Energy frontier probes of flavor specific scalar mediators

Mudit Rai (mur4@pitt.edu)

University of Pittsburgh

Collaborators : Brian Batell, Ayres Freitas, Ahmed Ismail & David Mckeen Phys. Rev. D **104**, 115032

Motivation

- New light scalar singlets feature prominently in SM extensions, e.g mediators to Dark sector.
- Apart from the Yukawa couplings, the new scalar-fermionic couplings tend to break the flavor symmetry, leading to possibility of new large FCNC's.
- A standard way to evade is via the MFV hypothesis, with new couplings $\propto Y_{u}, Y_{d}$.
- The flavor-specific hypothesis takes a different route by having couplings to only one flavor in the mass basis.

Motivation

- This is a technically natural, radiatively stable hypothesis
- We will explore two UV complete scenarios : VLQ & Heavy Higgs-like scalar.
- Focusing on an up-quark-specific model, we find that naturalness and experimental constraints in the UV theories are complementary to those in the EFT[1,2].
- In these models, the low energy and high energy frontiers can be linked via the interpolation of experimental constraints both ways.

EFT Review

Light scalar S with flavor-specific couplings :

$$\mathcal{L} \supset iar{\chi}(
ot\!\!/ D - m_\chi)\chi + rac{1}{2}\partial_\mu S\partial^\mu S - rac{1}{2}m_S^2S^2 - \left(g_\chi Sar{\chi}_L\chi_R + rac{c_S}{M}Sar{Q}_L U_R H_c + ext{h.c.}
ight)$$

In the up-specific hypothesis, the effective scalar up quark coupling is :

$$\mathcal{L}_S \supset -g_u S \bar{u} u, \qquad g_u = rac{c_S v}{\sqrt{2}M}$$

 EFT has implications for the naturalness of the light singlet scalar, flavor violation, and CP violation[1,2].

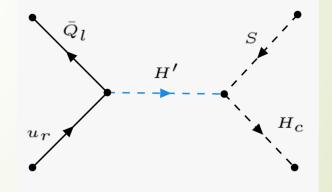
Scalar Completion: Model

We introduced a Higgs like scalar

$$\mathcal{L}_{\rm sd} = \mathcal{L}_{\rm SM} + \frac{1}{2} \partial_{\mu} S \partial^{\mu} S - \frac{1}{2} m_S^2 S^2 + (D_{\mu} H')^{\dagger} D^{\mu} H' - M^2 H'^{\dagger} H' - \left[y'_i^j \overline{Q_L^i} \, u_{Rj} \, H'_c + \kappa M \, S \, H^{\dagger} H' + \text{h.c.} \right] + \text{quartic scalar couplings},$$

The effective dim-5 operator would be :

$$\mathcal{L} \supset rac{\kappa \, {y'}_i^j}{M} S \, \overline{Q}_L^i \, u_{R\,j} \, H_c + ext{h.c.}$$



Model

We will rotate H, H' to the Higgs basis :

$$\hat{H} = \begin{pmatrix} G^+ \\ \frac{1}{\sqrt{2}}(v + \phi_1 + iG^0) \end{pmatrix}, \qquad \hat{H}' = \begin{pmatrix} H^+ \\ \frac{1}{\sqrt{2}}(\phi_2 + iA^0) \end{pmatrix}, \qquad S = v_S + \phi_3,$$

• The mixing angle is
$$\tan \beta = \frac{\nu'}{\nu_0} \ll 1$$

• Diagonalizing the mixed CP-even scalar fields φ_i will lead to mass eigenstates:

$$R^T\mathcal{M}_\phi^2R= ext{diag}\{m_h^2,m_{h'}^2,m_s^2\},$$

The CP even scalar masses are :

$$m_h^2 \simeq 2\lambda v^2, \qquad m_{h'}^2 \simeq M^2, \qquad m_s^2 \simeq m_S^2 - \kappa^2 v^2$$

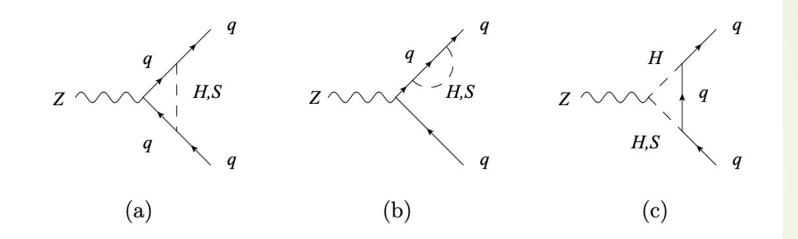
Decays

- The charged Goldstones A⁰ and H⁺ are approx. degenerate: $m_{A^0,H^{\pm}}^2 = M^2 \cos^2 \beta + (\mu'^2 - \kappa M v_s) \sin(2\beta) - \mu^2 \sin^4 \beta \approx M^2$
- The leading decays of heavier scalar are :

$$\begin{split} \Gamma(h' \to u \bar{u}) &= \Gamma(A^0 \to u \bar{u}) = \Gamma(H^+ \to u \bar{d}) \simeq \frac{3y'^2 M}{16\pi}, \\ \Gamma(h' \to sh) &= \Gamma(A^0 \to sZ) = \Gamma(H^+ \to sW^+) \simeq \frac{\kappa^2 M}{16\pi}. \end{split}$$

The decay for light scalar into an invisible DM will happen for $m_s > 2m_{\chi}$.

Scalar Phenomenology : EWPT

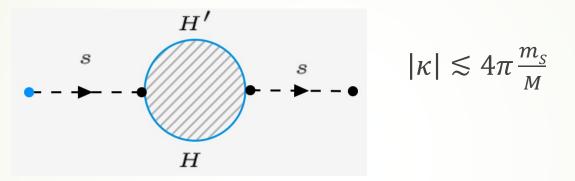


• EW precision bounds : Fixing M = 1 TeV, ms = 1 GeV,

- $R_l^{exp} R_l^{SM} = 0.83 \ (y' = \kappa = \sqrt{4\pi}), \ (\text{LEP}) \text{ excluded by current data} (\delta R_l = 0.034 \pm 0.025)$
- $R_l^{exp} R_l^{SM} = 5.5 \times 10^{-3} (y' = \kappa = 1)$, FCC-ee : expected $\delta R_l = 0.001$.

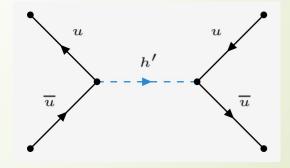
Scalar Phenomenology

Naturalness considerations :



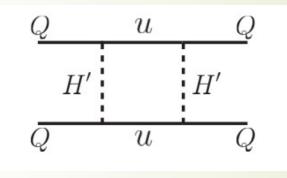
Another limit arises if we integrate out H' and put Higgs to its vev: $|\kappa| \lesssim \frac{m_s}{m_s}$

 Collider Phenomenology:
 Can be observed at LHC as Dijet resonances.



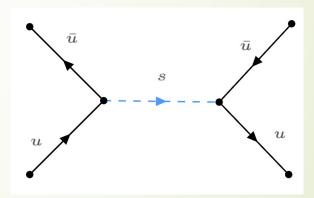
Scalar Phenomenology

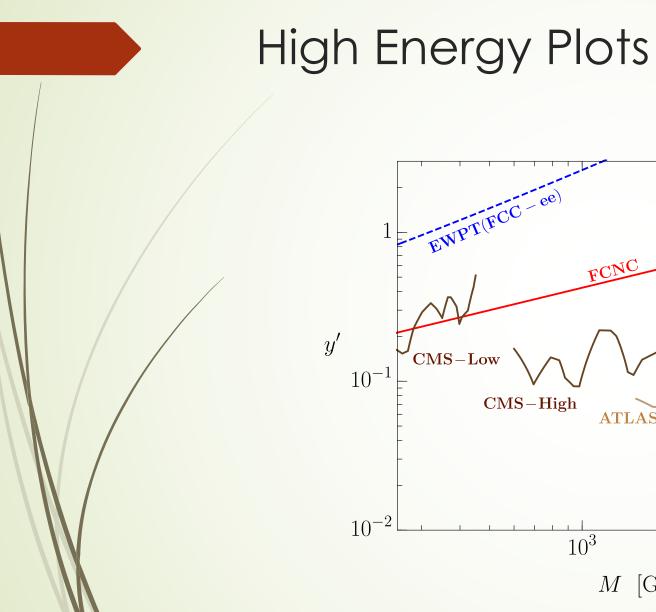
FCNC considerations (K⁰ system) : $y \lesssim 0.6 \left(\frac{M}{2 \,\text{TeV}}\right)^{1/2}$.

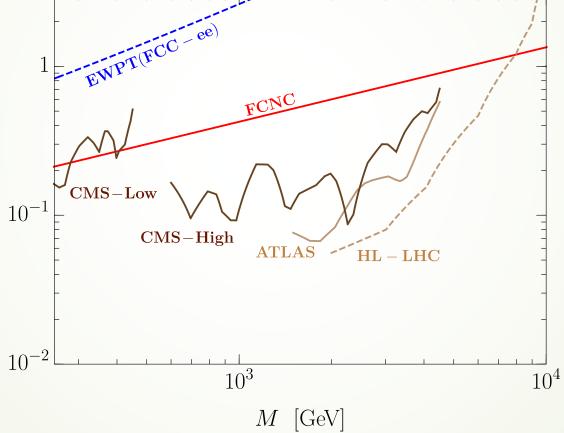


Neutron EDM from 4-up operator:

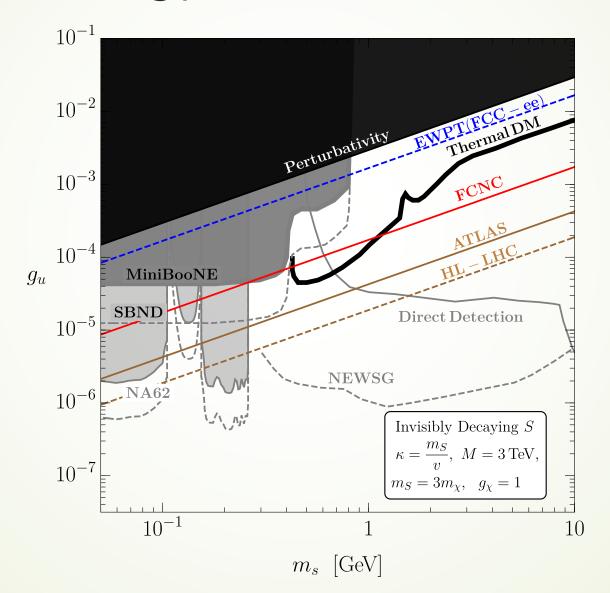
$$|g_u|\sqrt{\sin 2\phi_{\rm CP}} < 3 \times 10^{-6} \left(\frac{m_S}{1\,{\rm GeV}}\right)$$

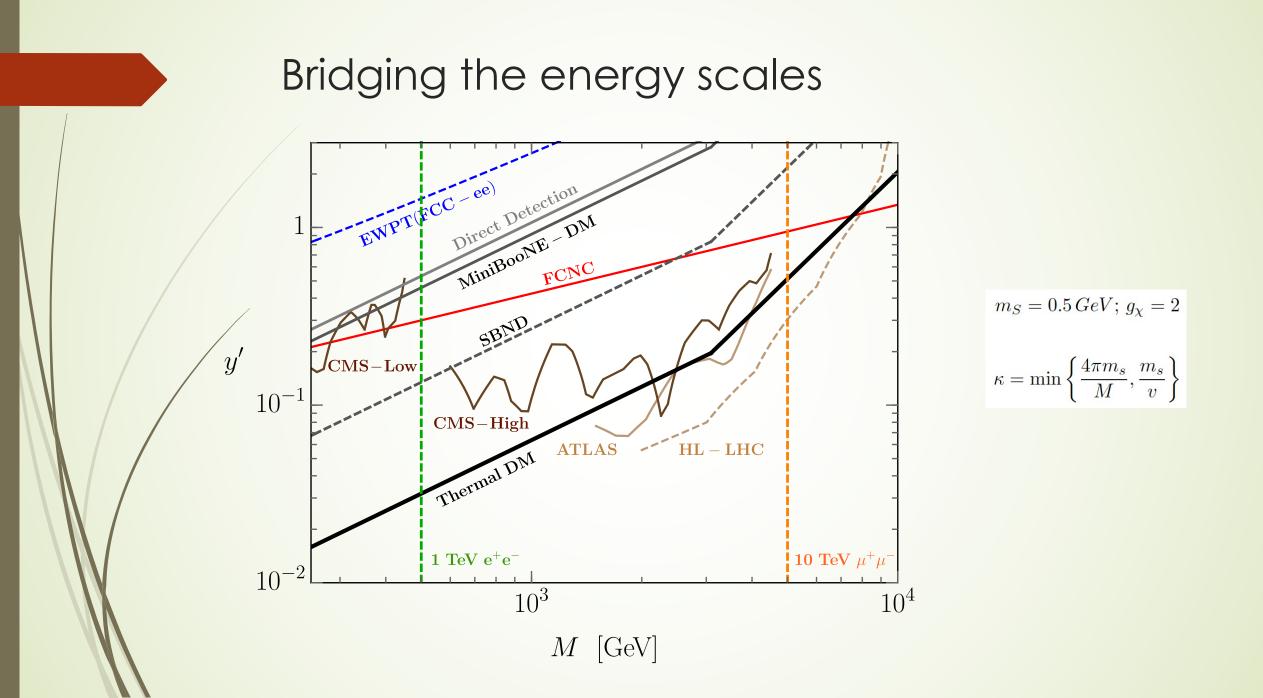






Low Energy DM Plot





Conclusions

- Light dark sectors are a particularly interesting realm of contemporary BSM phenomenology with promising precision, beam dump, and direct detection experiments on the horizon.
- The up-specific models provides an interesting complementary benchmark to Higgs-like scalars.
- Flavour-specific hypothesis can be applied easily to any of the quarks with minor modifications.
- UV completion of the previously studied EFT gives a wider picture and complimentary constraints on the parameter space.

Thank you!

VLQ completion : Model

• Consider a new RH-up quark like $U'(3,1,\frac{2}{3})$ s.t.:

$$-\mathcal{L} \supset M \,\overline{U}'_L U'_R + y_i \,\overline{Q}^i_L U'_R \,H_c + \lambda^i \,\overline{U}'_L u_{R\,i} \,S + \text{h.c.}$$

Integrating out VLQ gives :

$$-\mathcal{L} \supset rac{y_i \, \lambda^j}{M} S \, \overline{Q}_L^i \, u_{Rj} \, H_c + ext{h.c.}$$

$$(c_S)_i^j \equiv -y_i \, \lambda^j$$

After EWSB, there's a mass mixing b/w {u, U'} which can be diagonalized in the regime of {vy_u, λvs} ≪ yv < M via :</p>

$$u_L o \cos \theta \, u_L + \sin \theta \, U'_L, \qquad U'_L o \cos \theta \, U'_L - \sin \theta \, u_L,$$

 $\cos \theta = \frac{M}{M_{U'}}, \qquad \qquad \sin \theta = \frac{yv}{\sqrt{2} \, M_{U'}}.$

where $M_{U'} = \sqrt{M^2 + (yv)^2/2}$ is the physical mass of VLQ.

Appendix : VLQ : Decays

The decay widths for the VLQ are :

$$\begin{split} \Gamma(U' \to uS) &= \cos^2 \theta \; \frac{\lambda^2 \, m_{U'}}{32 \, \pi} \left(1 - \frac{m_S^2}{m_{U'}^2} \right)^2 \simeq \frac{\lambda^2 M}{32 \, \pi}, \\ \Gamma(U' \to uh) &= \sin^2 \theta \cos^2 \theta \; \frac{G_F \, m_{U'}^3}{16 \sqrt{2} \, \pi} \left(1 - \frac{m_h^2}{m_{U'}^2} \right)^2 \simeq \frac{y^2 M}{64 \, \pi}, \\ \Gamma(U' \to uZ) &= \sin^2 \theta \cos^2 \theta \; \frac{G_F \, m_{U'}^3}{16 \sqrt{2} \, \pi} \left(1 - \frac{m_Z^2}{m_{U'}^2} \right)^2 \left(1 + \frac{2m_Z^2}{m_{U'}^2} \right) \simeq \frac{y^2 M}{64 \, \pi}, \\ \Gamma(U' \to dW) &= \sin^2 \theta \; \frac{G_F \, m_{U'}^3}{8 \sqrt{2} \, \pi} \left(1 - \frac{m_W^2}{m_{U'}^2} \right)^2 \left(1 + \frac{2m_W^2}{m_{U'}^2} \right) \simeq \frac{y^2 M}{32 \, \pi}, \end{split}$$

The decay width for light scalar S :

$$\Gamma(S \to u\bar{u}) = \sin^2\theta \, \frac{\lambda^2 \, m_S}{8\pi} \simeq \frac{g_u^2 m_S}{8\pi}$$

Appendix : VLQ : Naturalness

- Naturalness considerations : From radiative sizes of terms generated by S,H up and U' interactions.
- Correction to scalar mass term at 1 loop :

$$\delta m_S^2 \sim \frac{\operatorname{Tr} \lambda^* \lambda}{16\pi^2} M^2 \implies \lambda^i \lesssim 4\pi \frac{m_S}{M}$$

Correction to Higgs mass at 1 loop :

$$\delta m_H^2 \sim \frac{\operatorname{Tr} y y^*}{16\pi^2} M^2 \quad \Rightarrow \quad y_i \lesssim 4\pi \frac{v}{M}$$

These two leads to an Naturalness bound on the EFT coupling :

$$g_u \lesssim \frac{16\pi^2}{\sqrt{2}} \, \frac{m_S v}{M^2} \approx (7 \times 10^{-4}) \left(\frac{m_S}{0.1 \,\text{GeV}}\right) \left(\frac{2 \,\text{TeV}}{M}\right)^2.$$

Appendix : VLQ : CKM considerations

There exists a tension between the SM theory and unitarity prediction for the top row CKM unitarity ("Cabbibo anomaly").

Current experimental bounds gives[3]:

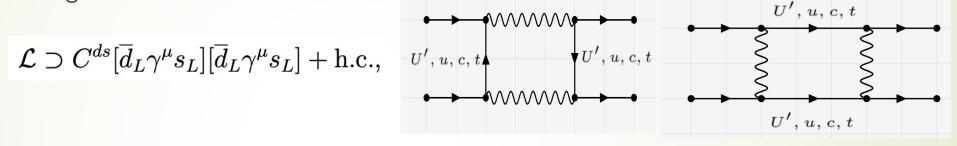
$$\left[\left. |V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 \right] \right|_{\exp} = 0.9985(3)_{V_{ud}}(4)_{V_{us}}$$

Requiring theory prediction to be within 3 σ gives sin $\theta \leq 0.055$, implying:

$$y \lesssim 0.6 \left(\frac{M}{2\,{\rm TeV}}\right). \label{eq:y_started}$$

Appendix : VLQ - FCNC bounds

FCNC considerations comes from the modification to Neutral Kaon mixing box diagrams :



- We get,
$$C^{ds} = -y^4 |V_{ud}^* V_{us}|^2/(128\pi^2 M^2)$$

• Current limits restrict : $\operatorname{Re}[C^{ds}] \lesssim (10^3 \,\mathrm{TeV})^{-2}$

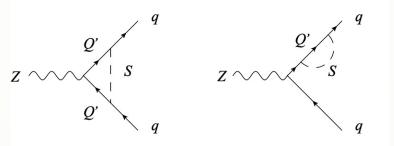
This can be translated as :

$$y \lesssim 0.6 \left(\frac{M}{2 \,\mathrm{TeV}}\right)^{1/2}.$$

Appendix : VLQ : EW Precision bounds

Heavy VLQ modifies the partial width of Z, R_l = $\frac{\Gamma[Z \rightarrow \text{had.}]}{\Gamma[Z \rightarrow \ell^+ \ell^-]}$.
 1. Tree-level shift through u-U' mixing is dominant.

2. Loops:



• Current data(LEP) is $\delta R_l^{exp} = 0.034 \pm 0.025$ leading to $\frac{yv}{M} < 0.063$.

• Future data(FCC-ee) will give $\delta R_l^{exp} = 0.001$ leading to $\frac{yv}{M} < 0.022$.

Appendix : VLQ: EDM bounds

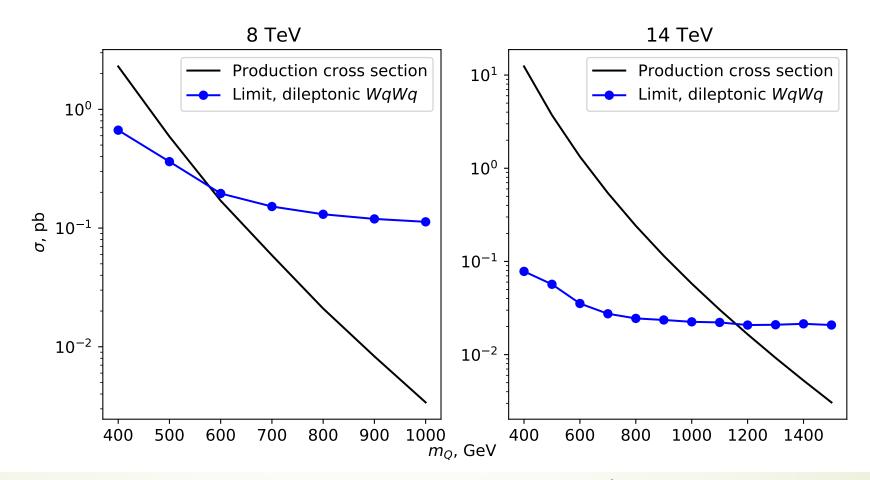
For complex M, y & λ , large nEDM can arise from effective CPV 4-quark operator :

$$\mathcal{L} \supset C'_u \,\overline{u} i \gamma^5 u \,\overline{u} u, \qquad C'_u = \frac{\operatorname{Re}(Y_{S \overline{u} u}) \operatorname{Im}(Y_{S \overline{u} u})}{m_S^2} \simeq -\frac{y^2 \lambda^2 v^2}{4M^2 m_S^2} \sin 2\phi_{\rm CP}$$

- Neutron EDM, in this terms gives, $d_n = 0.182 \, e \, C'_u \, {
 m GeV}$
- Experimentally, we have $|d_n| < 1.8 imes 10^{-26} e ~{
 m cm}$, thus leading to :

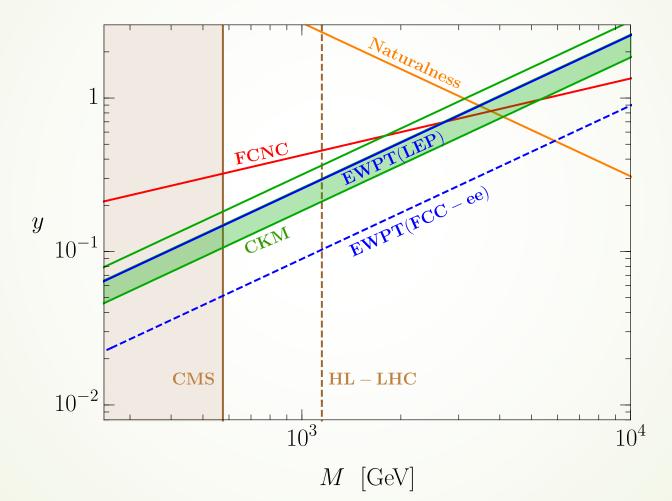
$$|g_u|\sqrt{\sin 2\phi_{\rm CP}} < 3 \times 10^{-6} \left(\frac{m_S}{1\,{\rm GeV}}\right)$$

VLQ : Collider Phenomenology

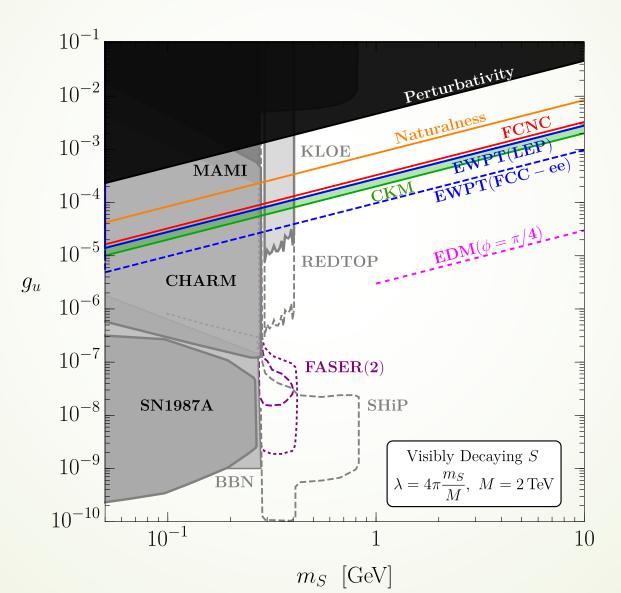


We consider pair production of U' and its decays $U' \rightarrow dW$. Assuming 20 fb⁻¹ at 8 TeV, the constraints is M > 575 GeV. At 14 TeV with 3000 fb⁻¹ luminosity, we get a constraint M > 1150 GeV. Analysis is close to that done by CMS.

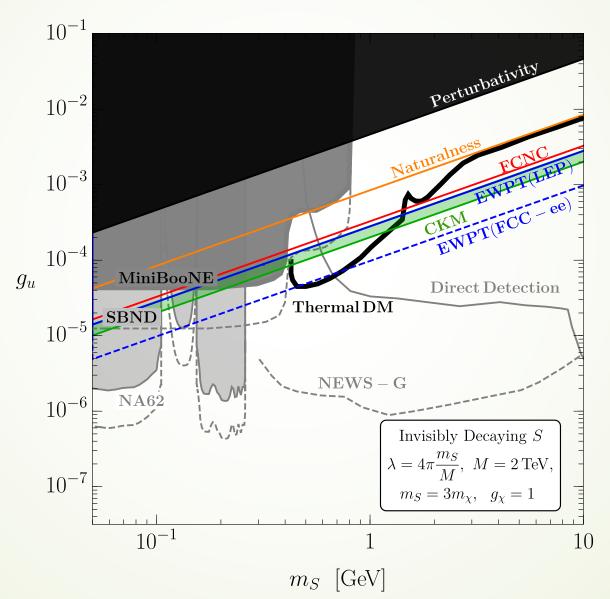
VLQ: Results (y vs M)



VLQ : Results (Visible Decay)



VLQ : Results (Invisible Decay)



EFT Structure Basics

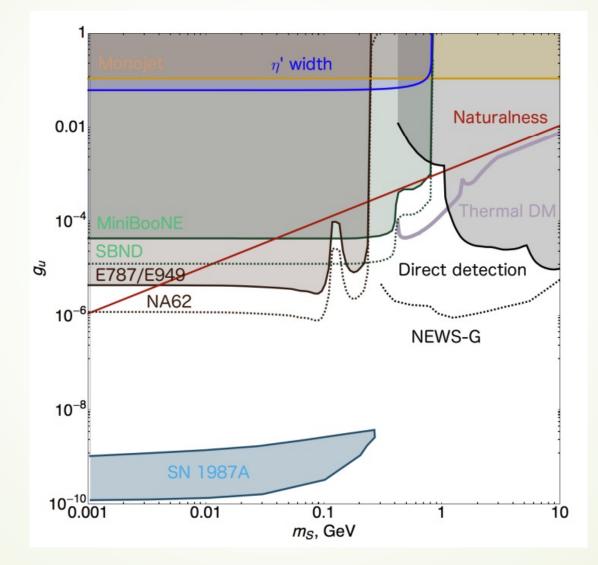
$$\mathcal{L}_S = \frac{1}{2} \partial_\mu S \partial^\mu S - \frac{1}{2} m_S^2 S^2 - \left(\frac{c_S}{M} S \bar{Q}_L U_R H_c + \text{h.c.}\right)$$

- In Flavor hypothesis, we have
 $c_S \propto \operatorname{diag}(1,0,0)$
 - This breaks the flavor symmetry as

 $U(3)_Q \times U(3)_U \to U(1)_u \times U(2)_{ctL} \times U(2)_{ctR}.$

In our flavor-specific scenario, in addition to the Yukawa couplings, we are adding a new flavor-breaking spurion c_s , and assume that it is aligned with the Yukawa coupling Y_u .

EFT Dark Sector Plot [2]



References

- Brian Batell, Ayres Freitas, Ahmed Ismail & David Mckeen, <u>1712.10022</u> [hepph]
- 2. Brian Batell, Ayres Freitas, Ahmed Ismail & David Mckeen, <u>1812.05103</u> [hepph]
- 3. P. A. Zyla et al. (Particle Data Group), PTEP 2020, 083C01 (2020).