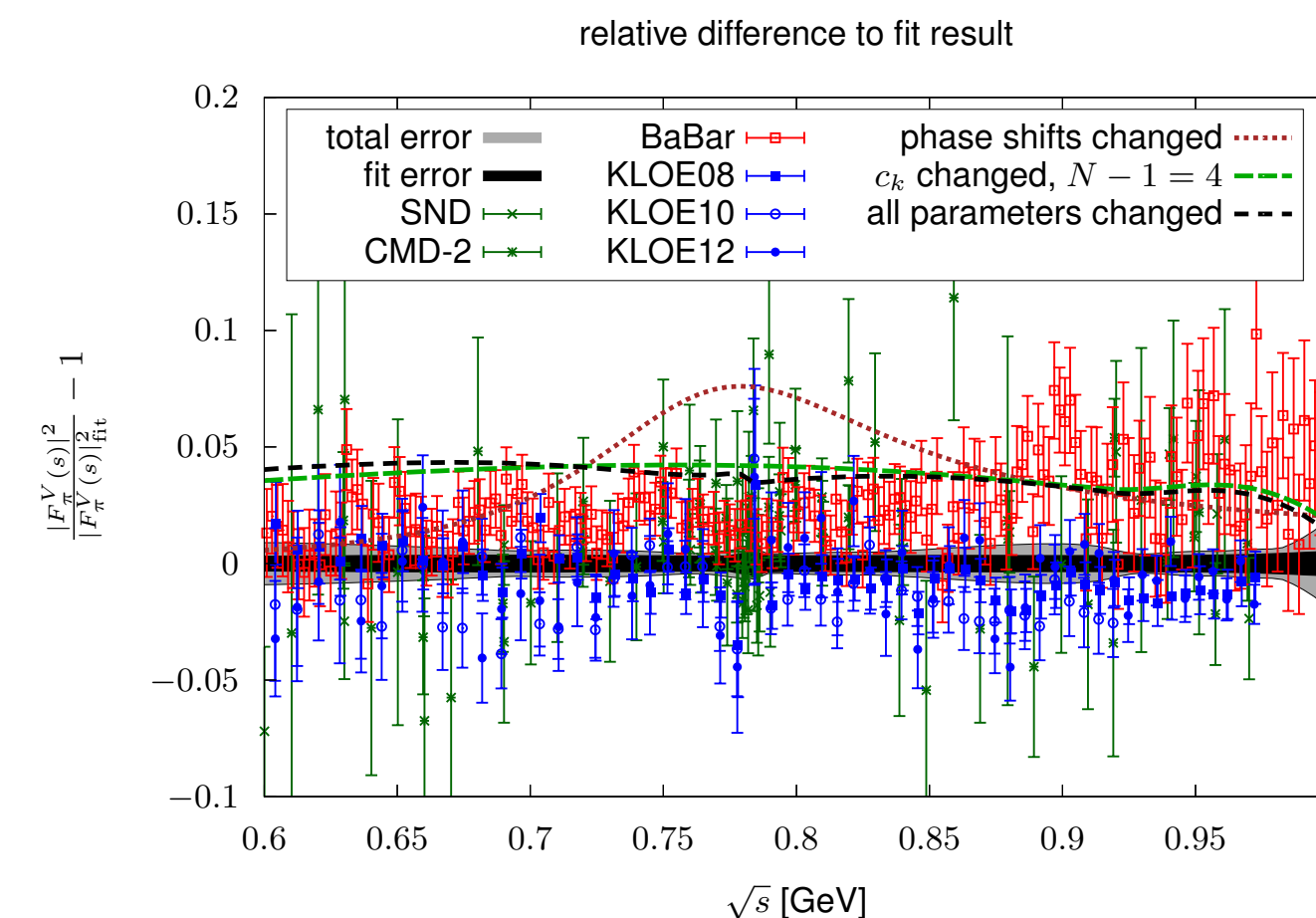
- requires **factor  $\sim 3$**

**improvement** over

$\chi$ QCD result:

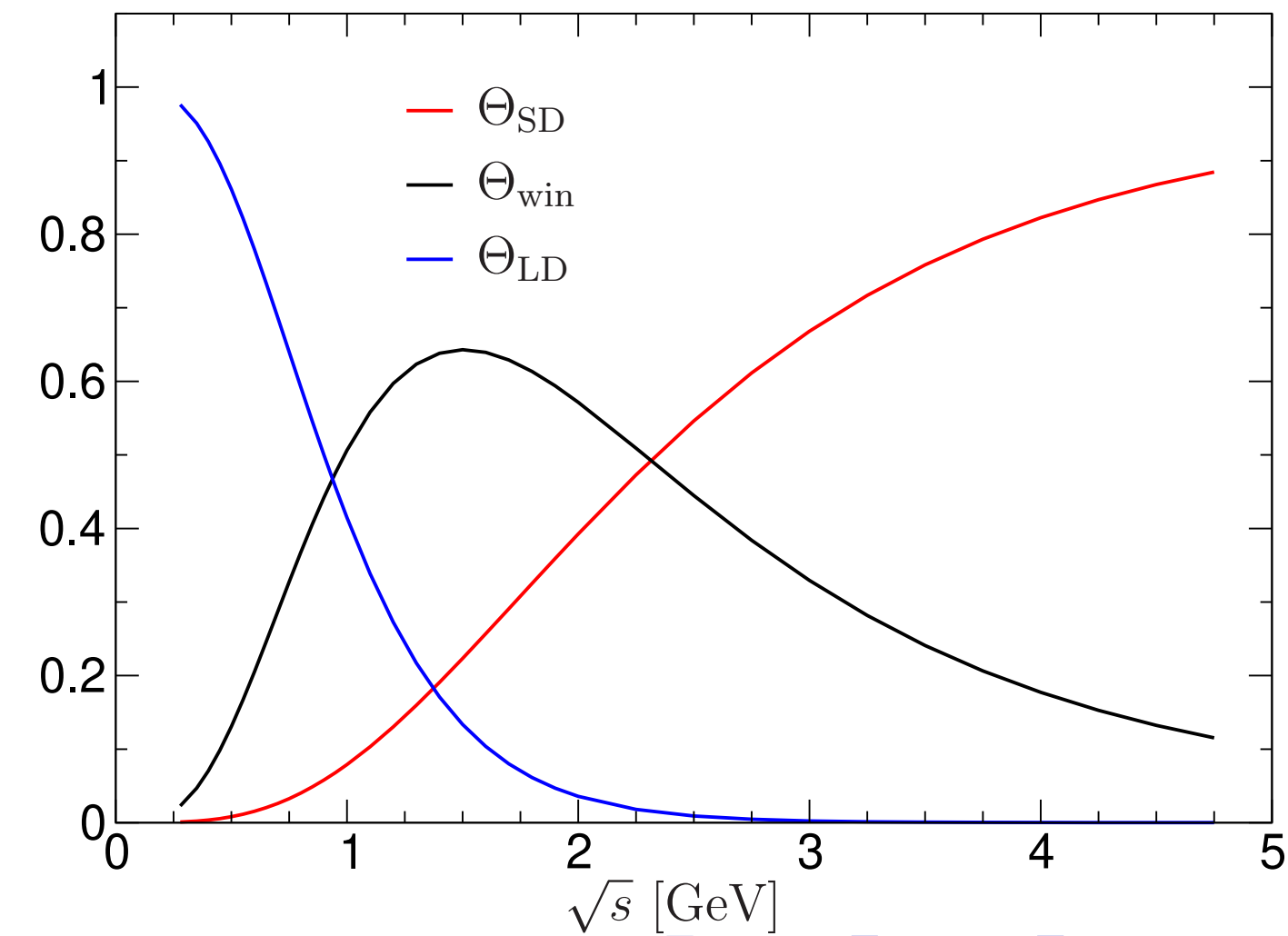
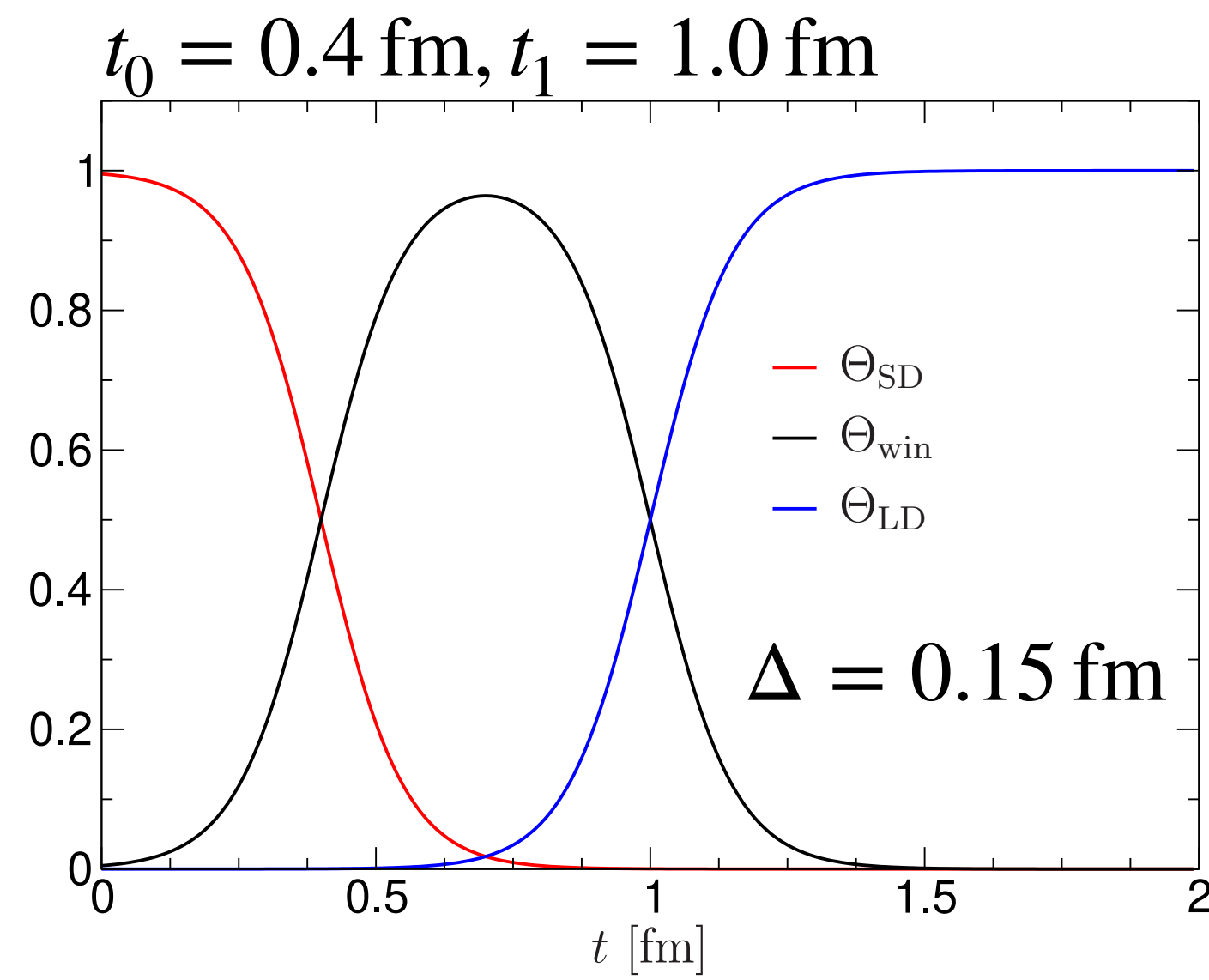
$$\langle r_{\pi}^2 \rangle = 0.433(9)(13) \text{ fm}^2$$

$\rightarrow$  arXiv:2006.05431 [hep-ph]



# Windows: Euclidean time vs $\sqrt{s}$

Martin Hoferichter @ Lattice HVP workshop



$[t_0, t_1]$ intermediate window	percentage captured of $\pi\pi$ channel $\leq 1 \text{ GeV}$		
	SD	intermediate	LD
$[0.4, 1.0] \text{ fm}$	3	28	69
$[1.0, 2.0] \text{ fm}$	31	51	18
$[1.0, 2.5] \text{ fm}$	31	61	9
$[1.0, 3.0] \text{ fm}$	31	65	4

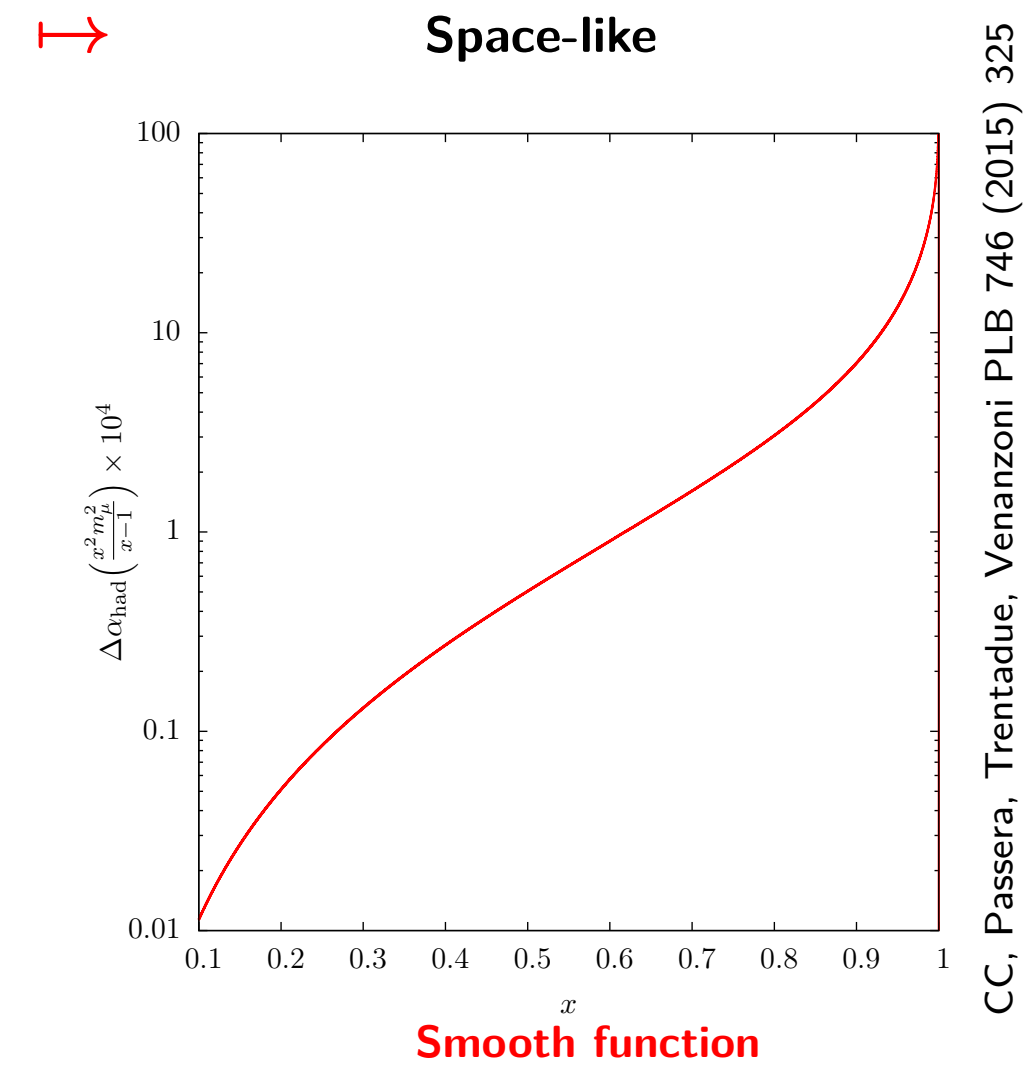
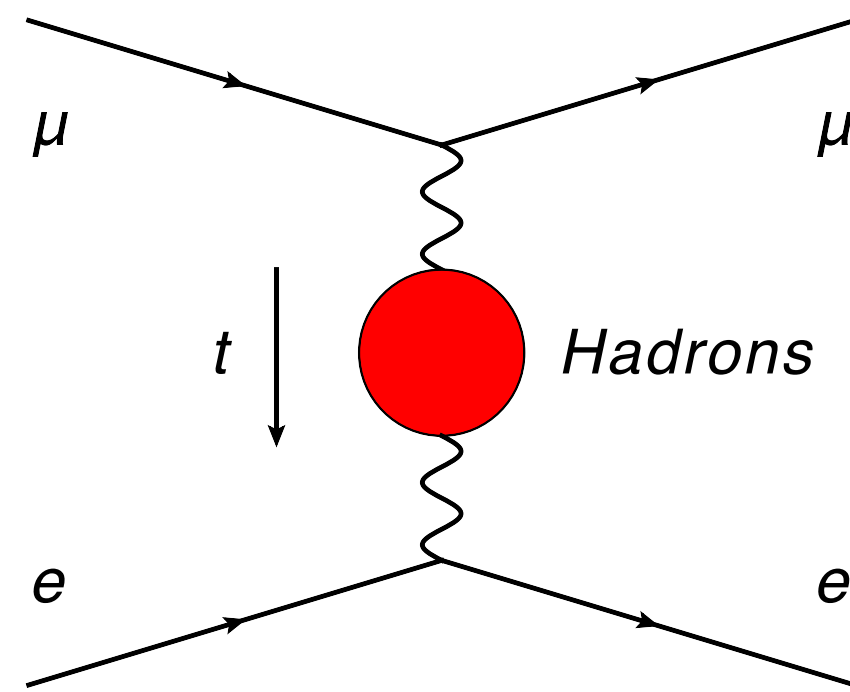
SD:  $[0, t_0]$   
LD:  $[t_1, \infty]$   
intermediate:  $[t_0, t_1]$

For intermediate window:  
 $\sim 30\%$  from  $\sigma(\pi\pi) \lesssim 1 \text{ GeV}$

# Hadronic vacuum polarization

$\mu$ -e elastic scattering to measure  $a_{\mu}^{\text{HVP}}$

LOI June 2019 [P. Banerjee et al, [arXiv:2004.13663](https://arxiv.org/abs/2004.13663), Eur.Phys.J.C 80 (2020)]



- use CERN M2 muon beam (150 GeV)
- Physics beyond colliders program @ CERN
- [LOI June 2019](https://arxiv.org/abs/2004.13663)
- pilot run in 2021
- full apparatus in 2023-2024

