

1 **Novel Sensors for Particle Tracking: A Contribution to the**
2 **Snowmass Community Planning Exercise of 2021**

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23 **ABSTRACT**

24 Four contemporary technologies are discussed in the context of their poten-
25 tial roles in particle tracking for future high energy physics applications. These
26 include sensors of the 3D configuration, in both diamond and silicon, submicron-
27 dimension pixels, and thin film detectors. Drivers of the technologies include
28 radiation hardness, excellent position and vertex resolution, simplified integra-
29 tion, and optimized power, cost, and material.

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1 Introduction

Research in particle tracking detectors for high energy physics application is underway with a goal of improving radiation hardness, achieving improved position and vertex resolution, simplifying integration, and optimizing power, cost, and material. The four technologies described here approach these goals in complementary ways.

2 Silicon Sensors in 3D Technology

Boscardin, Dalla Betta, Hoferkamp, Seidel, Sultan

Silicon sensors of the 3D technology [1] are employed in LHC experiments [2, 3] to provide radiation tolerant particle tracking at integrated fluences in the regime of $10^{16}n_{\text{eq}}/\text{cm}^2$. The decoupling of the depletion depth from the sensor thickness allows operation at bias voltages below breakdown despite very high integrated fluence, with significant savings on power dissipation, and the small inter-electrode distance suppresses the effect of radiation-induced charge trapping. The ATLAS IBL sensors, for example, are implemented in p-type with $230\mu\text{m}$ thickness and column electrodes of diameter approximately $10\mu\text{m}$, separated by approximately 62 microns. A slim edge of 200 microns is employed. Designs for application to the HL-LHC, where innermost tracking will be exposed over the course of 10 years to fluence $2.3 \times 10^{16}n_{\text{eq}}/\text{cm}^2$ [4], are more aggressive still, in anticipation of conditions in which the carrier lifetime will be reduced to 0.3 ns, corresponding to a mean free path of 30 microns. Up to 200 interactions per 25 ns bunch crossing are expected at the HL-LHC. Small-pitch 3D pixels ($25 \times 100 \mu\text{m}^2$ or $50 \times 50 \mu\text{m}^2$) have been developed to this purpose, with inter-electrode distances of approximately 30 microns [5] and a slim edge of 150 microns, and are currently in the pre-production phase for the ATLAS ITk.

Plans [6] for future facilities such as the FCC-hh anticipate a lifetime integrated luminosity of 30 ab^{-1} , predicting integrated fluence at the innermost tracking volume approaching $10^{18}n_{\text{eq}}/\text{cm}^2$. Estimates [7] of the pileup conditions are on the order of 1000 events per crossing. Continued development of silicon sensors of the 3D technology presents prospects both for restoration of signal loss in high radiation environments, and for separation of pileup signals by precision timing. Measurements [8] carried out on $50 \times 50 \mu\text{m}^2$ cell 3D sensors have shown signals with a full width of 5 ns, and a rise time of 1.5 ns, with a timing resolution of 30 - 180 ps (depending on the signal amplitude); this is a mode of operation comparable to that achieved by low gain avalanche detectors — but lacking gain — with the advantage of higher radiation tolerance and better fill factor. The standard column configuration of 3D has the disadvantages, however, that the electric and weighting fields are non-uniform, leading to position dependence of the pulse rise time; this is the limiting factor on the timing resolution. New geometries [9, 10] involving p-type trench electrodes spanning the entire length of the detector, separated by lines of segmented n-type electrodes for readout, promise improved uniformity and better timing resolution combined with further increased radiation tolerance. Nevertheless, at this time, trenched electrodes cause higher capacitance and introduce larger dead volumes within the substrate. Device optimizations, especially in terms of geometrical efficiency, remain to be carried out. In addition, this problem can be tackled at the system level by tilting the sensor plane with respect to the particle direction, so that a larger fraction of the charge is generated within the depleted

78 volume, and using multiple planes of sensors with an offset between the electrodes, so all
79 tracks would traverse several planes without crossing the electrodes [11].

80 3D columnar pixels with internal gain [12]-[14] offer an alternative approach to signal
81 restoration at high fluence. When implemented with very small inter-electrode separation,
82 approximately 15 microns or less, these devices can achieve controlled charge multiplication
83 at voltages on the order of 100 V, both before and after irradiation. Moderate gain values
84 can be achieved, sufficient to compensate the loss of charge signal due to irradiation of
85 these thin (approximately 100 μm) devices. Design optimization continues with a goal of
86 achieving uniform gain throughout the cell active volume, also benefiting from the wider
87 operating range that is possible due to increasing the breakdown voltage.

88 The goal of this research is to advance one or two 3D technologies in silicon for track-
89 ing particles, able to operate with adequate signal-to-noise ratio at fluences approaching
90 10^{18} $n_{\text{eq}}/\text{cm}^2$, and timing resolution on the order of 10 ps. Planned research activities
91 include TCAD simulations, process optimization and fabrication of several generations of
92 prototypes, and thorough characterization of the prototypes before and after irradiation to
93 extreme fluences.

94 3 3D Diamond Detectors

95 *Kagan, Trischuk*

96 By 2028, experiments operating at the HL-LHC must be prepared for an instantaneous
97 luminosity of $7.5 \times 10^{34}/\text{cm}^2/\text{s}$ and charge particle fluxes of GHz/cm^2 . After these doses, all
98 detector materials will be trap-limited, with the average drift distance a free charge carrier
99 travels before it gets trapped being below 50 μm [15]. 3D sensors reduce the drift distance
100 the charge carriers must travel to reach an electrode to much less than the sensor thickness.
101 This is particularly beneficial in detectors with a limited distance free charge carriers travel,
102 such as trap-dominated sensor materials like heavily irradiated silicon and pCVD diamond,
103 where the observed signal size is related to the mean free path divided by the drift distance.
104 Under these circumstances one gains radiation tolerance (larger signals) by keeping the drift
105 distance less than the mean free path. With the 3D geometrical structure, charge carriers
106 drift inside the bulk parallel to the surface over a typical drift distance of 25 – 100 μm
107 instead of perpendicular to the surface over a distance of 250 – 500 μm .

108 The RD42 collaboration has studied novel 3D detector designs in diamond, to extend
109 the radiation tolerance of diamond to fluences greater than 10^{17} hadrons/ cm^2 , exceeding
110 the HL-LHC doses. The detector design places column-like electrodes inside the detector
111 material using a 130 fs laser with a wavelength of 800 nm. After focusing to a 2 μm spot,
112 the laser has the energy density to convert diamond into an electrically resistive mixture
113 of different carbon phases [16]. A Spatial Light Modulator (SLM) [17] is used to correct
114 spherical aberrations during fabrication. This helps to achieve in 50 $\mu\text{m} \times 50 \mu\text{m}$ cells a
115 high column yield of $\geq 99.8\%$, a small column diameter of 2.6 μm , and a resistivity of
116 the columns of the order of 0.1 – 1 Ωcm . In this detector geometry, the drift distance an
117 electron-hole pair must travel to reach an electrode can be reduced below the mean free
118 path of an irradiated sensor without reducing the number of electron-hole pairs created. In
119 a detector with 25 $\mu\text{m} \times 25 \mu\text{m}$ cells, the maximum drift distance for charge carriers that
120 go into the saddle point region is 25 μm , and 17.5 μm for charge carriers that avoid the

121 saddle point. The goal of this research project is to create a detector that is essentially
122 immune to radiation doses at the level of 10^{17} hadrons/cm². Initial tests have shown that
123 after 3.5×10^{15} n/cm², the 3D geometry with $50 \mu\text{m} \times 50 \mu\text{m}$ cells has better than three
124 times less charge loss than a planar diamond detector after normalizing both unirradiated
125 devices to a relative charge of 1. Furthermore the charge in the unirradiated 3D device is
126 twice as large as that in the planar device. Thus, in addition to having twice the charge, the
127 3D device also has better than three times less damage, due to the shorter drift distance.
128 In order to achieve the 10^{17} hadrons/cm² goal, completion of the design of 3D diamond
129 devices with $25 \mu\text{m} \times 25 \mu\text{m}$ cells and testing of these devices after irradiation with 10^{17}
130 hadrons/cm² is proposed.

131 4 Beyond CMOS: Submicron Pixels for Vertexing

132 *Fourches, Renard, Barbier*

133 Development of a submicron position sensitive vertex detector for the future linear col-
134 lider experiments is proposed. Although improved relative to their predecessors, the present
135 vertexing pixel detectors at the Large Hadron Collider suffer from low position resolution.
136 The objective of vertex detectors is to enable the accurate secondary vertex determination
137 that is crucial for b-tagging [18, 19] in the case of high transverse momentum (p_T) events.
138 The heavy quark events are characterized by a relatively high lifetime that leads to a sec-
139 ondary vertex distinct from the interaction point [20]. For accurate track reconstruction,
140 it is necessary to improve point to point resolution well below the 5 micrometer limit. In
141 the framework of ILD, development of a pixel detector based on the DoTPiX structure is
142 proposed.

143 Accurate track reconstruction with a vertex detector is possible using a small pitch
144 detector, which in the case of the ILC can reduce the multiplicity (in which a pixel is
145 hit several times). This is crucial for the ILD where the readout of the detector is made
146 only after several bunches. With a track fit, displaced secondary vertices can be evaluated,
147 using an impact parameter technique to select the right track, and the analysis of the full
148 decay of the particle can be done, using all necessary jets. In addition, isolated tracks can
149 be tagged in order to reduce fake events. The vertex detectors implemented in the LHC
150 experiment are based on a hybrid design. The high particle rate at the LHC induces a
151 large dose in the detector where non-ionizing energy loss damages the detector material
152 and the electronic readout. Special techniques have been used to circumvent these effects
153 with the use of hardened processes [21] and adequately doped silicon pixel structures [22].
154 To accommodate the LHC beam crossing time, detectors use a triggered readout involving
155 a fast readout chip (ROC). The on-pixel electronics has to be elaborate to collect the
156 information of all the pixel hits' output when triggered. For the technologies available from
157 the late 1990's to the early 2000's, this requirement excluded small pitch pixel detectors.
158 Even with pitch of tens of microns, the number of channels (pixels) is of the order of tens
159 of millions in the inner vertex detectors.

160 The constraints are different for the ILC where more precise reconstruction is the ob-
161 jective. The advantage of vertex detectors with much improved resolution will be good
162 secondary vertex reconstruction with an accuracy of 0.5 micrometers (or in time at the
163 speed of light, of 1.6 fs). This cannot be matched by a timing procedure, which can only

164 estimate the position of the interaction point in the beam-crossing zone. This zone will
165 be reduced at the ILC compared with the LHC. Additionally, short-lived particles can be
166 tagged at this stage. Significant features of this proposal include the following.

167 The detectors close to the primary interaction point can detect low-mass charged par-
168 ticles that can escape the tracker due to the effect of the magnetic field on low-mass par-
169 ticles [23]. Tagging such particles can be established with a good vertex detector [24].
170 These particles can produce disappearing tracks. The energy of these long-lived particles
171 cannot be determined easily as they escape calorimetry. The only possibility, besides using
172 time-of-flight, is to add extra layers to the vertex detector to match the trajectory.

173 The operation of a vertex detector in a trigger-free mode means that many bunch cross-
174 ings will be combined (pile up) before being output and reset. This makes the use of very
175 small pixels necessary to avoid multiple hits in a single pixel. The pitch has to be reduced
176 to match these requirements, and only a fully monolithic pixel can be used for this purpose.

177 The reconstruction of tracks with relatively low p_T near the interaction point will be
178 easier with a pixel detector with large enough aspect (height/width) ratio. A small pitch
179 (less than 1 micrometer) with a height up to of 10 micrometers (the sensitive zone) opens
180 up such a possibility.

181 A pixel design has been proposed [25, 26]. The necessary simulations have been made
182 to assess the functionality of the proposed device. The next step is to find out what is the
183 best process to obtain the functionality and to reach some required specifications.

184 5 Thin Film Detectors

185 *Kim, Metcalfe, Sumant*

186 Nanoscience technologies are developing new cutting edge materials and devices for a
187 wide range of applications. HEP can take advantage of the many advances by looking
188 toward thin film fabrication techniques to implement a new type of particle detector. Thin
189 Film (TF) Detectors have the potential to be fully integrated, large area, low power, with
190 low dead material, and low cost. The present goal is to investigate potential research paths
191 using thin film technologies and to identify and characterize the performance benefits for
192 future particle experiments.

193 A new detector technology is proposed based on thin films that is aimed at dramatically
194 improving the precision of particle detectors by greatly reducing the mass of the detec-
195 tor [27]. Cleaner signatures of the particles from the primary collision will be obtained by
196 reducing those particles' interactions with dead material, which will improve reconstruc-
197 tion efficiencies and resolutions. Thin Film technologies could potentially replace the entire
198 detector including all the services. If a thin film detector could be printed in large areas
199 (square meters), it is estimated that the cost would be reduced to less than 1% of the
200 current cost. If the nuclear interaction length can be decreased by a factor of 10, then the
201 track reconstruction efficiencies would reach 99% and enable a host of new measurements
202 and searches.

203 Thin Film technology presents one possible solution to achieve these performance mile-
204 stones. TF transistors (TFs) were first conceived in the 1960s by Paul Weimer [28]. By
205 the 21st century, fabrication technologies had improved enough to make it competitive with

206 existing technologies. TFTs are the basis of technologies such as Liquid Crystal Display
207 (LCD) screens, solar cells, and light emitting diodes. It is a rapidly growing technology
208 area with a large market base and has corresponding investment in large scale fabrication
209 and industrialization. Ultimately, the broader interest of these technologies enables HEP
210 to leverage the investments in commercialization as well as the R & D into materials, tools,
211 and techniques.

212 Some of the advantages of TFs are optical transparency, mechanical flexibility, high
213 spatial resolution, large area coverage, and low cost relative to traditional silicon-based
214 semiconductor technology. TF technology uses crystalline growth techniques to layer ma-
215 terials. Monolithic sensors can be fabricated using layers of thin film materials for particle
216 detection with layers for amplification electronics. The advantages of a detector made with
217 this type of technology include single piece large area devices (on the order of a few square
218 meters), high resolution ($< 10 \mu\text{m}$), low cost (100 times less than that of Si-CMOS), low
219 mass, and high curvature for a cylindrical, edgeless design [29]-[31].

220 Fabrication processes such as chemical bath deposition and close-space sublimation on
221 a substrate material can produce thin films with a high degree of precision. Here, the
222 crystalline structure is grown in layers, avoiding drilling and etching techniques standard in
223 traditional silicon fabrication; consequently, TF processing is much less expensive.

224 Thin film electronics can be vertically integrated with a thin film sensor if the fabrication
225 techniques are compatible. This would allow vertical integration of sensor and pixel elec-
226 tronics. Further vertical integration using through-vias would enable signals to pass from
227 one layer to the next, thus enabling several levels of electronic processing. Typical front-
228 end ASIC functions could be integrated into the monolithic structure as well as higher-end
229 processing to perform functions such as data aggregation and region-of-interest processing.
230 Such processing would reduce the number of transmission lines integrated into a top layer
231 and further reduce the material inside the detector volume.

232 The transistor is the most basic unit that determines the power consumption in elec-
233 tronics. Complementary Metal Oxide Semiconductor (CMOS) is a low power technology
234 and is the current mainstay for most of commercial electronics. There are, however, many
235 types of transistor technologies that can outperform CMOS. Silicon Germanium (SiGe)
236 Heterojunction Bipolar Transistors (HBTs) are another class of transistor that typically
237 boast faster speeds and lower power consumption [32]-[34]. Fin-Field Effect Transistors
238 (FinFETs) are being pursued as the next ultra-low power technology and are manufactured
239 by companies such as IBM and Motorola. However, the most transformational transistor
240 is the Thin Film Transistor (TFT), which is breaking records in terms of size and power.
241 All of these technologies have the potential for reducing the power (and the copper in the
242 transmission lines) over the current technologies.

243 Thin Film Detectors have the potential to replace a wide range of detector types from
244 tracking to calorimetry. The present goals are to identify key areas of research within Thin
245 Film technologies, quantify the key requirements from different types of experiments, and
246 evaluate the potential physics impact.

247 **6 Conclusion**

248 Four contemporary technologies are under development for applications at future high en-
249 ergy physics experiments. Collaborators interested in joining any of these efforts are wel-
250 come.

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