

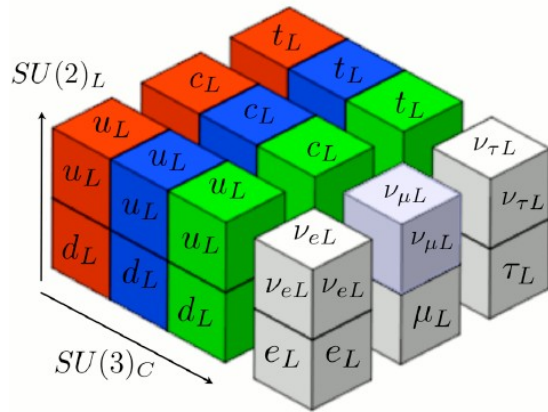


Modular symmetries and the flavor problem

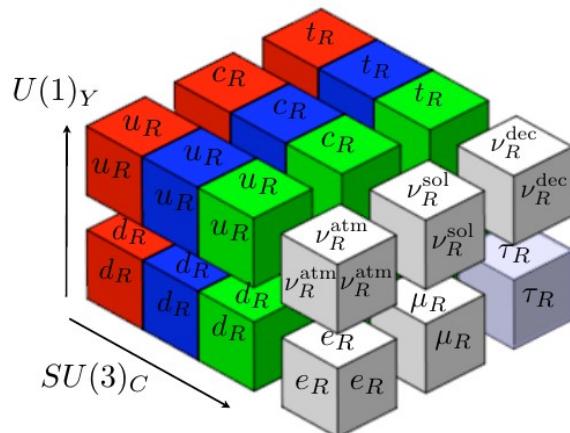
Davide Meloni
Dipartimento di Matematica e Fisica, Roma Tre

The Standard Model of Particle Physics

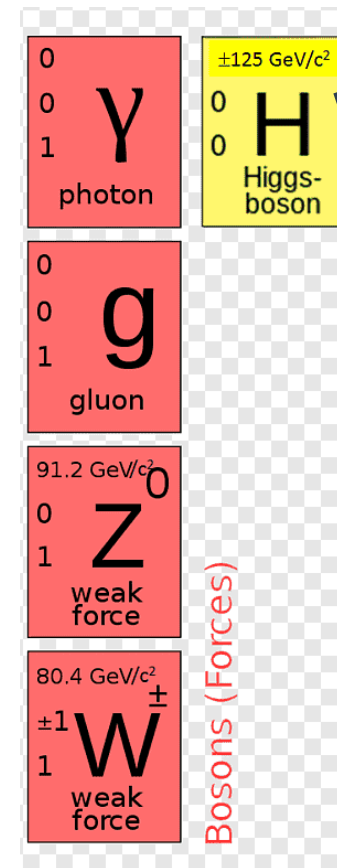
Left-handed



Right-handed



Gauge boson sector

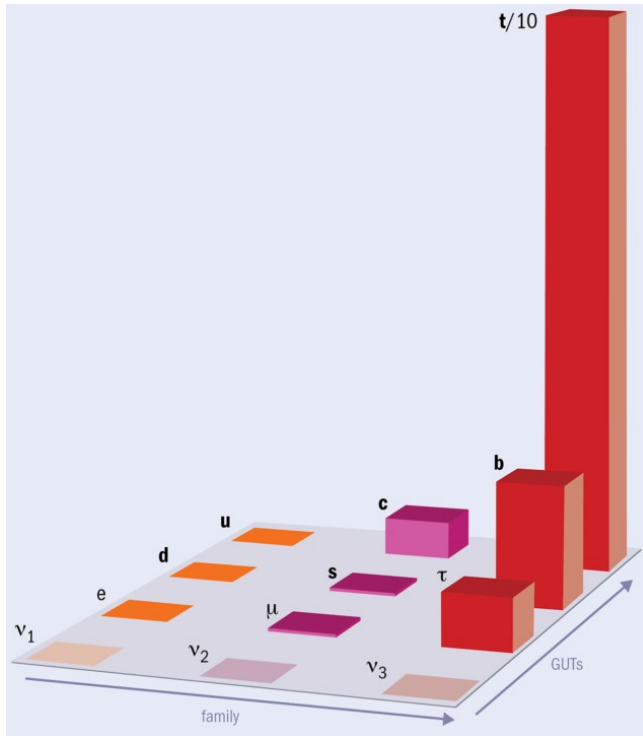


Scalar sector

Bosons (Forces)

The Flavor Problem

Mass hierarchies



$$m_d \ll m_s \ll m_b, \quad \frac{m_d}{m_s} = 5.02 \times 10^{-2},$$

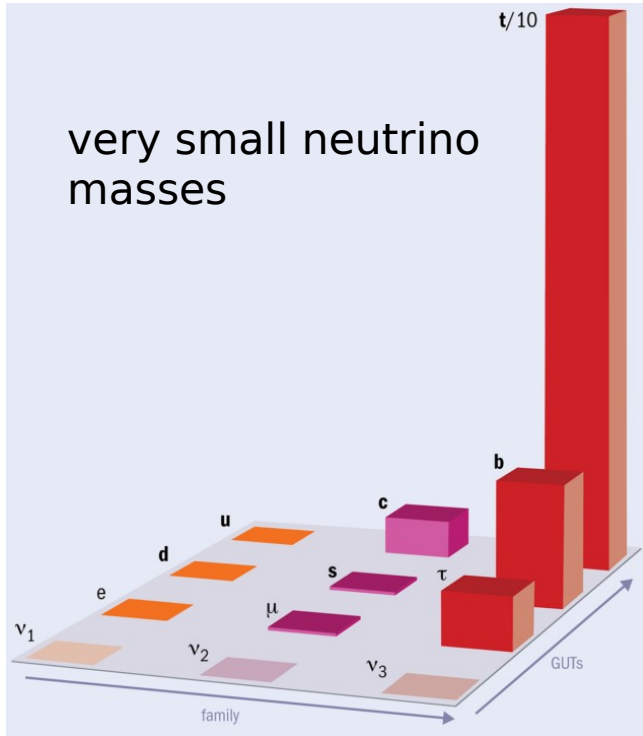
$$m_u \ll m_c \ll m_t, \quad \frac{m_u}{m_c} = 1.7 \times 10^{-3},$$

$$\frac{m_s}{m_b} = 2.22 \times 10^{-2}, \quad m_b = 4.18 \text{ GeV};$$

$$\frac{m_c}{m_t} = 7.3 \times 10^{-3}, \quad m_t = 172.9 \text{ GeV};$$

The Flavor Problem

Mass hierarchies



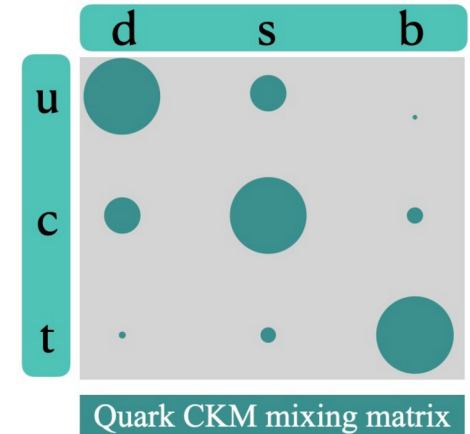
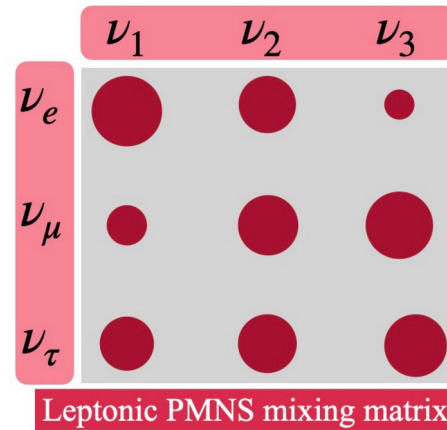
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Fermion mixing



all mixing are large but
the 13 element



almost a diagonal matrix

Suggested solutions

* Smallness of neutrino masses:

See-saw



$$\mathcal{M} = \begin{bmatrix} m_M^L & m_D \\ m_D & m_M^R \end{bmatrix}$$

$$m_{\text{light}} \sim \frac{m_D^2}{M_M^R}$$

No clue on mixing !

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- * Hierarchical Pattern

Froggatt-Nielsen mechanism

$$L \sim \overline{\Psi}_L H \Psi_R \left(\frac{\theta}{\Lambda} \right)^n \rightarrow e^{(-q_L + q_H + q_R + n * q_\theta)}$$

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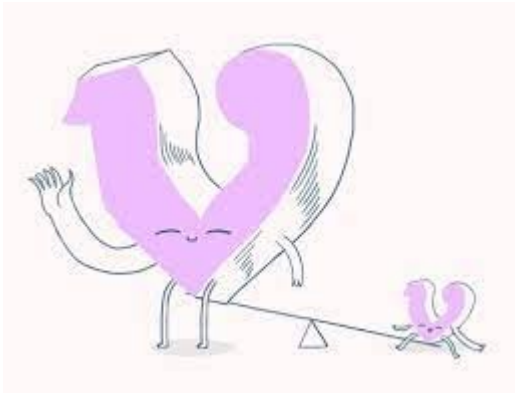
Too many O(1) coefficients

Works better for small mixing

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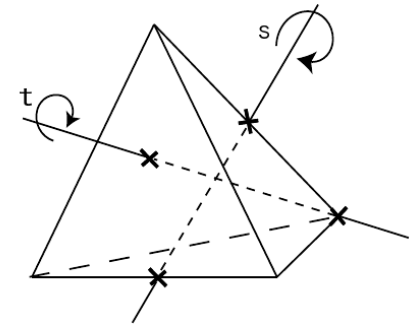
$$L \sim \overline{\Psi}_L H \Psi_R \left(\frac{\theta}{\Lambda} \right)^n$$

Too many O(1) coefficients

Works better for small mixing

- * mixing angles

elegant explanation:
non-Abelian
discrete flavour symmetries



Complicated scalar sector

Modular Symmetry

We start from

Feruglio, 1706.08749

$$\Gamma(N) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2, \mathbb{Z}), \begin{pmatrix} a & b \\ c & d \end{pmatrix} \equiv \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \pmod{N} \right\}$$

the group of 2x2 matrices with integer entries modulo N and determinant equals to one modulo N

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the group $\Gamma(N)$ acts on the complex variable τ ($\text{Im } \tau > 0$)

$$y\tau = \frac{a\tau + b}{c\tau + d}$$

Modular Symmetry

Important observation for $N=1$: a transformation characterized by parameters $\{a, b, c, d\}$ is identical to the one defined by $\{-a, -b, -c, -d\}$

$\Gamma(1)$ is isomorphic to $\mathbf{PSL}(2, \mathbf{Z}) = \mathbf{SL}(2, \mathbf{Z})/\{\pm 1\} = \Gamma$



inhomogeneous modular group (or simply **Modular Group**)

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In addition:

$$\bar{\Gamma}(2) = \Gamma(2)/\{1, -1\}$$



since 1 and -1 **cannot** be distinguished

$$\bar{\Gamma}(N) = \Gamma(N) \quad N > 2$$



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Finite Modular Group:

$$\Gamma_N = \frac{\bar{\Gamma}}{\bar{\Gamma}(N)}$$

Modular Symmetry

Generators of Γ_N : elements S and T satisfying

$$S^2=1, \quad (ST)^3=1, \quad T^N=1$$

$$S = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad T = \begin{pmatrix} 1 & 1 \\ 0 & -1 \end{pmatrix}$$

corresponding to:

$$\tau \xrightarrow{S} -\frac{1}{\tau}$$

$$\tau \xrightarrow{T} \tau+1$$

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relevant for model building:

for $N \leq 5$, the finite modular groups Γ_N are isomorphic to non-Abelian discrete groups

$$\Gamma_2 \simeq S_3 \quad \Gamma_3 \simeq A_4 \quad \Gamma_4 \simeq S_4 \quad \Gamma_5 \simeq A_5$$

Then the question is: why Modular Symmetry ?

Modular Forms

Modular Forms:

holomorphic functions of the complex variable τ with well-defined transformation properties under the group $\Gamma(N)$

$$f(\gamma\tau) = (c\tau + d)^{2k} f(\tau), \quad \gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma(N) \quad 2k = \text{weight}, \quad N = \text{level}$$

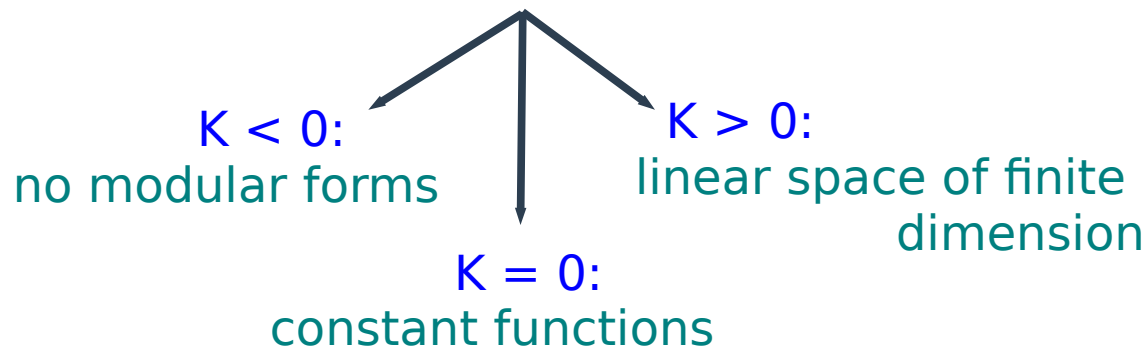
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N	$d_{2k}(\Gamma(N))$
2	$k + 1$
3	$2k + 1$
4	$4k + 1$
5	$10k + 1$
6	$12k$
7	$28k - 2$

R. C. Gunning, Lectures on Modular Forms, Princeton, New Jersey USA, Princeton University Press 1962

Model Building

Key points:

1. Modular forms of weight $2k$ and level $N \geq 2$ are invariant, up to the factor $(c\tau + d)^{2k}$ under $\Gamma(N)$ but they transform under Γ_N !

$$f_i(\gamma\tau) = (c\tau + d)^{2k} \rho(\gamma)_{ij} f_j(\tau)$$

representative element of Γ_N

unitary representation of Γ_N

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2. in addition, one assumes that the fields of the theory χ_i transforms non-trivially under Γ_N

$$\chi(x)_i \rightarrow (c\tau + d)^{-k_i} \rho(\gamma)_{ij} \chi(x)_j$$

not modular forms !
No restrictions on k_i

Model Building

Building blocks:

1. Modular forms and fields: $L_{eff} \in f(\tau) \times \phi^{(1)} \dots \phi^{(n)}$

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2. Invariance under modular transformation requires:

$$2k = \sum_i k_i$$
$$\rho_f \otimes \rho_{\chi_1} \otimes \dots \otimes \rho_{\chi_n} \supset I$$

To start playing the game:

Can someone give me the Modular Forms?

Model Building

Long list from **S.T. Petcov, Bethe Forum, University of Bonn, 04/05/2022**

For $(\Gamma_3 \simeq A_4)$, the generating (basis) modular forms of weight 2 were shown to form a 3 of A_4 (expressed in terms of log derivatives of Dedekind η -function η'/η of 4 different arguments).

F. Feruglio, arXiv:1706.08749

For $(\Gamma_2 \simeq S_3)$, the two basis modular forms of weight 2 were shown to form a 2 of S_3 (expressed in terms of η'/η of 3 different arguments).

T. Kobayashi, K. Tanaka, T.H. Tatsuishi, arXiv:1803.10391

For $(\Gamma_4 \simeq S_4)$, the 5 basis modular forms of weight 2 were shown to form a 2 and a 3' of S_4 (expressed in terms of η'/η of 6 different arguments).

J. Penedo, STP, arXiv:1806.11040

For $(\Gamma_5 \simeq A_5)$, the 11 basis modular forms of weight 2 were shown to form a 3, a 3' and a 5 of A_5 (expressed in terms of Jacobi theta function $\theta_3(z(\tau), t(\tau))$ for 12 different sets of $z(\tau), t(\tau)$).

P.P. Novichkov et al., arXiv:1812.02158; G.-J. Ding et al., arXiv:1903.12588

Multiplets of higher weight modular forms have been also constructed from tensor products of the lowest weight 2 multiplets:

i) for $N = 4$ (i.e., S_4), multiplets of weight 4 (weight $k \leq 10$) were derived in arXiv:1806.11040 (arXiv:1811.04933);

ii) for $N = 3$ (i.e., A_4) multiplets of weight $k \leq 6$ were found in arXiv:1706.08749;

iii) for $N = 5$ (i.e., A_5), multiplets of weight $k \leq 10$ were derived in arXiv:1812.02158.

Model Building

Constructing the Modular Forms

Crucial observation:

if $g(\tau) \rightarrow e^{i\alpha}(c\tau+d)^k g(\tau)$

then $\frac{d}{d\tau} \log[g(\tau)] \rightarrow (c\tau+d)^2 \frac{d}{d\tau} \log[g(\tau)] + \underbrace{k c (c\tau+d)}$

this term prevents of having a modular form of weight **2 k = 2**

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The inhomogeneous term can be removed if we combine several $f_i(\tau)$ with weights k_i

$$\frac{d}{d\tau} \sum_i \log[g_i(\tau)] \rightarrow (c\tau + d)^2 \frac{d}{d\tau} \sum_i \log[g_i(\tau)] + (\sum_i k_i) c (c\tau + d)$$

with $\sum_i k_i = 0$

A case study: $\Gamma_2 \sim S_3$

Let us find the functions $f(\tau)$!

The group S_3 contains $1 + 1' + 2$

N	$d_{2k}(\Gamma(N))$
2	$k + 1$

← two independent modular forms can fit into a doublet of S_3

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Dedekind eta functions $\eta(\tau) = q^{1/24} \prod_{n=1}^{\infty} (1 - q^n)$ $q \equiv e^{i2\pi\tau}$

S: $\eta(-1/\tau) = \sqrt{-i\tau} \eta(\tau)$, **T:** $\eta(\tau + 1) = e^{i\pi/12} \eta(\tau)$

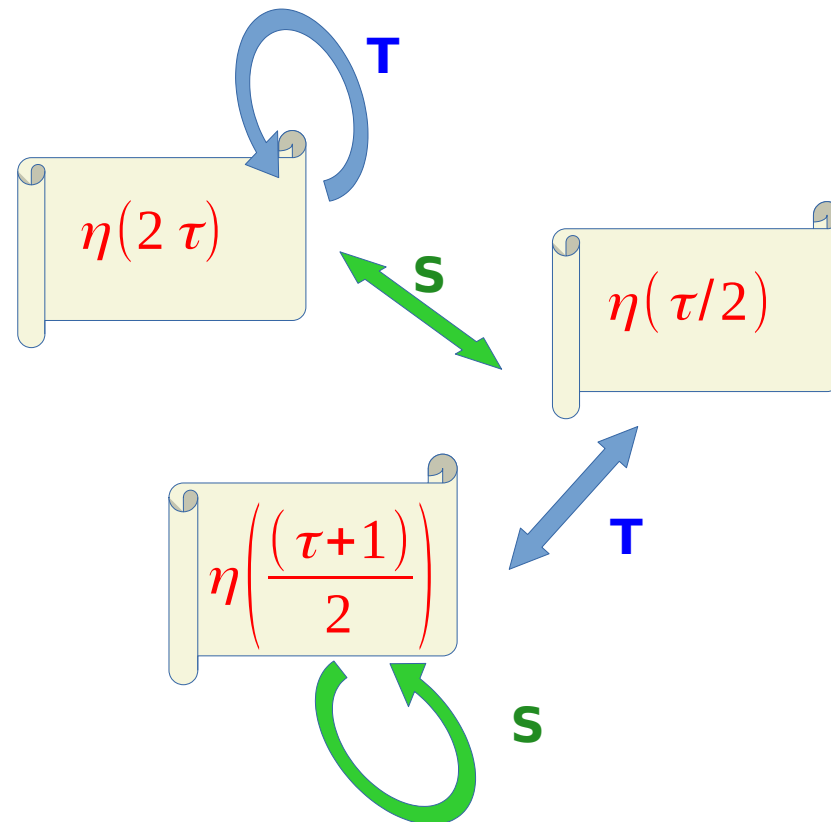


η^{24} is a modular form of weight 12

A case study: $\Gamma_2 \sim S_3$

Constructing the Modular Forms

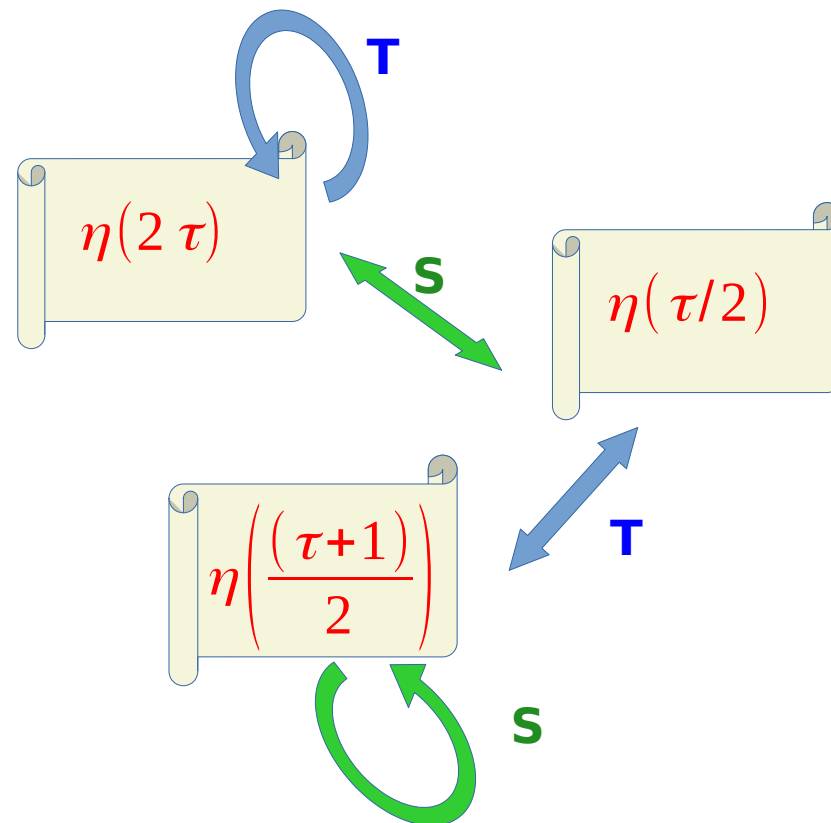
the system is closed under modular transformation



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Constructing the Modular Forms

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candidate modular form

$$Y(\alpha, \beta, \gamma) = \frac{d}{d\tau} \left[\alpha \log \eta(\tau/2) + \beta \log \eta\left(\frac{\tau+1}{2}\right) + \gamma \log \eta(2\tau) \right]$$

$$\alpha + \beta + \gamma = 0$$

A case study: $\Gamma_2 \sim S_3$

Constructing the Modular Forms

Equations to be satisfied:

$$\begin{pmatrix} Y_1(-1/\tau) \\ Y_2(-1/\tau) \end{pmatrix} = \tau^2 \rho(S) \begin{pmatrix} Y_1(\tau) \\ Y_2(\tau) \end{pmatrix}, \quad \begin{pmatrix} Y_1(\tau + 1) \\ Y_2(\tau + 1) \end{pmatrix} = \rho(T) \begin{pmatrix} Y_1(\tau) \\ Y_2(\tau) \end{pmatrix}$$



representation of generators

$$\rho(S) = \frac{1}{2} \begin{pmatrix} -1 & -\sqrt{3} \\ -\sqrt{3} & 1 \end{pmatrix}, \quad \rho(T) = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$
$$(\rho(S))^2 = \mathbb{I}, \quad (\rho(S)\rho(T))^3 = \mathbb{I}, \quad (\rho(T))^2 = \mathbb{I}$$

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$$(\rho(S))^2 = \mathbb{I}, \quad (\rho(S)\rho(T))^3 = \mathbb{I}, \quad (\rho(T))^2 = \mathbb{I}$$

$$Y_1(\alpha, \beta, \gamma) \sim Y(1, 1, -2)$$

$$Y_2(\alpha, \beta, \gamma) \sim Y(1, -1, 0)$$

$$\left. \begin{aligned} Y_1(\tau) &= \frac{i}{4\pi} \left(\frac{\eta'(\tau/2)}{\eta(\tau/2)} + \frac{\eta'((\tau+1)/2)}{\eta((\tau+1)/2)} - \frac{8\eta'(2\tau)}{\eta(2\tau)} \right) \\ Y_2(\tau) &= \frac{\sqrt{3}i}{4\pi} \left(\frac{\eta'(\tau/2)}{\eta(\tau/2)} - \frac{\eta'((\tau+1)/2)}{\eta((\tau+1)/2)} \right), \end{aligned} \right\} \text{doublet of } S_3: Y$$

A case study: $\Gamma_2 \sim S_3$

How to predict the Neutrino mass matrix (from the Weinberg operator, wrong path...)

For a satisfactory model, we ask:

1. small number of operators → ***predictability***
2. no new scalar fields beside Higgs(es) → ***symmetry breaking dictated by the vev of τ***

A case study: $\Gamma_2 \sim S_3$

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	S_3	$SU(2)$	k_i
$L_{e\mu} = (e, \mu)$	2	2	-1
L_τ	1	2	-1
H_u	1	2	0

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For a satisfactory model, we ask:

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using one power of Y (modular form of lowest weight)

	S_3	$SU(2)$	k_i
$L_{e\mu} = (e, \mu)$	2	2	-1
L_τ	1	2	-1
H_u	1	2	0

$$L = h_u^2 \left[a \left((L_{e\mu} L_{e\mu})_2, Y \right)_1 + b L_\tau (L_{e\mu} Y)_1 \right]$$



$$m_\nu = \begin{pmatrix} aY_2 & aY_1 & bY_1/2 \\ aY_1 & -aY_2 & bY_2/2 \\ bY_1/2 & bY_2/2 & 0 \end{pmatrix}$$

A case study: $\Gamma_2 \sim S_3$

How to predict the Neutrino mass matrix (from the Weinberg operator, wrong path...)

Mass matrix against the experimental data

$$m_\nu = \begin{pmatrix} aY_2 & aY_1 & bY_1/2 \\ aY_1 & -aY_2 & bY_2/2 \\ bY_1/2 & bY_2/2 & 0 \end{pmatrix}$$

$\sin^2 \theta_{12}/10^{-1}$	$2.97^{+0.17}_{-0.16}$
$\sin^2 \theta_{13}/10^{-2}$	$2.15^{+0.07}_{-0.07}$
$\sin^2 \theta_{23}/10^{-1}$	$4.25^{+0.21}_{-0.15}$
δ_{CP}/π	$1.38^{+0.23}_{-0.20}$
r	$2.92^{+0.10}_{-0.11} \times 10^{-2}$

5 observables, 2 complex parameters: a/b and τ → very difficult task!

large χ^2 of O(100) mainly driven by θ_{13}

Conclusions

Modular symmetries offer an alternative way for model building

Yukawa couplings dictated by modular forms

unified description of quarks and leptons

symmetry breaking by the vev of tau only

A lot to do:

mass hierarchy

more than one modulus

more pheno: leptogenesis, LFV...

Backup slides

Kahler potential

Under Γ :

$$\begin{cases} \tau \rightarrow \frac{a\tau + b}{c\tau + d} \\ \varphi^{(I)} \rightarrow (c\tau + d)^{-k_I} \rho^{(I)}(\gamma) \varphi^{(I)} \end{cases}$$

The invariance of the action requires the invariance of the superpotential $w(\Phi)$ and the invariance of the Kahler potential up to a Kahler transformation:

$$\begin{cases} w(\Phi) \rightarrow w(\Phi) \\ K(\Phi, \bar{\Phi}) \rightarrow K(\Phi, \bar{\Phi}) + f(\Phi) + f(\bar{\Phi}) \end{cases}$$

Kahler potential:

$$\sum_I (-i\tau + i\bar{\tau})^{-k_I} |\varphi^{(I)}|^2 \quad \rightarrow$$

modular invariant kinetic terms

$$\frac{h}{\langle -i\tau + i\bar{\tau} \rangle^2} \partial_\mu \bar{\tau} \partial^\mu \tau + \sum_I \frac{\partial_\mu \bar{\varphi}^{(I)} \partial^\mu \varphi^{(I)}}{\langle -i\tau + i\bar{\tau} \rangle^{k_I}}$$

Some definitions

a normal subgroup (also known as an invariant subgroup or self-conjugate subgroup) is a *subgroup* which is invariant under conjugation by members of the group of which it is a part:

a subgroup N of the group G is normal in G if and only if $(g n g^{-1}) \in N$ for all $g \in G$ and $n \in N$

$\Gamma(N)$, $N \geq 2$ are infinite normal subgroups of Γ , called *principal congruence subgroups*

the group $\Gamma(N)$ acts on the complex variable τ ($\text{Im } \tau > 0$)

$$\gamma \tau = \frac{a \tau + b}{c \tau + d}$$

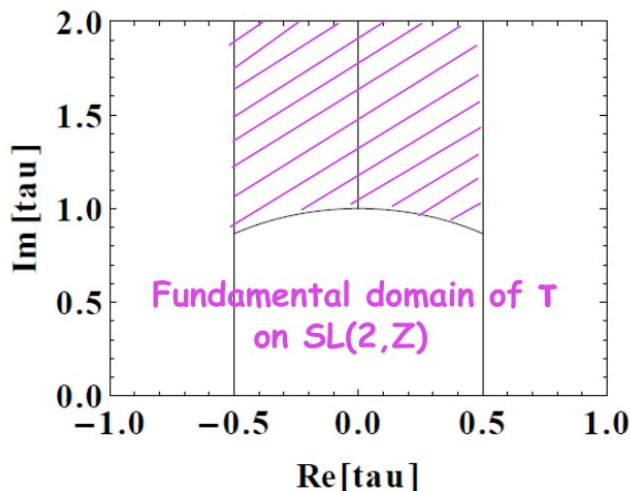
And it can be shown that the upper half-plane is mapped to itself under this action. The complex variable is henceforth restricted to have positive imaginary part

Some definitions

DEFINITION 0.2. A holomorphic function $f(z)$ on \mathbb{H} is a *modular form of level N and weight $2k$* if

- (a) $f(\alpha z) = (cz + d)^{2k} \cdot f(z)$, all $\alpha = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma(N)$;
- (b) $f(z)$ is “holomorphic at the cusps”.

Fundamental domain of τ on $SL(2, \mathbb{Z})$: connected open subset such that no two points of D are equivalent under $SL(2, \mathbb{Z})$



THEOREM 2.12. Let $D = \{z \in \mathbb{H} \mid |z| > 1, |\Re(z)| < 1/2\}$.

- (a) D is a fundamental domain for $\Gamma(1) = SL_2(\mathbb{Z})$; moreover, two elements z and z' of \bar{D} are equivalent under $\Gamma(1)$ if and only if
 - (i) $\Re(z) = \pm 1/2$ and $z' = z \pm 1$, (then $z' = Tz$ or $z = Tz'$), or
 - (ii) $|z| = 1$ and $z' = -1/z = Sz$.

A case study: $\Gamma_2 \sim S_3$

Constructing the Modular Forms

Under **T**: $Y(\alpha, \beta, \gamma) \rightarrow Y(\gamma, \beta, \alpha)$

Under **S**: $Y(\alpha, \beta, \gamma) \rightarrow \tau^2 Y(\gamma, \alpha, \beta)$

representation of generators

$$\rho(S) = \frac{1}{2} \begin{pmatrix} -1 & -\sqrt{3} \\ -\sqrt{3} & 1 \end{pmatrix}, \quad \rho(T) = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$(\rho(S))^2 = \mathbb{I}, \quad (\rho(S)\rho(T))^3 = \mathbb{I}, \quad (\rho(T))^2 = \mathbb{I}$$

A case study: $\Gamma_2 \sim S_3$

q-expansion of the Modular Forms

$$Y_1(\tau) = \frac{1}{8} + 3q + 3q^2 + 12q^3 + 3q^4 \cdots ,$$

$$Y_2(\tau) = \sqrt{3}q^{1/2}(1 + 4q + 6q^2 + 8q^3 \cdots).$$



$$Y_1(\tau) \gg Y_2(\tau)$$

for $\text{Im}(\tau) \gg 1$

Weinberg operators for $\Gamma_2 \sim S_3$

(1)

Neutrino mass matrices from the Weinberg operator

	S_3	$SU(2)$	k_i
$L_{e\mu} = (e, \mu)$	2	2	$k_{e\mu}$
L_τ	1	2	k_τ
H_u	1	2	0

Case a) $(L_{e\mu}^2)_1 \otimes (Y^2)_1, (Y^3)_1, \dots, (Y^n)_1 \Rightarrow -2k_{e\mu} + 2n = 0, \quad n=2\dots$

Case b) $(L_{e\mu}^2)_2 \otimes Y, (Y^2)_2, (Y^3)_2, \dots, (Y^n)_2 \Rightarrow -2k_{e\mu} + 2n = 0, \quad n=1\dots$

Case c) $(L_{e\mu} L_\tau)_2 \otimes Y, (Y^2)_2, (Y^3)_2, \dots, (Y^n)_2 \Rightarrow -k_{e\mu} - k_{e\tau} + 2n = 0, \quad n=1\dots$

Case d) $(L_\tau)^2 \otimes (Y^2)_1, (Y^3)_1, \dots, (Y^n)_1 \Rightarrow -2k_{e\tau} + 2n = 0, \quad n=2\dots$

Neutrino mass matrices from the Weinberg operator

(n=1)

Case b) $(L_{e\mu}^2)_2 \otimes Y, (Y^2)_2, (Y^3)_2, \dots, (Y^n)_2 \quad \Rightarrow \quad -2k_{e\mu} + 2n = 0, \quad n=1\dots$

Case c) $(L_{e\mu}L_{e\tau})_2 \otimes Y, (Y^2)_2, (Y^3)_2, \dots, (Y^n)_2 \quad \Rightarrow \quad -k_{e\mu} - k_{e\tau} + 2n = 0, \quad n=1\dots$

Solutions:

$[k_{e\mu}=1 \quad k_{e\tau}=0] \quad [k_{e\mu}=0 \quad k_{e\tau}=2] \quad [k_{e\mu}=1 \quad k_{e\tau}=1]$

$$m_\nu = \begin{pmatrix} bY_2 & bY_1 & cY_1/2 \\ bY_1 & -bY_2 & cY_2/2 \\ cY_1/2 & cY_2/2 & 0 \end{pmatrix}$$

Neutrino mass matrices from the Weinberg operator

(n=2)

Case a) $\rightarrow -2k_{e\mu} + 4 = 0$

Case c) $\rightarrow -k_{e\mu} - k_{e\tau} + 4 = 0$

Case b) $\rightarrow -2k_{e\mu} + 4 = 0$

Case d) $\rightarrow -2k_{e\tau} + 4 = 0$

Solutions:

$[k_{e\mu} = 2 \quad k_{e\tau} = 2] \quad [k_{e\mu} = 2 \quad k_{e\tau} \neq 2] \quad [k_{e\mu} \neq 2 \quad k_{e\tau} = 2]$

$$m_\nu = \begin{pmatrix} (a+b)y_1^2 + (a-b)y_2^2 & 2by_1y_2 & cy_1y_2 \\ * & (a-b)y_1^2 + (a+b)y_2^2 & 1/2c(y_1^2 - y_2^2) \\ * & * & d(y_1^2 + y_2^2) \end{pmatrix}$$

A case study: $\Gamma_2 \sim S_3$

Dedekind eta functions

$$\eta(\tau) = q^{1/24} \prod_{n=1}^{\infty} (1 - q^n) \quad q \equiv e^{i2\pi\tau}$$

Under **T**:

$$\left\{ \begin{array}{l} \eta(2\tau) \rightarrow e^{i\pi/6} \eta(2\tau) \\ \eta(\tau/2) \rightarrow \eta((\tau+1)/2) \\ \eta((\tau+1)/2) \rightarrow e^{i\pi/12} \eta(\tau/2) \end{array} \right.$$

Under **S**:

$$\left\{ \begin{array}{l} \eta(2\tau) \rightarrow \sqrt{-i\tau/2} \eta(\tau/2) \\ \eta(\tau/2) \rightarrow \sqrt{-2i\tau} \eta(2\tau) \\ \eta\left(\frac{(\tau+1)}{2}\right) \rightarrow e^{-i\pi/12} \sqrt{-i\tau(\sqrt{3}-i)} \eta\left(\frac{(\tau+1)}{2}\right) \end{array} \right.$$

Mod

$\text{Id}[a_, b_] := \{\{\text{Mod}[a, b], 0\}, \{0, \text{Mod}[a, b]\}\}$

$$\text{Id}[-1, 2] \quad \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\text{Id}[-1, 3] \quad \begin{pmatrix} 2 & 0 \\ 0 & 2 \end{pmatrix}$$