# Leptogenesis and Fundamental Symmetries of Nature

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# Baryon Number Asymmetry in SM

- Within the SM:
  - CP violation in quark sector not sufficient to explain the observed matterantimatter asymmetry of the Universe
- CP phase in quark sector:

Shaposhnikov, 1986; Farrar, Shaposhnikov, 1993

$$B \simeq \frac{\alpha_w^4 T^3}{s} \delta_{CP} \simeq 10^{-8} \delta_{CP} \qquad \delta_{CP} \simeq \frac{A_{CP}}{T_C^{12}} \simeq 10^{-20}$$

effects of CP violation suppressed by small quark mixing

$$A_{CP} = (m_t^2 - m_c^2)(m_c^2 - m_u^2)(m_u^2 - m_t^2)(m_b^2 - m_s^2)(m_s^2 - m_d^2)(m_d^2 - m_b^2) \cdot J$$

$$\longrightarrow B \sim 10^{-28}$$

too small to account for the observed

- Various Baryogenesis mechanisms (see Babu's talk)
- neutrino masses open up a new possibility

Fukugita, Yanagida, 1986

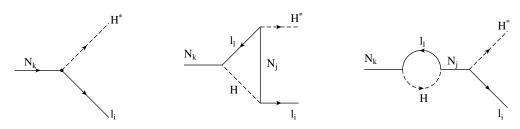
[For a review, see e.g. M.-C. C. TASI 2006 Lectures on Leptogenesis, hep-ph/0703087]

Leptogenesis

#### Leptogenesis

Fukugita, Yanagida, 1986

- Implemented in the context of seesaw mechanism
- out-of-equilibrium decays of RH neutrinos produce primordial lepton number asymmetry Luty, 1992; Covi, Roulet, Vissani, 1996; Flanz et al, 1996; Plumacher, 1997; Pilaftsis, 1997



$$\epsilon_1 = \frac{\sum_{\alpha} \left[ \Gamma(N_1 \to \ell_{\alpha} H) - \Gamma(N_1 \to \overline{\ell_{\alpha}} \overline{H}) \right]}{\sum_{\alpha} \left[ \Gamma(N_1 \to \ell_{\alpha} H) + \Gamma(N_1 \to \overline{\ell_{\alpha}} \overline{H}) \right]}$$

- sphaleron process convert ΔL → ΔB
- the asymmetry Buchmuller, Plumacher, 1998; Buchmuller, Di Bari, Plumacher, 2004

$$Y_B \simeq 10^{-2} \epsilon \kappa$$
  $\kappa$ : efficiency factor  $\sim (10^{-1} - 10^{-3})$   $Y_B = \frac{n_B - n_{\overline{B}}}{s} \sim 8.6 \times 10^{-11}$  (k: inverse decay  $\Delta L=1$ , scattering processes  $\Delta L=1$ , 2)

# Primordial $\Delta L$ from Heavy Neutrino Decay

[Animation Credit: Michael Ratz]

heavy neutrino decays generate lepton asymmetry

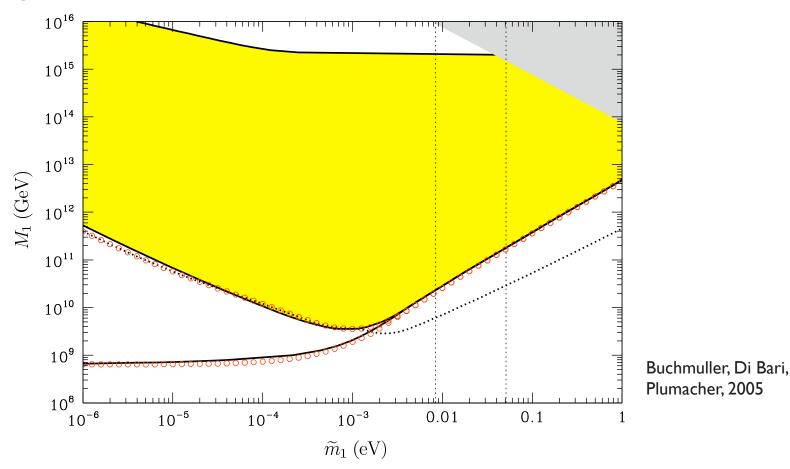
# Sphaleron Converting $\Delta L \rightarrow \Delta B$

[Animation Credit: Michael Ratz]

baryon asymmetry

# Bound on Light Neutrino Mass

- sufficient leptogenesis requires  $M_1 \geq 2 \times 10^9 \; {\rm GeV}$
- bound on light neutrino mass



#### Gravitino Problem

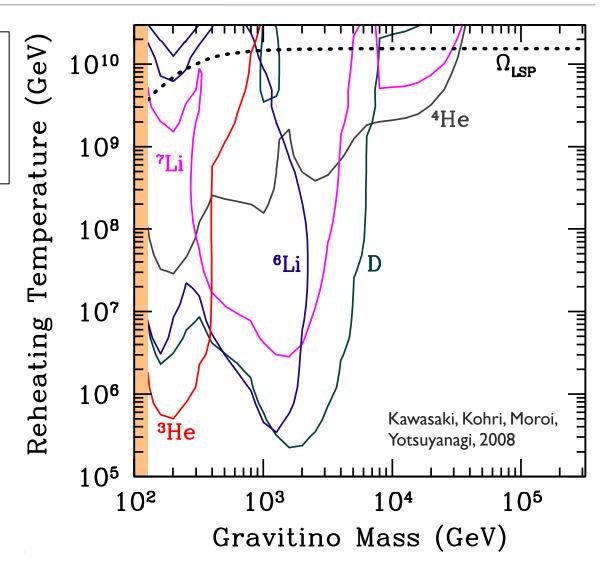
- Thermally produced RH neutrino N:
  - high reheating temperature needed:
    - $\Rightarrow$  T<sub>RH</sub> > M<sub>R</sub> > 2 x 10<sup>9</sup> GeV
- over-production of light state: gravitinos
- For gravitinos LSP:
  - DM constraint from WMAP
  - stringent bound on gluino mass for any given gravitino mass & TRH
- For unstable gravitinos:
  - long life time
  - decay during and after BBN ⇒ affect abundance of light elements

#### Gravitino Problem

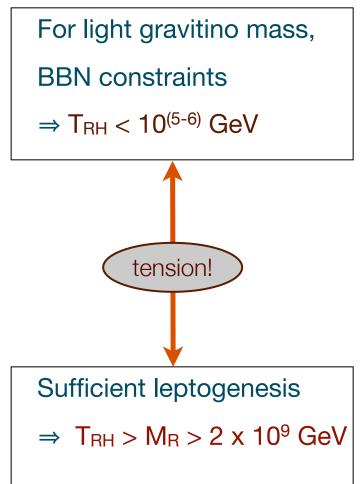
For light gravitino mass,

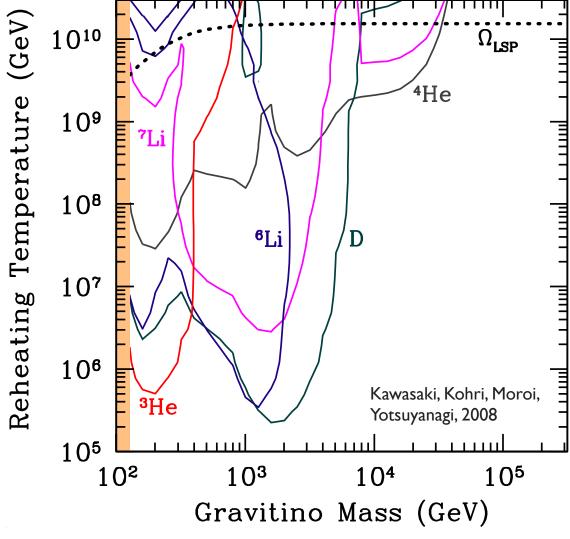
**BBN** constraints

 $\Rightarrow$  T<sub>RH</sub> < 10<sup>(5-6)</sup> GeV



#### Gravitino Problem

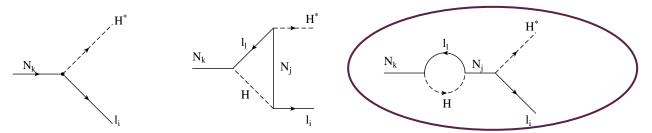




#### Alternatives: "Non-standard" Scenarios

- Possible ways to avoid the tension:
  - resonant enhancement in self-energy diagram ⇒ lowering M<sub>R</sub>, thus T<sub>RH</sub>
    - → resonant leptogenesis (near degenerate RH neutrinos) Pilaftsis, 1997

Recall: in standard leptogenesis:



self-energy diagram dominate for near degenerate RH neutrino masses,  $M_{1,2}$  enhanced O(1) asymmetry possible if

$$M_1 - M_2 \sim \frac{1}{2} \Gamma_{N_{1,2}}$$
, assuming  $\frac{Im(h_{\nu}h_{\nu}^{\dagger})_{12}^2}{(h_{\nu}h_{\nu}^{\dagger})_{11}(h_{\nu}h_{\nu}^{\dagger})_{22}} \sim 1$ 

leptogenesis possible even for low M<sub>1,2</sub>

Pilaftsis, Underwood, 2003

#### M.-C. Chen

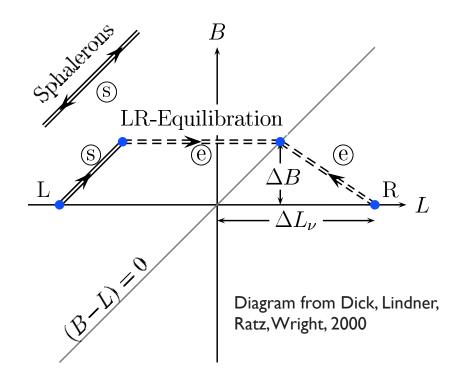
# Alternatives: "Leglands standard" Scenarios size of the natural SUSY breaking scale $\sqrt{\widetilde{m}^2} \sim \mathcal{O}(1)$ TeV, a genfar $_{ m st} R_{ m er}$ equiredepen the bresõnan eneconglition bl $R_H$ , of $_1 \sim \mathcal{O}(0.1)$ be obtained ible ways to avoid the tansians, $\begin{array}{c} & \sum_{f} \int_{0}^{\infty} [\textbf{retaxing}_{1} \textbf{retations}(\textbf{toetw}_{R_{1}}^{\dagger} \textbf{en} \textbf{p}) \textbf{pton number asymmetry and RH nu mass} \\ & \textbf{toe} \\ &$ ic tibet sween grange with a sufficient department of the bide to be a sign proattitimosiinactie (psydding fogen prixilogises softe compressiste pendence exting out the time integration the total CP asymmetry is $[^{42}M_R^{43}]$ , in thermal sis. This problem that be avoided if the relation has ween the reemperature And And Spring States of the second and the second sec the difference between the poccupation numbers of bosons and the difference between the poccupation numbers of bosons and the value of his head of his dedineutrinoshthroughatherinflatom2degay [48]. on solves the horizon and flatness problem, and it accounts for the density flucty attorned same that the inflat other ay production of RH neutrinos) r of lightest $\mathbb{R}^{H^8}$ neutrinos, $\Phi \to N_1 + N_1$ . For this decay to occur, on mass $m_{\Phi}$ (has solve) greater than $12M_1$ . For (simplicity, let fix) Hamaguchi, Yanagida, 2002

me that the decay modes into  $N_{2,3}$  are energetically forbidden.  $l_{\tilde{\nu}}$  in the first line is the density of the lightest sneutrino in equiuced. Which in the first line is the decay then subsequently decays into  $H + \ell_L$  in units of entropy density, and is given by,  $d_{\tilde{\nu}_R} = 45 \zeta(3)/(\pi^* g_*)$ ;

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K. Dick, M. Lindner, M. Ratz, D. Wright, 2000; H. Murayama, A. Pierce, 2002

- Leptogenesis possible even when neutrinos are Dirac particles
- small Dirac mass through suppressed Yukawa coupling
- Characteristics of Sphaleron effects:
  - only left-handed fields couple to sphalerons
  - sphalerons change (B+L) but not (B-L)
  - sphaleron effects in equilibrium for T > Tew
- If L stored in RH fermions can survive below EW phase transition, net lepton number can be generated even with L=0 initially



left – handed lepton asymmetry

right – handed lepton asymmetry

[Animation Credit: Michael Ratz]

left – handed
lepton
asymmetry sphalerons

right – handed
lepton
lepton
asymmetry
asymmetry

[Animation Credit: Michael Ratz]

- for neutrinos: LH equilibration can occur at late time (  $T_{eq} \ll T_{EW}$  ) because of their much suppressed masses ( $m_D < 10~{\rm keV}$  )
- Naturally small Dirac neutrino mass?
- Two examples:
  - non-anomalous U(1) family symmetry M.-C.C., J. Huang, W. Shepherd (2011)
    - gives realistic quark and lepton masses and mixing patterns
    - naturally small Dirac neutrino masses due to higher dimensional operators
    - primordial asymmetry by U(1) flavor higgs decay
  - discrete R-symmetries M.-C.C., M. Ratz, C. Staudt, P. Vaudrevange, to appear
    - satisfy all anomaly cancellation conditions a la Green-Schwarz mechanism
    - automatically suppressed the mu term, thus solving the mu problem in MSSM
    - automatically suppressed the Dirac neutrino masses

# Testing Leptogenesis?

- Sakharov Conditions:
  - out-of-equilibrium
    - expanding Universe
    - smallness of neutrino masses

Leptogenesis with Majorana neutrino: heavy field decay

Dirac Leptogenesis: late equilibration temperature

- Baryon Number Violation
  - abound in many extensions of the SM
  - neutrinoless double beta decay
    - Leptogenesis with Majorana (if observed) or Dirac (if not observed) neutrinos
- CP violation
  - → Long baseline neutrino oscillation experiments

#### Connection to Low Energy Observables

Leptogenesis ↔ Low Energy Observables
• Seesaw Lagrangian at high energy (in the presence of RH neutrinos)

6 mixing angles + 6 physical phases

Low energy effective Lagrangian (after integrating out RH neutrinos)

3 mixing angles + 3 physical phases

presence of low energy leptonic CPV (neutrino oscillation, neutrinoless double beta decay)

high energy → low energy: numbers of mixing angles and CP phases reduced by half



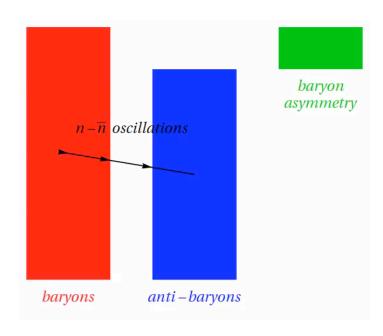
- No model independent connection
- Statement is weakened when the so-called flavor effects are taken into account (relevant if leptogenesis at T < 10<sup>12</sup> GeV)
- BUT, in certain models, connection can be established even without the flavor

#### Connection in Specific Models

- models for neutrino masses:
  - additional symmetries
  - reduce the number of parameters ⇒ connection can be established
- rank-2 mass matrix (may be realized by symmetry)
  - models with 2 RH neutrinos (2 x 3 seesaw) Kuchimanchi & Mohapatra, 2002
  - sign of baryon asymmetry ↔ sign of CPV in v oscillation Frampton, Glashow, Yanagida, 2002
- all CP come from a single source
  - models with spontaneous CP violation:
    - minimal LR model: only 1 physical leptonic CP phase M.-.C.C, Mahanthappa, 2005
    - SM + vectorial quarks + singlet scalar Branco, Parada, Rebelo, 2003
    - SCPV in SO(10): <126>B-L complex Achiman, 2004, 2008
  - SUSY SU(5) x T' Model: M.-.C.C, Mahanthappa, 2009
    - group theoretical origin of CP violation ⇒ only low energy lepton phases ≠ 0

#### Neutron-Antineutron Oscillation

- Neutrino Experiments → "archeological" evidence for leptogenesis
- n-nbar oscillation searches → complementarity test of leptogenesis (baryogenesis) mechanisms
  - constrain the scale of leptogenesis
- observation of neutron antineutron oscillation.
  - new physics with  $\Delta B = 2$  at  $10^{(5-6)}$  GeV
  - erasure of matter-antimatter generated at high scale, e.g. standard leptogenesis
- Low scale leptogenesis scenarios preferred:
  - Dirac Leptogenesis
  - Resonance Leptogenesis
  - .....



[Animation Credit: Michael Ratz]

#### Conclusions

- origin of matter: one of the great mysteries in particle physics and cosmology
- leptogenesis: an appealing baryogenesis mechanism connected to neutrino physics
- various leptogenesis mechanisms:
  - standard leptogenesis: gravitino problem, incompatible with SUSY
  - resonance leptogenesis
  - Dirac leptogenesis
- While there is no model-independent way to test leptogenesis, searches at neutrino experiments (leptonic CPV, neutrino-less double beta decay) can provide supports for/distinguish among the mechanisms
- neutron-antineutron oscillation: complementarity test
  - if observed ⇒ low scale leptogenesis scenarios preferred