THREE FLAVOR NEUTRINO OSCILLATION RESULTS FROM NOvA

[robust measurements in an uncertain world]



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Joint Experimental-Theoretical Seminar Fermilab September 18, 2020



Preface: "an uncertain world"

[this slide reflects the opinions of the author only]

Science happens in a context.

Acknowledge NOvA collaborators & Fermilab community for whom the last few months have been even less "business as usual":

- Members of the Black community
- Those who have lost family/close friends/etc. to COVID-19



[Tyler LaRiviere/Chicago Sun-Times]



Let's listen to & support them. (There's work to be done!)

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Plan for this talk

- Why study neutrino oscillations?
- NOvA: observing oscillations with a long-baseline experiment
- Interpretations: events into probabilities
- Inferences: NOvA oscillation results

Part I: why study neutrino oscillations?

I suppose that I tend to be optimistic about the future of physics. And nothing makes me more optimistic than the discovery of broken symmetries.

—Steven Weinberg (Nobel Lecture 1979)

Why study neutrino oscillations?

Despite our relatively complete picture of the fundamental particles & interactions, many questions remain.

Neutrino oscillations help us approach many of them...



Why study neutrino oscillations?

- v oscillations probe <u>symmetries</u>:
 - Do the neutrino flavors mix in a predictable way?
 - Are the neutrino masses distributed in a 'regular' fashion? ... in a fashion resembling the other fermions?
 - Do neutrinos and antineutrinos oscillate the same way?
 - Are there exactly 3 v flavors, like the charged leptons and quarks?



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Does this help us understand "generations?"

Maybe insights into neutrino mass generation mechanism?

If not, can we learn more about other matterantimatter differences? CP violation in leptons??

If not... where do the extras come from?

Oscillations arise from transitions from one neutrino flavor eigenstate to another.

$$P(\mathbf{v}_{\mathbf{o}} \to \mathbf{v}_{\mathbf{b}}) = \left| \sum_{j} U_{\beta j}^{*} e^{-i \frac{m_{j}^{2} L}{2E}} U_{\alpha j} \right|^{2}$$



Oscillations depend on two essential ingredients:

$$P(v_{\alpha} \to v_{\beta}) = \left[\begin{array}{c} \sum_{j} U_{\beta j}^{*} e^{-i \frac{m_{j}^{2} L}{2E}} U_{\alpha j} \\ \sum_{j} U_{\mu 1}^{*} U_{\mu 2} & U_{\mu 3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{array} \right] \left[\begin{array}{c} v_{1} \\ v_{2} \\ v_{3} \end{array} \right]$$

PMNS matrix ("mixing" matrix: connects flavor to mass states) must have off-diagonal elements

Oscillations depend on two essential ingredients:



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Focus today on questions addressed by <u>3-flavor long-</u> <u>baseline accelerator</u> <u>neutrino oscillations</u>:

> E_v ~ few GeV L ~ 100s of km



 Is there a symmetry governing the ordering of the lepton mass states?
 Is the most electron-like state the lightest one, like with the charged leptons?



Most poorly known parameters: θ_{23} (~5%), δ_{CP} (weak constraints)



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Part II: Observing neutrinos over a long baseline with NOvA

$$P(v_{\alpha} \to v_{\beta}) = \left| \sum_{j} U_{\beta j}^{*} e^{-i \frac{m_{j}^{2}L}{2E}} U_{\alpha j} \right|^{2}$$



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$$P(v_{\alpha} \to v_{\beta}) = \left| \sum_{j} U_{\beta j}^{*} e^{-i \frac{m_{j}^{2}L}{2E}} U_{\alpha j} \right|^{2}$$

$$P_{\nu_{\mu} \to \nu_{\mu}}^{(-)} \approx 1 - \sin^2 2 \theta_{23} \sin^2 \left(\Delta m_{32}^2 \frac{L}{4E} \right)$$





$$P(\mathbf{v}_{\alpha} \rightarrow \mathbf{v}_{\beta}) = \left| \sum_{j} U_{\beta j}^{*} e^{-i\frac{m_{j}^{2}L}{2E}} U_{\alpha j} \right|^{2}$$

$$\mathbf{v}_{\mu} \operatorname{disappearance} \quad \mathbf{v}_{e} \operatorname{appearance}$$

$$P_{\nabla_{\mu}^{\downarrow} \rightarrow \nabla_{\mu}^{\downarrow}} \approx 1 - \sin^{2} 2 \theta_{23} \sin^{2} \left(\Delta m_{32}^{2} \frac{L}{4E} \right) \quad P_{\nabla_{\mu}^{\downarrow} \rightarrow \nabla_{e}^{\downarrow}} \approx \sin^{2} 2 \theta_{13} \sin^{2} \theta_{23} \sin^{2} \frac{(A-1)\Delta}{(A-1)^{2}}$$

$$+ 2 \alpha \sin \theta_{13} \sin 2 \theta_{12} \sin 2 \theta_{23}$$

$$(1) \quad \mathbf{v}_{e} \rightarrow \Delta m_{32}^{2}$$

$$(2) \quad \mathbf{v}_{e} \rightarrow \theta_{23}$$

$$(3) \quad \mathbf{v}_{e} \rightarrow \delta_{CP}$$

$$\alpha = \frac{\Delta m_{31}^{2}}{\Delta m_{31}^{2}} \quad \Delta = \Delta m_{31}^{2} \frac{L}{4E} \quad A = \mp G_{I} N_{e} \frac{L}{\sqrt{2\Delta}} \quad (2)$$





Where do neutrinos come from?



K. ENGMAN/SCIENCE 345, 6204

Where do neutrinos come from?



Where do neutrinos come from?



MW-capable target (installed 2019)



MW-capable horn (installing during 2020 shutdown)

Working towards 900+ kW

- With 2019-2020 improvements to NuMI beamline components, complex will support ~800 KW
- Early PIP-II improvements to Booster will allow 900+ KW with faster cycle times

Detecting neutrinos







- <u>Near Detector</u>: 300 ton, 1 km from source (FNAL)
 - 100m underground, 20K channels
- Far Detector: 14 kton, 810 km from source (Ash River, MN)
 - On the surface, 3m concrete+barite overburden; 344K channels

Detecting neutrinos





- Good energy resolution for muons, electromagnetic & hadron showers:
 - Mostly (65%) active detector
 - Radiation length ~ 40 cm → 6 samples per radiation length

Detectors differ mainly in size (otherwise functionally identical)

The task:

Get from this...

... to this



FD sits on the surface \rightarrow ~150 KHz cosmics

 v_{e} candidate



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- Workhorse tool: convolutional neural network (CNN) called CVN
 - Technique borrowed from computer vision community:
 - Learns topological "features"
 - Eventually mapped onto desired output categories
 - Performs neutrino event classification; also important part of cosmic rejection
 - Updated for 2020 (3rd edition!)
 - Significantly (~×3) faster network architecture (modified MobileNet \vee 2)
 - Slightly better physics performance
 - Further reading: <u>JINST 11, P09001</u>



v_e candidates near detector edges are recovered into "peripheral" sample using tighter PID cut & dedicated cosmic BDT





Goal #2: Measuring E,

Strategy: divide and conquer



Goal #2: Measuring $E_v(v_\mu)$

Strategy: divide and conquer



Goal #2: Measuring $E_v(v_e)$

Strategy: divide and conquer


What do we observe?



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Part III: Interpretations

Best fits

The <u>principle</u> of finding the values of the oscillation parameters most compatible with the data is straightforward...



(use a fitter to find the values that minimize the data-prediction difference)

Making predictions

... but there's a lot hiding under the hood of the predicted specta



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Neutrino interaction modeling



- For 2020: upgrade to GENIE 3.0.6
 - Introduces choices of prepackaged collections of models, often with tuning to data
 - We choose a "theory-driven" set of models w/ GENIE collaboration's tune to free-nucleon data*
- Challenges always arise treating <u>nuclear</u> dynamics of neutrino interactions
 - Low-Q² suppression of quasielastic scattering relative to free nucleon
 - Multinucleon knockout (2p2h, ...)
 - Reinteraction of hadrons after primary scatter (FSI)

We apply custom tuning in two places

* We call our "tune" **N1810j_0211a.** It is built by starting with G1810b_0211a and substituting the Z-expansion QE axial form factor for the dipole one. This combination was not available in the 3.0.6 release, but it may be available in future versions.

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QE	Multinucleon	RES	DIS	FSI
València 1p1h w/ Z-expansion axial form factor	València MEC	Berger-Sehgal	Bodek-Yang	hN semi- classical cascade

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Multinucleon knockout





"2p2h" Knock out two nucleons with an elastic-like interaction.

Lots of recent progress on theory, but no model in GENIE describes extant data well

Employ fits to NOvA ND data in the meantime

Central value prediction + uncertainties based on fits to ND data



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Final-state interactions

• FSI model choice: "hN 2018"

- More rigorous theoretical foundation than older "hA" effective model
- Challenge: not directly reweightable
- Some tuning required...
 - <u>Use BDT reweighting</u> technique adapted from DUNE (see overflow)
 - <u>Adjust central value</u> to agree better with pion scattering data at low energies where most relevant for NOvA
 - <u>Construct uncertainty bands</u> in same spirit as work from T2K [Phys. Rev. D99, 052007]
- 5-10% unc. on pion kinematics
 - \rightarrow Ultimately subdominant for calorimetric
 - E_{ν} reco. used in NOvA



Making predictions

... but there's a lot hiding under the hood of the predicted specta



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Detector response



- Detector response is largest syst
 - Data-MC discrepancies in proton candidate response at 5% level
- NOvA Test Beam program underway
 - Should allow direct measurement of hadron responses
 - Expected to constrain detector energy scale uncertainty

Extrapolating ND → FD mitigates both "known" and "unknown" effects



Extrapolating ND → FD mitigates both "known" and "unknown" effects



This technique applied both to variations from <u>systematic uncertainties</u> ("known unknowns") and <u>central value</u> ("unknown unknowns")

Extrapolating ND → FD mitigates both "known" and "unknown" effects



Extrapolating ND → FD mitigates both "known" and "unknown" effects



Extrapolating in quartiles of $E_{_{had}}/E_{_{v}}$ matches the resolutions between detectors

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Extrapolating ND → FD mitigates both "known" and "unknown" effects



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<u>New in 2020</u>: extrapolating in sub-ranges of lepton $|\rho_{T}|$ enables matching the acceptance between detectors

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Extrapolating ND → FD mitigates both "known" and "unknown" effects



New in 2020: extrapolating in sub-ranges of lepton |ρ_T| enables matching the acceptance between detectors FNAL JETP Seminar / Sept. 18, 2020 J. Wolcott / Tufts University

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Applications of extrapolation





- New $|p_T|$ extrapolation improves analysis robustness
 - 30% reduction in cross section uncertainties vs. previous analyses
 - Slight increase from lepton reconstruction syst. (but well understood)
 - Overall decrease of 5-10%
- Dominated by detector energy scale uncertainties

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Part IV: Inferences

The data to be interpreted



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Oscillation results: (θ_{23} , Δm_{32}^2)



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Oscillation results: (θ_{23} , Δm_{32}^2)



Oscillation results: (δ_{CP} , θ_{23})

NOvA Preliminary 0.7 0.6 $\sin^2 \theta_{23}$ 0.5 0.4 <u>≤</u>2σ $\leq 3\sigma$ + Best Fit ≤1σ NH-0.3 $\frac{\pi}{2}$ <u>3π</u> 2 π 2π $\boldsymbol{\delta}_{\text{CP}}$ 0.7 Best fit: • $\sin^2 \theta_{23} = 0.57^{+0.03}_{-0.04}$ 0.6 • $\Delta m_{32}^2 = (+2.41 \pm 0.07) \times 10^{-3} eV^2/c^4 (NH)$ $\sin^2 \theta_{23}$ • $\delta_{CP} = 0.82 \pi$ 0.5 0.4 $\leq 2\sigma \leq 3\sigma$ l≤1σ 0.3 IH [All contours and significances calculated using $\frac{\pi}{2}$ $\frac{3\pi}{2}$ 2π 0 π $\boldsymbol{\delta}_{\text{CP}}$ Feldman-Cousins method thanks to NERSC] 64 [note: $\sin^2\theta_{13} = 0.085 \pm 0.005$ (from PDG avg. of reactor data)]

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Oscillation results: (δ_{CP} , θ_{23})



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NOvA Preliminary 60 -NOvA FD $sin^2 2\theta_{13} = 0.085$ 13.60. 10²⁰ POT-equiv (v) Total events - antineutrino beam 12.50 · 10²⁰ POT (⊽) UO LO NH δ_{CD} ★ 2020 best fit ο δ_{CP}= 0 • δ_{CP}= π/2 ⁻□ δ_{CP}= π ■ δ_{CP}= 3π/2 60 80 100 20 40 120 Total events - neutrino beam



[All contours and significances calculated using Feldman-Cousins method thanks to NERSC]

13.6×10²⁰ POT equiv v + 12.5×10²⁰ POT ⊽ NOvA FD 5 **NOvA Preliminary** NOvA Preliminary - - - NH Lower octant - NH Upper octant 60 -NOvA FD Significance (0) $sin^2 2\theta_{13} = 0.085$ --- IH Lower octant 13.60. 10²⁰ POT-equiv (v) Total events - antineutrino beam IH Upper octant 12.50 · 10²⁰ POT (v) 3 IH UO <u>3π</u> 2 $\frac{\pi}{2}$ π 2π $\boldsymbol{\delta}_{\mathsf{CP}}$ LO • Hie-oct- δ_{CP} combinations that produce 'asymmetric' v_e - v_e NH appearance are disfavored $[\circ \delta_{CP} = 0$ • $\delta_{CP} = \pi/2$ \star 2020 best fit $\neg \Box \delta_{CP} = \pi \quad \bullet \delta_{CP} = 3\pi/2$ 60 80 100 20 40 120 Total events - neutrino beam

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- Hie-oct- δ_{CP} combinations that produce 'asymmetric' $v_e \mathchar`-v_e$ appearance are disfavored
- Combinations that include some "cancellation" are preferred
 - There are such combinations for either hierarchy or octant \rightarrow no strong preferences for hierarchy (or octant)

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Vs. other measurements

[n.b.: not yet updated for NEUTRINO 2020 results from all other expt's]



"atmospheric" parameters

allowed values of δ_{CP}

Vs. other measurements



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Looking ahead

- Will resume in neutrino mode after summer shutdown
 - Run plan: 50:50 v: \overline{v}
 - NOvA is expected to run until 2025
 - Beam improvements an important part of story!
- Good sensitivity to resolution of hierarchy at full exposure
 - 4-5 σ for δ_{CP} =3 $\pi/2$
 - ≥3σ for 30-50% of δ_{CP} values (depending on θ₂₃ & true hierarchy)
 → current measurements already beginning to show power
- Anticipating reductions in detector uncertainties that should further improve analysis robustness


Summary

- With 13.6x10²⁰ POT neutrino + 12.5x10²⁰ POT antineutrino beam exposure, NOvA reports:
 - Precision measurements of atmospheric parameters:
 - ∆m²₃₂ = (2.41±0.07)×10⁻³ eV² (2.9%)
 - $\sin^2 \theta_{23} = 0.57^{+0.04}_{-0.03} \ (\sim 6\%)$
 - Constraints on strongly asymmetric v_e - \overline{v}_e appearance PMNS solutions:
 - (IH, $\delta_{CP} = \pi/2$) excluded at >3 σ
 - (NH, $\delta_{\mbox{\tiny CP}}\mbox{=} 3\pi/2$) disfavored at ~2\sigma
 - Progress towards answering the "deep questions"!
- With continued running through 2025, NOvA anticipates:
 - 230 sensitivity to mass hierarchy determination for 30-50% of δ_{CP} values
 - Input from NOvA Test Beam program, neutrino interactions community to further improve robustness to systematics



Overflow



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Reweighting using BDTs

To avoid having to fully resimulate v scattering to apply tunes, we train BDTs using truth quantities to build reweights for each variation.

$$f_{BDT} = \alpha_1 + \dots + \alpha_N + \dots + \alpha_N$$

$$\approx \frac{1}{N} \sum_{\substack{tree \\ i=1}}^{N} \alpha_i \Theta \left(\vec{x} - \vec{x}_0^i \right)$$

$$\stackrel{\text{event trained values "trained values "trained tree i}}{\stackrel{\text{event trained tree i}}{\stackrel{\text{trained tree i}}{\stackrel{\text{tree i}}}}}}}}}}}}}}}}$$

with $\vec{\chi} = (\# \text{ hadrons, hadron KE, ...})$

We use a binary logistic loss as the training objective:

$$L_{\log} = \sum_{\substack{\text{training}\\ \text{evt } n}} -y_n \ln \hat{y}_n + (y_n - 1) \ln (1 - \hat{y}_n)$$

[Technique inspired by J. Phys. Conf. Series 762, 012036; built on work by C. Vilela for DUNE]

The desired weights for an event are:



The weighted nominal distributions adequately reproduce the simulated variations. 76

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FSI tuning & uncertainties

• FSI model choice: "hN"

- Propagates hadrons through nucleus in finite steps
- Interaction probabilities simulated according to Oset quantum model
- More rigorous foundation than older "hA" effective model (hA applies hadron scattering data directly to FSI and ... hopes for the best)

• Challenge: hN not directly reweightable

 \rightarrow Addressed with novel BDT reweighting technique adapted from DUNE (see also J. Phys. Conf. Series 762, 012036)

• Tuning:

- <u>Adjust central value</u> to agree better with pion scattering data at low energies where most relevant for NOvA
- <u>Construct uncertainty bands</u> in same spirit as work from T2K [Phys. Rev. D99, 052007]



PID validation: MRE, MRDiF

- Validating selection performance:
 - Remove muons in both data & MC
 - Compare PID efficiencies relative to a preselection
- ND: muon removal, electron addition (MRE):
 - Begin with $v_{\!\scriptscriptstyle \mu}$ CC candidates
 - Replace the removed muon with simulated electron of same kinematics
- FD: muon removal from decay-in-flight (MRDiF)
 - Begin with cosmic-ray muons that decay in flight to electrons
 - Remove muon part, study electron shower



v_e bkgds: beam v_e constraints



v_e bkgds: NC/ v_μ CC constraints



A priori sensitivities



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