Neutrinos in cosmology:

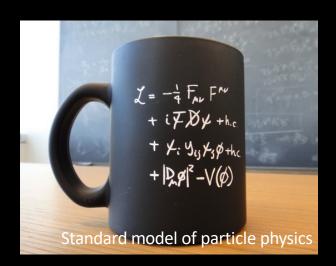
A match made in the Heavens

Yvonne Y. Y. Wong, UNSW Sydney

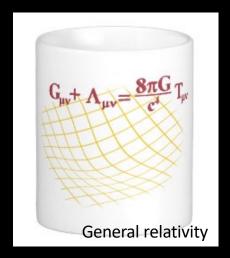
Snowmass neutrino colloquium IV, May 4, 2022

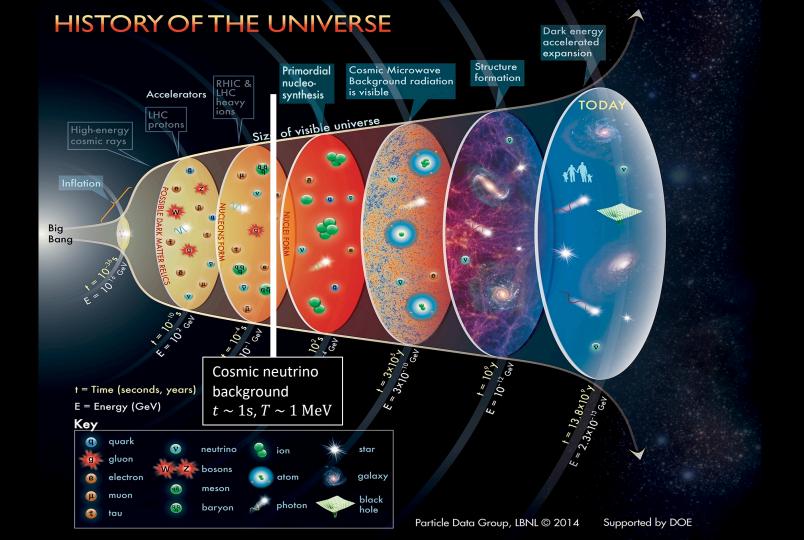
An unlikely partnership?

Neutrino = one of the lightest and most weakly-interacting known particles



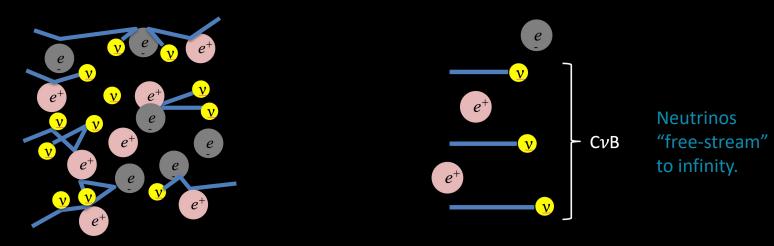
Cosmology = gravitation on the largest observable scales





Formation of the CνB...

The CvB is formed when neutrinos decouple from the cosmic plasma.

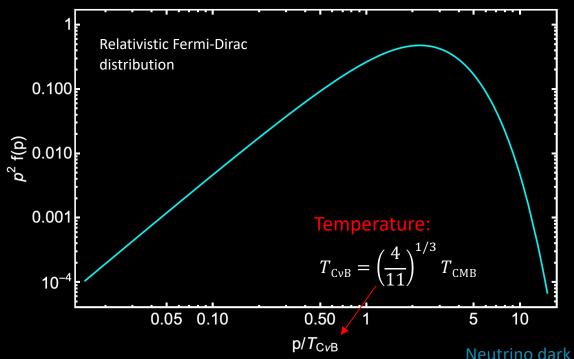


Above $T \sim 1$ MeV, even weakly-interacting neutrinos can be produced, scatter off e^+e^- and other neutrinos, and attain thermodynamic equilibrium

Below $T \sim 1$ MeV, expansion dilutes plasma, and reduces interaction rate: the universe becomes transparent to neutrinos.

The cosmic neutrino background...

Standard model predictions



Per family of Number density: +antineutrinos

$$n_{\text{CvB}} \simeq 110 \text{ cm}^{-3}$$

Energy density:

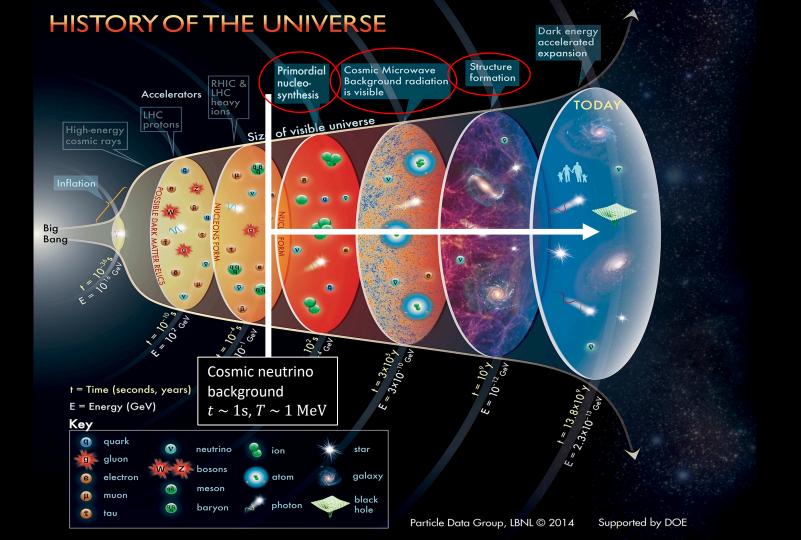
Relativistic (if $T_{\text{CVB}} \gg m_{\nu}$):

$$\rho_{\text{CvB}} \simeq \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} \rho_{\text{CMB}}$$

Non-rel (if $T_{\text{CVB}} \ll m_{\nu}$):

$$\Omega_{\text{CvB}} \simeq \frac{m_{\nu}}{93 \, h^2 \, \text{eV}}$$

Neutrino dark matter



What can cosmology do for neutrino physics?

Precision cosmological observations allow us to infer the properties of the cosmic neutrino background, from which to determine:

- Absolute neutrino mass scale, $\sum m_{\nu}$
- Number of neutrino families, $N_{\rm eff}$
 - Deviations from SM prediction of $N_{\rm eff} \approx 3$
 - e.g., test for the existence of light sterile states

"Standard" tests (even a raison d'être)

- Neutrino decay/lifetime, τ_0
- Non-standard neutrino interactions
 - Self, neutrino-dark matter, neutrino-dark energy
 - **—** ..

More exotic, but of growing interest

From the theoretical perspective:

- Origin of dark matter = keV sterile neutrinos as a dark matter candidate
- Origin of the matter-antimatter asymmetry = leptogenesis linked to neutrino mass generation

More directly, neutrino experiments can also help to pin down parameters of the $C\nu B$.

Allow us to gain more precise and accurate information about the other stuff in the universe.

1. What can cosmology do for neutrino physics?

Cosmological observables...

+ Supernova Ia, local H₀, etc.

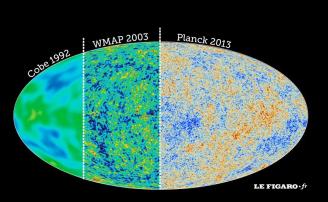
(No direct neutrino effects)

Light element abundances from primordial nucleosynthesis

Exoergic Direction $(0,\gamma)$ $(0,\gamma)$

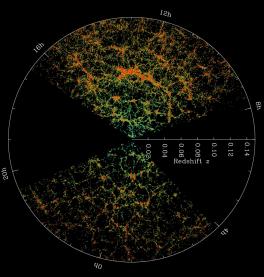
 $N_{\rm eff}$ (expansion rate)

Cosmic microwave background anisotropies



N_{eff} (expansion rate) Interactions (free-streaming) Lifetime (free-streaming)

Large-scale matter distribution

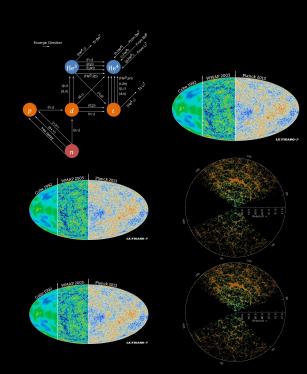


 $\sum m_{
u}$ (perturbation growth)

What do these probes really probe?

They may look different, but ultimately the information contained is

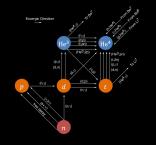
- Universal expansion rate at different times
 - How much matter, radiation, "in-between" (e.g., neutrinos), vacuum energy, etc.
- Growth of fluctuations under gravity
 - Kinematic properties and interactions of the various types of stuff in the universe; good for neutrino physics
- Distance measurements
 - Spatial geometry, dark energy; not directly relevant for neutrino physics but has indirect effects on inference

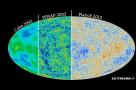


What do these probes really probe?

They may look different, but ultimately the information contained is

- Universal expansion rate at different times
 - How much matter, radiation, "in-between" (e.g., neutrinos), vacuum energy, etc.





Neutrinos & the expansion rate...

The Hubble expansion rate depends on the energy content of the universe:

$$H^2(a(t)) = H_0^2(\Omega_m a^{-3} + \Omega_r a^{-4} + \Omega_\Lambda + \Omega_k a^{-2} + \cdots)$$
 Scale factor Matter Radiation Cosmological constant curvature Reutrinos = radiation at early times

Neutrinos = radiation at early times = matter at late times

Neutrinos & the expansion rate...

The Hubble expansion rate depends on the energy content of the universe:

$$H^2(a(t)) = H_0^2(\Omega_m a^{-3} + \Omega_r a^{-4} + \Omega_\Lambda + \Omega_k a^{-2} + \cdots)$$
 Scale factor Matter Radiation Cosmological Spatial constant curvature
$$\rho_{\rm CMB} + \sum \rho_{\rm C\nu B} = \left[1 + N_{\rm eff} \times \frac{7}{8} \left(\frac{4}{11}\right)^{4/3}\right] \rho_{\rm CMB}$$

$$N_{\rm eff}^{\rm SM} = 3.0440 \pm 0.0002 \quad \text{For 3 SM families, includes}$$

Bennett et al. 2020, 2021;

Froustey, Pitrou & Volpe, 2020

For 3 SM families, includes m_e/T corrections, non-instantaneous decoupling, finite-temperature QED, and neutrino oscillations.

Neutrinos & the expansion rate...

The Hubble expansion rate depends on the energy content of the universe:

$$H^2(a(t)) = H_0^2(\Omega_m a^{-3} + \Omega_r a^{-4} + \Omega_\Lambda + \Omega_k a^{-2} + \cdots)$$
 Scale factor Matter Radiation Cosmological constant curvature
$$\rho_{\text{Other}} + \rho_{\text{CMB}} + \sum \rho_{\text{CvB}} = \left[1 + N_{\text{eff}} \times \frac{7}{8} \left(\frac{4}{11}\right)^{4/3}\right] \rho_{\text{CMB}}$$

Any relativistic, feebly-interacting, thermalised particle species will look like a neutrino cosmologically, e.g., light sterile neutrinos, thermal axions, etc.

 $N_{\rm eff}^{\rm SM} = 3.0440 \pm 0.0002$ Bennett et al, 2020, 2021; Froustey, Pitrou & Volpe, 2020

For 3 SM families, includes m_e/T corrections, non-instantaneous decoupling, finite-temperature QED, and neutrino oscillations.

Nucleosynthesis & N_{eff}...

Constraining $N_{\rm eff}$ with the primordial elemental abundances has a long history.

COSMOLOGICAL LIMITS TO THE NUMBER OF MASSIVE LEPTONS

Gary STEIGMAN

National Radio Astronomy Observatory and Yale University, USA

David N. SCHRAMM

University of Chicago, Enrico Fermi Institute (LASR), 933 E 56th, Chicago, Ill. 60637, USA

James E. GUNN

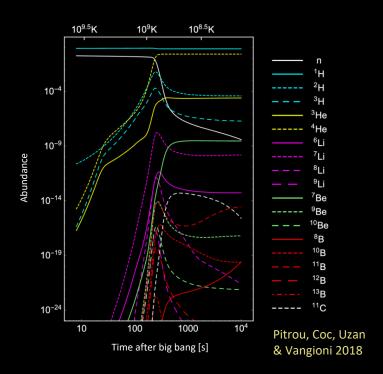
University of Chicago and California Institute of Technology, USA

Received 29 November 1976

If massive leptons exist, their associated neutrinos would have been copiously produced in the early stages of the hot, big bang cosmology. These neutrinos would have contributed to the total energy density and would have had the effect of speeding up the expansion of the universe. The effect of the speed-up on primordial nucleosynthesis is to produce a higher abundance of 4He. It is shown that observational limits to the primordial abundance of 4He lead to



the constraint that the total number of types of heavy lepton must be less than or equal to 5.



How much of these elements is produced depends on how fast the universe expands.

Nucleosynthesis & N_{eff}...

Constraining $N_{\rm eff}$ with the primordial elemental abundances has a long history.

COSMOLOGICAL LIMITS TO THE NUMBER OF MASSIVE LEPTONS

Gary STEIGMAN

National Radio Astronomy Observatory and Yale University, USA

David N. SCHRAMM

University of Chicago, Enrico Fermi Institute (LASR), 933 E 56th, Chicago, Ill. 60637, USA

James E. GUNN

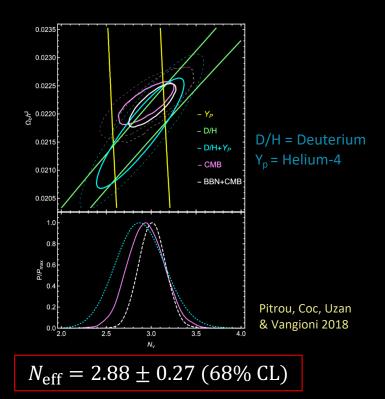
University of Chicago and California Institute of Technology, USA

Received 29 November 1976

If massive leptons exist, their associated neutrinos would have been copiously produced in the early stages of the hot, big bang cosmology. These neutrinos would have contributed to the total energy density and would have had the effect of speeding up the expansion of the universe. The effect of the speed-up on primordial nucleosynthesis is to produce a higher abundance of 4He. It is shown that observational limits to the primordial abundance of 4He. It is shown that observational limits to the primordial abundance of 4He. It is shown that observational limits to the primordial abundance of 4He lead to



the constraint that the total number of types of heavy lepton must be less than or equal to 5.

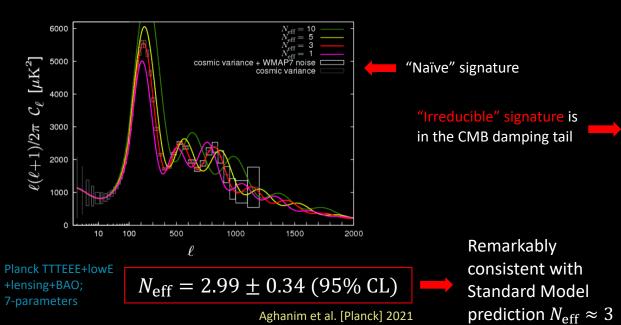


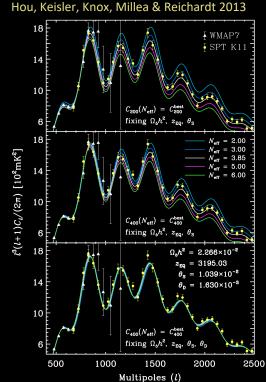
At face value a strong statement against thermalised light sterile neutrinos ($N_{\rm eff}=4$).

CMB anisotropies & N_{eff}...

 $N_{\rm eff}$ also affects the expansion rate at recombination.

Observable in the CMB temperature power spectrum



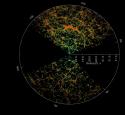


What do these probes really probe?

They may look different, but ultimately the information contained is

- Growth of fluctuations under gravity
 - Kinematic properties and interactions of the various types of stuff in the universe; good for neutrino physics



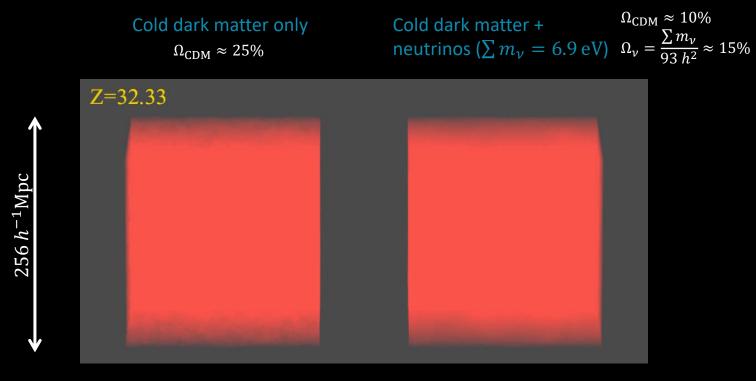




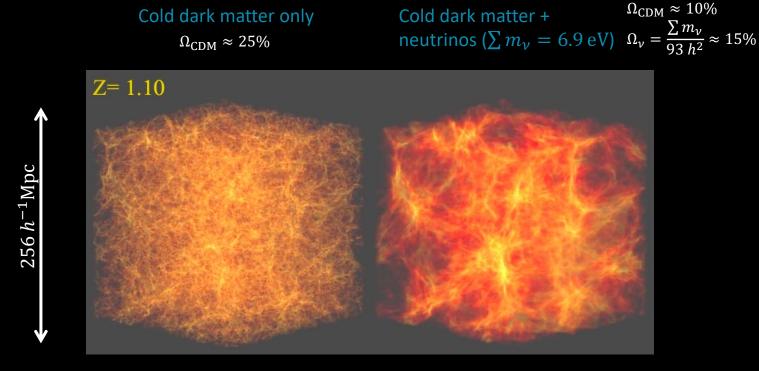
Neutrino masses, $\sum m_{\nu}$ (large-scale structure) Neutrino decay/lifetime, τ_0 (CMB)

Non-standard neutrino interactions (CMB)

Neutrino masses & large-scale structure...



Neutrino masses & large-scale structure...

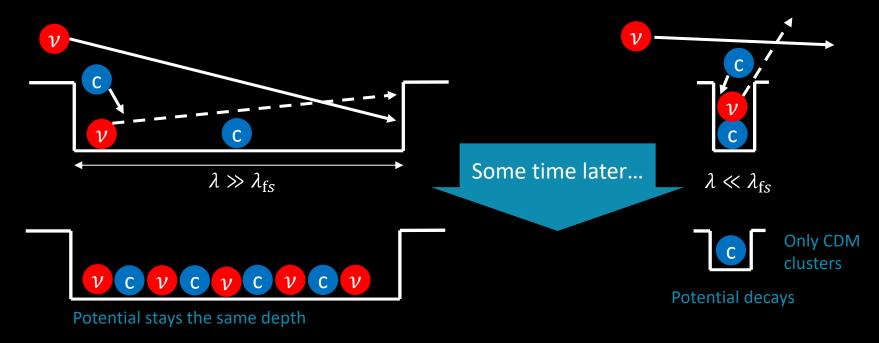


Why? Free-streaming suppression...

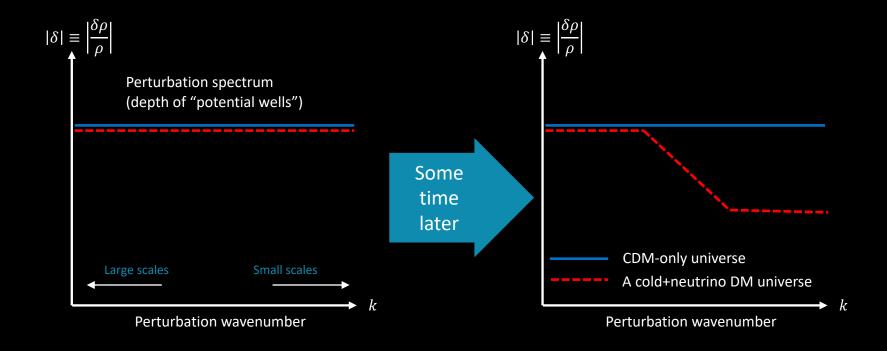
Neutrino thermal motion prevents efficient clustering on small length scales.

Free-streaming scale:
$$\lambda_{\rm fs} \equiv \sqrt{\frac{8\pi^2 v_{\rm thermal}^2}{3\Omega_m H^2}} \approx 4.2 \sqrt{\frac{1+z}{\Omega_{m,0}} \left(\frac{\rm eV}{m_\nu}\right)} \ h^{-1} \ {\rm Mpc}$$

A neutrino and a cold DM particle encounter 2 gravitational potential wells of different physical sizes in an expanding universe:

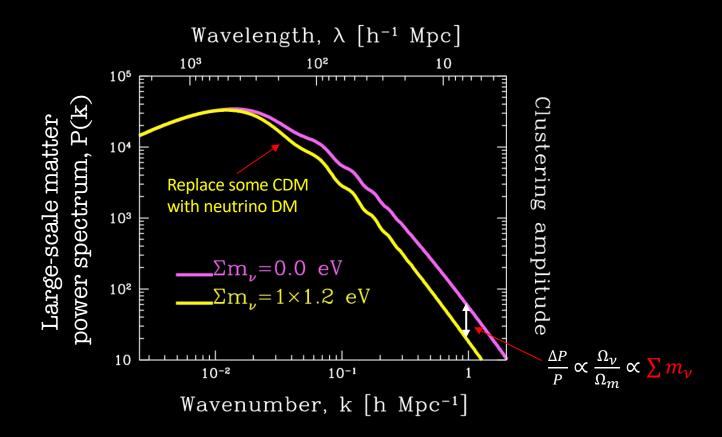


Free-streaming induces gravitational potential decay on length scales $\lambda \ll \lambda_{FS}$.

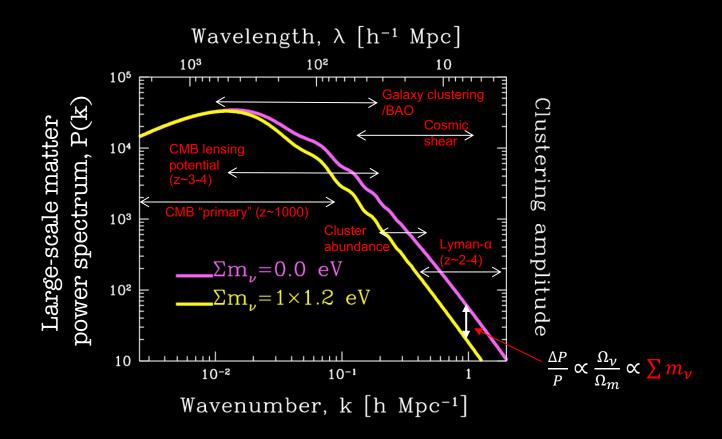


The presence of neutrino dark matter induces a step-like feature in the spectrum of gravitational potential wells.

A real matter power spectrum calculated from linear perturbation theory...

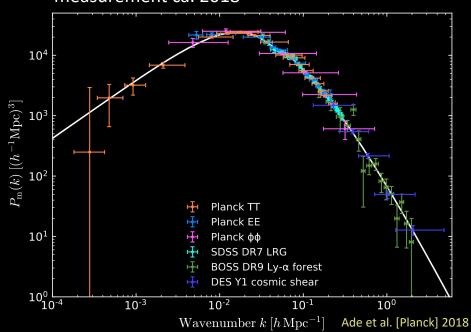


A real matter power spectrum calculated from linear perturbation theory...



Neutrino mass from cosmology...

Large-scale matter power spectrum measurement ca. 2018



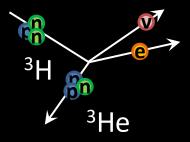
Planck TTTEEE+lowE+lensing+BAO; 7-parameters

$$\sum m_{\nu} < 0.12 \text{ eV } (95\% \text{ CL})$$

Aghanim et al. [Planck] 2021

At face value a factor of 30 tighter than current lab bound from KATRIN, $\sum m_{\nu} < 3 \text{ eV}$.

Aker et al. [KATRIN] 2019





Future cosmological probes...

	46	
P	e	

ESA Euclid

2024

 1σ sensitivity to $\sum m_{\nu}$ 1σ sensitivity to $N_{\rm eff}$

0.011 - 0.02 eV

0.05



LSST

2024

0.015 eV

0.05



CMB-S4

2027

0.015 eV

0.02 - 0.04

Minimum $\sum m_{\nu} = 0.06 \text{ eV}$ From neutrino oscillations (assuming normal mass ordering)

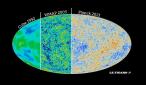


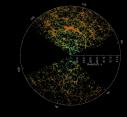
Detection of the absolute neutrino mass may be possible!

What do these probes really probe?

They may look different, but ultimately the information contained is

- Growth of fluctuations under gravity
 - Kinematic properties and interactions of the various types of stuff in the universe; good for neutrino physics





Neutrino masses, $\sum m_{\nu}$ (large-scale structure)

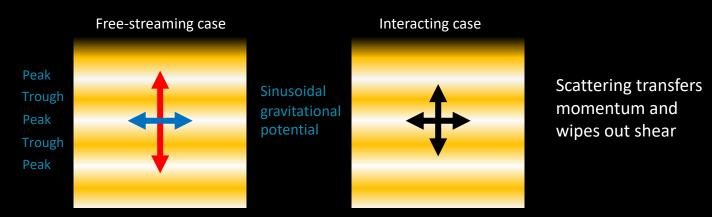


Non-standard neutrino interactions (CMB)

Neutrino free-streaming & the CMB...

Standard neutrinos free-stream.

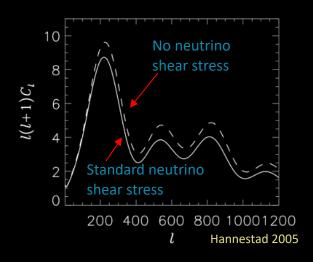
- Free-streaming in a spatially inhomogeneous background induces shear stress
- Conversely, interactions transfer momentum and wipe to out shear.

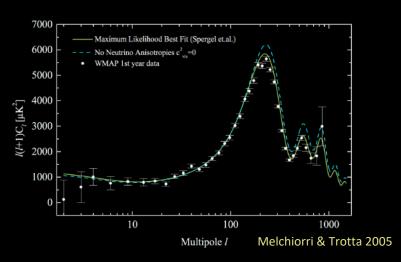


- Neutrino shear stress (or lack thereof) leave distinct imprints on the spacetime metric.
 - Affects the evolution of CMB perturbations; observable in the TT spectrum.

Neutrino free-streaming & the CMB...

That CMB prefers neutrino shear stress to no shear stress is well known.



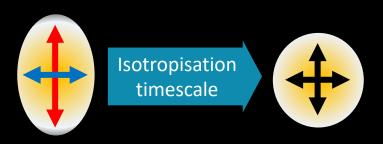


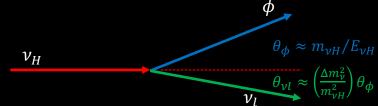
- The trick is in translating this preference to constraints on the fundamental parameters of a non-standard neutrino interaction → What is the isotropisation timescale?
 - Isotropisation timescale should not be longer than the CMB timescale (400k years).

Isotropisation from relativistic (inverse) decay...

Consider $\nu_H \rightarrow \nu_l + \phi$ and its inverse process.

Isotropisation timescale = How long it takes for decay and inverse decay to wipe out the momentum anisotropy in a fluid element.



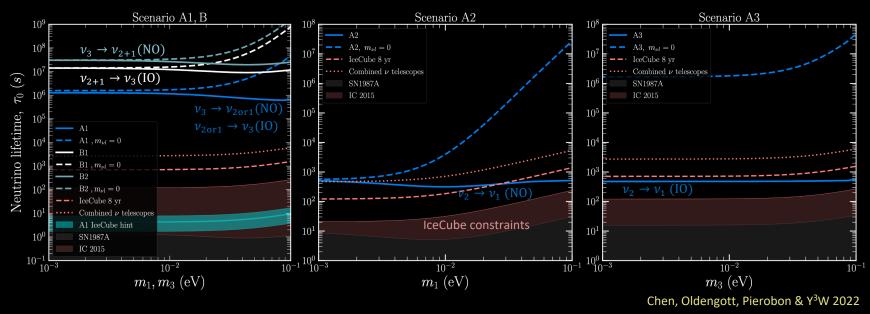


In relativistic decay, the decay products are beamed. Inverse decay can also only happen when the daughter particles satisfy strict momentum/angular requirements.

$$ightharpoonup$$
 Isotropisation is a looooong process: $T_{\rm isotropise} \sim \left(\theta_\phi \theta_{vl}\right)^{-2} \gamma_{vH} \, \tau_{\rm rest}$

CMB lower bounds on the neutrino lifetime...

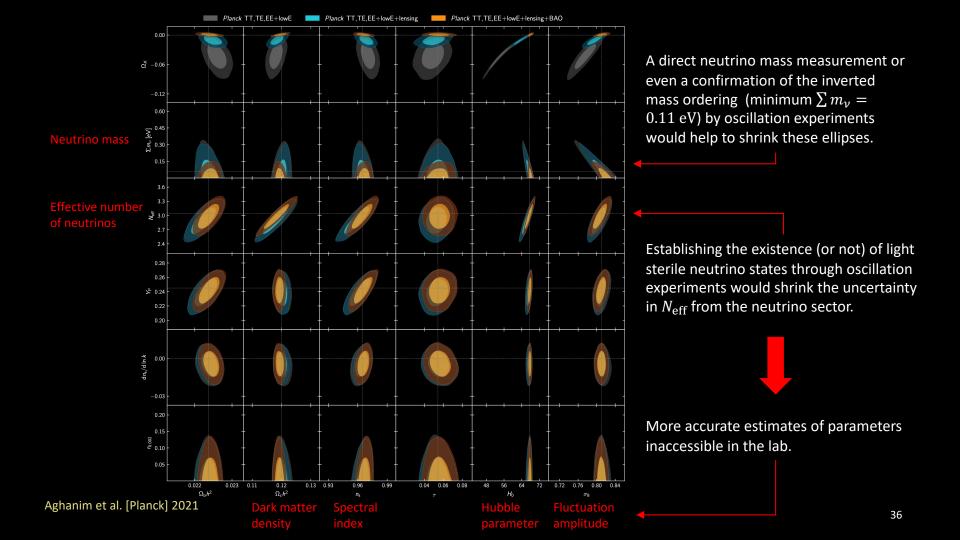
Mass-spectrum consistent constraints on invisible neutrino decay $\nu_H \rightarrow \nu_l + \phi$.



In some scenarios, neutrino telescopes and CMB probe the same parameter space.

Parameter estimation from cosmological observations is based on statistical inference.

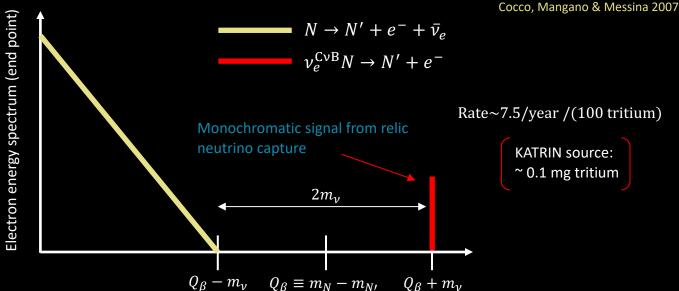
- Observations don't actually measure the dark matter density or the dark energy equation of state, and much less anything about inflation.
- Inference always assumes a model:
 - The less the model uncertainty, the more precise and accurate the parameter estimates.
- Neutrinos are unique in that they are the only cosmologically significant component that has a
 precise prediction within the Standard Model and whose properties can be independently
 measured in a laboratory.
 - Eliminating uncertainty in the neutrino sector will help us pin down other cosmological parameter inaccessible in the laboratory.



Ultimate prize = Direct detection of the CvB itself

• Best idea uses the β -decay end-point spectrum \rightarrow Goes hand-in-hand with direct neutrino mass detection.

Weinberg 1962

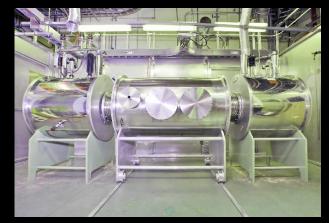


Ultimate prize = Direct detection of the CvB itself

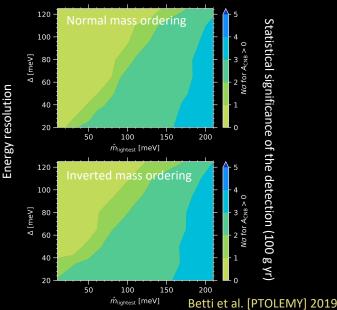
• Best idea uses the β -decay end-point spectrum \rightarrow Goes hand-in-hand with direct neutrino

mass detection.

Ptolemy Experiment



The PTOLEMY prototype at the Princeton Plasma Physics Laboratory



Summary...

- The existence of a cosmic neutrino background is a fundamental prediction of SM+FLRW cosmology.
 - Precision cosmological observations have allowed us to infer the properties of this background, from which to determine neutrino properties.
 - e.g., masses, effective number of neutrinos, non-standard interactions, lifetime.
- Conversely, better determination of neutrino properties in laboratory experiments will allow us to eliminate some model uncertainty in the cosmological parameter inference exercise.
 - More precise and accurate constraints on the dark matter density, dark energy properties, inflationary physics, and other cosmological physics inaccessible in the laboratory.