

Nucleon & Nuclear Structure Inputs to Beyond-SM Searches

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SNOWMASS Rare Processes & Precision Frontier
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- Nucleon Charges & Form Factors
 - Beyond-SM Contributions to neutron decays*
 - Dark Matter searches*
 - Lepton Flavor Violation*
- Nucleon/Nuclear Electric Dipole Moments
 - nEDM from quark EDMs*
 - Intrinsic nucleon moments from θ_{QCD}*
 - and quark & gluon chromoEDMs*
- Baryon number violation
 - Proton decay*
 - Neutron-antineutron oscillations*
- Neutrinoless double-beta decay
 - Short- and Long-range Lepton Number Violation*
 - Charged Pion Transition Amplitudes*
 - Pion Annihilation Amplitudes*

Nucleon ($\bar{q}\Gamma q$) M.E's: Neutrinos & β -decay

$$\langle p + q | \bar{q} \gamma_\mu \gamma_5 q | p \rangle = \bar{u}' \left[G_A \gamma_\mu + \tilde{G}_P \frac{q_\mu}{M} \right] \gamma_5 u$$

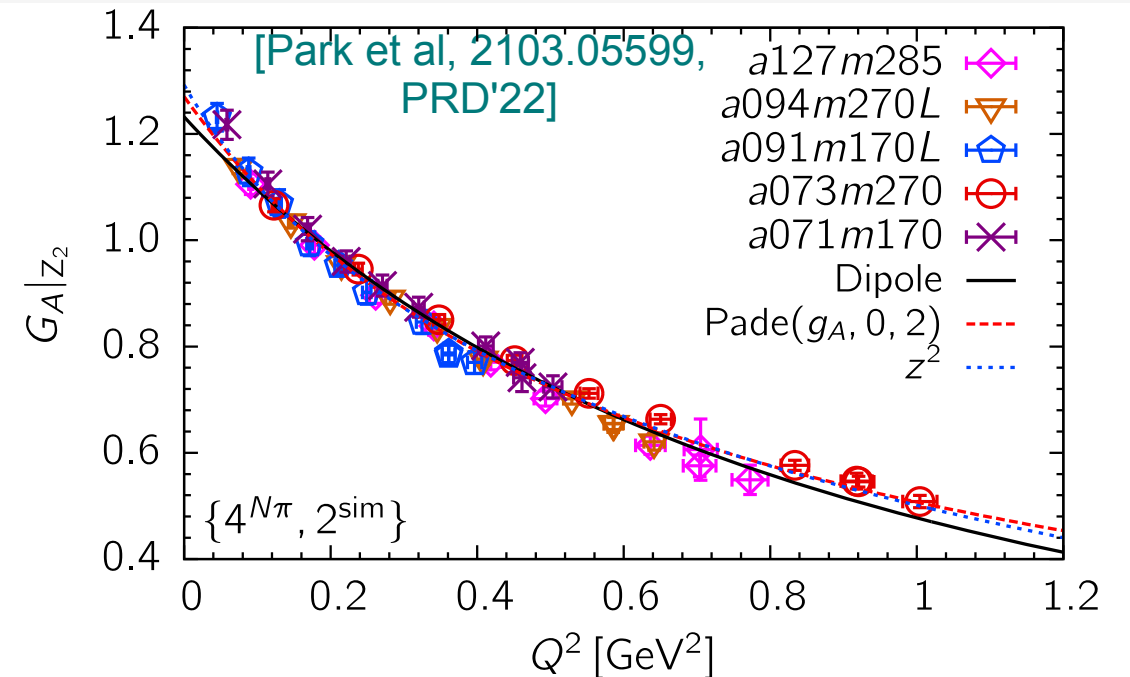
$$\langle p + q | \bar{q} \gamma_5 q | p \rangle = G_P \bar{u}' \gamma_5 u$$

$$\langle p | \bar{q} q | p \rangle = g_S \bar{u}' u$$

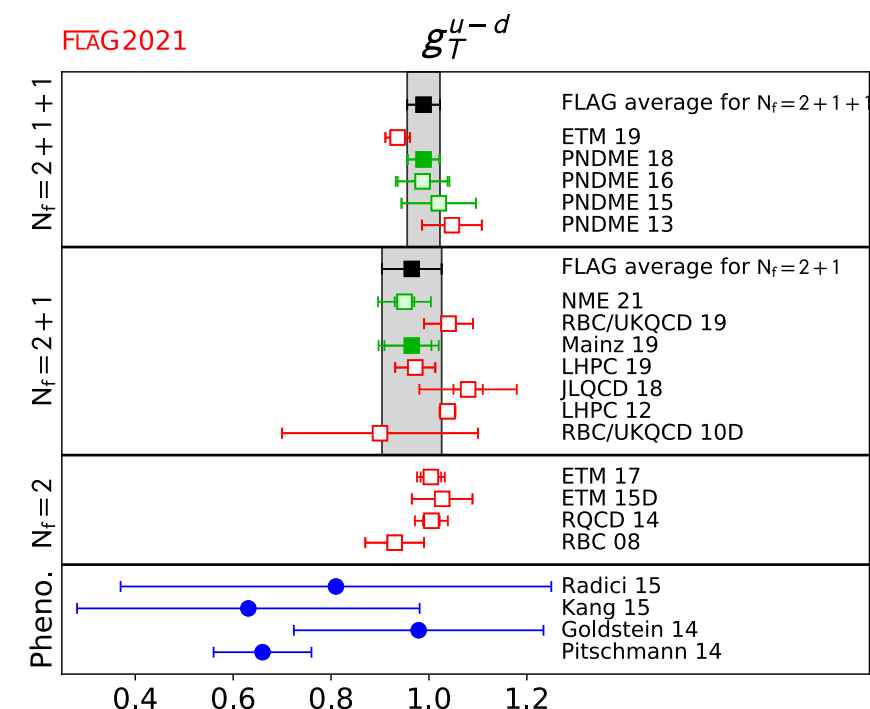
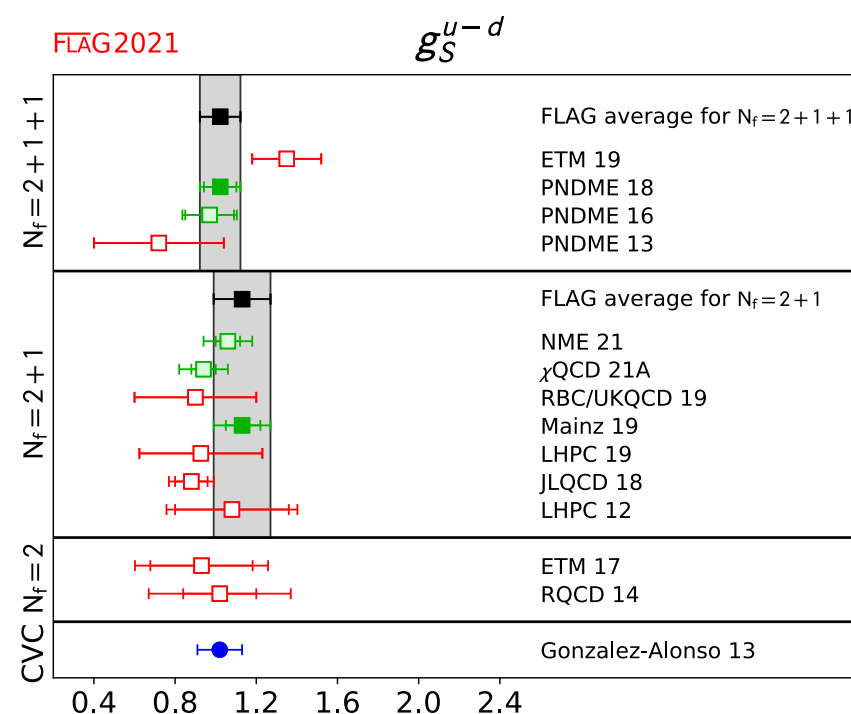
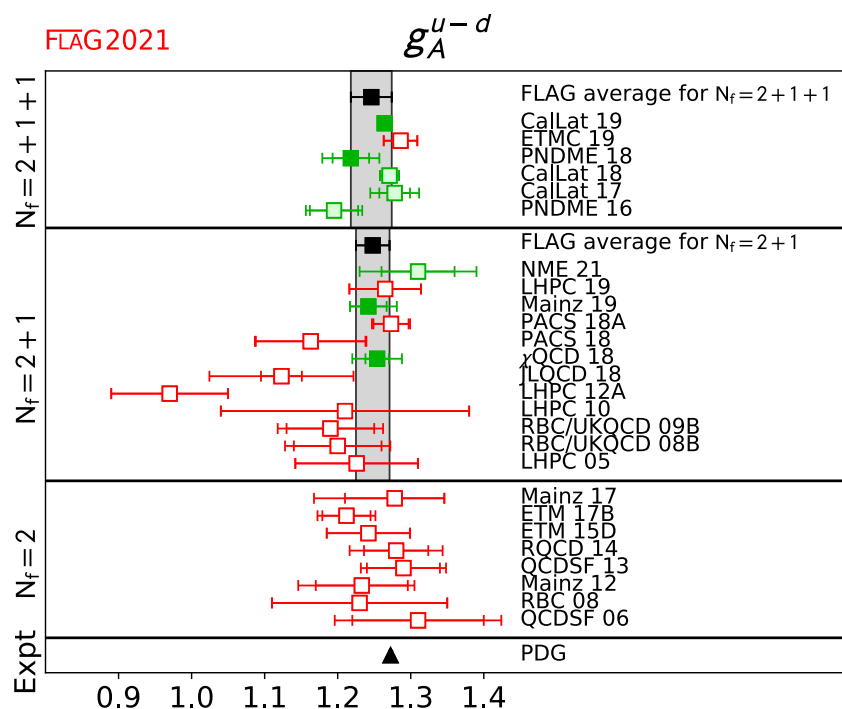
$$\langle p | \bar{q} \sigma_{\mu\nu} q | p \rangle = g_T \bar{u}' \sigma_{\mu\nu} u$$

- Axial form factors $G_A(Q^2)$: QE neutrino scattering
- isovector g_A : neutron β -decay
- g_S, g_T : sensitivity of β -decay to Beyond-SM physics

[Y.Aoki et al, FLAG'21 2111.09849]



- **Ongoing controversy on excited-state contamination in single-nucleon ME: e.g. $N\pi$ states required for PCAC**
- **Much harder problem in light nuclei**



Nucleon ($\bar{q}\Gamma q$) M.E's: DM, LFV, CPv Sensitivity

- DM direct detection:
scalar, fermion DM EFT couplings

$$\sim g_{A,S,P,T}, \langle GG \rangle, \langle G\tilde{G} \rangle$$

- Lepton Flavor violation
 $\mu \rightarrow e$ conversion EFT couplings

$$\sim G_{V,A,S,P,T}(Q^2 = -(m_\mu)^2), \langle GG \rangle$$

- g_T : CP violation in nucleon EDM from quark EDM

Example: scalar dark matter couplings

[Bishara et al 1707.06998]

$$Q_{1,q}^{(6)} = (\varphi^* i \overleftrightarrow{\partial}_\mu \varphi) (\bar{q} \gamma^\mu q),$$

$$Q_{3,q}^{(6)} = m_q (\varphi^* \varphi) (\bar{q} q),$$

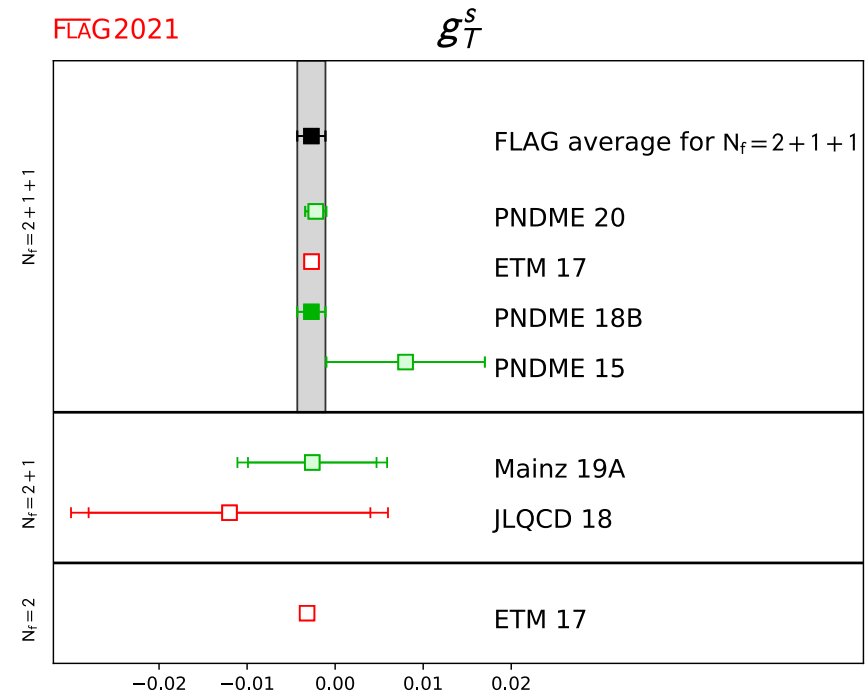
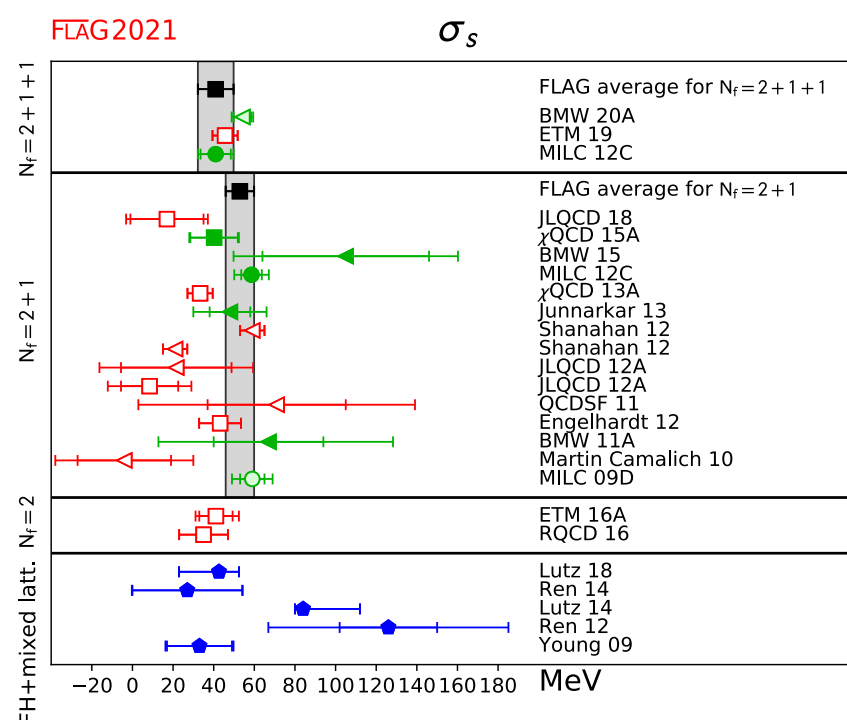
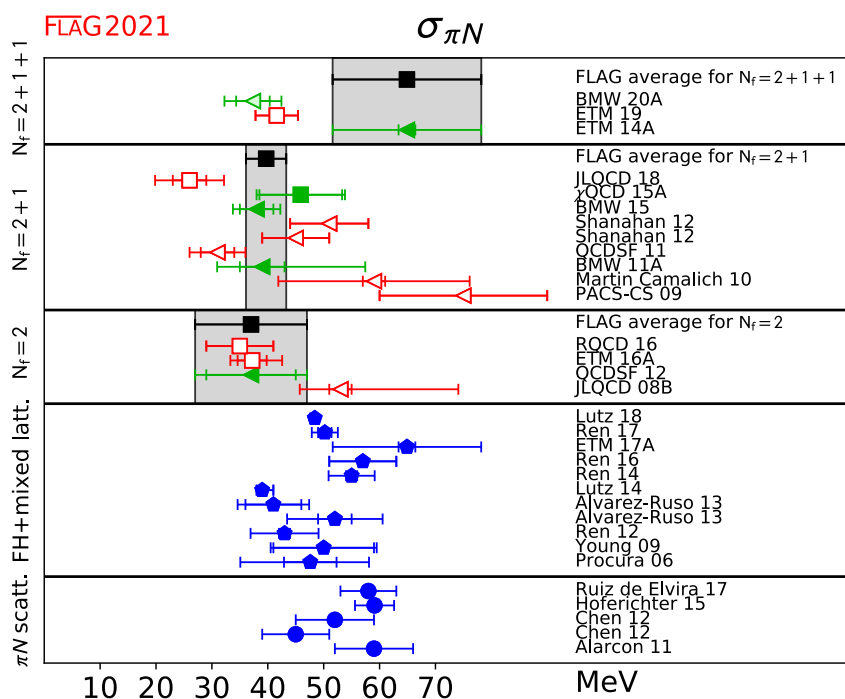
$$Q_5^{(6)} = \frac{\alpha_s}{12\pi} (\varphi^* \varphi) G^{a\mu\nu} G_{\mu\nu}^a,$$

$$Q_{2,q}^{(6)} = (\varphi^* i \overleftrightarrow{\partial}_\mu \varphi) (\bar{q} \gamma^\mu \gamma_5 q),$$

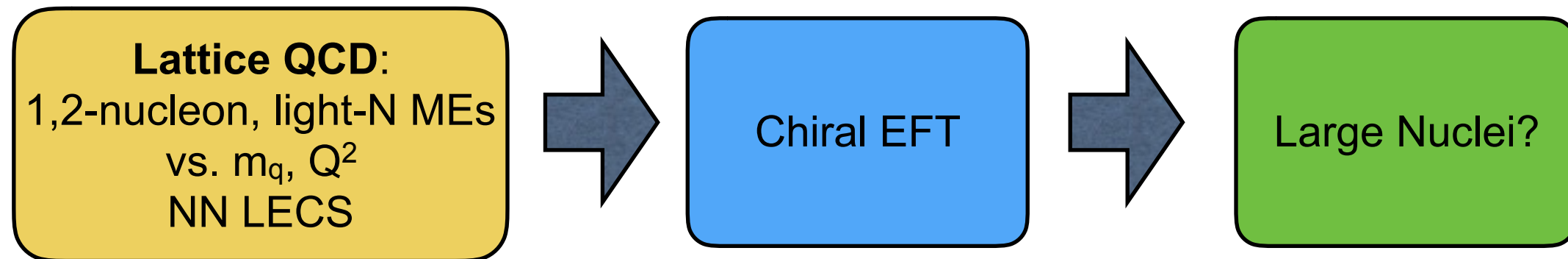
$$Q_{4,q}^{(6)} = m_q (\varphi^* \varphi) (\bar{q} i \gamma_5 q),$$

$$Q_6^{(6)} = \frac{\alpha_s}{8\pi} (\varphi^* \varphi) G^{a\mu\nu} \tilde{G}_{\mu\nu}^a.$$

[Y.Aoki et al, FLAG'21 2111.09849]



Nuclear M.E's: Summary / Challenges



Single-nucleon matrix elements

- Quark matrix elements $(\bar{q}\Gamma q)$: **precision results available**
 - persistent doubt about excited states
- Gluon matrix elements $\langle GG \rangle$, $\langle G\tilde{G} \rangle \sim$ **straightforward (albeit expensive)**
 - Noise(Gluon observables) \gg Noise (Quark observables)
 - mixing with isoscalar quark densities
- Momentum dependence (form factors): **straightforward**
 - large V_3 for small momentum transfer $Q^2 \lesssim 0.1 \text{ GeV}^2$

Two-nucleon and light-nuclei MEs

- at the physical point: **challenging**
 - dense spectrum of excitations
 - multiple-volume dependence needed

2-body amplitude of local interaction:

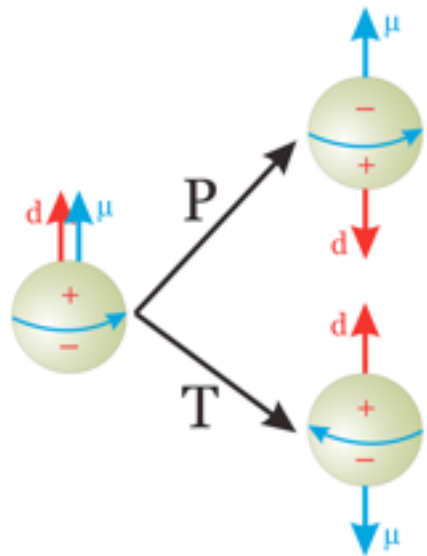
$$|\mathcal{A}_{V \rightarrow \infty}|^2 \propto \left(q \frac{\partial \phi(q)}{\partial q} + k \frac{\partial \delta(k)}{\partial k} \right) |\mathcal{A}_{V=L^3}|^2$$

$\delta(k)$ = scattering phase

$\phi(q)$ = lattice zeta fcn.

[Lellouch, Luscher, '01]

P, CP(T) Violation: Nucleon/Nuclear EDM



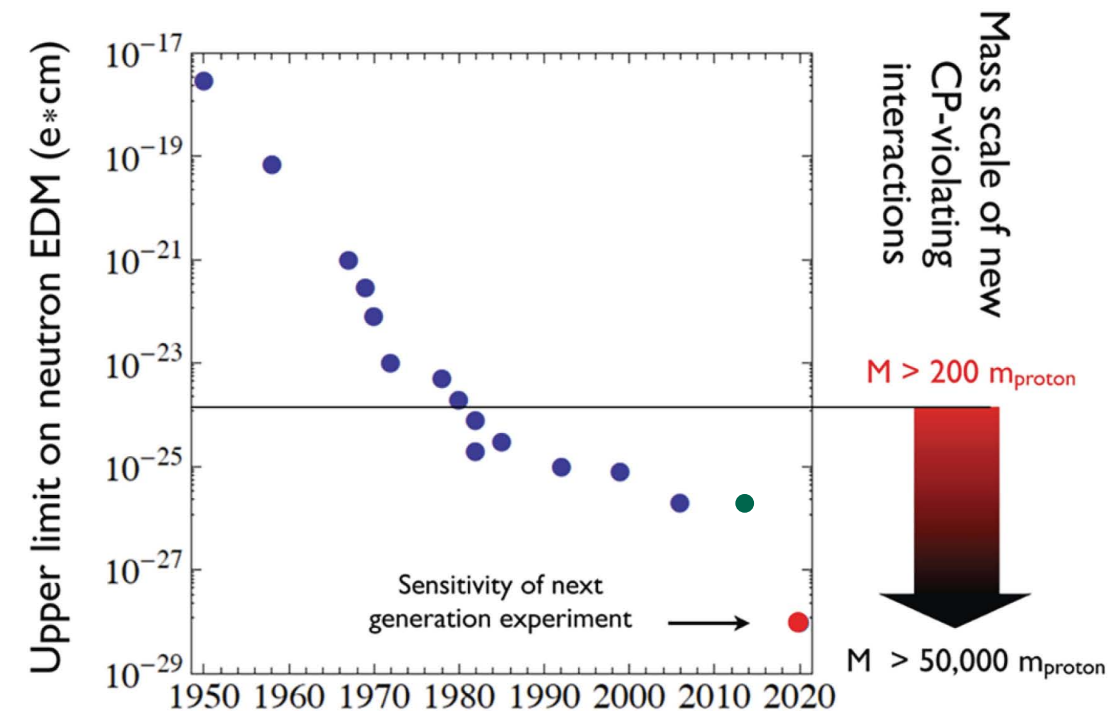
$$\vec{d}_N = d_N \frac{\vec{S}}{S}$$

$$\mathcal{L}_{int} = eA_\mu^{\text{em}} \gamma^\mu \quad (\text{P,T-even})$$

$$+ eA_\mu^{\text{em}} \mathcal{A}^\mu \quad (\text{P,T-odd})$$

The most sensitive probe of non-CKM CPv:

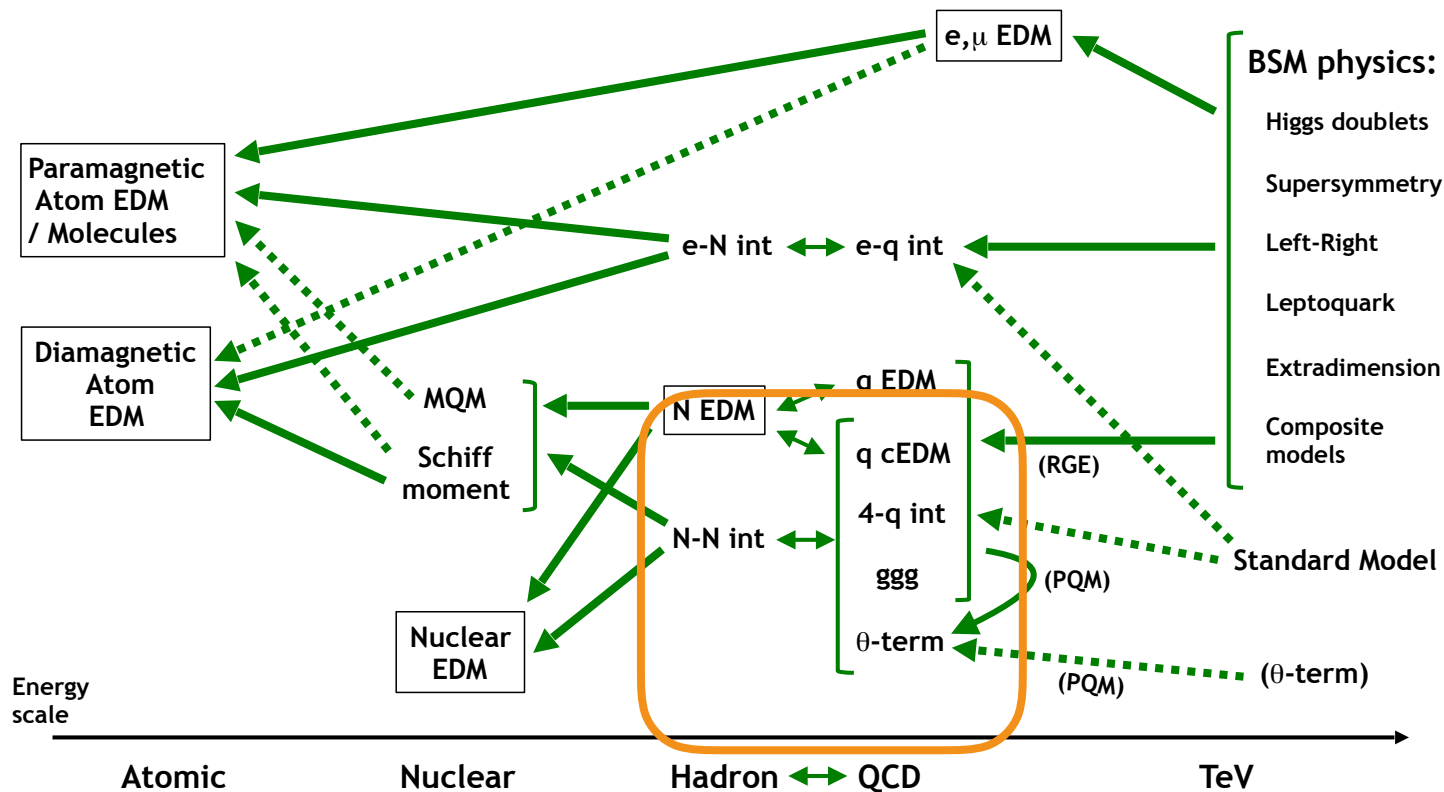
- Any signal $> 10^{-5} \cdot (\text{current bound}) \rightarrow \text{discovery}$
- θ_{QCD} -induced EDM : Strong-CP problem
- Prerequisite for Baryogenesis (non- θ_{QCD} EDM)



	10^{-28} e cm
CURRENT LIMIT	<300
Spallation Source @ORNL	< 5
Ultracold Neutrons @LANL	~30
PSI EDM	<50 (I), <5 (II)
ILL PNPI	<10
Munich FRMII	< 5
RCMP TRIUMF	<50 (I), <5 (II)
JPARC	< 5
Standard Model (CKM)	< 0.001

[B.Filippone '16]

Sources of Nucleon and Nuclei EDMs



[Yamanaka et al, EPJA53:54 (2017)]

- Effective quark-gluon CPv interactions organized by dimension
[Engel, Ramsey-Musolf, van Kolck, Prog.Part.Nucl.Phys. 71:21 (2013)]

$$\mathcal{L}_{eff} = \sum_i \frac{C_i}{[\Lambda_{(i)}]^{d_i-4}} \mathcal{O}_i^{[d_i]}$$

$$d=4 : \theta_{QCD}$$

$$d=5(6) : \text{quark EDM, chromo-EDM}$$

$$d=6 : 4\text{-fermion CPv, 3-gluon (Weinberg)}$$

- Lattice QCD: CPv at hadronic scale (Nucleon EDM and πN CPv interactions)

$$d_{n,p} = d_{n,p}^\theta \theta_{QCD} + d_{n,p}^{cEDM} c_{cEDM} + \dots$$

Lattice methods for computing nEDM

- Background electric field

$$\mathcal{H} = -\vec{d}_N \cdot \vec{E}$$

- CPv form factor $F_3(Q^2 \rightarrow 0)$

$$\langle p+q | J^\mu | p \rangle_{CP} = \bar{u}_{p'} \left[F_1 \gamma^\mu + (F_2 + iF_3 \gamma_5) \frac{\sigma^{\mu\nu} q_\nu}{2m_N} \right] u_p$$

Dirac
Pauli (anomalous magnetic dipole)
Electric dipole

Nucleon EDM from Quark EDMs

Quark EDM \rightarrow nucleon EDM: "tensor charge"

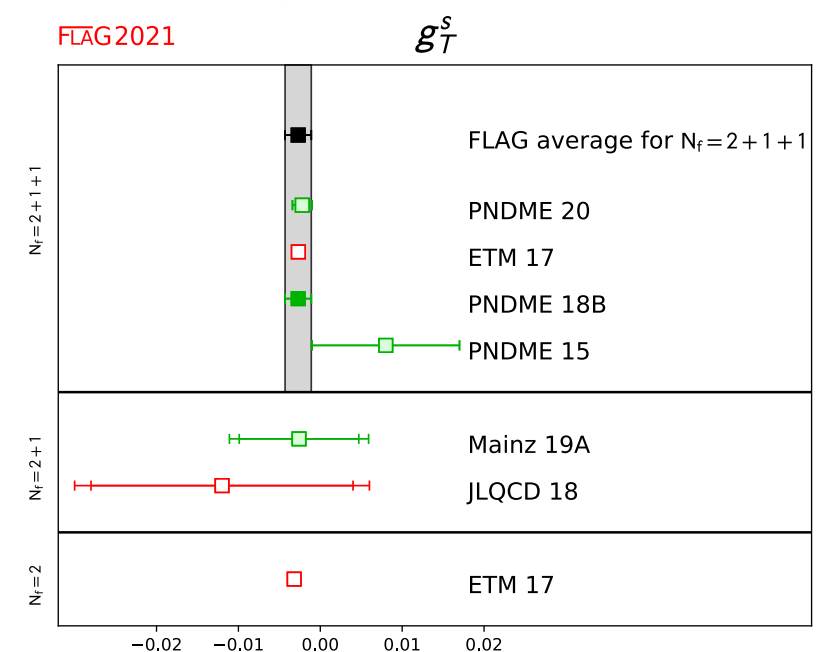
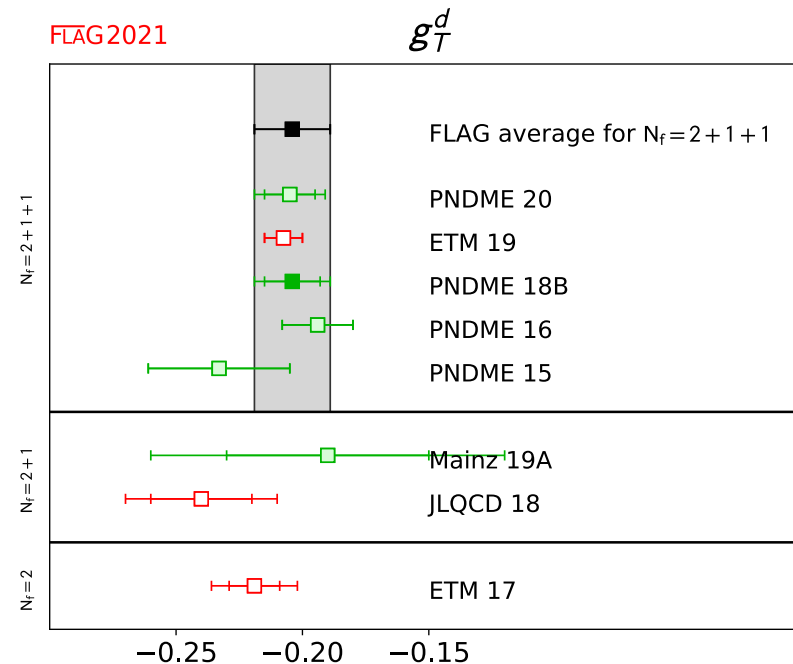
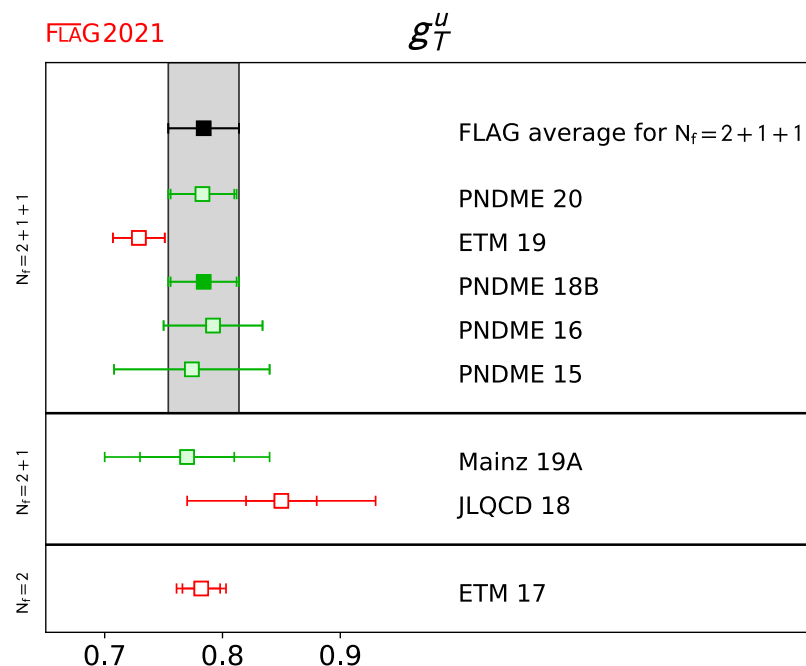
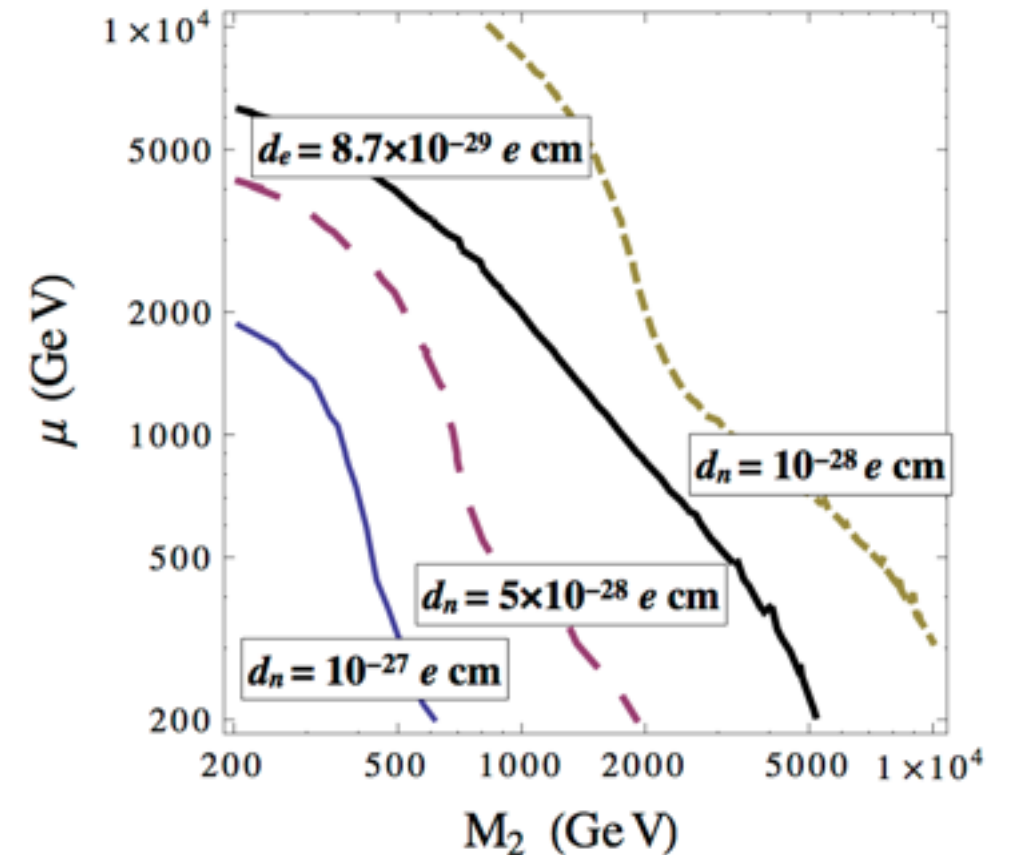
$$\delta\mathcal{L}_{\text{CPV}} \supset -\frac{ie}{2} \sum_{f=u,d,s,e} d_f \bar{f} \sigma_{\mu\nu} \gamma_5 F^{\mu\nu} f$$

$$d_N = g_T^u d_u + g_T^d d_d + g_T^s d_s$$

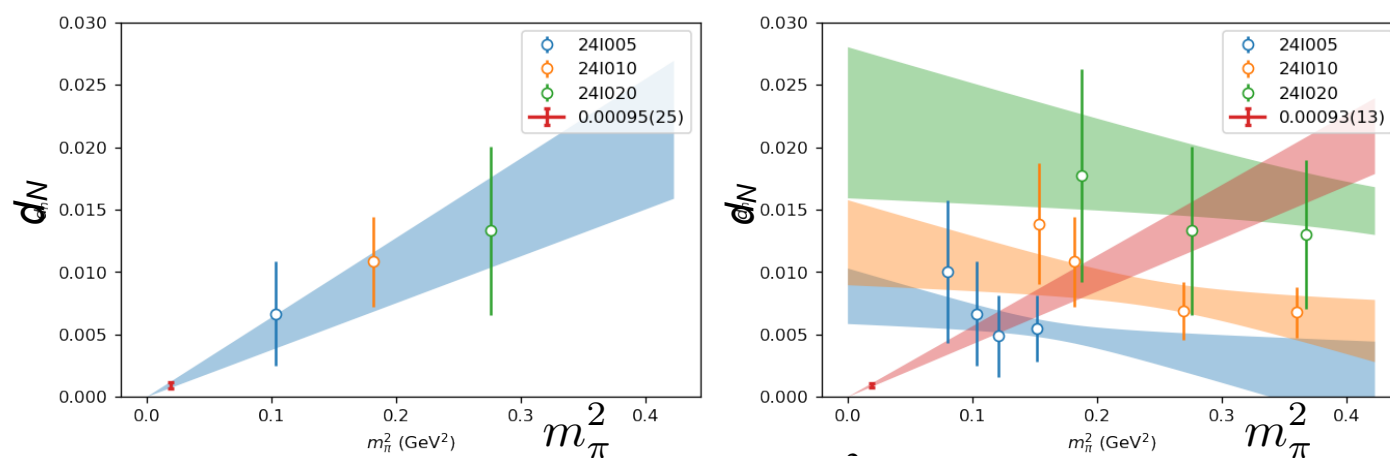
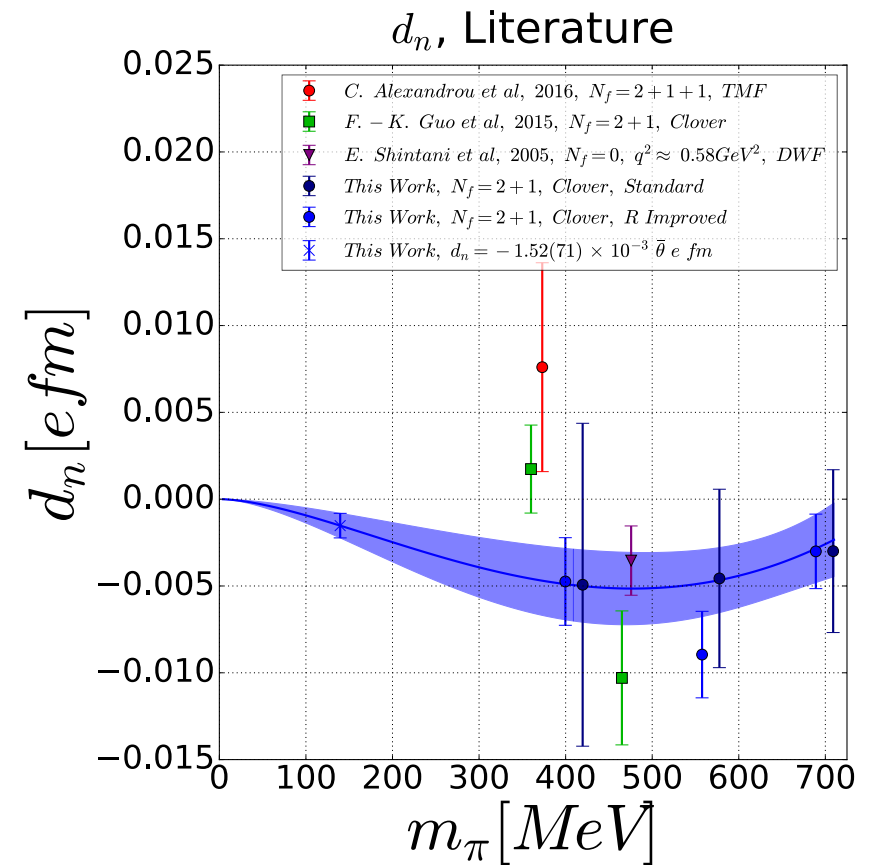
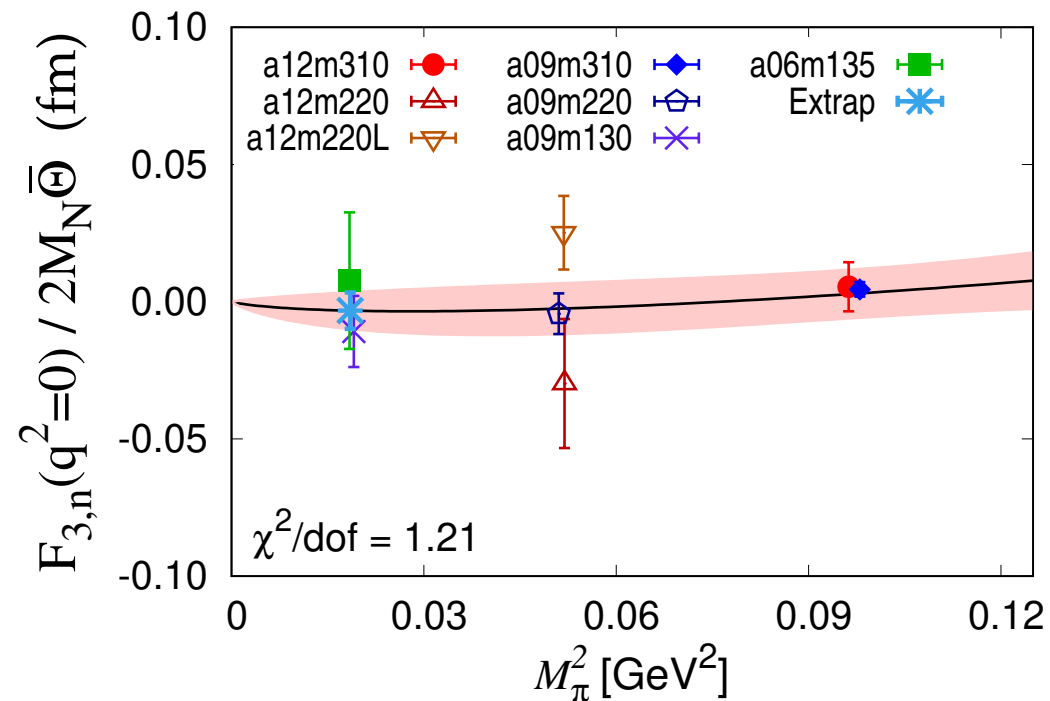
$$\langle N | \bar{q} \sigma^{\mu\nu} q | N \rangle = d_T^q \bar{u} \sigma^{\mu\nu} u$$

[Y.Aoki et al, FLAG'21 2111.09849]

Constraints on split-SUSY model
[Bhattacharya et al, PRL115:212002 (2019)]



Nucleon EDM from θ_{QCD} -Term



DW + Overlap quarks: [K.-F.Liu et al, PRELIMINARY]:

- partial- V_4 sampling of Q
- partial quenching to explore m_q dependence

Challenges:

- effect of θ_{QCD} vanishes at $m_q \rightarrow m_q^{\text{phys}}$
chiral symmetry important
- noise in global $Q = \int F \tilde{F}$ grows with V_4
- $Q^2 \rightarrow 0$ extrapolation(*)

Topological Charge with Gradient Flow

[M.Luscher, JHEP08:071; 1006.4518]

Gradient flow: covariant $4D$ -diffusion of quantum fields with "G.F." time t_{GF} :

$$\frac{d}{dt_{GF}} B_\mu(t_{GF}) = D_\mu G_{\mu\nu}(t_{GF}), \quad B_\mu(0) = A_\mu$$

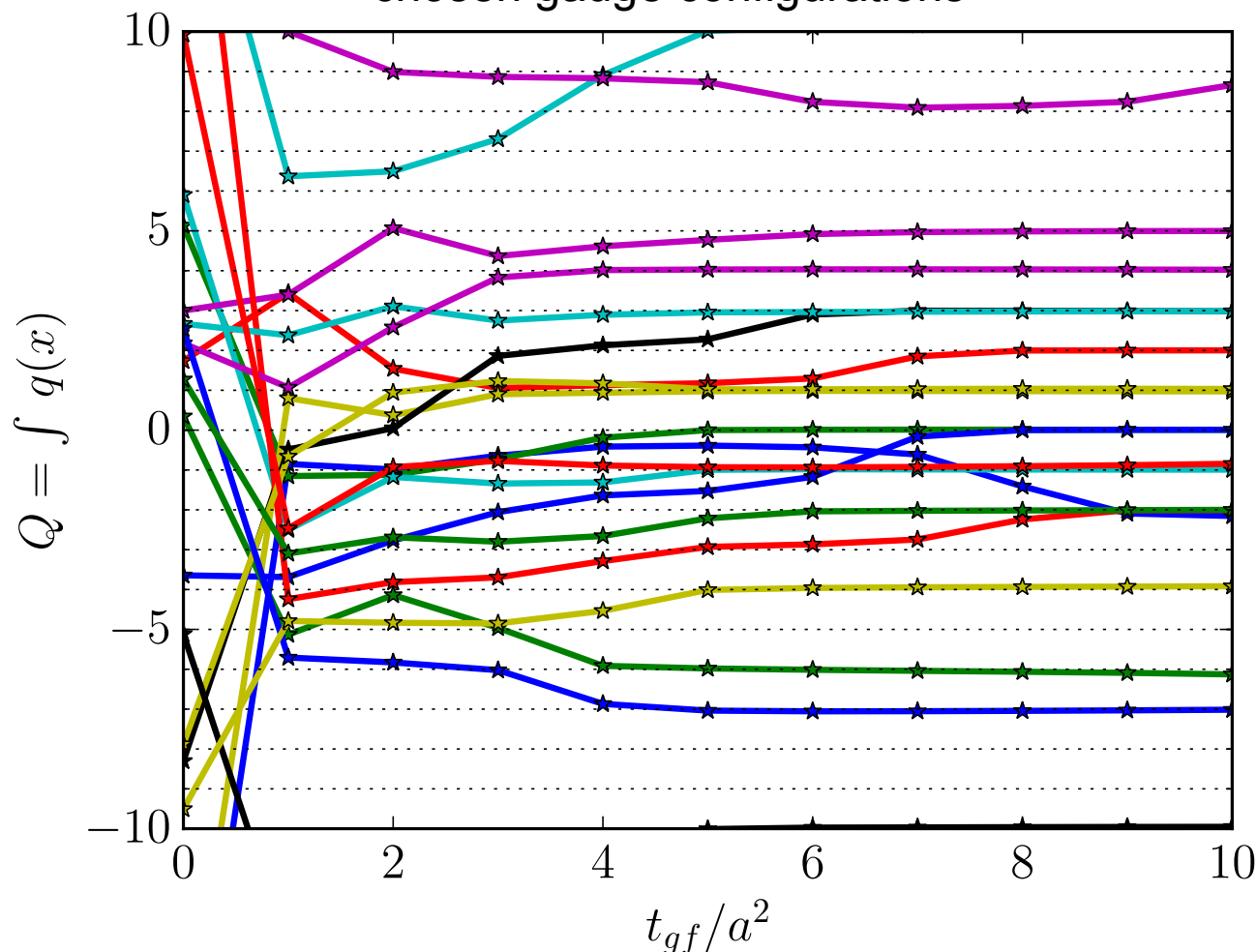
Tree-level:

$$B_\mu(x, t_{GF}) \propto \int d^4y \exp \left[-\frac{(x-y)^2}{4t_{GF}} \right] A_\mu(y)$$

Gradient-flowed topological charge:

$$\tilde{Q}(t_{GF}) = \int d^4x \frac{g^2}{32\pi^2} \left[G_{\mu\nu} \tilde{G}_{\mu\nu} \right] \Big|_{t_{GF}}$$

total top. charge on 20 randomly chosen gauge configurations



- effective scale $\Lambda_{UV} \rightarrow (t_{GF})^{-1/2}$
 \Rightarrow renormalization prescription
- smooth fields (reduce $|G_{\mu\nu}|$)
 \Rightarrow continuous "cooling"
- remove $G_{\mu\nu}$ dislocations
 \Rightarrow separation of top. sectors

[M.Luscher, JHEP08:071; 1006.4518]

Topological Charge with Gradient Flow

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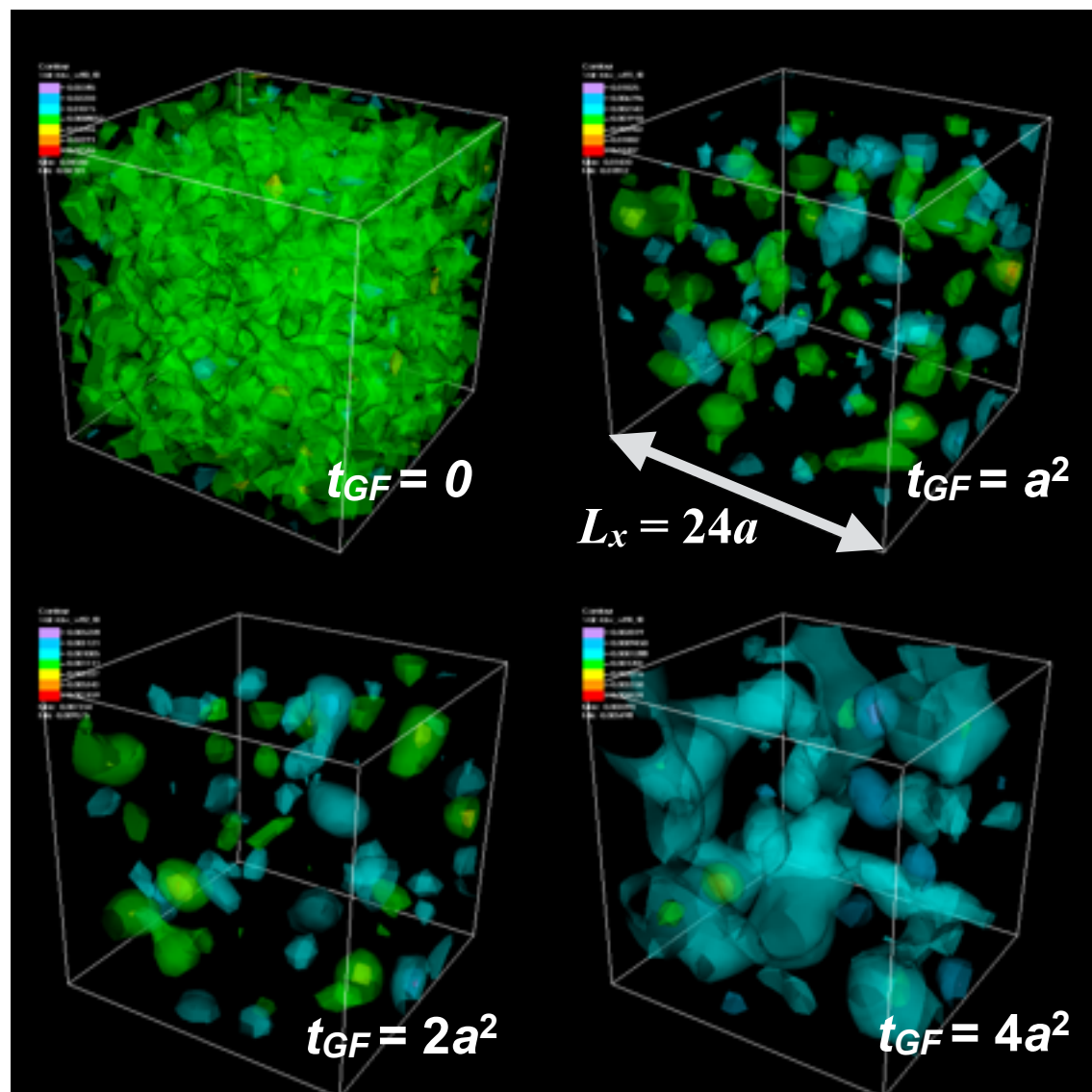
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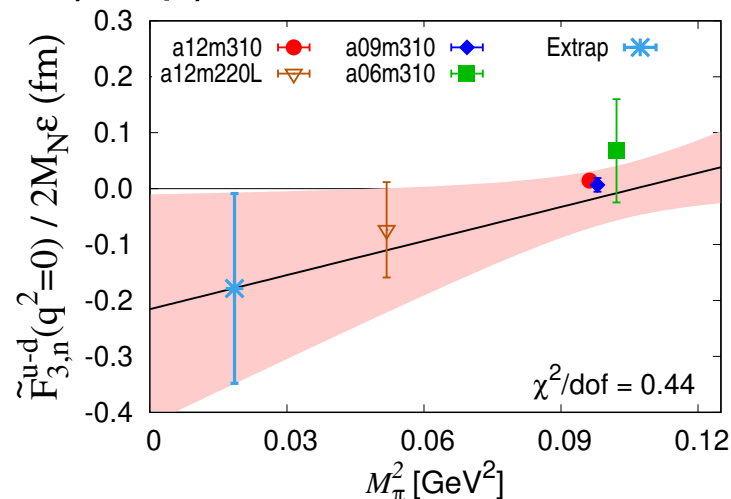
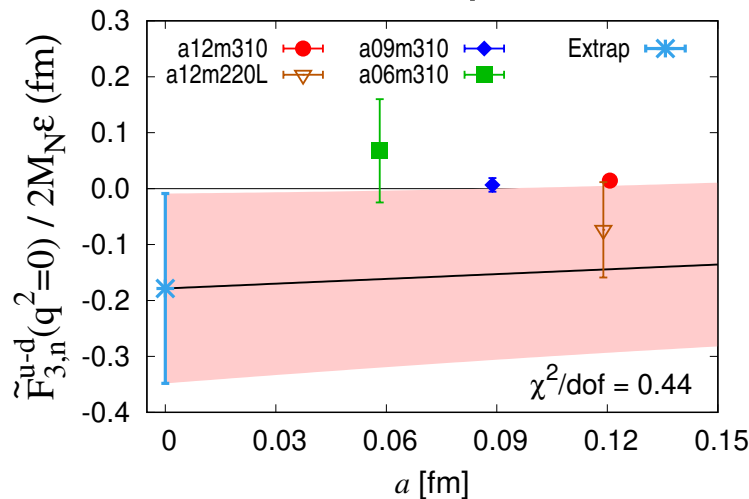
[M.Luscher, JHEP08:071; 1006.4518]



$$\Leftarrow \tilde{q}(t_{GF}) = \frac{g^2}{32\pi^2} \left[G_{\mu\nu} \tilde{G}_{\mu\nu} \right] \Big|_{t_{GF}} \propto \tilde{\mathbf{E}} \cdot \tilde{\mathbf{H}}$$

Nucleon EDM from Quark and Gluon cEDM

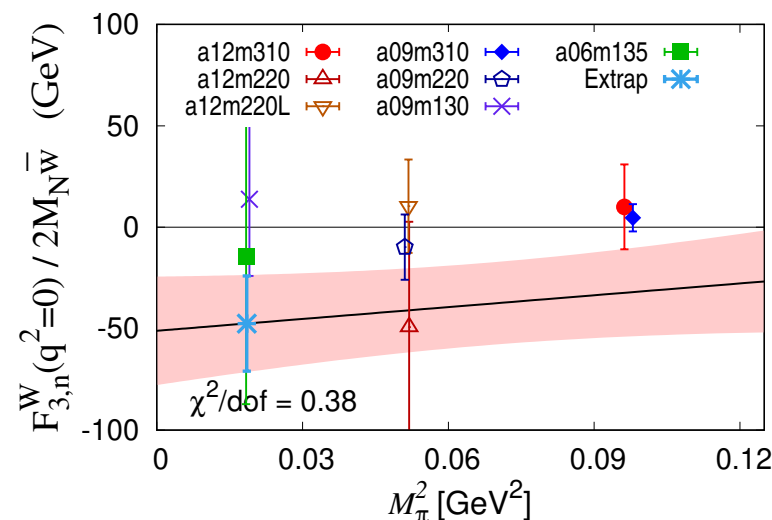
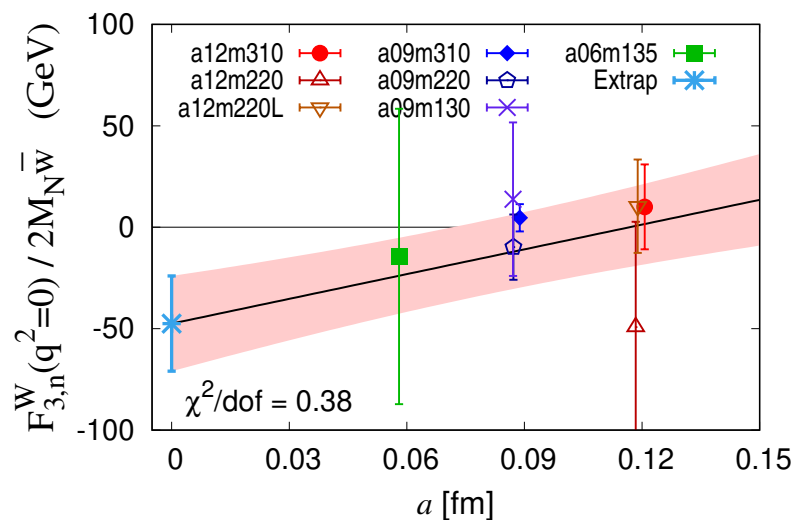
- dim-5 CPv: quark cEDM ($\bar{q} (\sigma \cdot F) \gamma_5 q$) :



Challenges:

- renormalization & mixing with lower-dim operators (e.g. dim-4 $F\tilde{F} \rightarrow \theta_{\text{QCD}}\text{-nEDM}$)
- quark cEDM (dim-5): quark-disconnected contractions for vector current and cEDM
- Weinberg term (dim-6): gluon noise

- dim-6 CPv: Weinberg term ($\sim F \cdot F \cdot \tilde{F}$):



Wilson-Clover quarks

[T.Bhattacharya, Lattice 2021; 2203.03746]

Renormalization & mixing with Gradient Flow:

[Mereghetti et al 2111.11449; J.Kim et al 2106.07633]

- novel nonperturbative scheme
- small t_{GF} expansion \Rightarrow matching to perturbative
- avoid power divergences

Nucleon&Nuclear EDMs: Summary/Challenges

Single-nucleon EDM

- θ_{QCD} -induced nEDM : ongoing work; some extrapolated results
 - potentially **challenging** at the physical point
- Higher-dim effective operators
 - quark/gluon cEDM mixing/renormalization tractable with Gradient flow
 - very noisy / weak "signal" -> **straightforward** with improved statistics
 - 4-quark CPv just starting; potentially **challenging** due to quark loops
 - depend on precision of θ_{QCD} -induced nEDM due to mixing

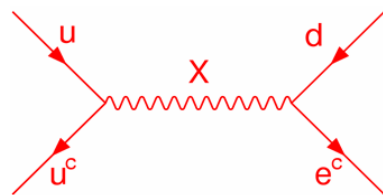
Multi-nucleon EDM and CPv NN interaction

- **extremely challenging**, will need EFTs

Baryon Number Violation: Proton Decay

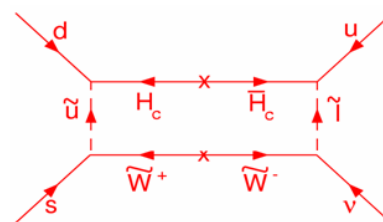
● Proton decay: $\Delta B=1$ violation

- Baryon number = accidental "symmetry" of SM; violated by sphalerons
- Probe scales inaccessible to colliders: Limits on GUT, extra dim., etc
- Limits on stability of nuclear matter



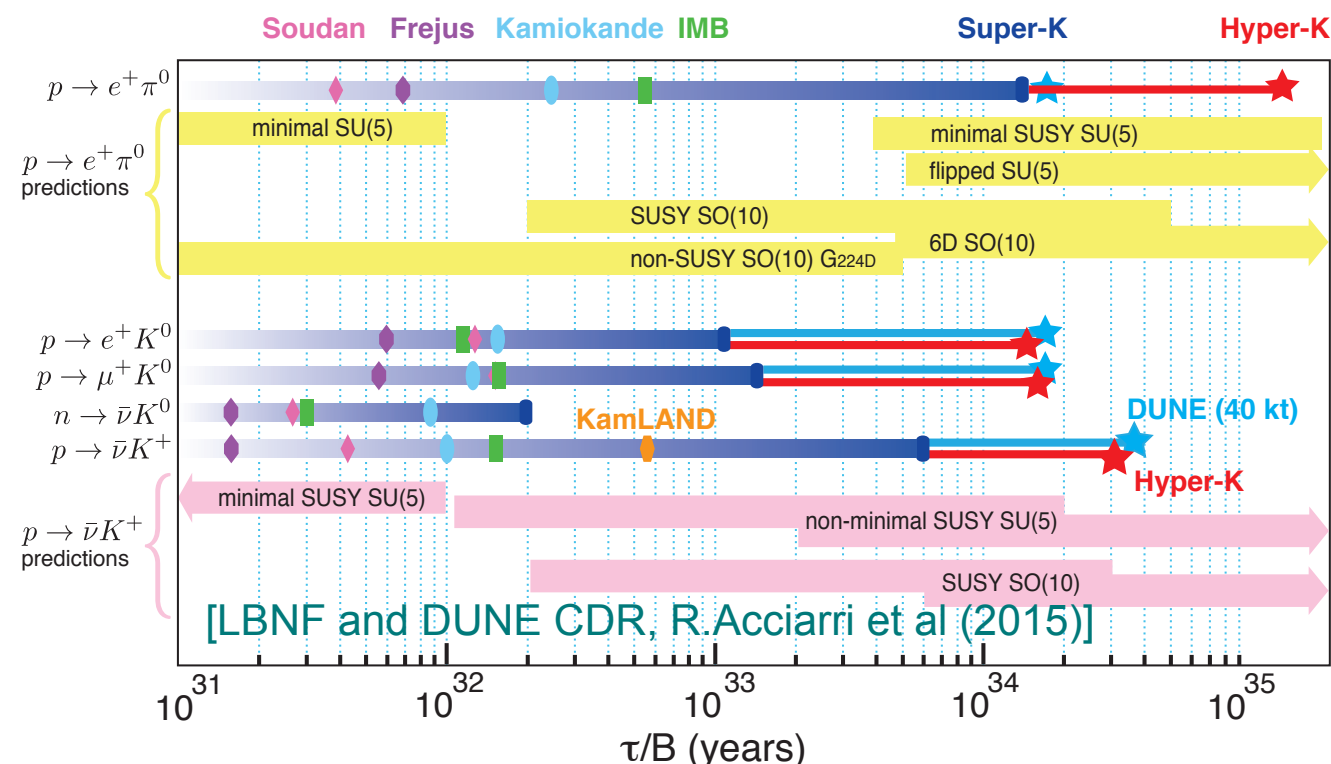
ordinary GUT

- min. $SU(5)$ ruled out by $\tau(p \rightarrow e^+\pi^0)$
- $SO(10)$ probed by next-gen exp.



SUSY GUT [Sakai, Yanagida '82; Weinberg '82]

- min.SUSY- $SU(5)$ ruled out by $\tau(p \rightarrow \bar{\nu}K^+)$
- SUSY- $SO(10)$ probed by next-gen exp.



- $\tau(p \rightarrow e^+\pi^0) \gtrsim 1.6 \cdot 10^{34}$,
 $\tau(p \rightarrow \bar{\nu}K^+) \gtrsim 5.9 \cdot 10^{33}$ [Super-K]
- Expect x10 limits from Hyper-K, DUNE
- DUNE: sensitive to $p \rightarrow \bar{\nu}K^+$ (SUSY GUT)

Alt. explanation of proton stability?

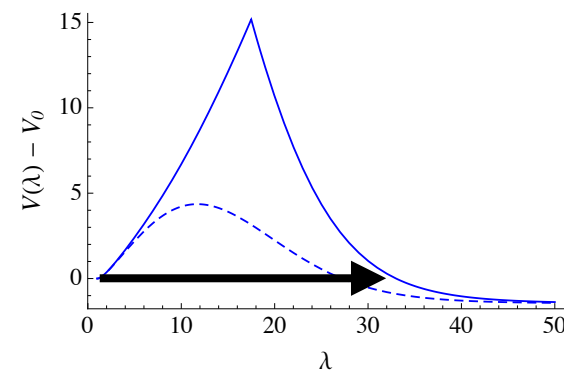
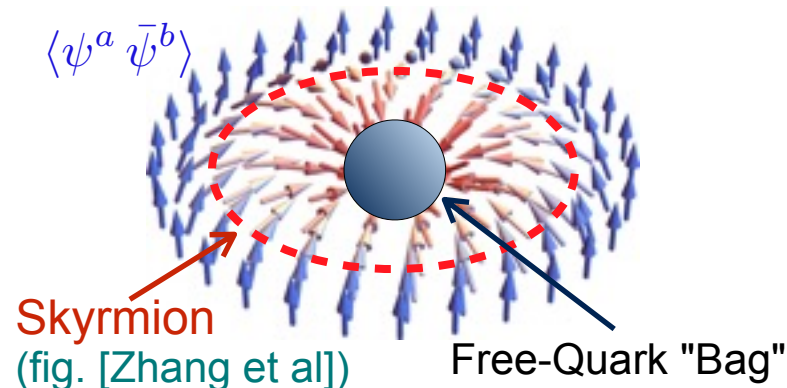
Multi-meson decay channels?

Proton Decay Matrix Elements

Is proton inherently stable?

Conjecture [A.Martin, G.Stavenga '12]

Topological stability of "Chiral Bag" proton :



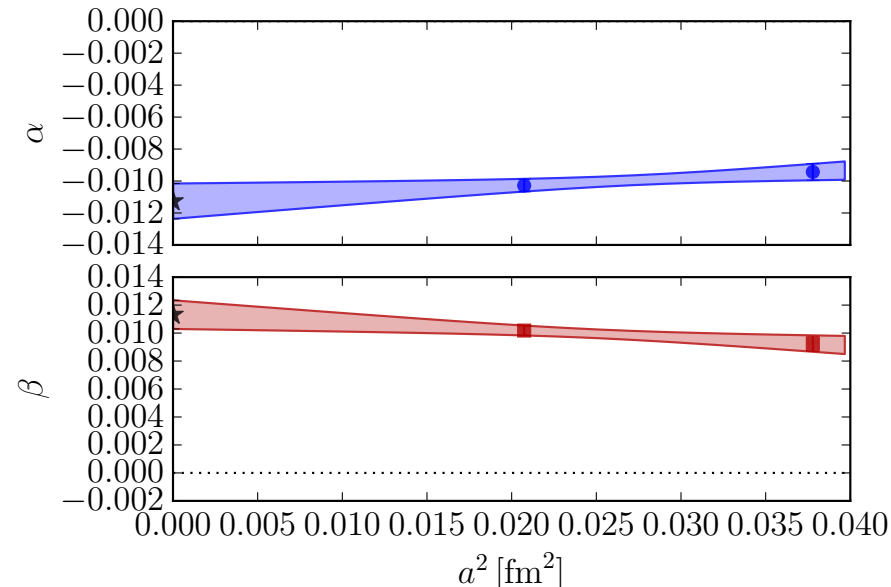
Lattice calculations with chirally-symmetric quarks:

- Prev. at $m_\pi \gtrsim 300$ MeV
[S.Aoki et al (2000)]
[Y.Aoki et al (2006), (2013), (2017)]
- Physical quarks
[J.Yoo, PRD'22]
- NEXT: $p \rightarrow \pi\pi$, $p \rightarrow \pi K$

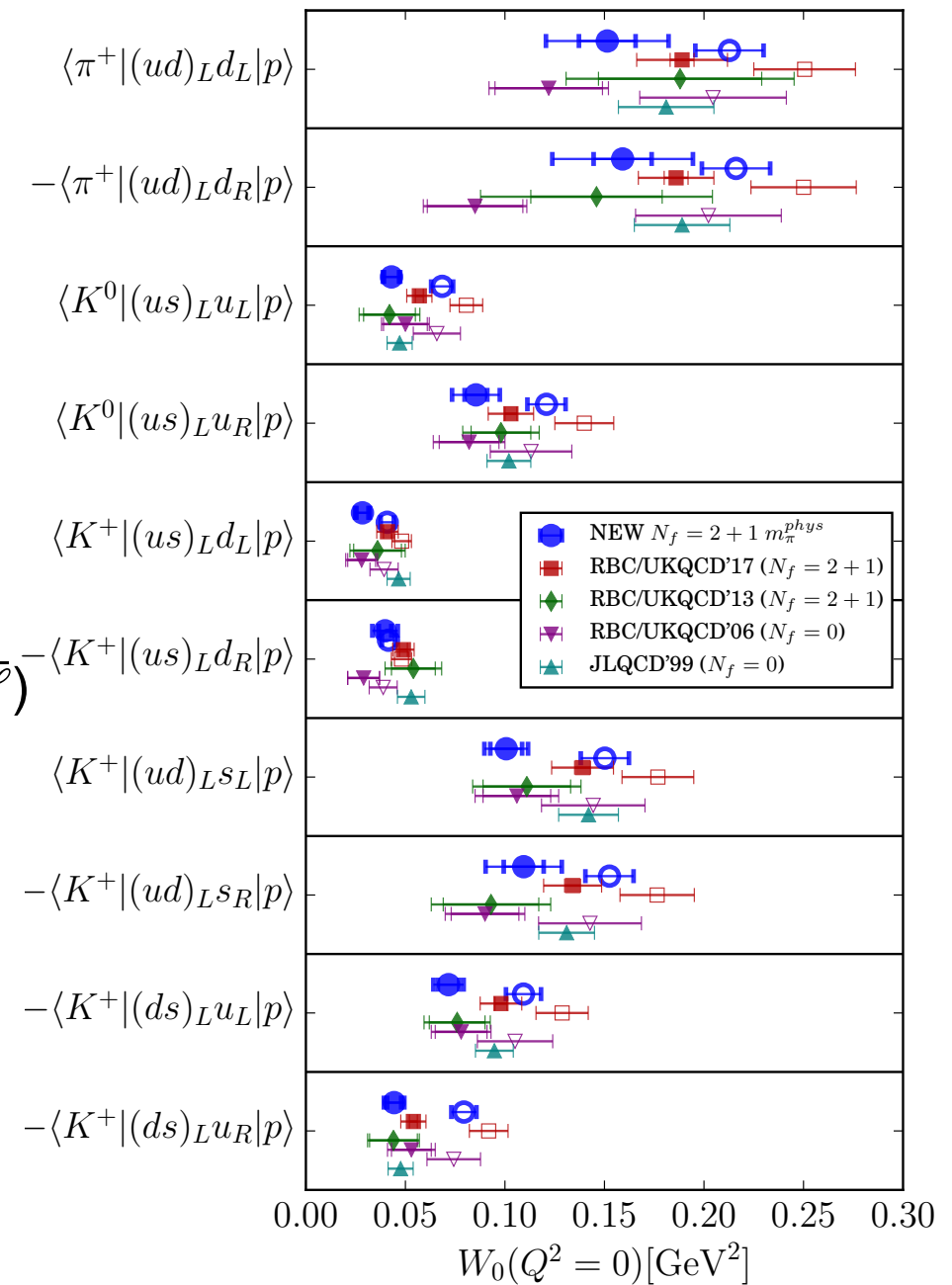
Nucleon decay constants
($p \rightarrow 3\ell$ decays; LO ChPT $p \rightarrow K\bar{\ell}$, $\pi\bar{\ell}$)

$$\langle \text{vac} | \epsilon^{abc} (\bar{u}^{aC} d^b)_R u_L^c | N \rangle = \alpha P_L U_N$$

$$\langle \text{vac} | \epsilon^{abc} (\bar{u}^{aC} d^b)_L u_L^c | N \rangle = \beta P_L U_N$$



Nucleon-to-meson amplitudes
($p \rightarrow \pi\bar{\ell}$, $K\bar{\ell}$, decays)



\Rightarrow NO SUPPRESSION
at physical quark masses

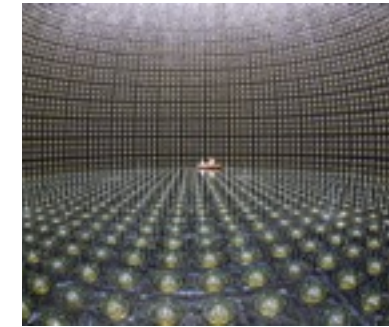
Baryon Number Violation : Neutron Oscillation

● $N\bar{N}$ oscillation: $\Delta B = 2 = \Delta(B-L)$ violation

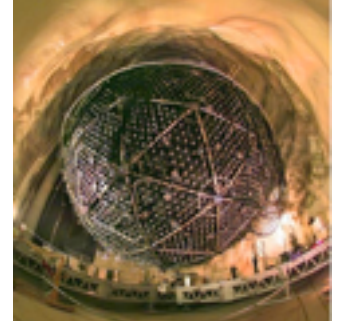
- $\tau(^{56}\text{Fe}) \gtrsim 0.72 \cdot 10^{32} \text{ yr}$
 $\Rightarrow \tau_{N\bar{N}} \gtrsim 1.4 \cdot 10^8 \text{ s}$ [Soudan]
- $\tau(^{16}\text{O}) \gtrsim 1.77 \cdot 10^{32} \text{ yr}$
 $\Rightarrow \tau_{N\bar{N}} \gtrsim 3.3 \cdot 10^8 \text{ s}$ [Super-K]
- $\tau(^2\text{H}) \gtrsim 0.54 \cdot 10^{32} \text{ yr}$
 $\Rightarrow \tau_{N\bar{N}} \gtrsim 1.96 \cdot 10^8 \text{ s}$ [SNO]
- Quasi-free reactor neutrons
 $\tau_{N\bar{N}} \gtrsim 10^8 \text{ s}$ [ILL'94]



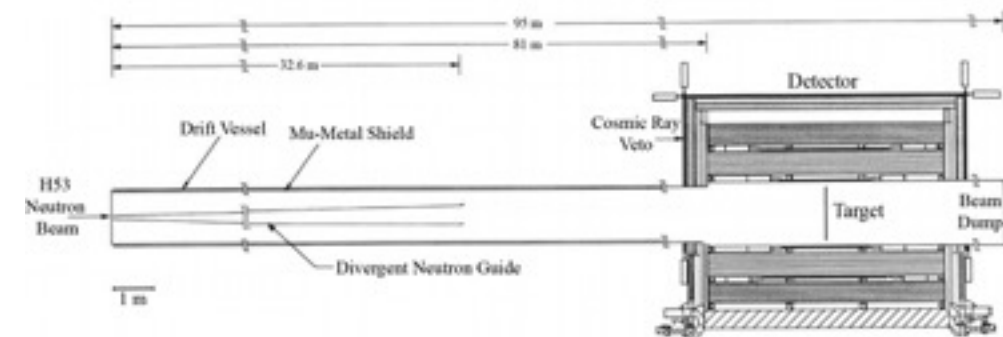
Soudan



Super Kamiokande



SNO



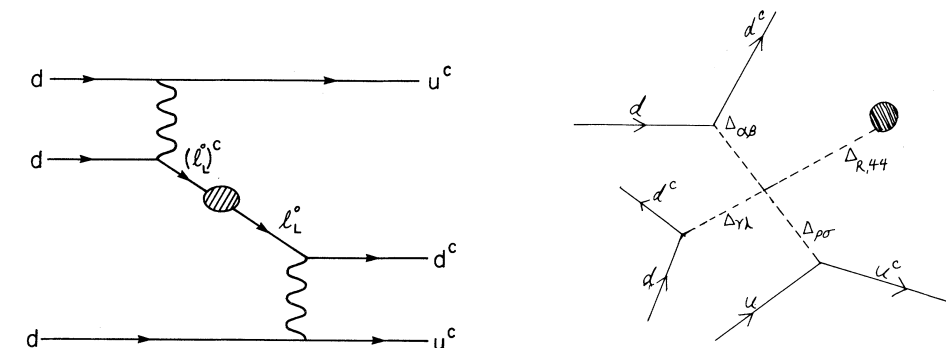
- GUT + massive Majorana lepton?
[T.K.Kuo, S.T.Love, PRL45:93 (1980)]
- partial unification and $(B-L)$ viol.?
[R.N.Mohapatra, R.E.Marshak, PRL44:1316 (1980)]

$$\delta m = -\langle \bar{n} | \int d^4x \mathcal{L}_{\text{eff}} | n \rangle = -\sum_i \frac{c_i}{M_X^5} \langle \bar{n} | \mathcal{O}_i^{6q} | n \rangle$$

oscillation time
 $\tau_{n\bar{n}} = (2\delta m)^{-1}$

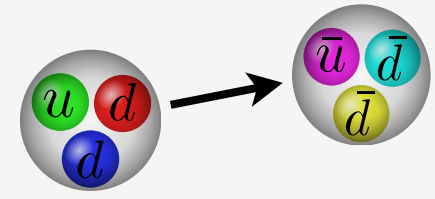
$\sim (\text{BSM scale})^{-5}$
 $M_X \gtrsim (200-300) \text{ TeV}$

$N\text{-}\bar{N}$ amplitude



$$\mathcal{L}_{\text{eff}} = \sum_i [c_i \mathcal{O}_i^{6q} + \text{h.c.}]$$

Lattice QCD Result: Enhanced $N \leftrightarrow \bar{N}$

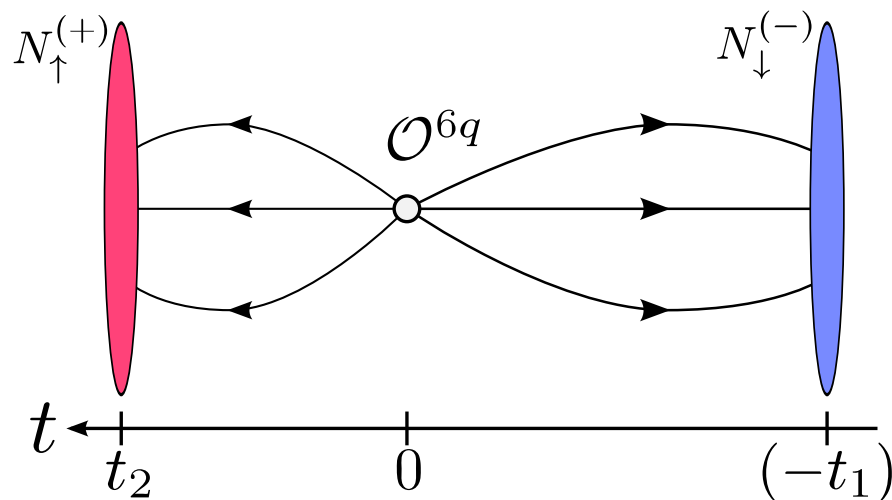


Lattice QCD with physical-mass, chiral-symmetric quarks:

x(5-10) larger N - N bar oscillation vs. nucleon Bag model

[E.Rinaldi, S.S., M.Wagman, et al, PRD99:074510 (2019)]

[E.Rinaldi, S.S., M.Wagman, et al, PRL122:162(2018)]



	$\mathcal{O}^{\overline{MS}}(2 \text{ GeV})$	Bag "A"	$\frac{\text{LQCD}}{\text{Bag "A"}}$	Bag "B"	$\frac{\text{LQCD}}{\text{Bag "B"}}$
$[(RRR)_3]$	0	0	—	0	—
$[(RRR)_1]$	45.4(5.6)	8.190	5.5	6.660	6.8
$[R_1(LL)_0]$	44.0(4.1)	7.230	6.1	6.090	7.2
$[(RR)_1 L_0]$	-66.6(7.7)	-9.540	7.0	-8.160	8.1
$[(RR)_2 L_1]^{(1)}$	-2.12(26)	1.260	-1.7	-0.666	3.2
$[(RR)_2 L_1]^{(2)}$	0.531(64)	-0.314	-1.7	0.167	3.2
$[(RR)_2 L_1]^{(3)}$	-1.06(13)	0.630	-1.7	-0.330	3.2
	$[10^{-5} \text{ GeV}^{-6}]$	$[10^{-5} \text{ GeV}^{-6}]$		$[10^{-5} \text{ GeV}^{-6}]$	

MIT Bag model results from [S.Rao, R.Shrock, PLB116:238 (1982)]

Next steps: non-quasi-free oscillation

- full systematic UQ : finite volume, continuum limit
- "crossed" 2-neutron annihilation amplitudes $\langle \text{vac} | \mathcal{O}^{6q} | nn \rangle$
- Nuclear medium effects

Neutrinoless Double-Beta Decay

$0\nu 2\beta$: Experimental window into

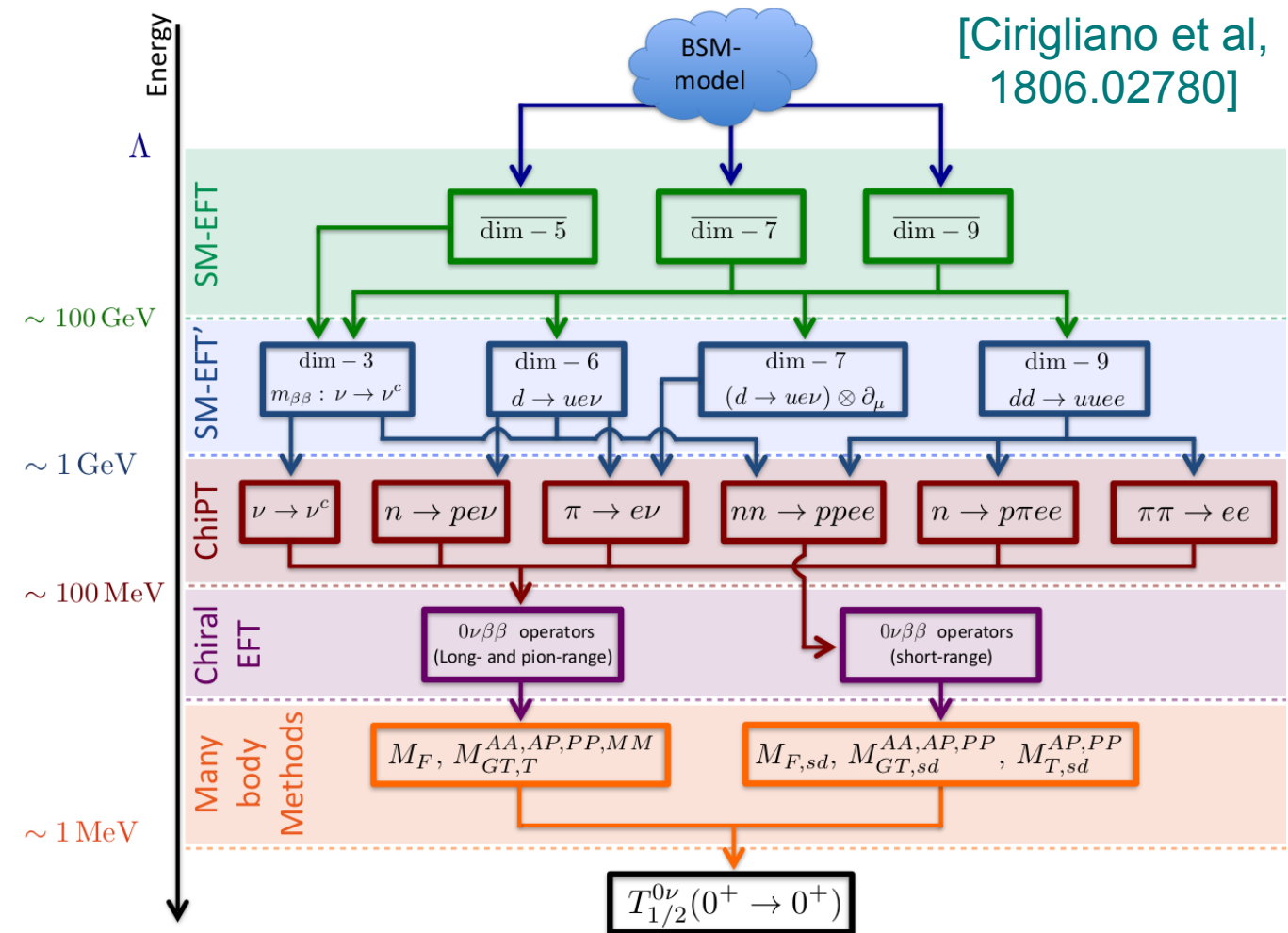
- Neutrino mass mechanism
- mass hierarchy
- leptogenesis (LNV \Rightarrow BNV thru spalerons)
- scale of BSM

Experiments

- current : NEMO3, KamLAND-Zen, EXO-200, Majorana, GERDA, CUORE, CUPID
- Next-gen: x100 lifetime constraints (e.g. ton-range nEXO (1t ^{136}Xe))

LNV mechanisms:

- light Majorana neutrino? – only in $0\nu 2\beta$
- low-scale seesaw?
- BSM at TeV scale?
- neutrion flavor models



$$|m_{\beta\beta}| = |m_i U_{ei}^2| \leq 10^{-2} eV$$

$$\Gamma^{0\nu\beta\beta} \propto |m_{\beta\beta}|^2 |M^{0\nu}|^2$$

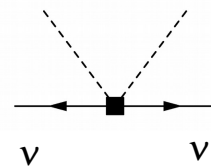
nuclear matrix elements
lattice QCD + EFT + many-body

EFTs for $0\nu 2\beta$

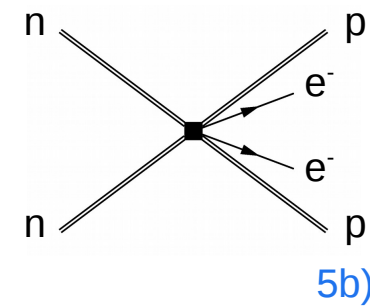
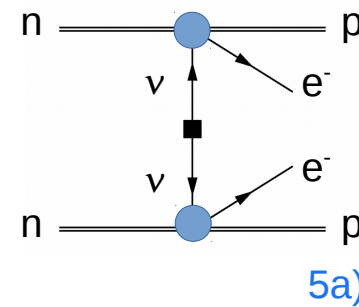
[2203.12169]

SM EFT

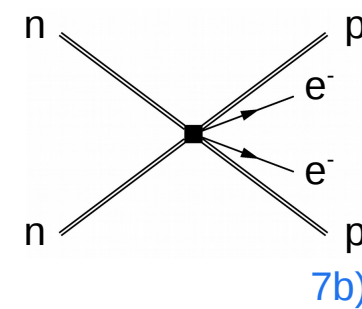
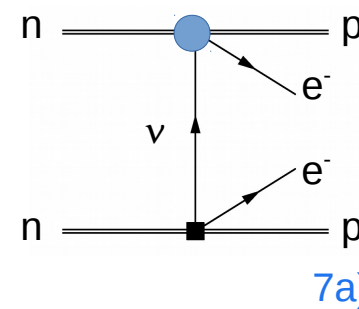
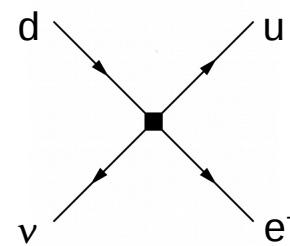
Dim 5



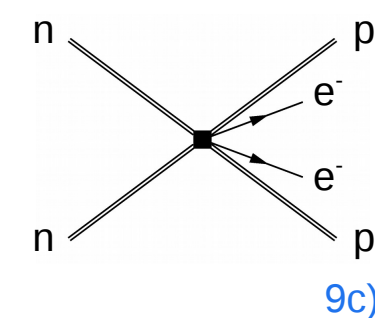
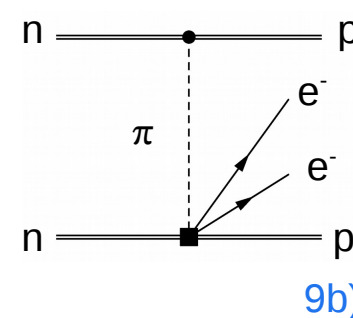
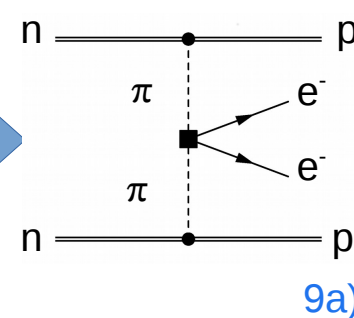
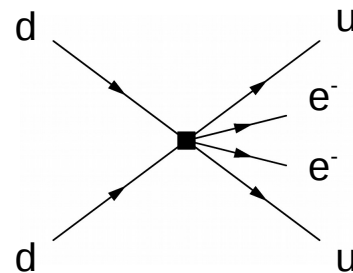
Chiral EFT



Dim 7



Dim 9

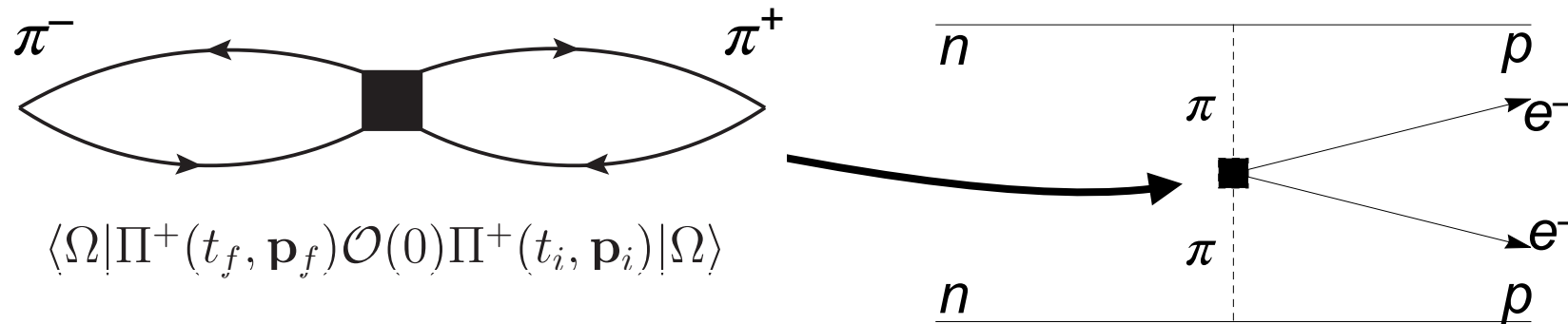


LVN operators in SM EFT: dim-5,7,9 [Prezeau et al, '03]

- dim-5, dim-7 \Rightarrow Long-range LVN
- dim-9 \Rightarrow Short-range, π -range LVN
- dim-5,7,9 \Rightarrow short-range contact contractions in Chiral EFT

Short-Range LNV in $\pi^- \rightarrow \pi^+ ee$

$\langle \pi^+ | \mathcal{O}^{4q} | \pi^- \rangle$ [Nicholson et al, PRL'18, 1805.02634]



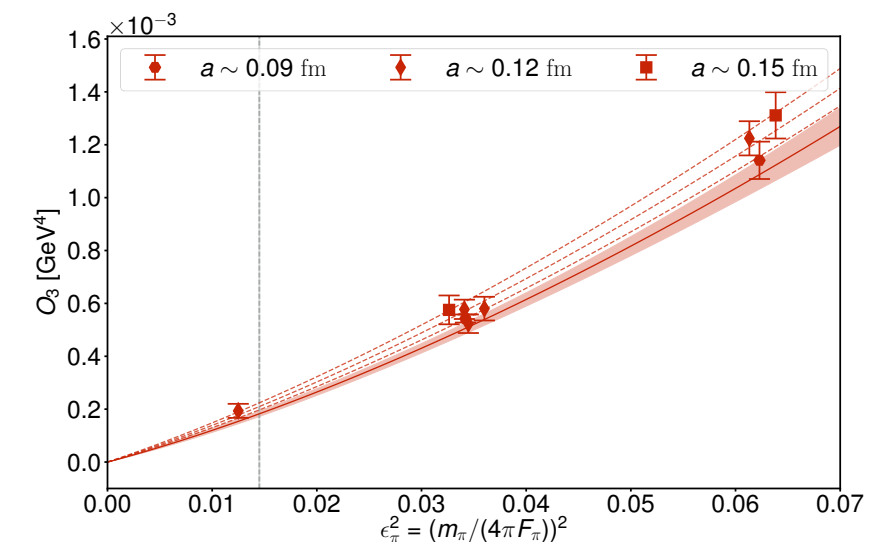
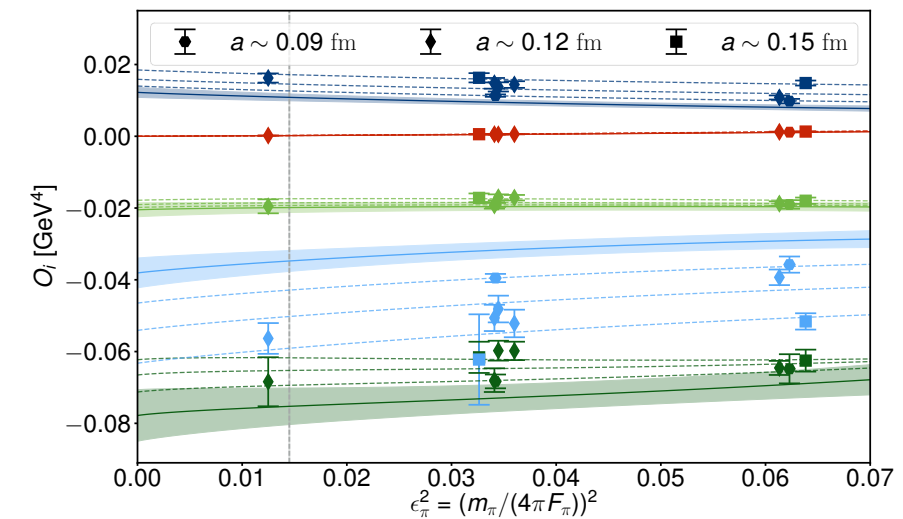
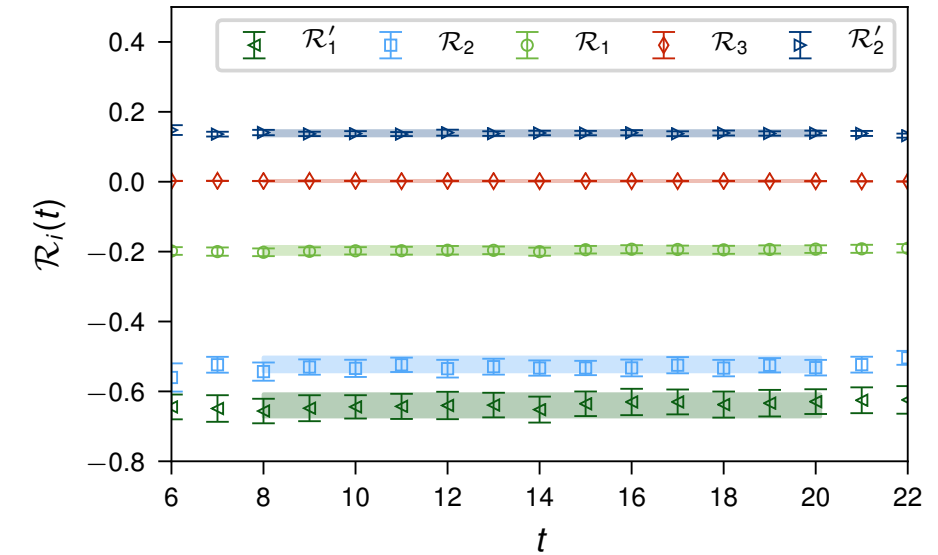
$$\mathcal{O}_{1+}^{++} = (\bar{q}_L \tau^+ \gamma_\mu q_L)(\bar{q}_R \tau^+ \gamma_\mu q_R)$$

$$\mathcal{O}_{2+}^{++} = (\bar{q}_L \tau^+ q_R)(\bar{q}_R \tau^+ q_L) + \{L \leftrightarrow R\}$$

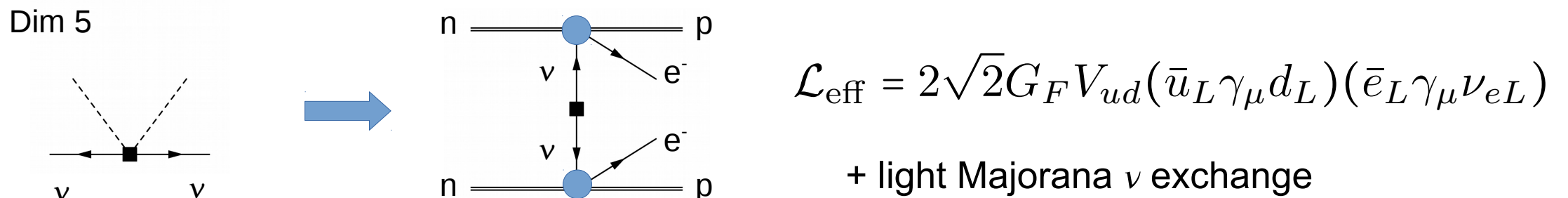
$$\mathcal{O}_{3+}^{++} = (\bar{q}_L \tau^+ \gamma_\mu q_L)(\bar{q}_L \tau^+ \gamma_\mu q_L) + \{L \leftrightarrow R\}$$

&& color-mixed combinations $\mathcal{O}_1', \mathcal{O}_2'$

Short-Range LNV $\sim (M_{\text{BSM}})^{-1}$
comparable in See-Saw models to
Long-Range LNV $\sim (\text{chirality flip}) \sim m_{\beta\beta} \sim (M_{\text{BSM}})^{-1}$



Long-Range LNV on a Lattice



Challenges in $nn \rightarrow ppe^-e^-$

- nn, pp scattering / low-lying states on present lattices
- nonlocal operator in Euclidean time:
light intermediate states $E < 2m_N$ are exp. diverging with Euc.time

First step: $\pi^- \rightarrow \pi^+ e^- e^-$

- $1 \rightarrow 1$ amplitude, large exc. gap
- component of EFT matching

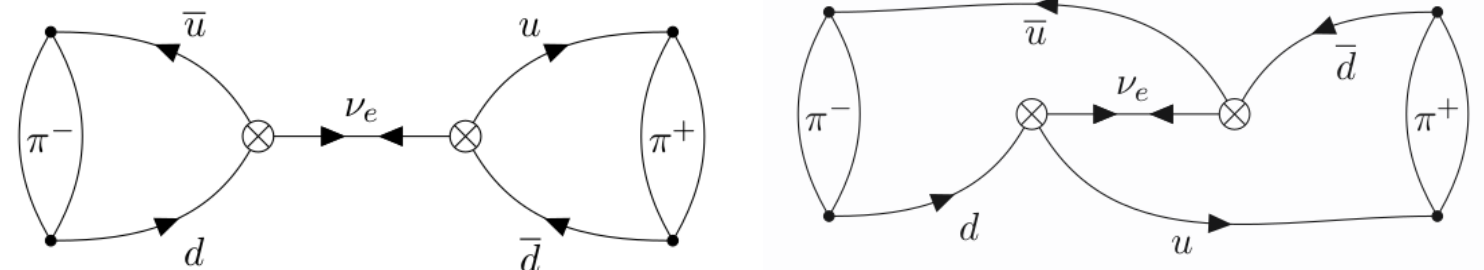


Fig. [Detmold, Murphy 2004.07404]

$$C_{\pi^- \rightarrow \pi^+ e^- e^-}(T, t_x, t_y) = \sum_{\vec{x}, \vec{y} \text{ reg. at } |x-y| \sim O(a)} S_\nu(x-y) \langle \pi^+(T) \mathcal{T} \{ j_{L\alpha}(x) j_{L\beta}(y) \} (\pi^-)^\dagger \rangle$$

Long-Range LNV: $\pi^- \rightarrow \pi^+ e^- e^-$ (1)

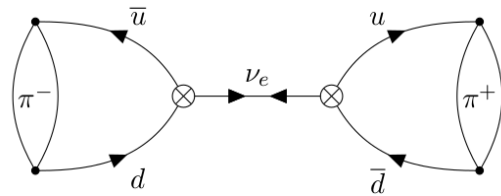
- $\langle \pi^+ e^- e^- | S_{NL} | \pi^- \rangle$ with reg. photon correlator
[Detmold, Murphy; 1811.05554; 2004.07404]

$$\sum_{\Delta \leq t_{x,y} \leq T-\Delta} \frac{C_{\pi^- \pi^+}(T, t_x, t_y)}{C_\pi(T)} \approx$$

$$\approx (T - 2\Delta) \cdot \boxed{M^{0\nu}} + O\left(\frac{e^{-(|\vec{q}| + E_k - m_\pi)T} - 1}{|\vec{q}| + E_k - m_\pi}\right)$$

~ linear fit

E_k = energy of $\langle k | j_L | \pi \rangle$ state
vac. $E_k < m_\pi$ treated separately



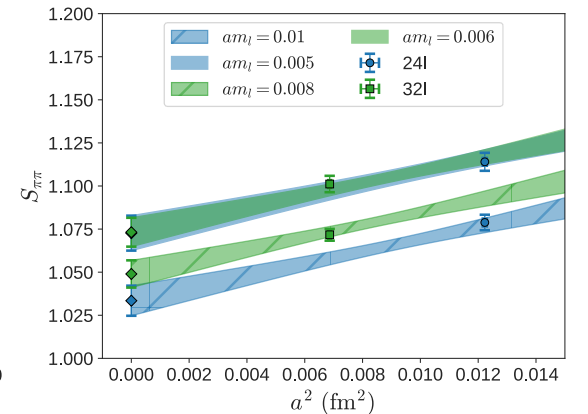
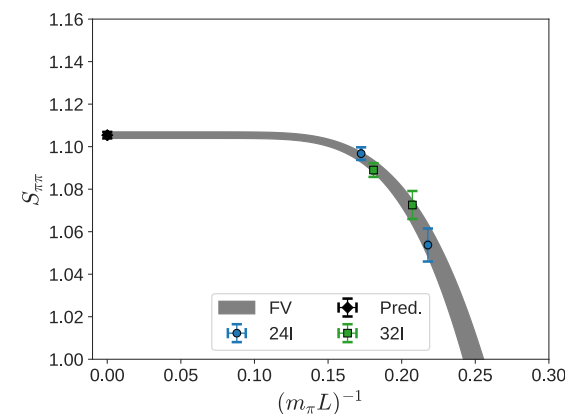
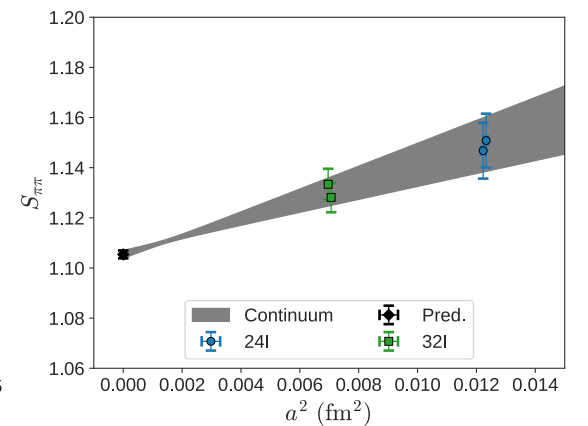
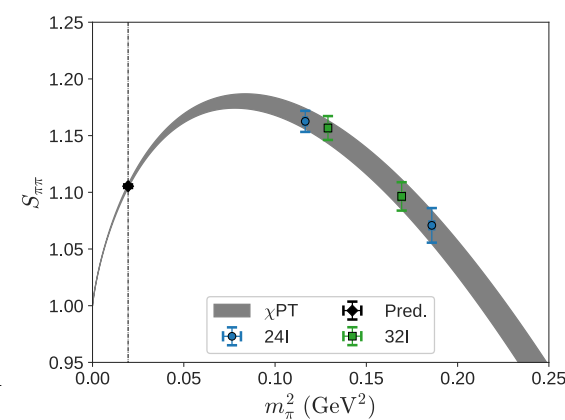
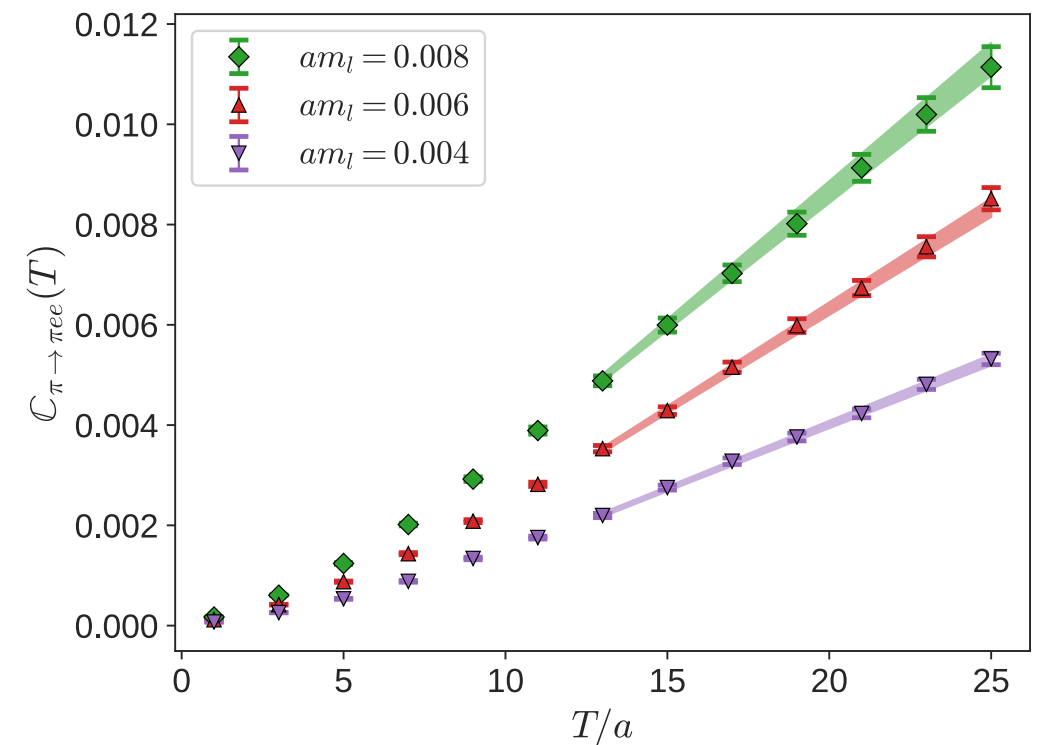
- Chiral, FVE, continuum limit fit

$$S_{\pi\pi} = 1 + \underbrace{\frac{m_\pi^2}{8\pi^2 f_\pi^2} \left(3 \log \left(\frac{\mu^2}{m_\pi^2} \right) + 6 + \frac{5}{6} g_\nu^{\pi\pi}(\mu) \right)}_{\text{NLO } \chi\text{PT}} + \underbrace{c_{FV}^{\text{NLO}} \frac{e^{-m_\pi L}}{(m_\pi L)^{3/2}}}_{\text{FV}} + \underbrace{c_a a^2}_{\text{Continuum}}$$

$$g_\nu^{\pi\pi}(770 \text{ MeV}) = -10.78(12)_{\text{stat}}(4)_{\text{fit}}(50)_{\text{FV}}(9)_{\chi\text{PT}},$$

$$S_{\pi\pi} = 1.1054(14)_{\text{stat}}(6)_{\text{fit}}(61)_{\text{FV}}(10)_{\chi\text{PT}},$$

$$M^{0\nu} = 0.01880(6)_{\text{stat}}(2)_{\text{fit}}(10)_{\text{FV}}(2)_{\chi\text{PT}} \text{ GeV}^2$$



Long-Range LNV: $\pi^- \rightarrow \pi^+ e^- e^-$ (2)

- $\langle \pi^+ e^- e^- | S_{NL} | \pi^- \rangle$ with Infinite-Volume Reconstruction
[Tuo, Feng, Jin, PRD, 1909.13525]

$$\sum_{t, \vec{x}} S(t, \vec{x}) H(t, \vec{x})$$

lattice
hadronic tensor

$$\approx \sum_{|t| < t_s, \vec{x}} S(x) H(x)$$

$$+ 2 \sum_{\vec{x}} L(t_s, \vec{x}) H(t_s, \vec{x})$$

long-range extrap.
based on 1-state dominance

- FVE, Long-range effects examined

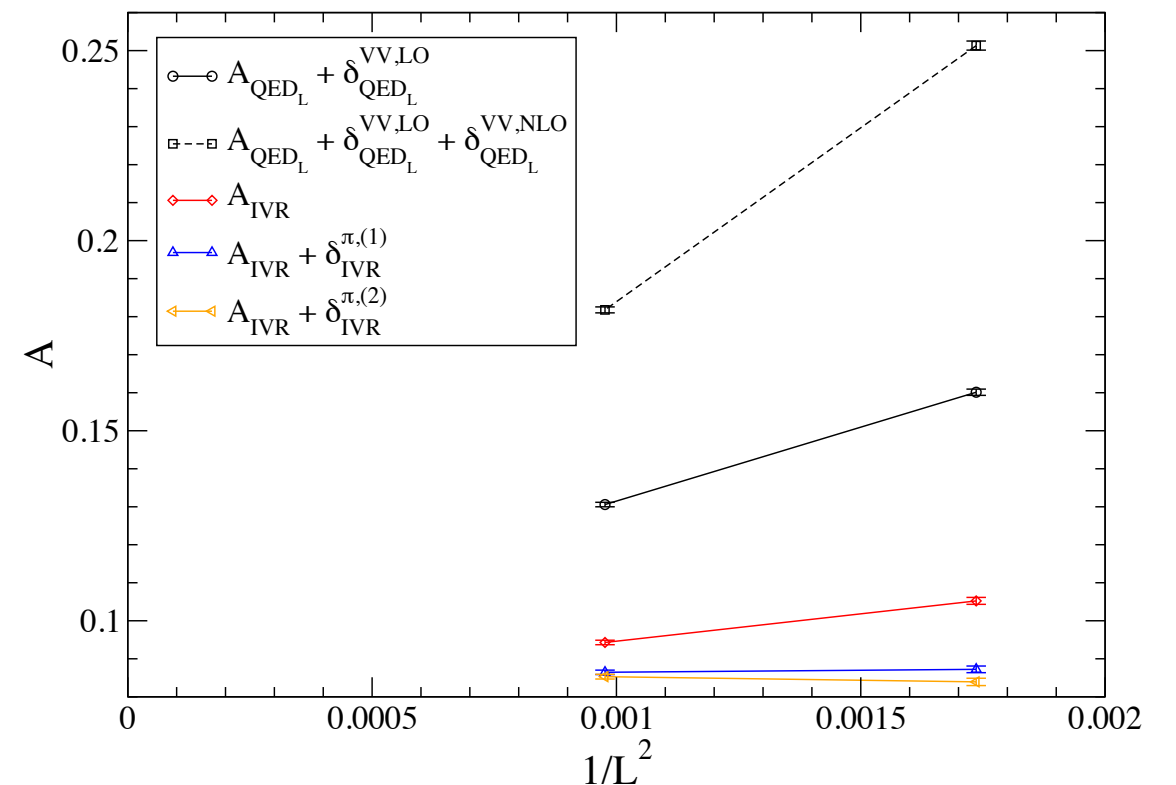
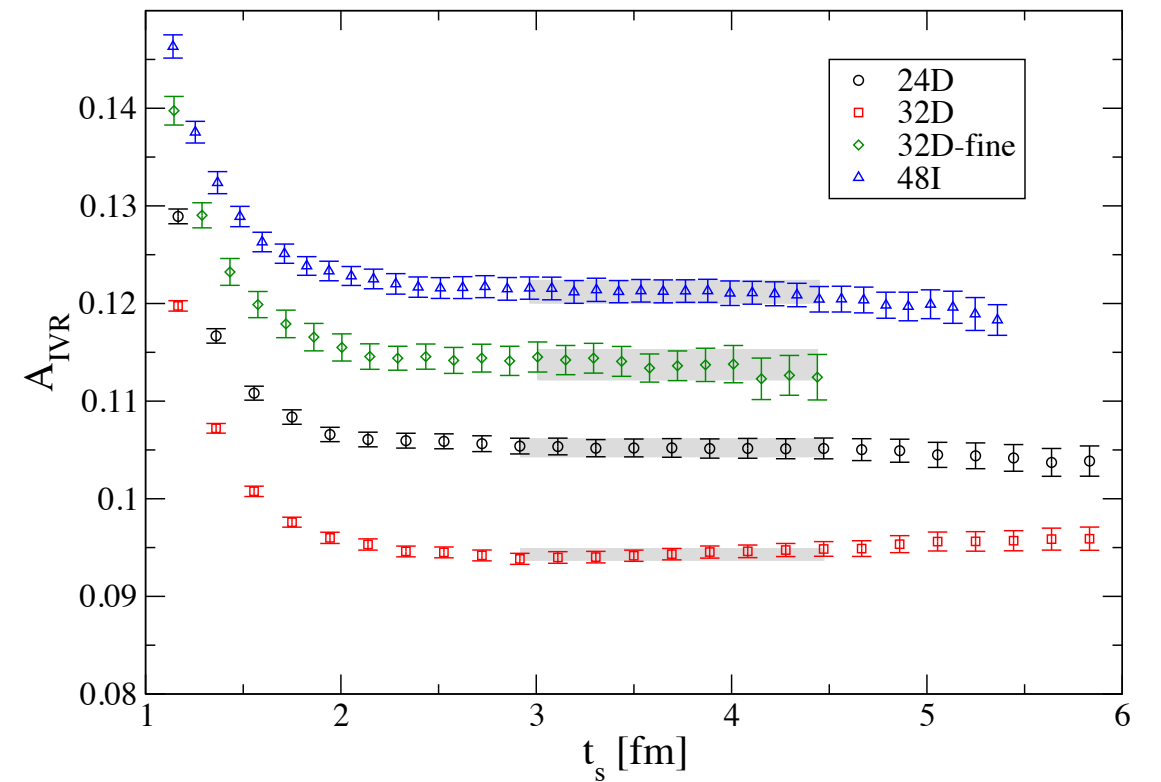
$$g_\nu^{\pi\pi}(\mu) \Big|_{\mu=m_\rho} = -10.89(28)(33)_L(66)_a$$

$$S_{\pi\pi} = 1.1045(34)_{\text{stat}}(74)_{\text{sys}}.$$

- in agreement with [Detmold, Murphy; 2004.07404]

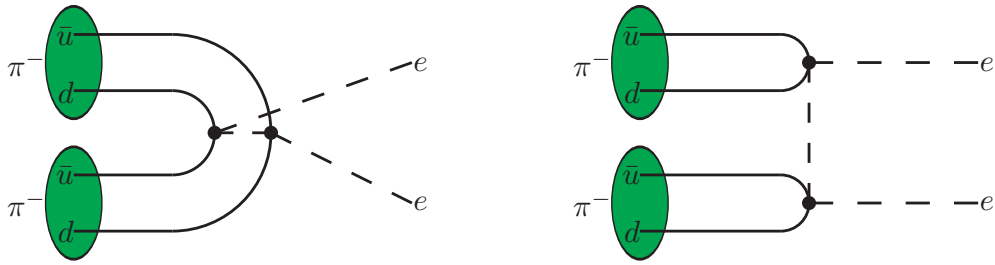
$$g_\nu^{\pi\pi}(770 \text{ MeV}) = -10.78(12)_{\text{stat}}(4)_{\text{fit}}(50)_{\text{FV}}(9)_{\chi\text{PT}},$$

$$S_{\pi\pi} = 1.1054(14)_{\text{stat}}(6)_{\text{fit}}(61)_{\text{FV}}(10)_{\chi\text{PT}},$$



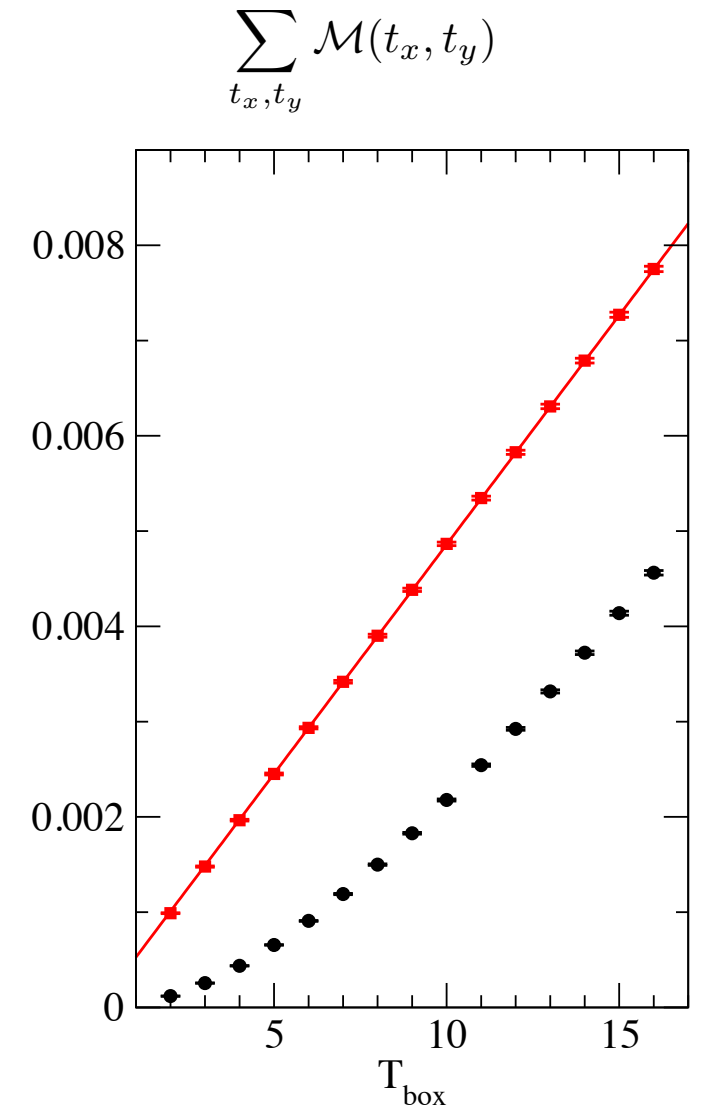
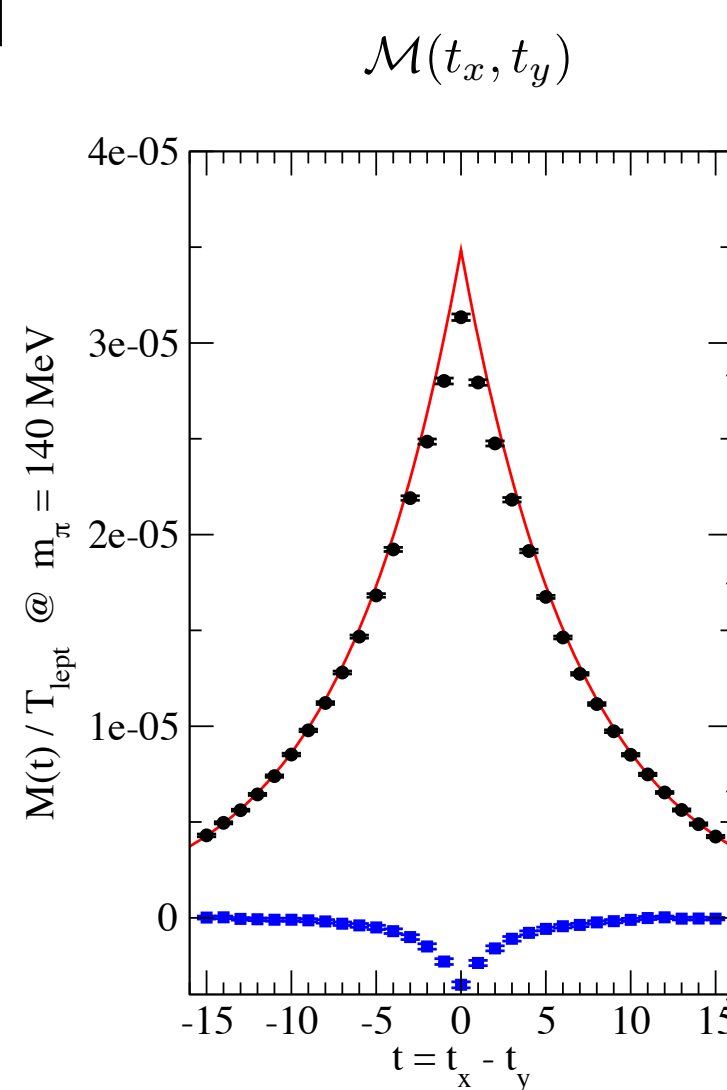
Long-Range LNV: $\pi^-\pi^-\rightarrow e^-e^-$

● $\langle ee|S_{\text{NL}}|\pi\pi\rangle$ [Feng et al, PRL, 1809.10511]



$$\mathcal{M}(T, t_x, t_y) = \frac{C_{\pi\pi\rightarrow ee}(T, t_x, t_y)}{C_{\pi\pi}(T)}$$

● (*) $\pi\bar{\nu}$ state treated analytically



$$\left. \frac{\mathcal{A}(\pi\pi \rightarrow ee)}{F_\pi^2 T_{\text{lept}}} \right|_{m_\pi=420 \text{ MeV}} = 1.517(13),$$

$$\left. \frac{\mathcal{A}(\pi\pi \rightarrow ee)}{F_\pi^2 T_{\text{lept}}} \right|_{m_\pi=140 \text{ MeV}} = 1.820(6).$$

NLO ChPT

$$g_\nu^{\pi\pi}(m_\rho) \Big|_{m_\pi=420 \text{ MeV}} = -8.50(9)_{\text{stat}}$$

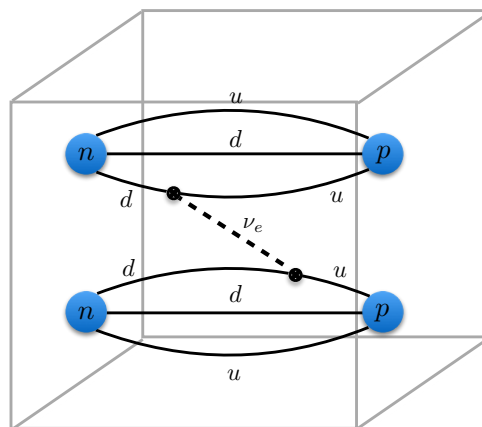
$$g_\nu^{\pi\pi}(m_\rho) \Big|_{m_\pi=140 \text{ MeV}} = -11.96(31)_{\text{stat}}$$

[Detmold, Murphy; 2004.07404]

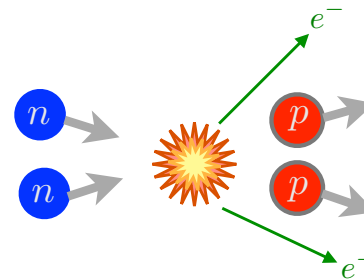
Towards $nn \rightarrow ppe^-e^-$ Amplitudes

Short-range LNV

- Methodology straightforward [R.Briceño, M.Hansen'15] analogous to NN Parity-violation [CalLat]
- Main computing challenge: nn , pp spectroscopy & scattering states on a lattice with physical quarks



?



[Z.Davoudi, ECT* '19]

2-body amplitude of local interaction:

$$|\mathcal{A}_{V \rightarrow \infty}|^2 \propto \left(q \frac{\partial \phi(q)}{\partial q} + k \frac{\partial \delta(k)}{\partial k} \right) |\mathcal{A}_{V=L^3}|^2$$

$\delta(k)$ = scattering phase

$\phi(q)$ = lattice zeta fcn.

[Lellouch, Luscher, '01]

Long-range LNV : complicated by bi-locality of $nn \rightarrow pp$

- Can be matched to pion-less EFT \Rightarrow short-range $(g_\nu)^{NN}$ [Z.Davoudi, S.V.Kadam, 2012.02083]
- On-shell intermediate states : may need addl. calculations to constrain $(n)n \rightarrow (n)p\nu$ amplitudes

Matching to EFTs beyond LO : combination of $nn \rightarrow pp$, $\pi^+ \rightarrow \pi^-$, $n \rightarrow p\pi$ LECs

SUMMARY

- Nucleon charges & Dark Matter scattering
Remarkable progress sys/stat precision for single nucleons
- Nucleon/nuclear Electric Dipole Moments
 θ_{QCD} is challenging at the physical point
Higher-dim operators: renormalization&mixing tractable
Nuclear EDMs: $CP\nu$ πN , NN coupling
- Baryon number violation operators
Single-nucleon decay/oscillation straightforward
Annihilation, in-medium effects challenging as $0\nu 2\beta$
- Neutrinoless double-beta decay
Depend critically on
 - progress in NN spectrum & scattering
 - Effective many-body Theory