

Separating ν and anti- ν components with non-magnetized detectors

Joe Grange
New Perspectives 2012

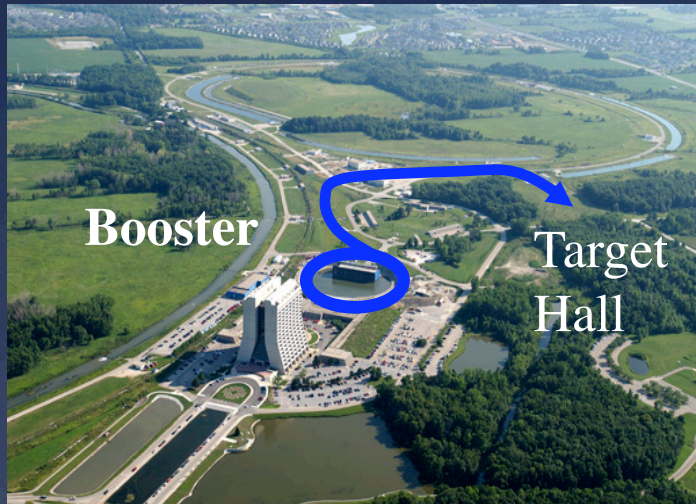


Outline

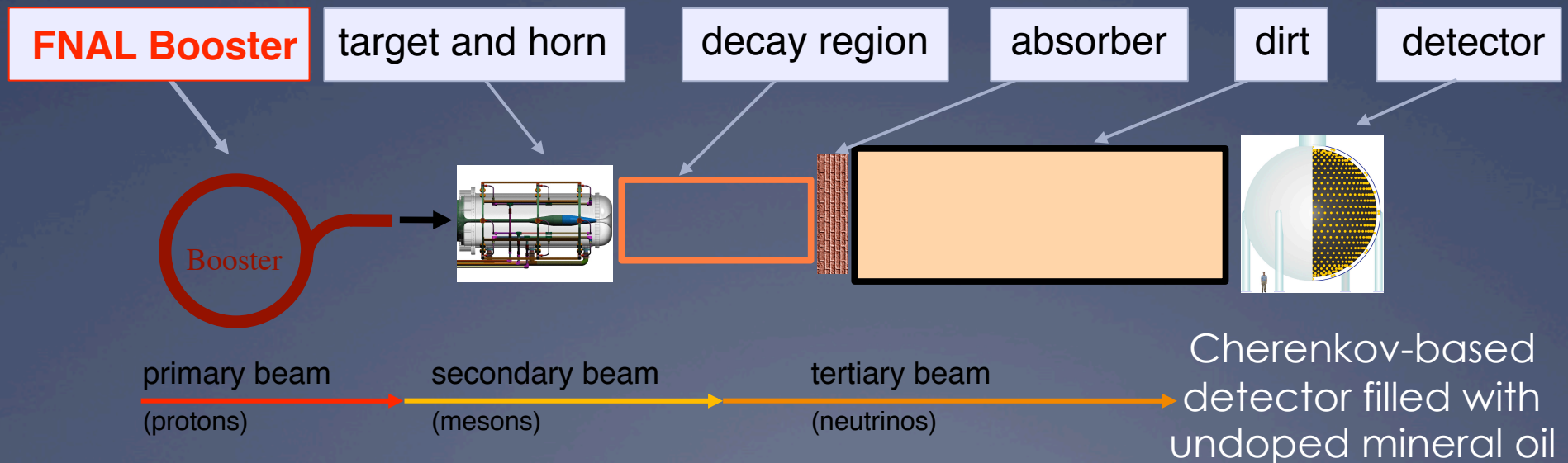
1. Booster Neutrino Beam (BNB)
2. Three measurements of ν_μ flux in BNB $\bar{\nu}_\mu$ beam
3. Community interest, conclusions

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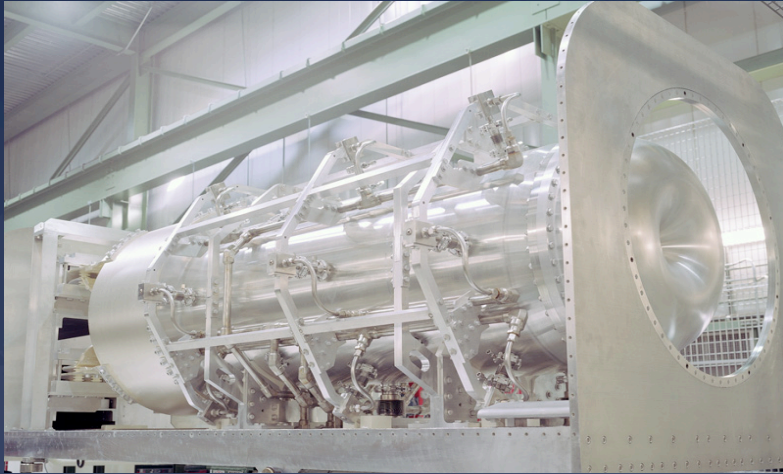
Booster Neutrino Beam



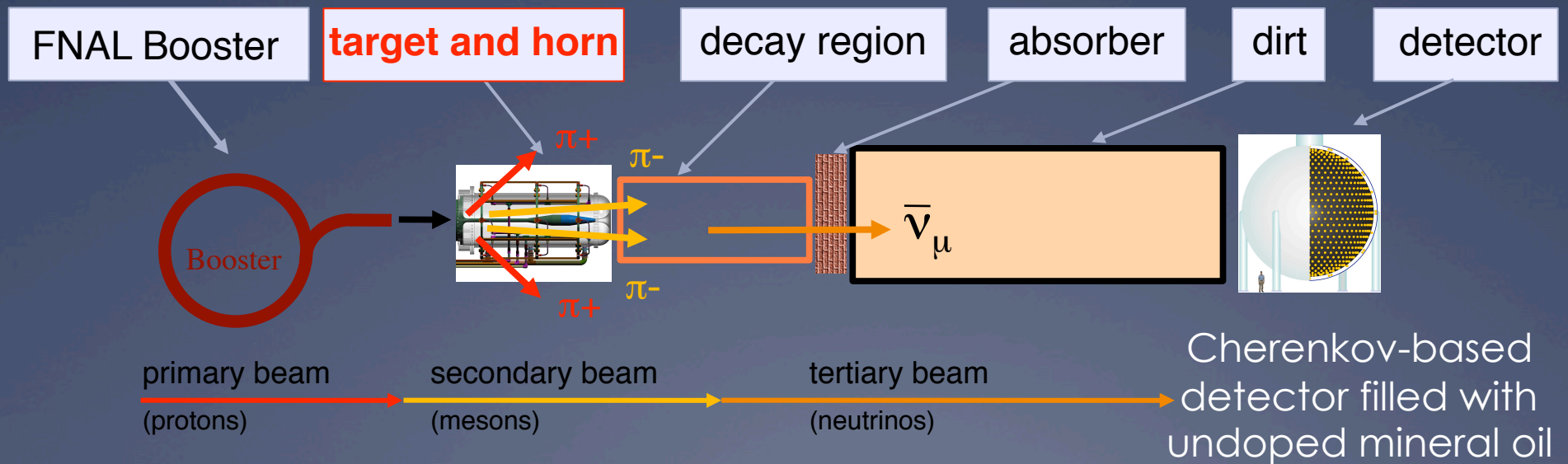
8.9 GeV/c momentum protons
extracted from Booster, steered
toward a Beryllium target in
bunches of 5×10^{12} at a maximum
rate of 5 Hz



Booster Neutrino Beam

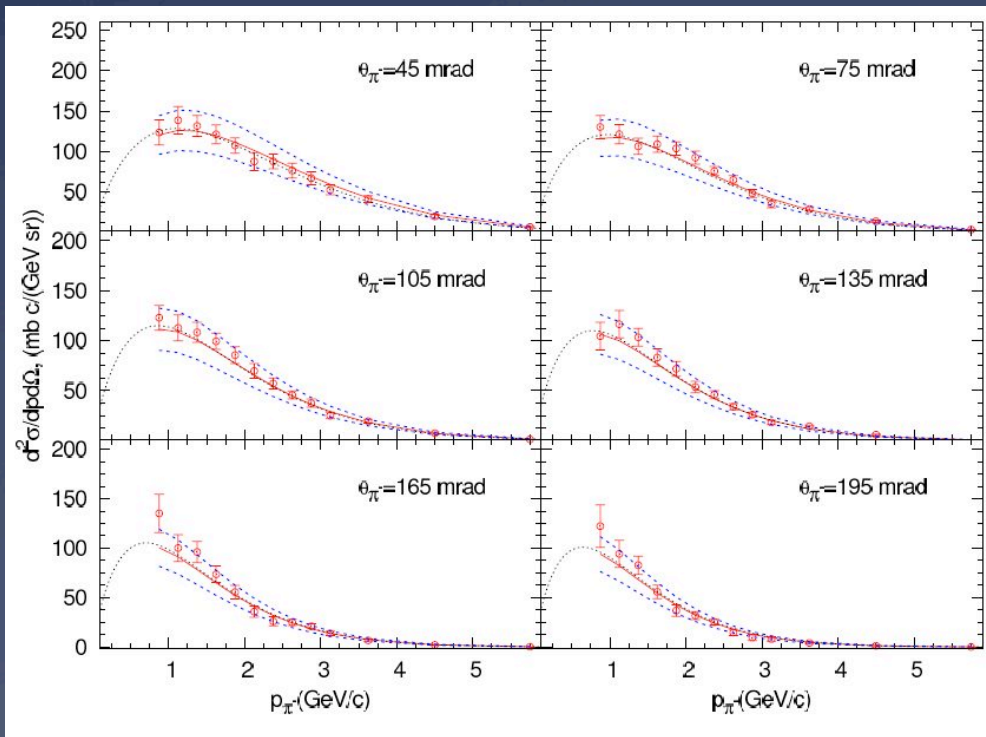


Magnetic horn with reversible polarity focuses either neutrino or anti-neutrino parent mesons ("neutrino" vs "anti-neutrino" mode)

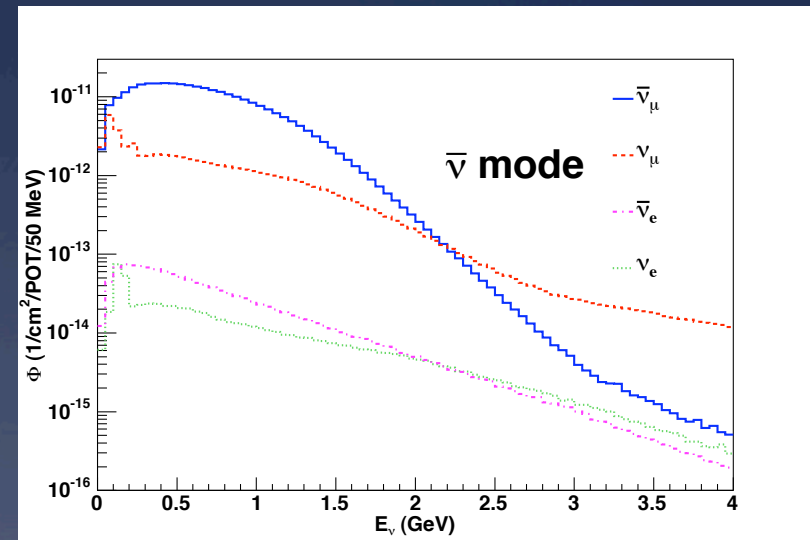


MiniBooNE Flux

- * Flux prediction based exclusively on external data - no *in situ* tuning



HARP collaboration,
Eur. Phys. J. C52 29 (2007)

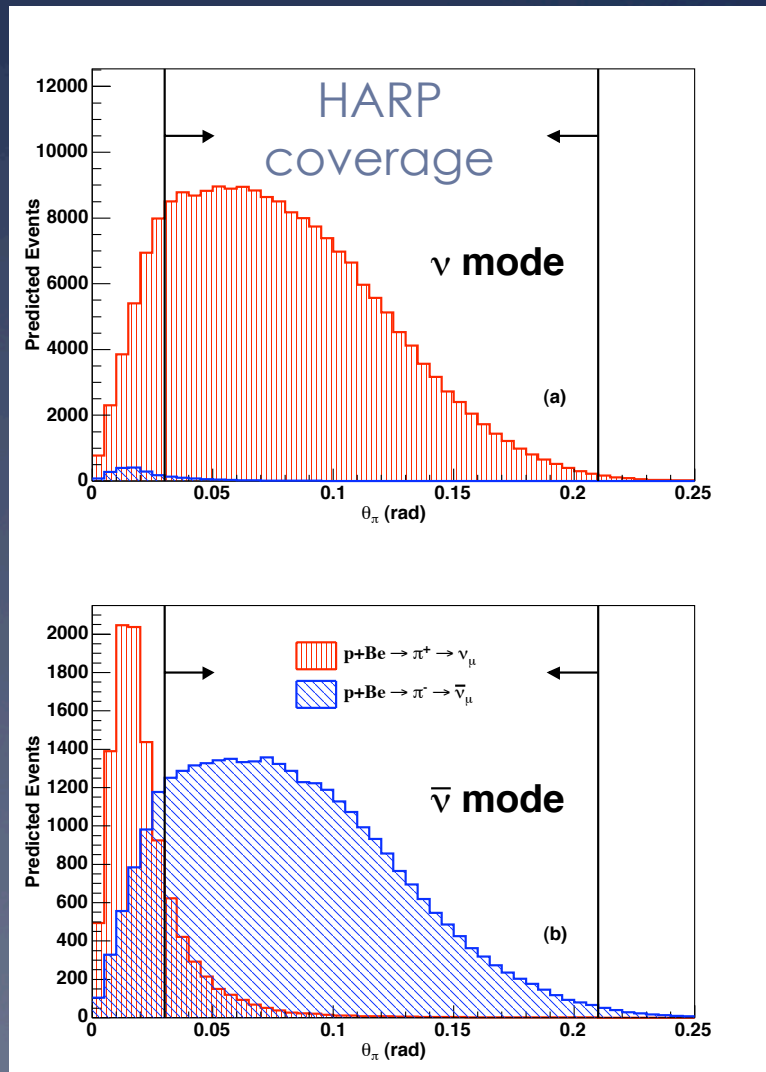


MiniBooNE collaboration,
Phys. Rev. D79, 072002 (2009)

- * Dedicated pion production data taken by HARP experiment to predict neutrino flux at MiniBooNE
- * A spline fit to these data brings flux uncertainty to $\sim 9\%$

MiniBooNE Flux

- * ~9% errors only true for pions produced in HARP-covered phase space
- * Due to large proton background, pion production below 30 mrad not reported
- * While not a serious issue for neutrino mode (top plot), severe complication for anti-neutrino mode (bottom)



MiniBooNE Flux

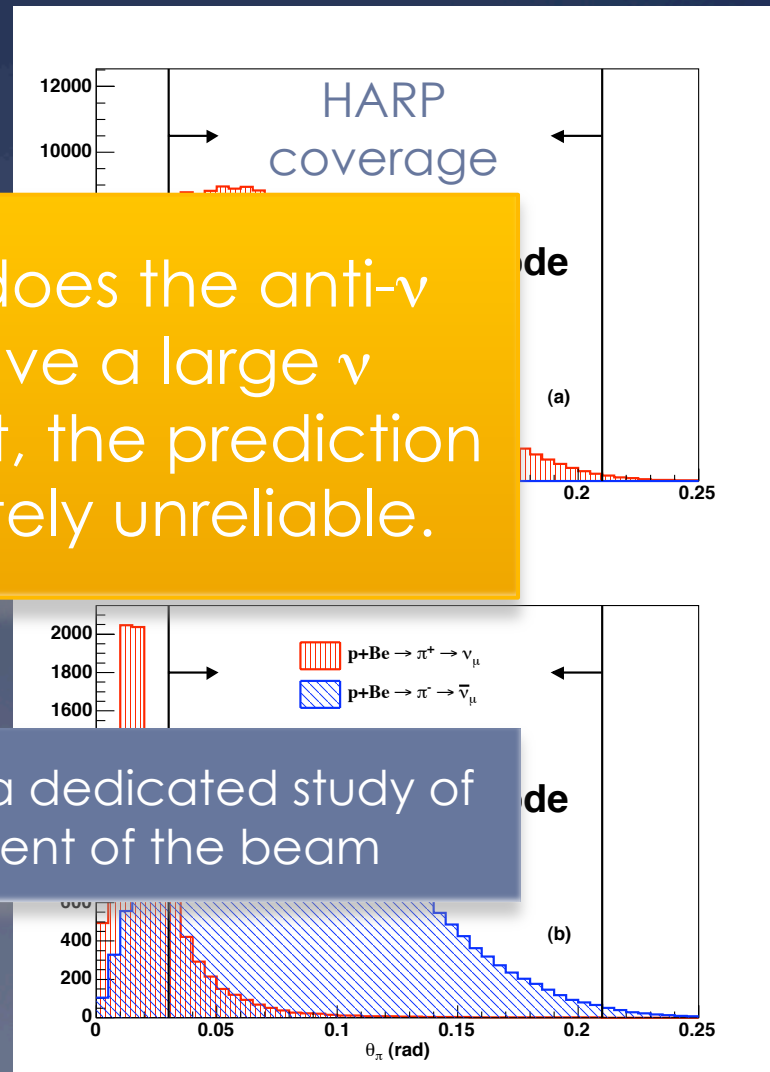
- * ~9% errors only true for pions produced in HARP-coverage space

- * Due to large background production 30 mrad not

- * While not a serious issue for neutrino (plot), severe complication neutrino mode (bottom)

Not only does the anti- ν data have a large ν component, the prediction is completely unreliable.

This motivates a dedicated study of the ν_μ content of the beam



1. Booster Neutrino Beam (BNB)
2. Three measurements of ν_μ flux in BNB $\bar{\nu}_\mu$ beam
3. Community interest, conclusions

Wrong-sign measurements

- * Three independent and complementary measurements of the wrong-sign background:
 1. Fitting the angular distribution of the CCQE sample for the neutrino and anti-neutrino content
 2. Comparing predicted to observed event rates in the $\text{CC}\pi^+$ sample
 3. Measuring how often muon decay electrons are produced (exploits μ^- nuclear capture)

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First measurement of the ν_μ content of a $\bar{\nu}_\mu$ beam using a non-magnetized detector.

Phys. Rev. D81: 072005 (2011)

Wrong-sign measurements

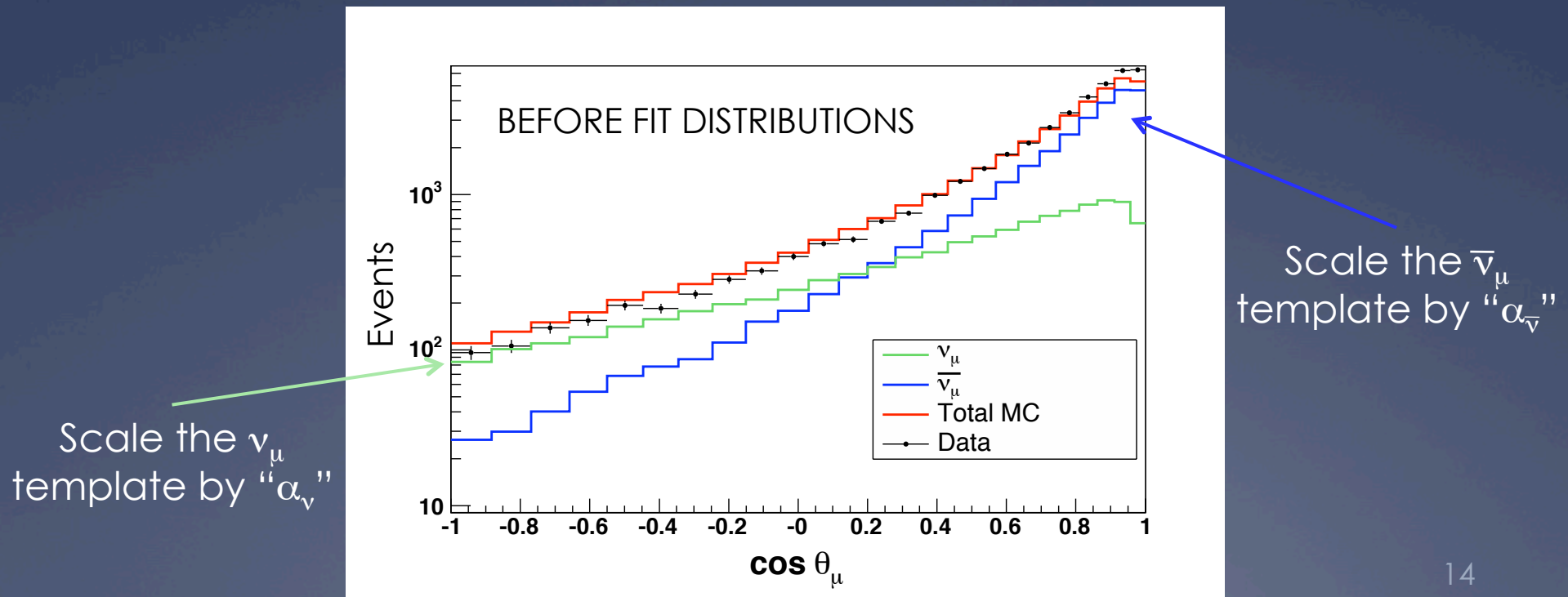
- * General strategy: isolate samples sensitive to the ν_μ beam content, apply the measured cross sections from neutrino mode (CCQE, CC π^+)
 - * *Crucial* application of BoONE-measured ν_μ σ 's
- * The level of data-simulation agreement then reflects the accuracy of the ν_μ flux prediction

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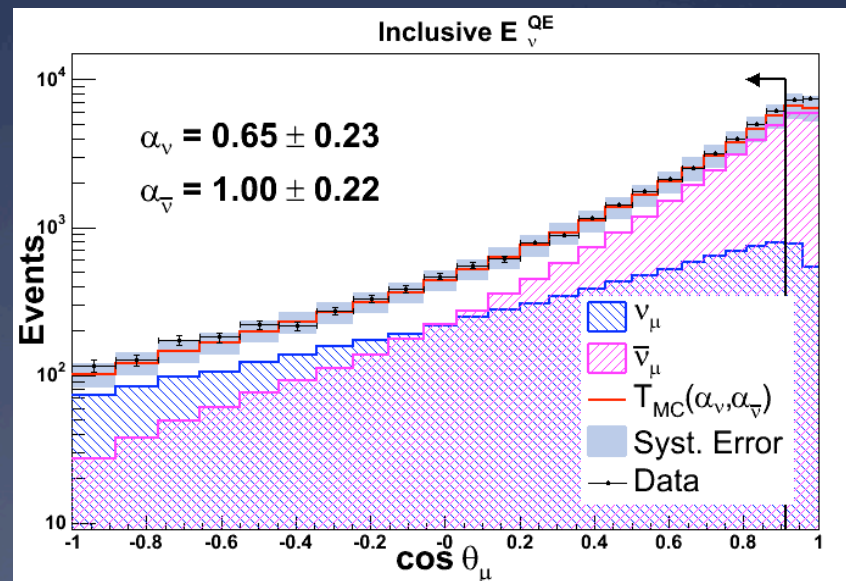
Fitting the outgoing muon angular distribution

- * We form a linear combination of the neutrino and anti-neutrino content to compare with CCQE data:



Fitting the outgoing muon angular distribution

- * Results indicate the ν_μ flux is over-predicted by $\sim 30\%$
- * Fit also performed in bins of reconstructed energy; consistent results indicate flux spectrum shape is well modeled

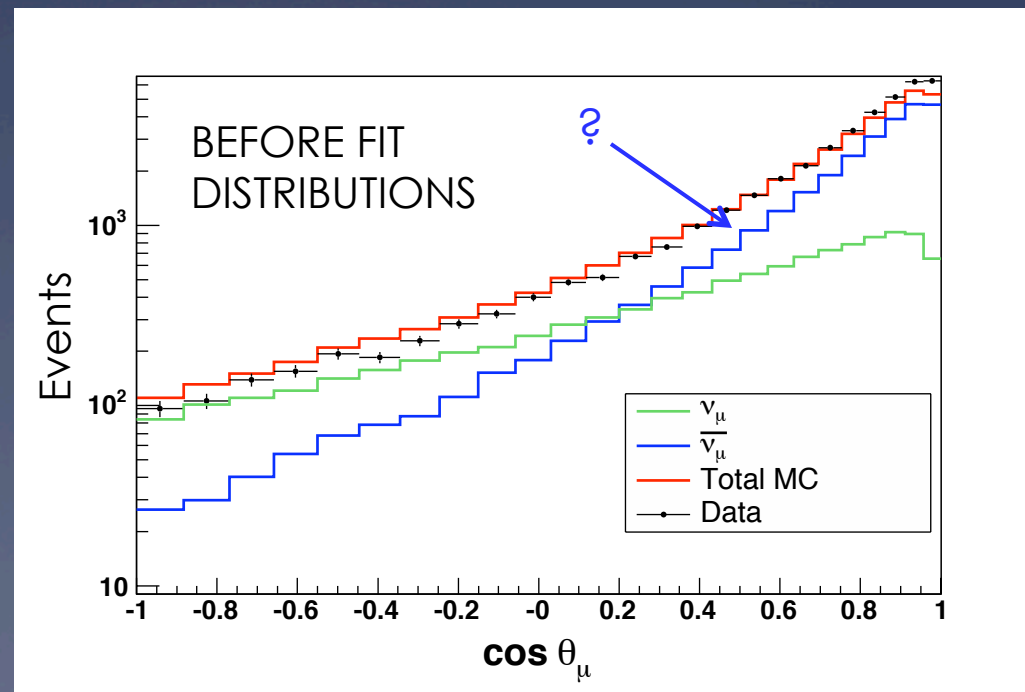


$E_\nu^{QE}(\text{MeV})$	α_ν	$\alpha_{\bar{\nu}}$
< 600	0.65 ± 0.22	0.98 ± 0.18
$600 - 900$	0.61 ± 0.20	1.05 ± 0.19
> 900	0.64 ± 0.20	1.18 ± 0.21
Inclusive	0.65 ± 0.23	1.00 ± 0.22

Model dependence

- * Though the ν_μ CCQE scattering template is known (from our measurement), the result is correlated to the (unknown) anti- ν_μ distribution and therefore biased

- * In the future, thanks to current expt's, σ 's will be much better known and this technique could be very powerful



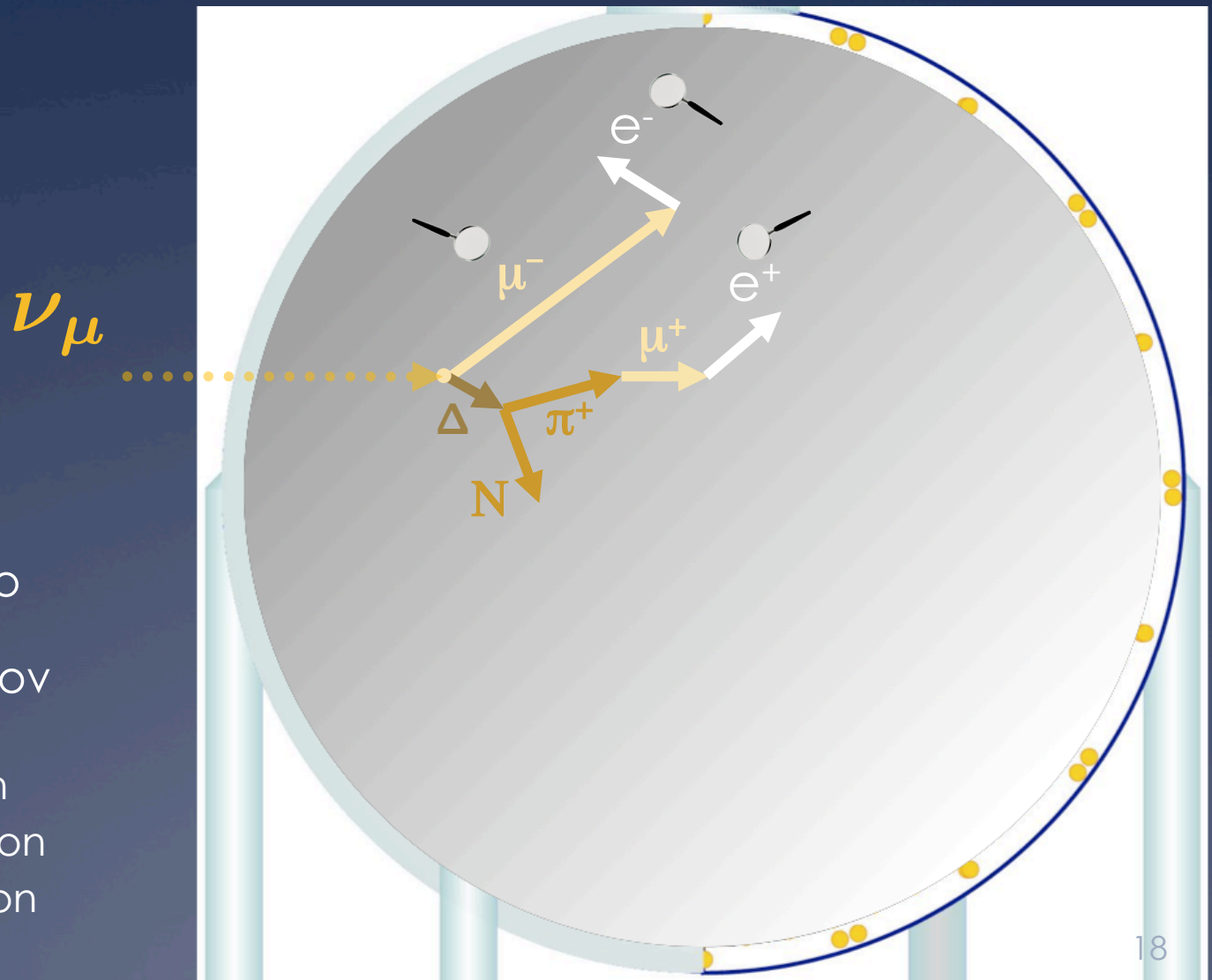
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CC π^+ sample formation

* The neutrino induced resonance channel leads to three leptons above Cherenkov threshold

1. Primary muon
2. Decay electron
3. Decay positron



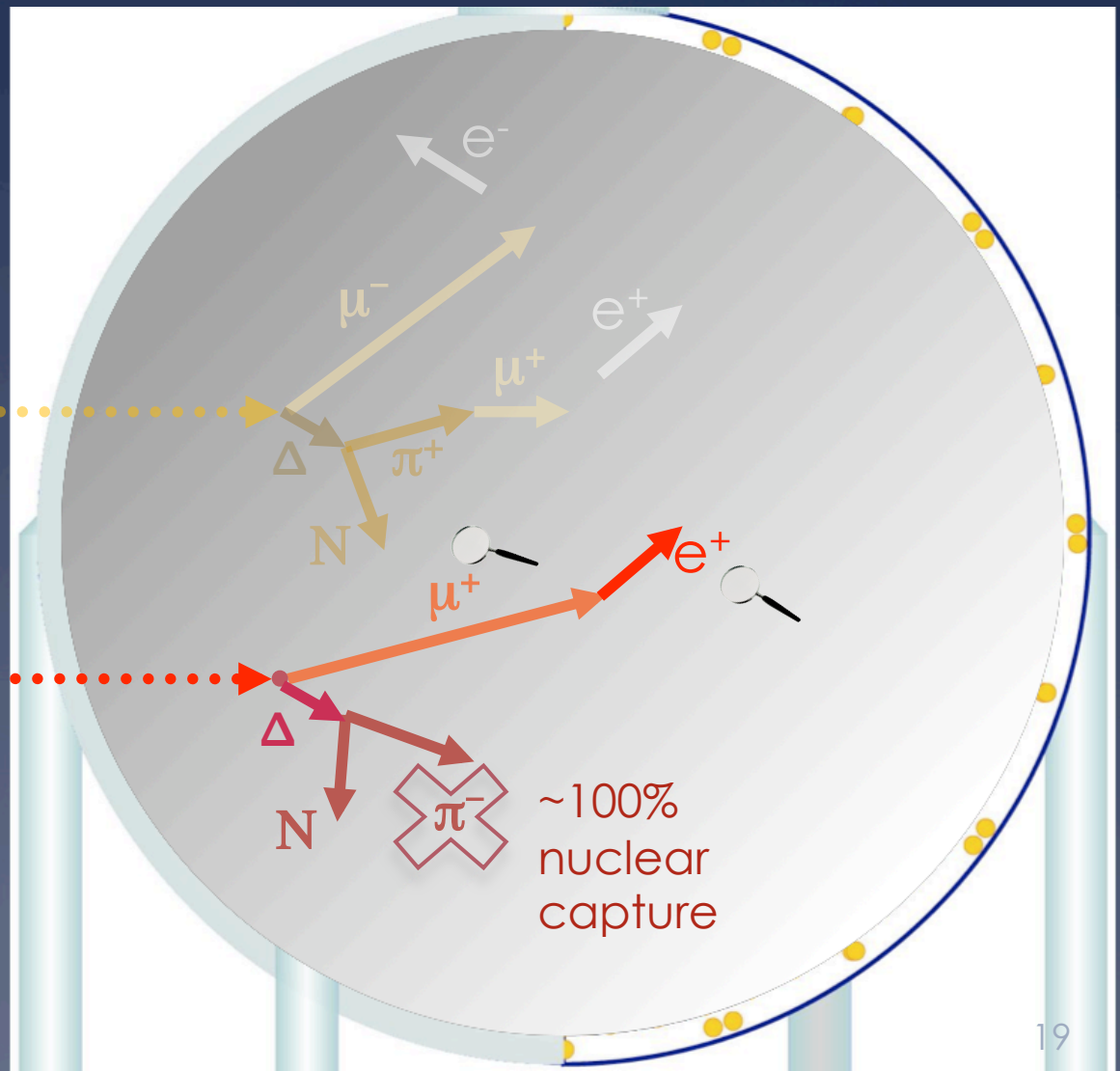
CC π^+ sample formation

* Due to nuclear π^- capture, the corresponding anti-neutrino interaction has only two:

1. Primary muon
2. Decay positron

ν_μ

$\bar{\nu}_\mu$



CC π^+ ν_μ flux measurement

- * With the simple requirement of two decay electrons subsequent to the primary muon, we isolate a sample that is $\sim 80\%$ neutrino-induced.
- * Data/simulation ratios in bins of reconstructed energy indicate the neutrino flux is over-predicted in normalization, while the spectrum shape is consistent with the prediction

E_ν^Δ (MeV)	ν_μ Φ scale
600 - 700	0.65 ± 0.10
700 - 800	0.79 ± 0.10
800 - 900	0.81 ± 0.10
900 - 1000	0.88 ± 0.11
1000 - 1200	0.74 ± 0.10
1200 - 2400	0.73 ± 0.15
Inclusive	0.76 ± 0.11

Wrong-sign measurements

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 1. Fitting the angular distribution of the CCQE sample for the neutrino and anti-neutrino content
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μ^- capture measurement

- * We isolate a $> 90\%$ CC sample for both μ^- -only and μ^+e samples
- * CC events typically observe both μ^+e - two reasons why we may not observe the decay electron:
 1. Michel electron detection efficiency
 2. μ^- nuclear capture (ν_μ CC events only)

μ^- capture measurement

- * By requiring $(\mu\text{-only}/\mu+e)^{\text{data}} = (\mu\text{-only}/\mu+e)^{\text{MC}}$ and normalization to agree in the $\mu+e$ sample we can calculate a ν_μ flux scale α_ν and a rate scale $\alpha_{\bar{\nu}}$

$$\frac{\mu}{\mu + e}^{\text{data}} = \left(\frac{\alpha_\nu \nu^\mu + \alpha_{\bar{\nu}} \bar{\nu}^\mu}{\alpha_\nu \nu^{\mu+e} + \alpha_{\bar{\nu}} \bar{\nu}^{\mu+e}} \right)^{\text{MC}}$$

Predicted neutrino content in the $\mu+e$ sample, for example

μ^- capture measurement

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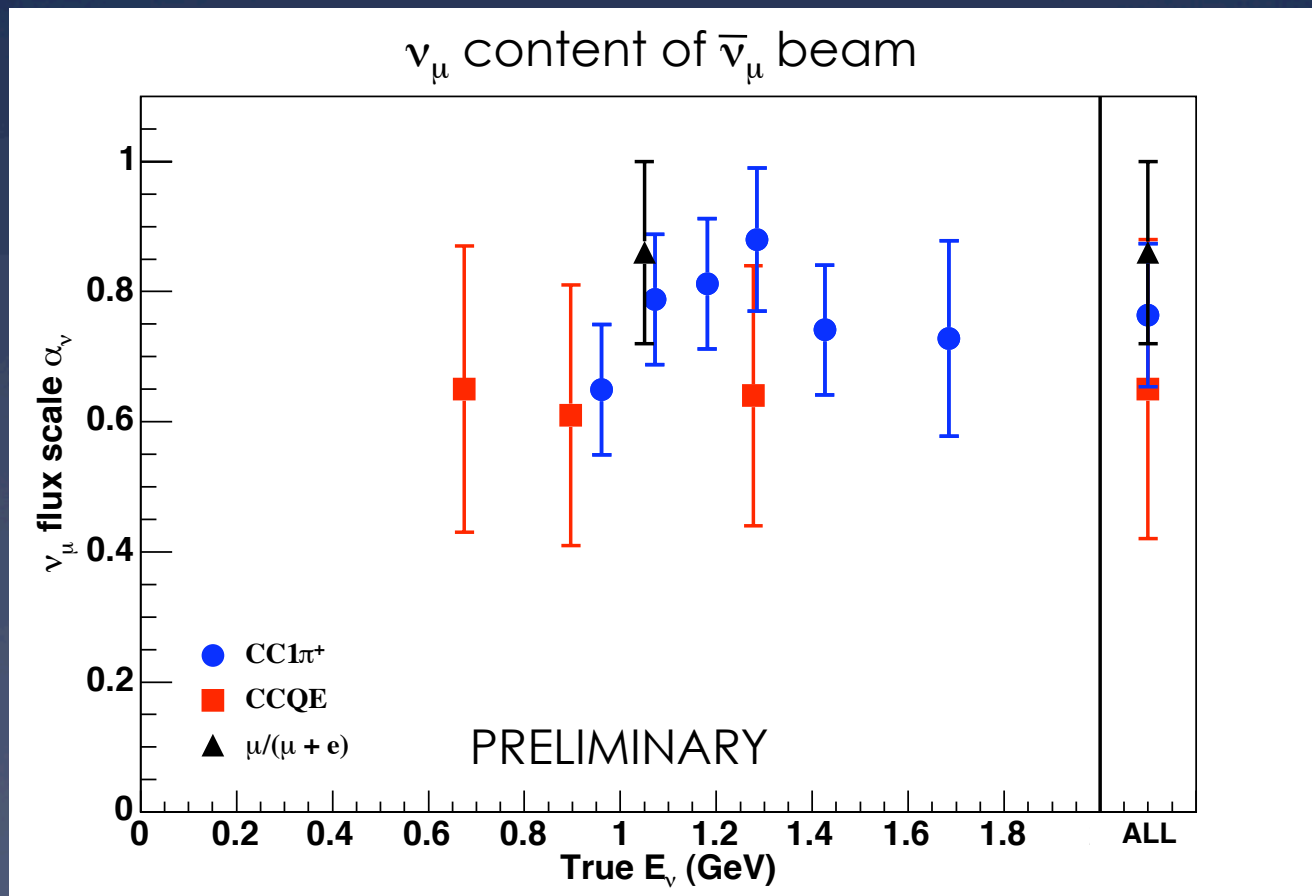
$$\frac{\mu}{\mu + e}^{\text{data}} = \left(\frac{\alpha_\nu \nu^\mu + \alpha_{\bar{\nu}} \bar{\nu}^\mu}{\alpha_\nu \nu^{\mu+e} + \alpha_{\bar{\nu}} \bar{\nu}^{\mu+e}} \right)^{\text{MC}}$$

Results:

$$\alpha_\nu = 0.86 \pm 0.14$$
$$\alpha_{\bar{\nu}} = 1.09 \pm 0.23$$

PRELIMINARY

Neutrino flux measurement summary



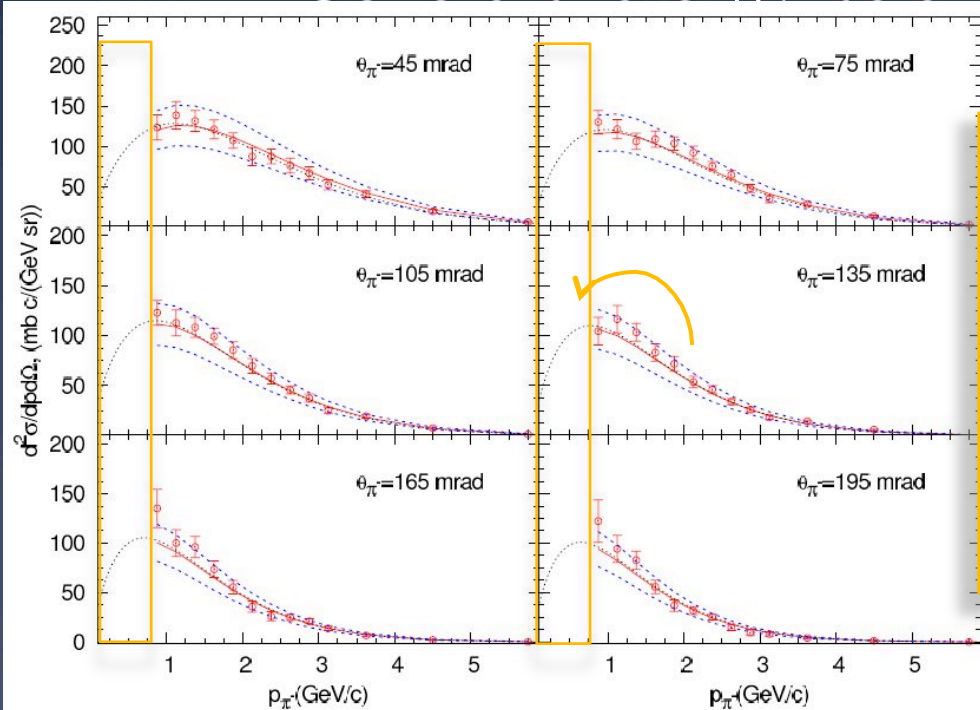
- * Discrepancy with prediction appears to be in normalization only - flux shape is well modeled

Strategy revisited

- * General strategy: isolate samples sensitive to the ν_μ beam content, apply the measured cross sections from neutrino mode (CCQE, CC π^+)
 - * *Crucial* application of BoONE-measured ν_μ σ 's
- * The level of data-simulation agreement then reflects the accuracy of the ν_μ flux prediction

Strategy revisited

* General strategy: isolate samples sensitive to the
the measured cross



Takes hadro-production
data, uses it to place similar
errors on the flux region
not measured!

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Who else cares?

- * Anyone using anti- ν beams without B-fields!
 - * Nova
 - * T2K far detector
- * LBNE: yesterday we heard “preferred reconfiguration” is FD at Homestake without near detector. If no B-field, μ^- capture technique could be very powerful in WS discrimination (argon: ~75% capture, carbon: ~8%!)
 - * almost event-by-event discrimination without B-field
- * Minerva: can get powerful stat increases if use μ 's stopped in main detector

Conclusions

- * Though MiniBooNE is unmagnetized, model-independent statistical techniques measure the ν_μ content in the ν_μ beam to $\sim 15\%$ uncertainty
- * This is the first demonstration of a set of techniques that could well be used in the near future for CP-violation, mass hierarchy and σ measurements

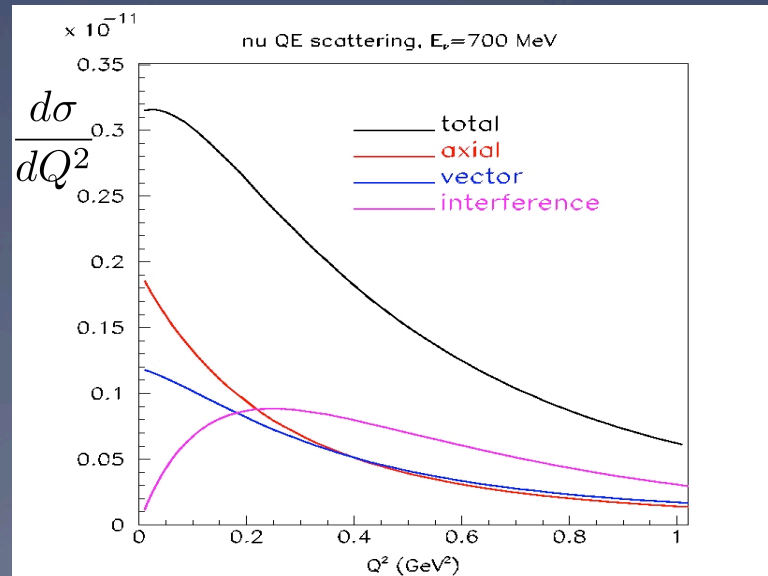
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Fitting the outgoing muon angular distribution

- * Neutrino vs anti-neutrino CCQE cross sections differ exclusively by an interference term that changes sign between the two

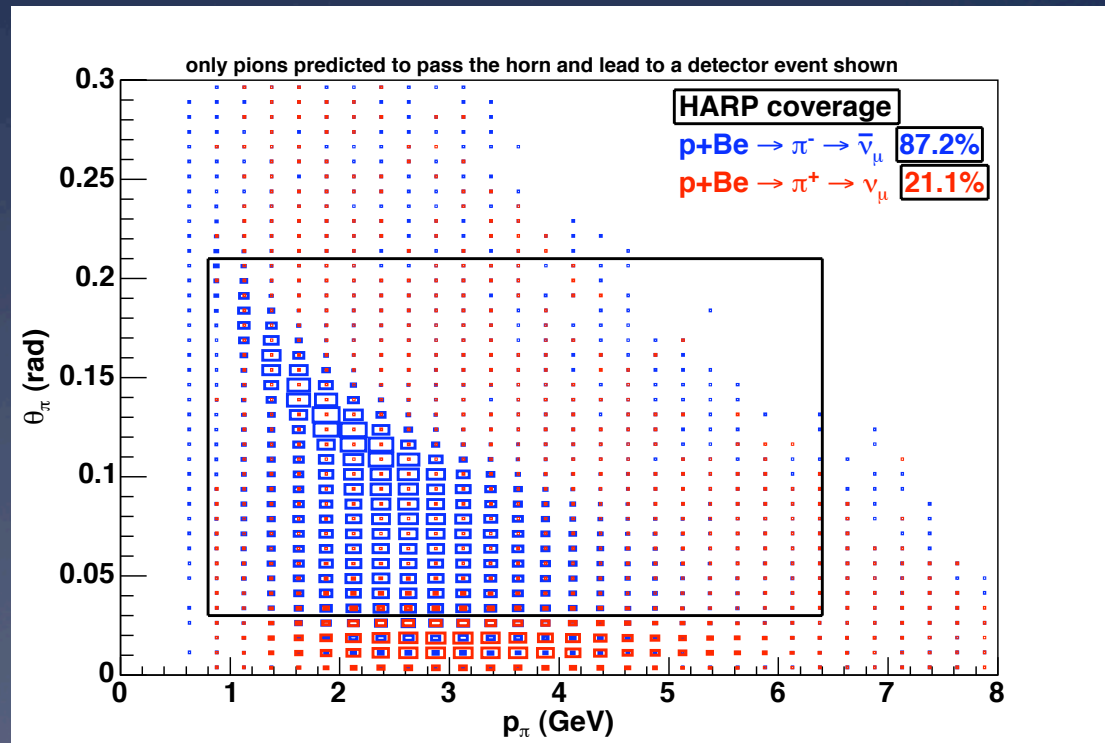
$$\frac{d\sigma}{dQ^2} = \frac{M^2 G_F^2 |V_{ud}|^2}{8\pi E_\nu^2} \left[A(Q^2) \boxed{\pm} B(Q^2) \left(\frac{s-u}{M^2} \right) + C(Q^2) \left(\frac{s-u}{M^2} \right)^2 \right]$$

- * The divergence is more pronounced at higher Q^2 , which is strongly correlated with backward scattering muons



How wrong signs contribute to flux

- * Wrong-sign pions escape magnetic deflection and contribute to the anti-neutrino beam via low angle production



- * In anti-neutrino mode low-angle production is a *crucial* flux region and we do not have a reliable prediction

This motivates a dedicated study of ν_μ content of the beam

Why so different?

- * Cross section: at MiniBooNE energies ($E_\nu \sim 1$ GeV), neutrino cross section ~ 3 x higher than anti-neutrino

$$\frac{d\sigma}{dQ^2} = \frac{M^2 G_F^2 |V_{ud}|^2}{8\pi E_\nu^2} \left[A(Q^2) \boxed{\pm} B(Q^2) \left(\frac{s-u}{M^2} \right) + C(Q^2) \left(\frac{s-u}{M^2} \right)^2 \right]$$

- * Flux: leading particle effect creates ~ 2 x as many π^+ as π^-



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