Benchmarking intranuclear cascade models for neutrino scattering with relativistic optical potentials

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A direct comparison of RDWIA with relativistic optical potential (ROP) with intranuclear cascade (INC) model of the NEUT generator

In short:

The ROP removes nucleons that suffer inelastic FSI

The INC explicitly models inelastic FSI

Kinematic cuts on the INC results can be used to define an event sample that is equivalent to the ROP results

INC and ROP are found compatible for large energy only!

Broader context:

- Most precision experiments that measure interactions with nuclei apply stringent kinematic cuts to isolate observables of interest, minimize FSI, etc... e.g. exclusive (e,e'p), elastic p-A
- Theory to calculate/interpret these observables successfully are often specific to the problem/experimental setup at hand
- In accelerator-based neutrino experiments an exclusive measurement does not exist. Measurements are semi-inclusive! This means (in principle) all coupled final-states have to be described consistently → Currently done with INC

Some open questions:

- What is the minimum needed to describe a specific CC1p + X signal ? (depends on flux, kinematics, cuts)
- What is the 'correct' input for a cascade model ?

See e.g. [R. Gonzalez-Jimenez et al. PRC 105, 025502] [J.W. Van Orden et al. PRC 100, 044620] [J.M. Franco Patino Arxiv:2207.02086 (2022)] (Talk NuFACT) [Isaacson et al. Arxiv:2205.06378 (2022)] [B. Bourgouille et al. JHEP4 (2021) 004]

Intra-nuclear cascade model (INC) in NEUT

- 1.) Nucleon propagates in straight lines
- 2.) Check for interaction based on density and nucleon-nucleon CS
- 3.) Pauli-blocking: Reaction products must be above p_{Fermi}
- 4.) Track the created particles on the way and propagate them

Proton-Carbon cross-sections 300 თ **(mb)** 250 200 N-N inelastic 150 100 N-N elastic 50 500 1000 2000 1500 2500 proton kinetic energy (MeV)

Why ?

Treat the involved coupled-channels problem posed by FSI

Tradeoffs:

Factorized and classical approach Expected to break down at low Energy Eur. Phys. J. Spec. Top. (2021) 230:4469–4481



Coupled channels problem, separate out elastic channel

$$(H_{00} - E)|\phi_0\rangle = -V_{0j}|\phi_j\rangle \ (j > 0) \bullet$$
$$(H_{ij} - E)|\phi_j\rangle = -V_{j0}|\phi_0\rangle \ (i, j > 0) \bullet$$

$$H_{ij} = H^{free}\delta_{ij} + V_{ij}^{nA} + \epsilon_i \delta_{ij}$$



Coupled channels problem \rightarrow one-body problem with optical potential

$$\begin{aligned} (H_{00} - E)|\phi_0\rangle &= -V_{0j}|\phi_j\rangle \ (j > 0) \bullet \\ (H_{ij} - E)|\phi_j\rangle &= -V_{j0}|\phi_0\rangle \ (i, j > 0) \bullet \\ & \downarrow \\ \left[H^{free} + V_{00}^{nA} + V_{0j}\frac{1}{E - H_{ij} + i\eta}V_{j0} - E\right]|\phi_0\rangle \bullet \end{aligned}$$



$$\begin{aligned} (H_{00} - E) |\phi_0\rangle &= -V_{0j} |\phi_j\rangle \ (j > 0) \bullet \\ (H_{ij} - E) |\phi_j\rangle &= -V_{j0} |\phi_0\rangle \ (i, j > 0) \bullet \\ \downarrow \\ \left[H^{free} + V_{00}^{nA} + V_{0j} \frac{1}{E - H_{ij} + i\eta} V_{j0} - E \right] |\phi_0\rangle \bullet \\ &\approx \left[H^{free} + \mathcal{V}^{opt} - E \right] |\phi_0\rangle \end{aligned}$$

Empirical relativistic optical potential (ROP)

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Global Dirac phenomenology for proton-nucleus elastic scattering

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Target	T_p (MeV)	σ_R (mb)				
		EDAI-fit		EDAD-fit		Reference
			fit 1	fit 2	fit 3	
¹² C	29.00	420.2	435.5	433.1	422.7	[6]
	30.30	415.9	429.0	425.6	414.2	[7]
	49.00	358.8	363.0	348.4	327.7	[6]
	49.48	357.4	361.8	347.0	326.1	[8]
	61.40	323.3	335.6	317.0	294.8	[9]
	65.00	313.5	329.0	309.7	287.4	[10]
	122.00	202.2	269.0	254.4	230.5	[11]
	160.00	177.8	252.3	246.4	215.2	[11]
	200.00	177.6	243.0	243.9	205.0	[11-13]
	300.00	201.1	233.0	235.4	194.9	[14]
	398.00	215.8	227.4	218.6	199.1	[15]
	494.00	227.2	223.7	203.0	211.6	[16]
	797.50	238.4	235.3	209.9	250.0	[17,18]
	1040.00	198.6	259.4	243.8	232.2	[19.20]

(Received 31 August 1992)



10⁵

Energy-dependent potentials fit to *elastic* proton-nucleus scattering

N–A scattering with INC and ROP

N-A data is used as benchmark/input to INC



$$\sigma_{tot} = \sigma_{reac} + \sigma_{el}$$

 σ_{el} not modeled in (NEUT) INC

 σ_{tot} Obtained in ROP with optical theorem





[Dytman et al. PRD104, 053006]

Exclusive electron scattering with ROP



Incoming energy is known Never satisfied for accelerator neutrinos! [M. Leuschner et al. PRC49, 955 (1994)]

Independent particle model (IPM)



angular momentum Energy

(No momentum eigenstates)

Experimental data implies smaller occupations than IPM → Correlations **and FSI**





RDWIA with ROP for exclusive (e,e'p)





The optical potential removes nucleons that suffer inelastic FSI \rightarrow changes E_m

FIG. 11. The reduced cross section (σ_{red}) of the ${}^{16}O(e,e'p)$ reaction as a function of the recoil momentum p_m for the transitions to the $1/2^-$ ground state and to the $3/2^-$ excited state of ${}^{15}N$, in

The ROP also changes the observed p_m spectrum!

p_m [MeV/c]

Direct comparison of INC and ROP

1. Input to NEUT INC

Events according to **unfactorized** fluxfolded five-fold differential nucleon knockout cross section

2. Kinematic cuts on NEUT result

Select events with E_m corresponding to shell-model peaks

Signal with only 'elastic' FSI That does not change $\rm E_{\rm m}$

Can be directly compared to ROP



NEUT Cascade with rROP input



Simplified E_m contains recoil energy

$$E_m = E_i - E_l - T_p$$

T2K flux-folded calculations

- rROP+NEUT moves strength from shell model peaks to larger E_m
- Resulting strength of shell model peaks agrees with ROP predictions
- FSI leads to mostly large ~50 MeV shifts in missing energy

 $T_{\mbox{\tiny p}}$ distributions INC and ROP



The INC agrees with the ROP for high T_n

At $T_p < 100$ differences go up to 100% !

In NEUT low energy nucleons don't

Large model dependence in INC's at low E



The INC agrees with the ROP for high T_{p}

At $T_p < 100$ differences go up to 100% !

In NEUT low energy nucleons don't undergo FSI at all!

 \rightarrow Will influence a.o. neutron rates

How do other INC perform ? You can use these results!



Important issue: the input to the INC

1. Input to NEUT INC

Events according to **unfactorized** fluxfolded five-fold differential nucleon knockout cross section

(computationally non-trivial work by RGJ)



RDWIA calculations with real final-state potentials

- Describe *inclusive* (e,e') [PRC 100 045501 (2019)]
- Similar to SuSAv2 [PRC 101 015503 (2020)]
- Not limited to inclusive CS! [PRC 105 025502 (2022)]
 [Arxiv:2207.02086 (2022)]

= J.M. Franco-Patino's talk



4 August 2022

NuFACT 2022

17/29

Important issue: the input to the INC

1. Input to INC includes some FSI!

FSI treatment is important to reproduce the inclusive cross section In RDWIA, CRPA, effective SF, LFG+RPA, ... some FSI included! Cannot always (ever?) separate 'initial' and 'final'!



Inclusive cross sections in generators



Implementation of the CRPA model in the GENIE event generator and analysis of nuclear effects in low-energy transfer neutrino-nucleus interactions

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 (Dated: November 2, 2021)

Implementation of the SuSAv2-MEC 1p1h and 2p2h models in GENIE and analysis of nuclear effects in T2K measurements

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 ⁴University of Tokyo, Institute for Cosmic Ray Research, Research Center for Cosmic Neutrinos, Kashiwa, Japan (Dated: February 21, 2020)

Implementation strategy

- Include inclusive responses
- Add nucleon final state in GENIE based on momentum distribution
- SuSAv2, LFG+RPA (Valencia), CRPA, SuSA-MEC, ...

Inclusive cross sections in generators

The idealized version of the event generator:



The reality:

(part of) the initial-state, the 'extra effects' and the FSI are included

It is not possible to obtain the hadron information from inclusive CS!!

Nucleon variables in GENIE

Input to the generator is inclusive cross section:

$$\frac{\mathrm{d}\sigma(E_{\nu})}{\mathrm{d}E_{l}\mathrm{d}\cos\theta_{l}} = G^{2}\frac{k_{l}}{E_{\nu}}L_{\mu\nu}\int\mathrm{d}\Omega_{N}\sum_{n,\kappa}H_{n,\kappa}^{\mu\nu}(\omega,q,\Omega_{N},E_{n,\kappa})$$

Lost nucleon information \rightarrow Need to generate it in GENIE

1. Draw initial nucleon \mathbf{p}_{m} from p² n(p) (e.g. LFG)

2. Compute
$$E_m^2 = p_m^2 + M_N^2$$

3.
$$E_N = E_m + \omega - E_b(q)$$

4.
$$k_N^2 = E_N^{2-} M_N^2$$

1. $|\mathbf{p}_m + \mathbf{q}| \neq k_N = \sqrt{E_N^2 - M_N^2}$

$$\mathbf{\rightarrow } \mathbf{k}_{N} = \frac{k_{N}}{|\mathbf{p}_{m} + \mathbf{q}|} \ (\mathbf{p}_{m} + \mathbf{q})$$





Nucleon variables in GENIE and RDWIA

Comparison of the nucleon kinematics

- The full 5-fold RDWIA calculation
- The GENIE algorithm for 'adding' the nucleon



Nucleon variables in GENIE and RDWIA

Comparison of the nucleon kinematics

- The full 5-fold RDWIA calculation
- The GENIE algorithm for 'adding' the nucleon



23/29

Conclusions

- The ROP and INC approaches both use nucleon-nucleus scattering to constrain FSI, albeit in very different ways.
- A consistent comparison of the NEUT INC and optical potential approaches shows that there is quantitative agreement only at large kinetic energies.
- For small kinetic energy the differences are up to 100% !!
- The ROP should be more reliable in this comparison, but the true answer is unknown!
- Results of the generator will depend crucially on the input to the INC, but current implementations (necessarily) use unrealistic approximations
- Unfactorized Events for flux-averaged signals over the whole phase space can be generated
 → You can use these for validation/error estimation/... of your own INC/simulation/...

Transverse Kinematic Imbalance



Because $p_p > 450$ MeV low energy differences are not seen

Effect of non-elastic FSI visible in $P_{_{T}} and \; \alpha_{_{T}}$

Large non-QE 'background' not separable from FSI effects

Non-QE from [Bourguille et al. JHEP04(2021)004]

(e,e'p) results for CLAS e4nu kinematics

Fixed-E experiment (e,e'p) in CLAS

Range of nucleon and lepton kinematics

<u>Restrict E_m to probe</u> <u>specific interaction</u> <u>mechanisms</u>



 T_p (GeV)

 ${ ilde E}_m < 25 {
m ~MeV}$ — No 2p2h, RES, inelastic FSI

But with full range of ω , q, T_N to study FSI 351.4ROP 301.2 $d\sigma/dT_p ~(\mu {\rm b/GeV})$ E = 2.2 GeVRPWIA+NEUT ····· $d\sigma/dT_p~(\mu {
m b/GeV})$ 25 $\tilde{E}_m < 25 \text{ MeV}$ rROP+NEUT 1 200.8E = 1.1 GeV150.6 $\tilde{E}_m < 25 \text{ MeV}$ 100.450.20 0 0.20.050.10.150.250.30 0.20.30.10.40.50

 T_p (GeV)

Degeneracy of interaction channels

Is reduced, but not removed by fixing incoming energy

Is removed by restricting the energy/momentum of residual system

Restricted kinematics	$l + A \to l' + N + B$	Unrestricted kinematics	
exclusive (e,e'p)	Fixed-E experiment (v _ı ,l'p) or (e,e'p)	Accelerator (v _ı ,l'p)	
Fixed nucleon and lepton kinematics	Range of nucleon and lepton kinematics	Range of nucleon and lepton kinomatics	
Severely restrict E _m	<u>specific interaction</u> <u>mechanisms</u>	No E _m restriction	
Learn nuclear structure and Reduce FSI and kinematic effects		possible	

E_m (& p_m) : (e,e'p)



Figures from talk by C. Giusti (INT Workshop INT-18-1a)

Relativistic mean field with nlsw interaction

$$\mathcal{L} = \overline{\Psi} \left(i \gamma_{\mu} \partial^{\mu} - M \right) \Psi + \frac{1}{2} \left(\partial_{\mu} \sigma \partial^{\mu} \sigma - m_{\sigma}^{2} \sigma^{2} \right) - U(\sigma)$$

$$- \frac{1}{4} \Omega_{\mu\nu} \Omega^{\mu\nu} + \frac{1}{2} m_{\omega}^{2} \omega_{\mu} \omega^{\mu} - \frac{1}{4} \mathbf{R}_{\mu\nu} \mathbf{R}^{\mu\nu} + \frac{1}{2} m_{\rho}^{2} \rho_{\mu} \rho^{\mu} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$

$$- g_{\sigma} \overline{\Psi} \sigma \Psi - g_{\omega} \overline{\Psi} \gamma_{\mu} \omega^{\mu} \Psi - g_{\rho} \overline{\Psi} \gamma_{\mu} \tau \rho^{\mu} \Psi - g_{e} \frac{1 + \tau_{3}}{2} \overline{\Psi} \gamma_{\mu} A^{\mu} \Psi .$$

Extension of the original σ-ω Walecka model (Ann. Phys.83,491 (1974)).

$$\left[\hat{\alpha} \cdot \hat{\mathbf{p}} + \beta \left(m_N + S\left(r\right)\right) - \left(E - V\left(r\right)\right)\right] \psi = 0,$$

Where:

$$S(r) = g_{\sigma}\sigma(r)$$

$$V(r) = g_{\omega}\omega^{0}(r) + g_{\rho}\tau_{3}\rho_{3}^{0}(r) + e\frac{1+\tau_{3}}{2}A^{0}(r).$$

 $U(\sigma) = \frac{1}{3}g_2\sigma^3 + \frac{1}{4}g_3\sigma^4$

$$\begin{aligned} \left(\nabla^2 - m_{\sigma}^2\right)\sigma(r) &= g_{\sigma}\rho_s(r) + g_2\sigma^2(r) + g_3\sigma^3(r),\\ \left(\nabla^2 - m_{\omega}^2\right)\omega^0(r) &= g_{\omega}\rho_B(r),\\ \left(\nabla^2 - m_{\rho}^2\right)\rho_3^0(r) &= g_{\rho}\rho_I(r),\\ \nabla^2\sigma(r) &= -e\rho_e(r), \end{aligned}$$

Main approximations:

1) Mean-field
approximation:
$$\sigma \rightarrow \langle \sigma \rangle$$
 $\rho_{\mu} \rightarrow \langle \rho_{\mu} \rangle$
2) Static limit:
 $\partial^{0}\omega_{0} = \partial^{0}\rho_{0} = \partial^{0}\sigma = 0$ $\omega_{\mu} = \delta_{\mu0}\omega_{0}$, $\rho_{\mu} = \delta_{\mu0}\rho_{0}$
3) Spherical symmetry for finite
 $\omega_{0} = \omega_{0}(r)$ $\rho_{0} = \rho_{0}(r)$ $\sigma = \sigma(r)$