



Quantum Monte Carlo calculations of electron Scattering from ^{12}C in the Short-Time Approximation

Fermilab Theory Division seminar

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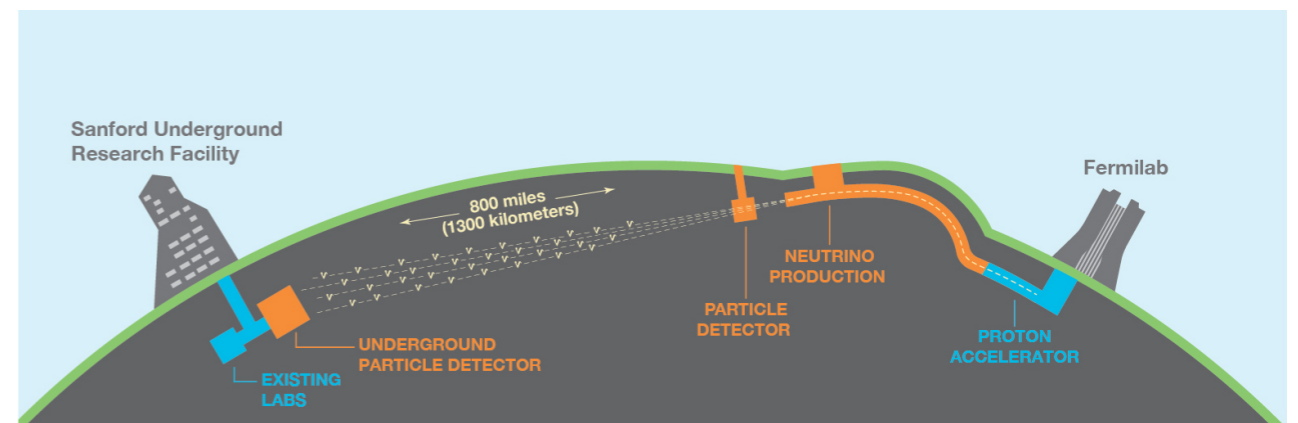
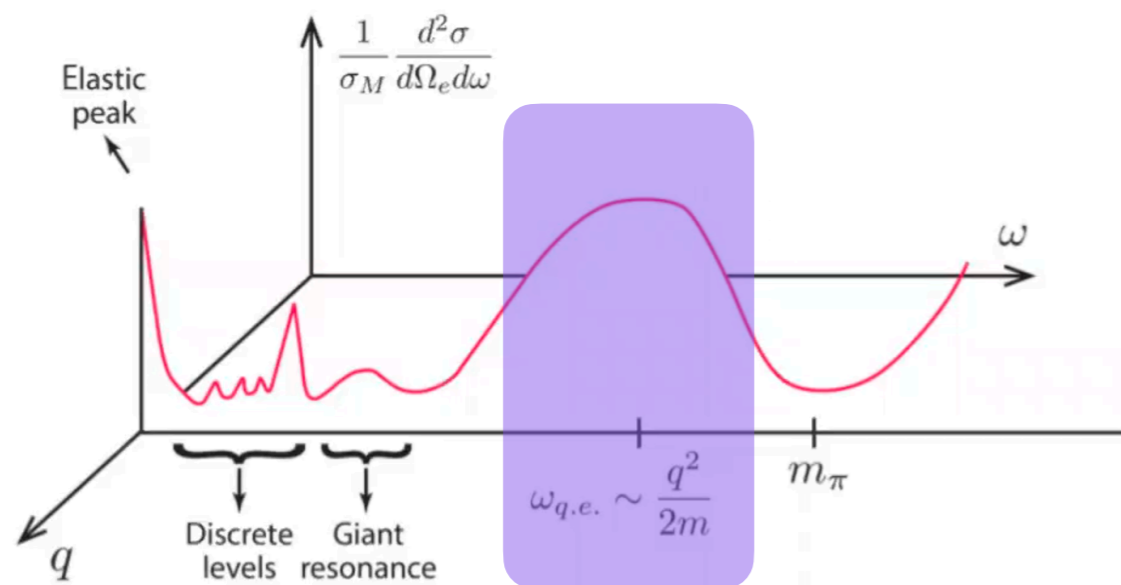
Maria Piarulli and Saori Pastore

 Washington University in St. Louis



Electron-nucleus scattering

Theoretical understanding of nuclear effects is extremely important for experimental programs



Lepton-nucleus cross sections $\omega \sim 10^2$ MeV

Ab-initio description of nuclei



- Nuclear interaction
- Electroweak interaction of leptons with nucleons
- Computational method

Many-body nuclear interaction



Many-body Nuclear Hamiltonian: Argonne v_{18} + Urbana IX

$$H = \sum_i T_i + \sum_{i < j} v_{ij} + \sum_{i < j < k} v_{ijk}$$

Quantum Monte Carlo methods:

Use nuclear wave functions that minimize the expectation value of E

$$E_V = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} \geq E_0$$

The evaluation is performed using Metropolis sampling

Nuclear Wave Functions



Variational wave function for nucleus in J state

$$|\psi\rangle = \mathcal{S} \prod_{i<j}^A \left[1 + U_{ij} + \sum_{k \neq i,j}^A U_{ijk} \right] \left[\prod_{i<j} f_c(r_{ij}) \right] |\Phi(JMTT_3)\rangle$$

Two-body spin- and isospin-dependent correlations

$$U_{ij} = \sum_p f^p(r_{ij}) O_{ij}^p$$

$$O_{ij}^p = [1, \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j, S_{ij}] \otimes [1, \boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j]$$

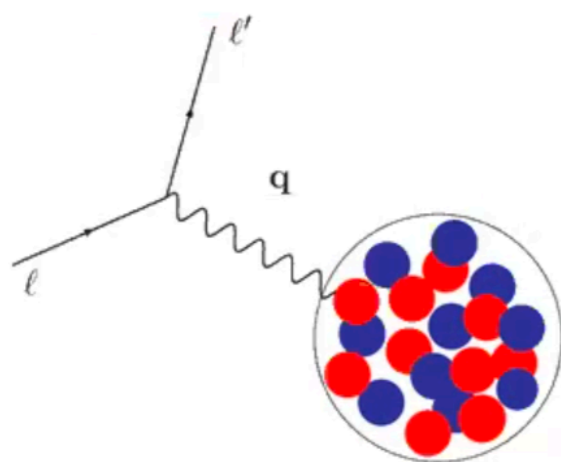
$$U_{ijk} = \epsilon v_{ijk}(\bar{r}_{ij}, \bar{r}_{jk}, \bar{r}_{ki})$$



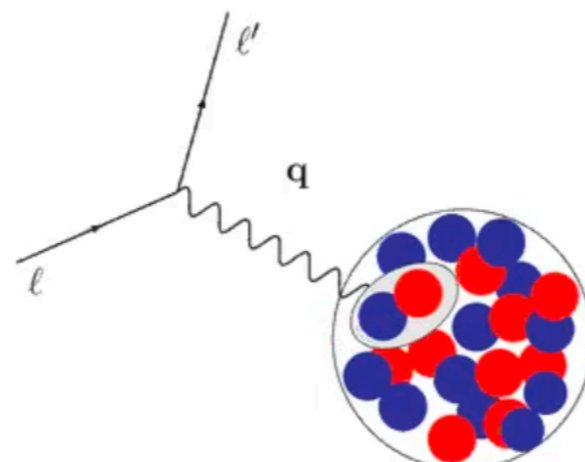
Electromagnetic interactions

Phenomenological Hamiltonian for NN and NNN

The interaction with external probes is described in terms on one- and two-body charge and current operators



one-body



two-body

Charge operators

$$\rho = \sum_{i=1}^A \rho_i + \sum_{i<j} \rho_{ij} + \dots$$

Current operators

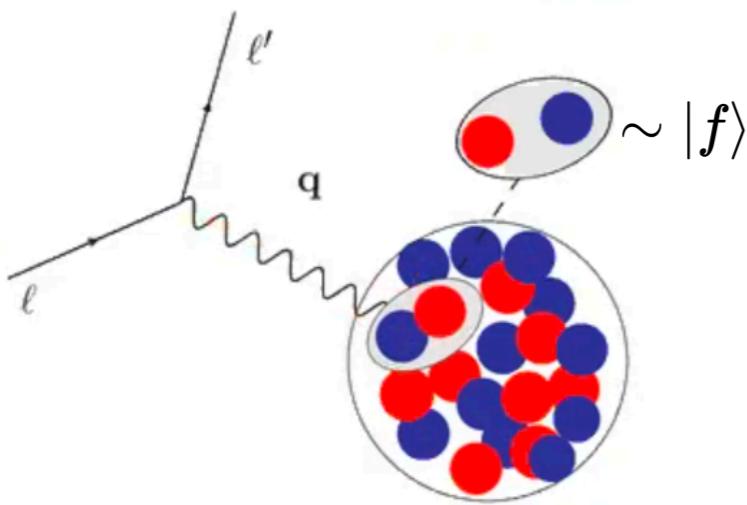
$$\mathbf{j} = \sum_{i=1}^A \mathbf{j}_i + \sum_{i<j} \mathbf{j}_{ij} + \dots$$



Short-time approximation

Pastore et al. PRC101(2020)044612

Quasielastic scattering cross sections are expressed in terms of response function



Response functions

$$R_{\alpha}(q, \omega) = \sum_f \delta(\omega + E_0 - E_f) |\langle f | O_{\alpha}(\mathbf{q}) | 0 \rangle|^2$$

$$R_{\alpha}(q, \omega) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{i(\omega + E_i)t} \langle \Psi_i | O_{\alpha}^{\dagger}(\mathbf{q}) e^{-iHt} O_{\alpha}(\mathbf{q}) | \Psi_i \rangle$$

The sum over all final states is replaced by a two nucleon propagator

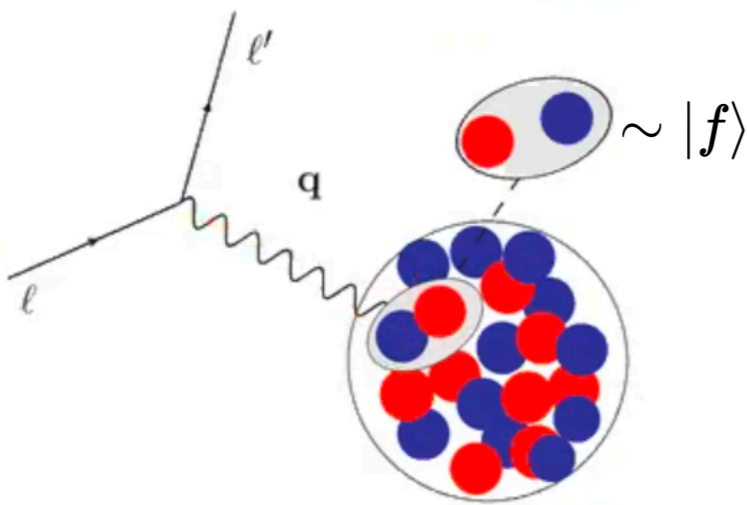
$$\begin{aligned} O^{\dagger} e^{-iHt} O &= \left(\sum_i O_i^{\dagger} + \sum_{i < j} O_{ij}^{\dagger} \right) e^{-iHt} \left(\sum_{i'} O_{i'} + \sum_{i' < j'} O_{i'j'} \right) \\ &= \sum_i O_i^{\dagger} e^{-iHt} O_i + \sum_{i \neq j} O_i^{\dagger} e^{-iHt} O_j \\ &\quad + \sum_{i \neq j} \left(O_i^{\dagger} e^{-iHt} O_{ij} + O_{ij}^{\dagger} e^{-iHt} O_i \right. \\ &\quad \left. + O_{ij}^{\dagger} e^{-iHt} O_{ij} \right) + \dots \end{aligned}$$



Short-time approximation

Pastore et al. PRC101(2020)044612

Quasielastic scattering cross sections are expressed in terms of response function



Response functions

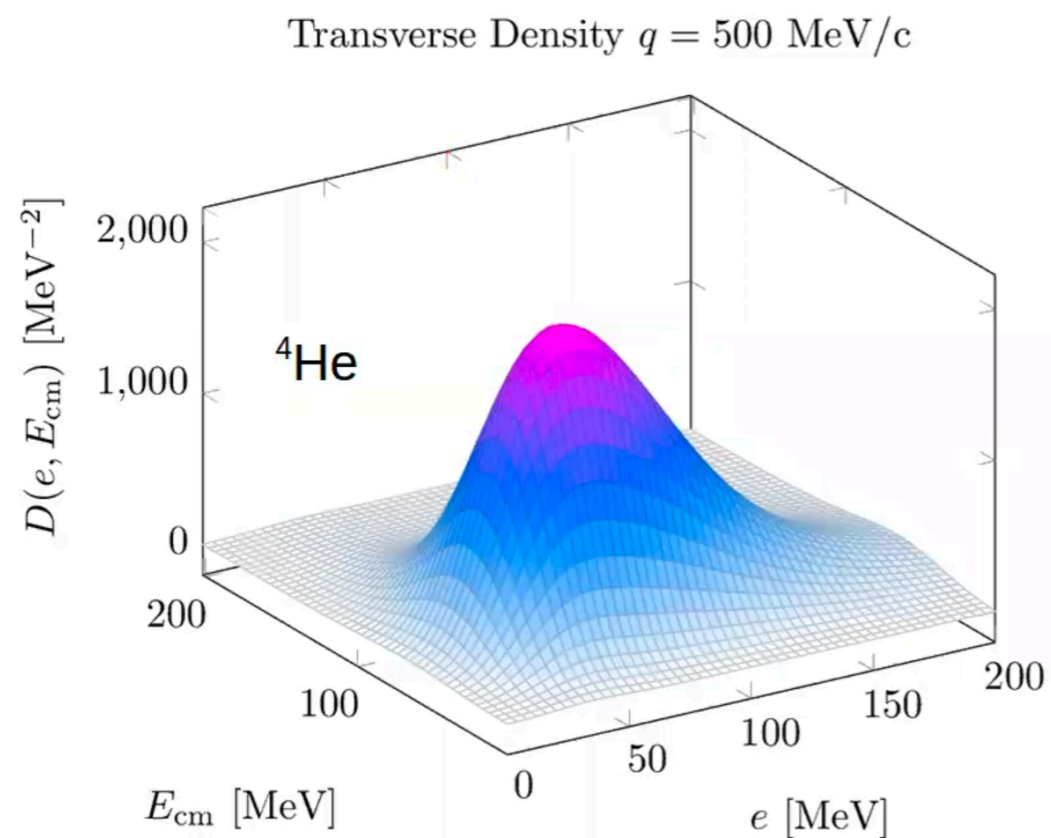
$$R_{\alpha}(q, \omega) = \sum_f \delta(\omega + E_0 - E_f) |\langle f | O_{\alpha}(\mathbf{q}) | 0 \rangle|^2$$

Response densities

$$R^{\text{STA}}(q, \omega) \sim \int \delta(\omega + E_0 - E_f) de dE_{cm} \mathcal{D}(e, E_{cm}; q)$$

STA: scattering of a correlated pair of nucleons inside a nucleus

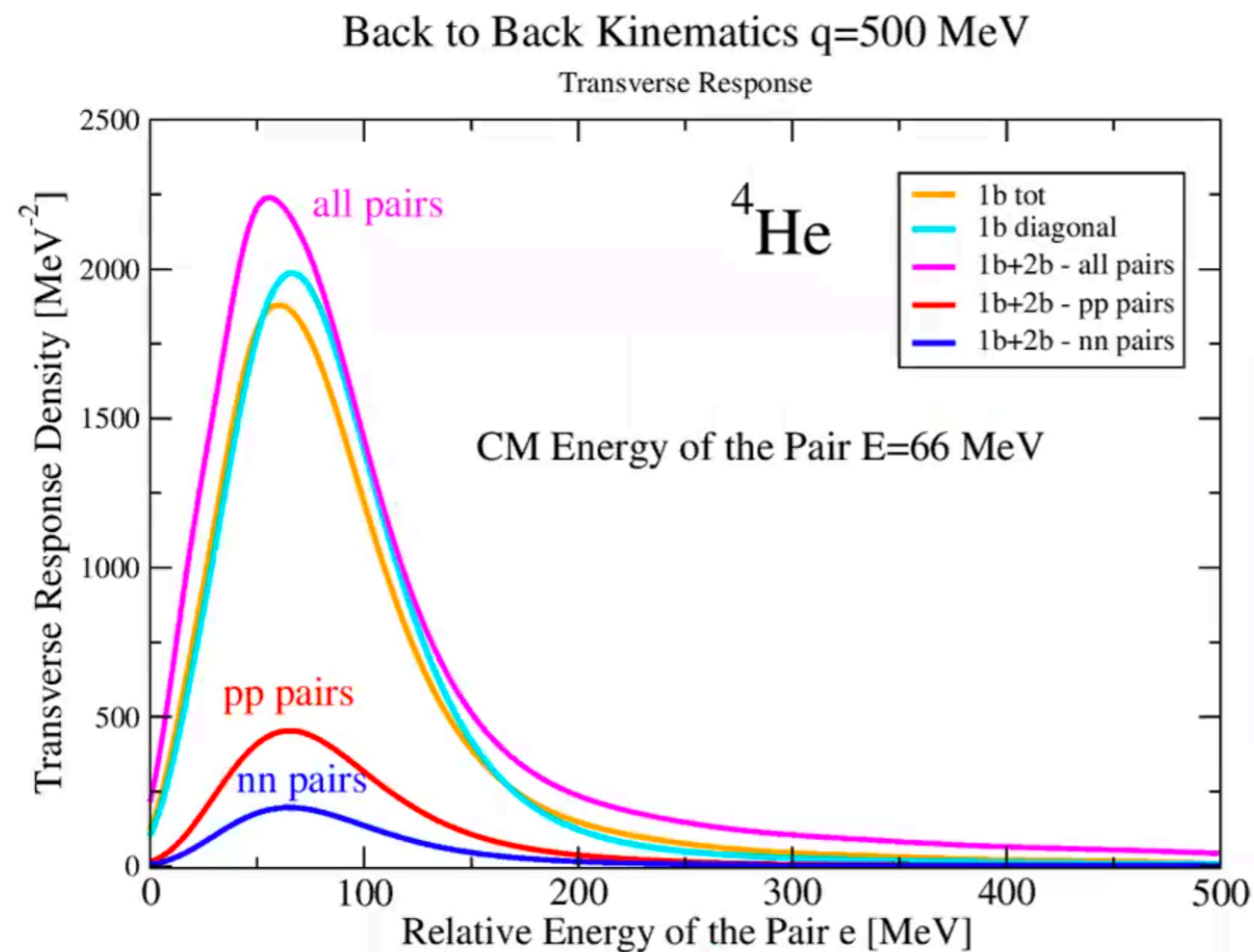
Transverse response density



Electron scattering from ^4He :

- Response density as a function of (E, e)
- Give access to particular kinematics for the struck nucleon pair

Back-to-back kinematic



We can select a particular kinematic, and assess the contributions from different particle identities

Benchmark



L.A., S. Pastore, N. Rocco, et al. arXiv:2108.10824

- We benchmarked three different methods based on the same description of nuclear dynamics of the initial target state
- Compared to the experimental data for the longitudinal and transverse electromagnetic response functions of ^3He , and the inclusive cross sections of both ^3He and ^3H
- Comparing the results allows for a precise quantification of the uncertainties inherent to factorization schemes

Benchmark



L.A., S. Pastore, N. Rocco, et al. arXiv:2108.10824

Green's function Monte Carlo

$$|\Psi_0\rangle \propto \lim_{\tau \rightarrow \infty} \exp[-(H - E_0)\tau] |\Psi_T\rangle$$

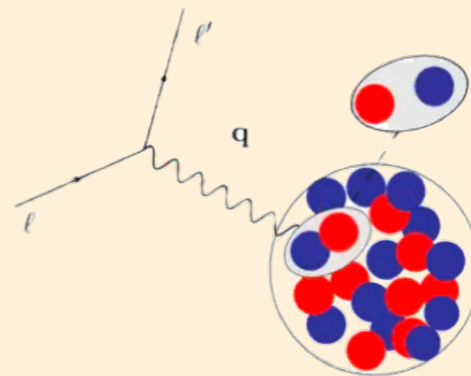
$$E_\alpha(\mathbf{q}, \tau) = \int_{\omega_{\text{th}}}^{\infty} d\omega e^{-\omega\tau} R_\alpha(\mathbf{q}, \omega), \quad \alpha = L, T$$

$$E_\alpha(\mathbf{q}, \tau) = \langle \Psi_0 | J_\alpha^\dagger(\mathbf{q}) e^{-(H-E_0)\tau} J_\alpha(\mathbf{q}) | \Psi_0 \rangle - |F_\alpha(\mathbf{q})|^2 e^{-\omega_{el}\tau}$$

Start-time approximation

$$R_\alpha(\mathbf{q}, \omega) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{i(\omega+E_0)t} \times \langle \Psi_0 | J_\alpha^\dagger(\mathbf{q}) e^{-iHt} J_\alpha(\mathbf{q}) | \Psi_0 \rangle$$

$$J^\dagger e^{-iHt} J = \sum_i J_i^\dagger e^{-iHt} J_i + \sum_{i \neq j} J_i^\dagger e^{-iHt} J_j + \sum_{i \neq j} (J_i^\dagger e^{-iHt} J_{ij} + J_{ij}^\dagger e^{-iHt} J_i + J_{ij}^\dagger e^{-iHt} J_{ij}) + \dots$$



Spectral function

$$|\Psi_f\rangle = |\mathbf{p}\rangle \otimes |\Psi_n^{A-1}\rangle$$

$$R_\alpha(\mathbf{q}, \omega) = \sum_{\tau_k=p,n} \int \frac{d^3k}{(2\pi)^3} dE [P_{\tau_k}(\mathbf{k}, E) \times \frac{m_N^2}{e(\mathbf{k})e(\mathbf{k}+\mathbf{q})} \sum_i \langle k | j_{i,\alpha}^\dagger | k + \mathbf{q} \rangle \langle p | j_{i,\alpha} | k \rangle \times \delta(\tilde{\omega} + e(\mathbf{k}) - e(\mathbf{k} + \mathbf{q}))]$$

Spectral function P, probability distribution of removing a nucleon with momentum k

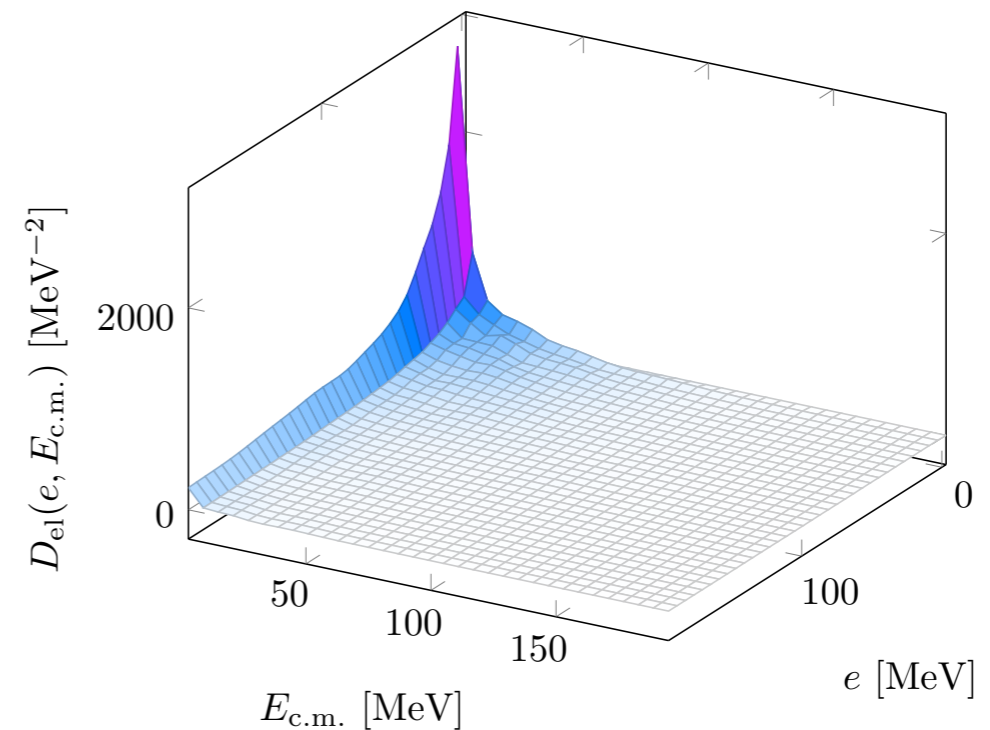
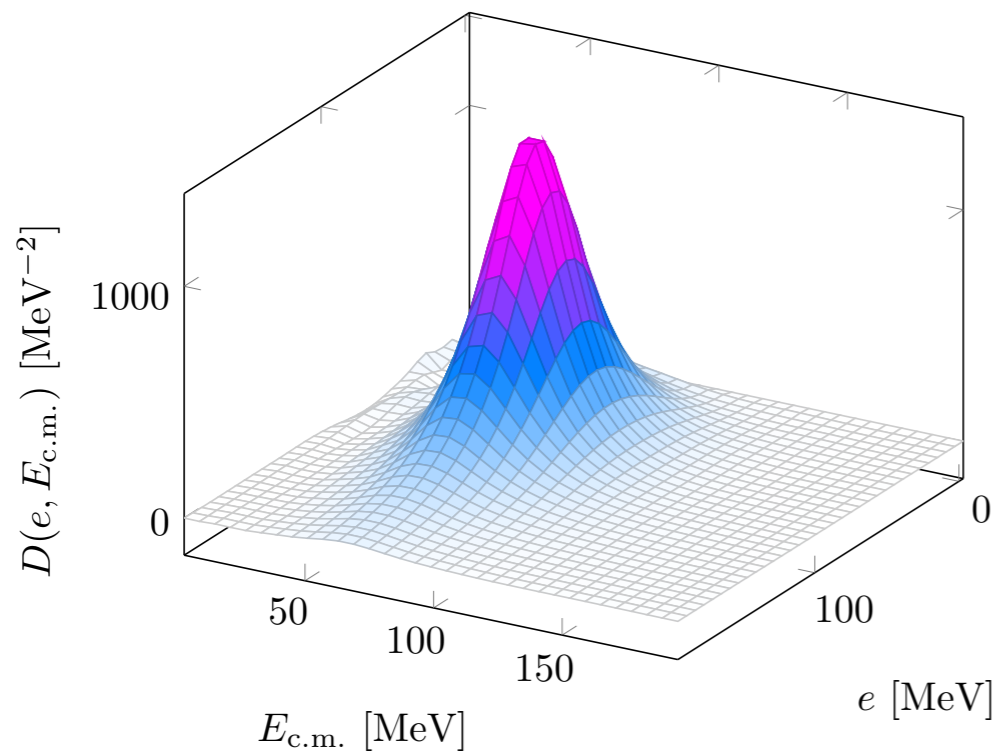


Elastic contribution

$$R^{\text{STA}}(q, \omega) \sim \int \delta(\omega + E_0 - E_f) de dE_{\text{cm}} \mathcal{D}(e, E_{\text{cm}}; q)$$

$$\mathcal{D}(e, E_{\text{cm}}) - \mathcal{D}_{\text{el}}(e, E_{\text{cm}})$$

$$\mathcal{D}_{\text{el}}(\mathbf{q}, \mathbf{p}', \mathbf{P}') = |\langle \Psi_0 | J(\mathbf{q}) | \Psi_0 \rangle|^2 \times \sum_{\beta} \langle \Psi_0 | \Psi_2(\mathbf{p}', \mathbf{P}', \beta) \rangle \langle \Psi_2(\mathbf{p}', \mathbf{P}', \beta) | \Psi_0 \rangle$$

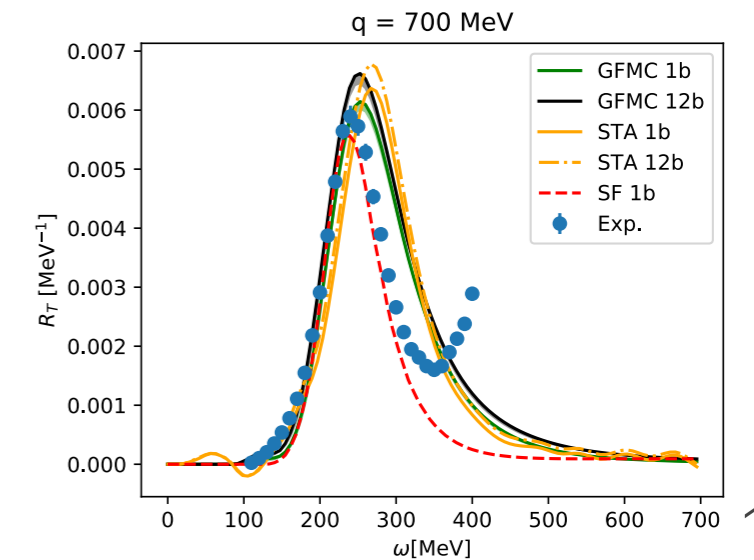
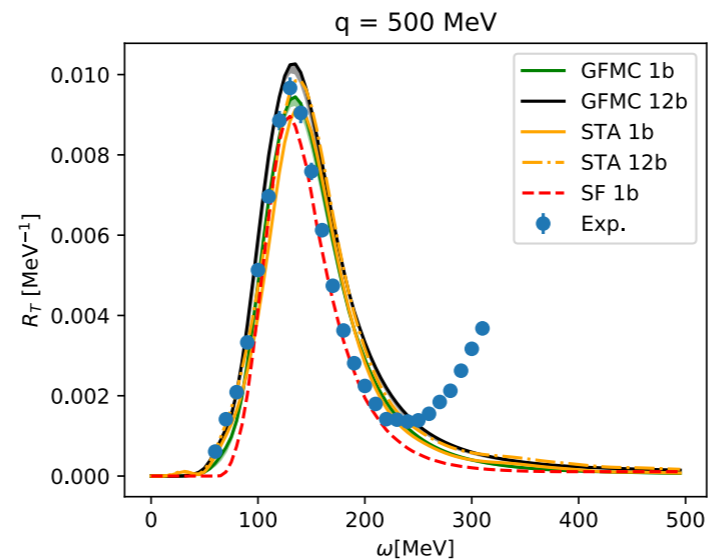
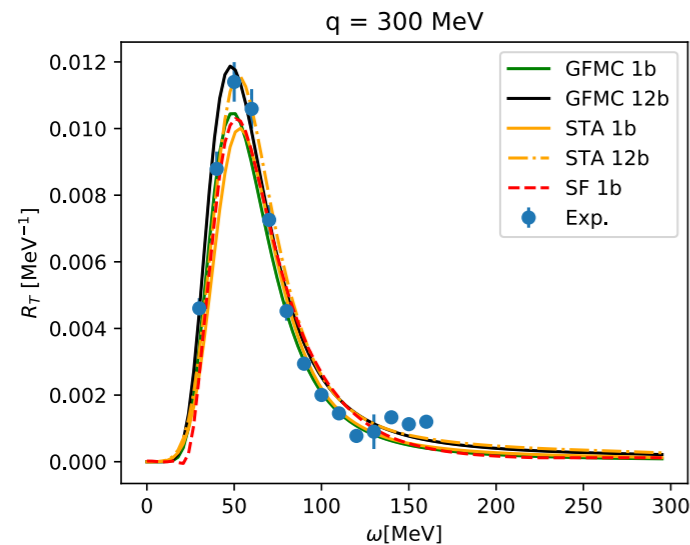
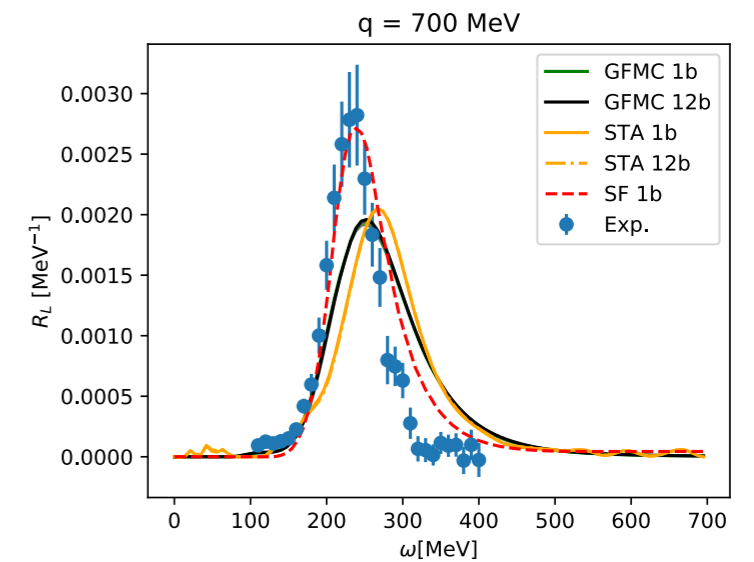
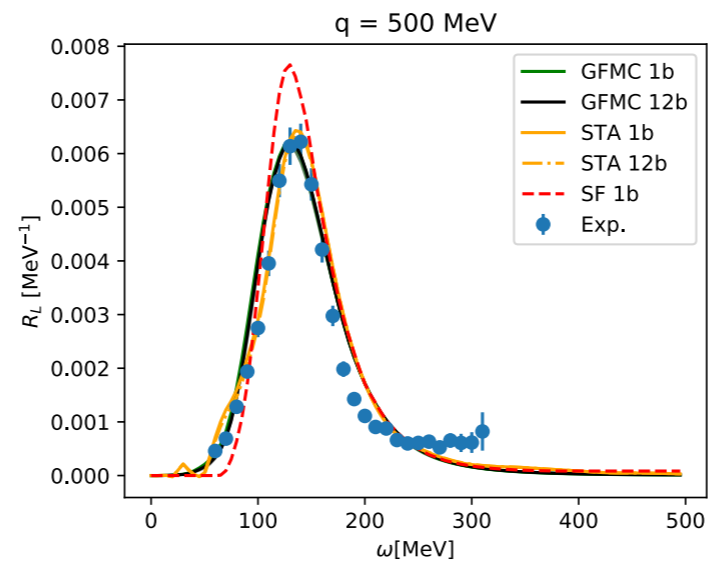
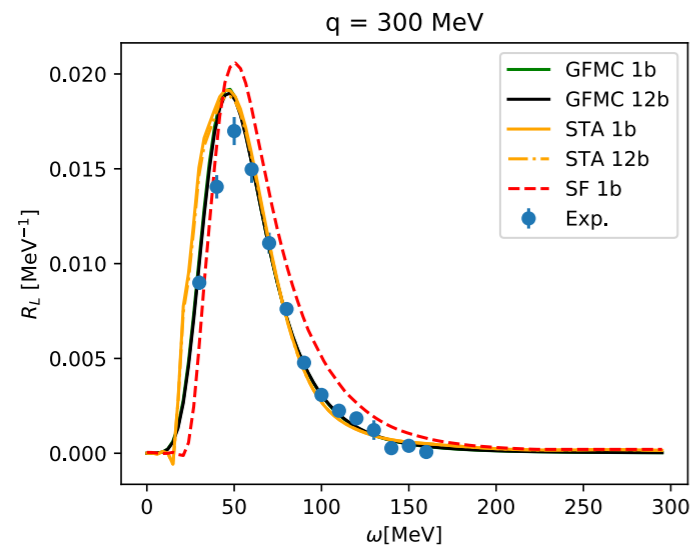


^3H Longitudinal response at 300 MeV

Benchmark



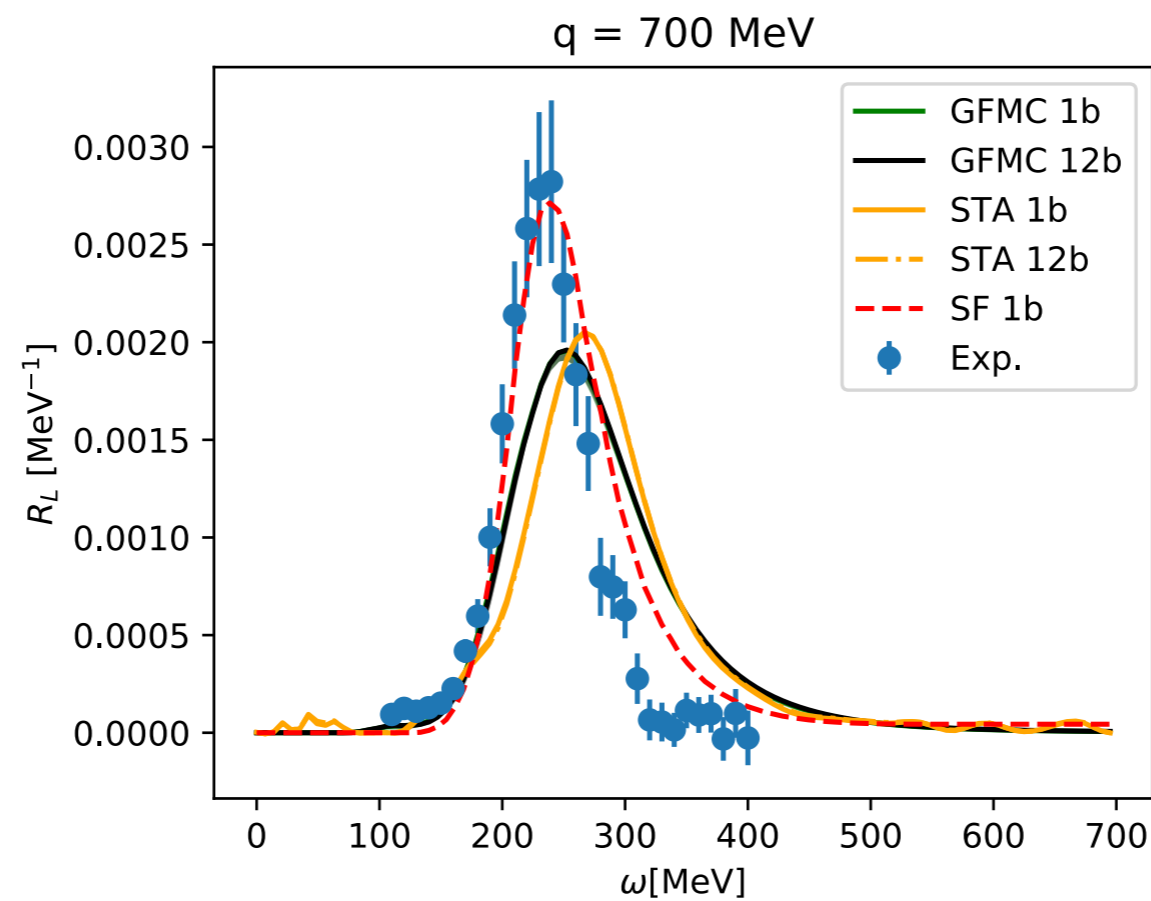
Longitudinal and transverse response function in ^3He



Relativistic corrections



Necessary to include relativistic correction at higher momentum q

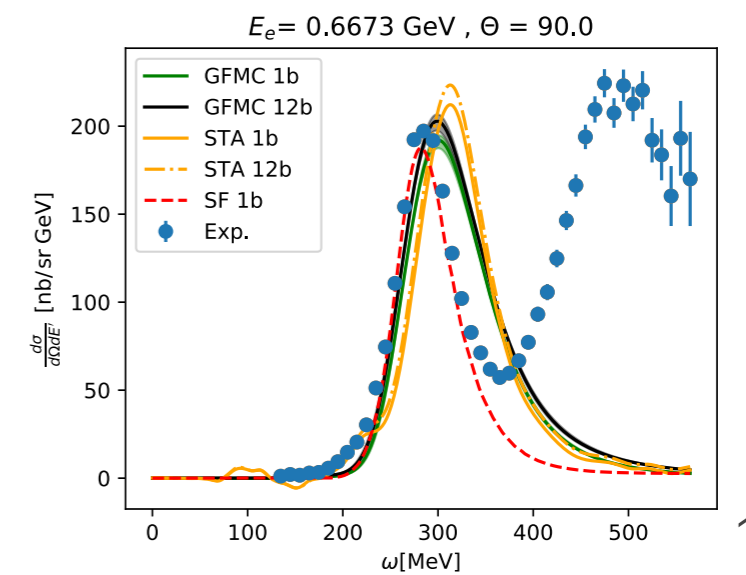
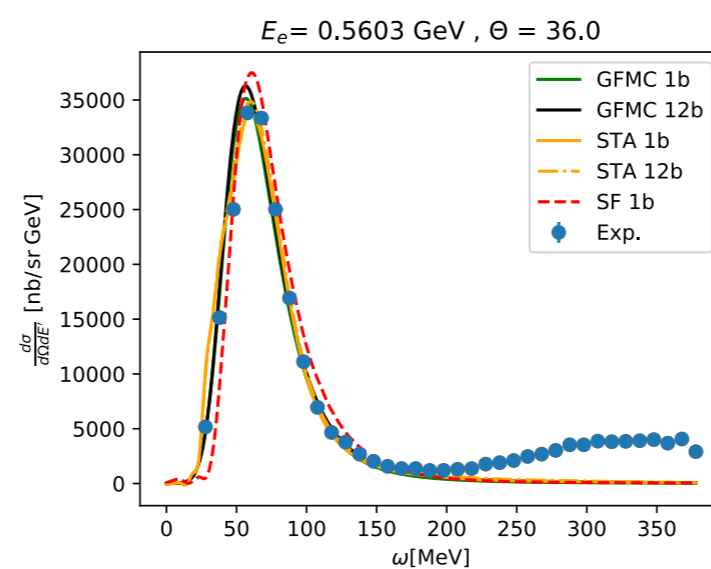
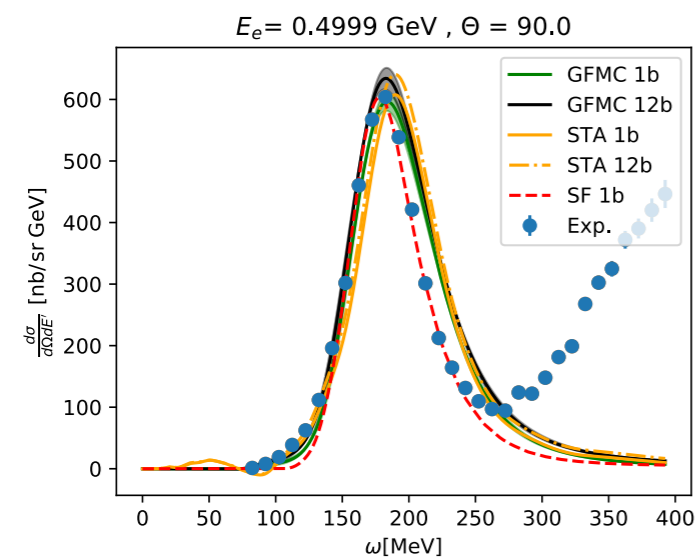
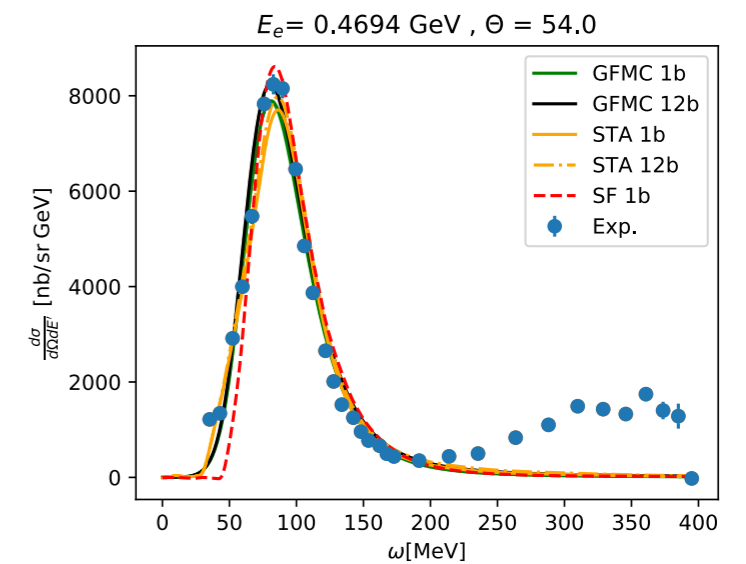
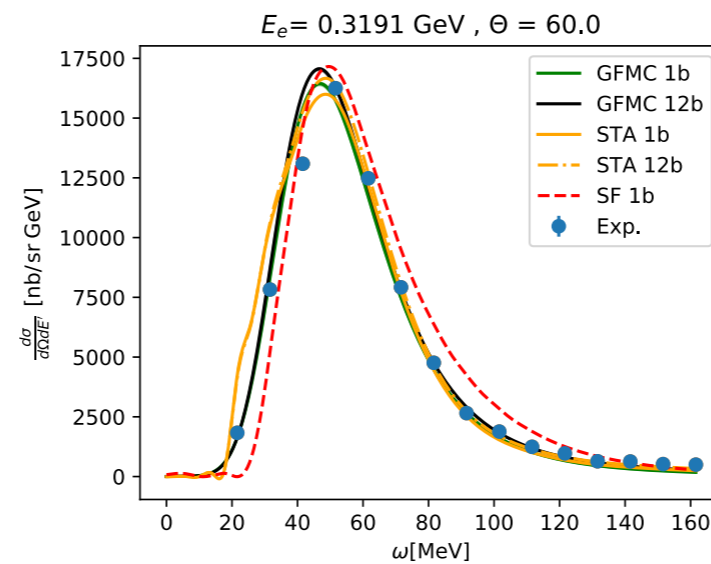
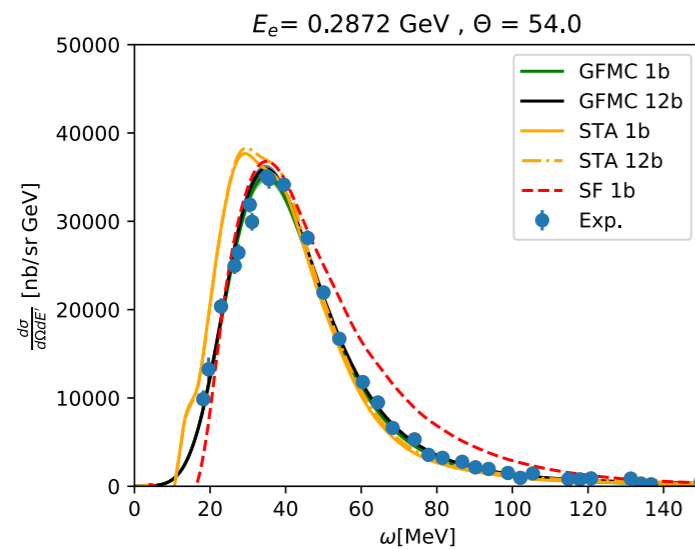


L.A., S. Pastore, N. Rocco, et al. PRC105(2022)014002

Cross sections



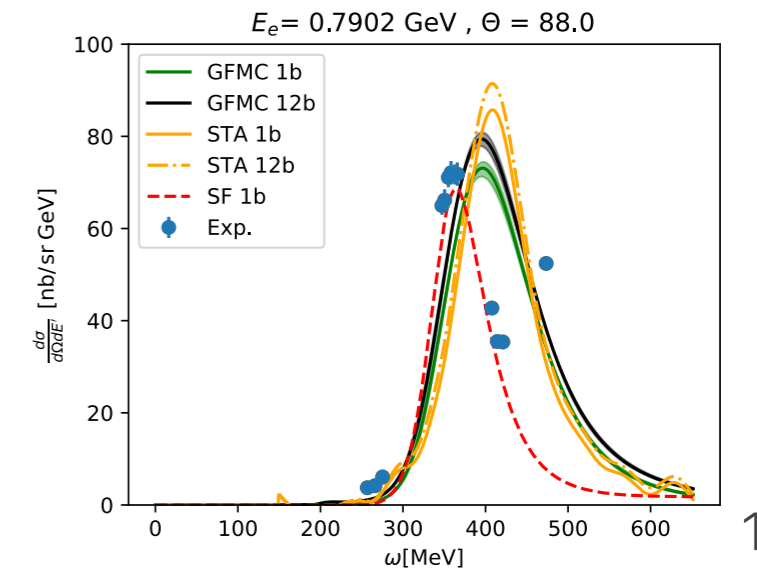
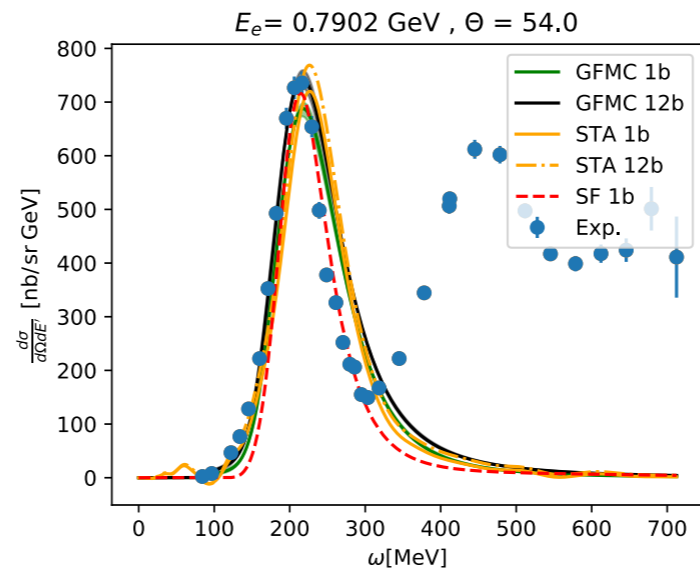
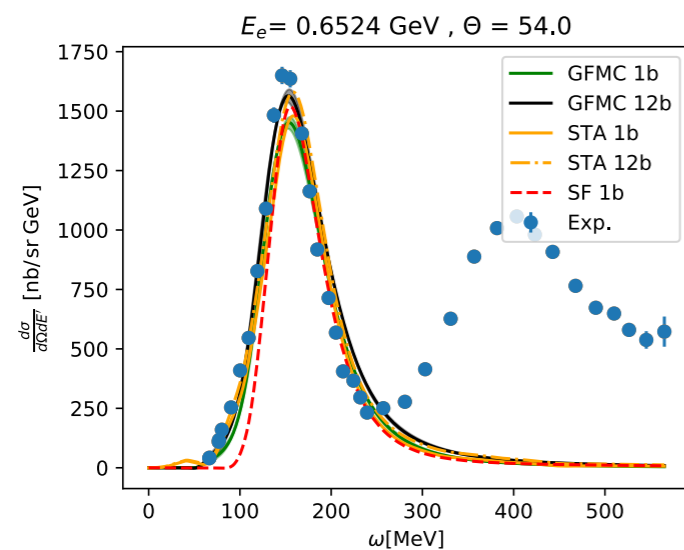
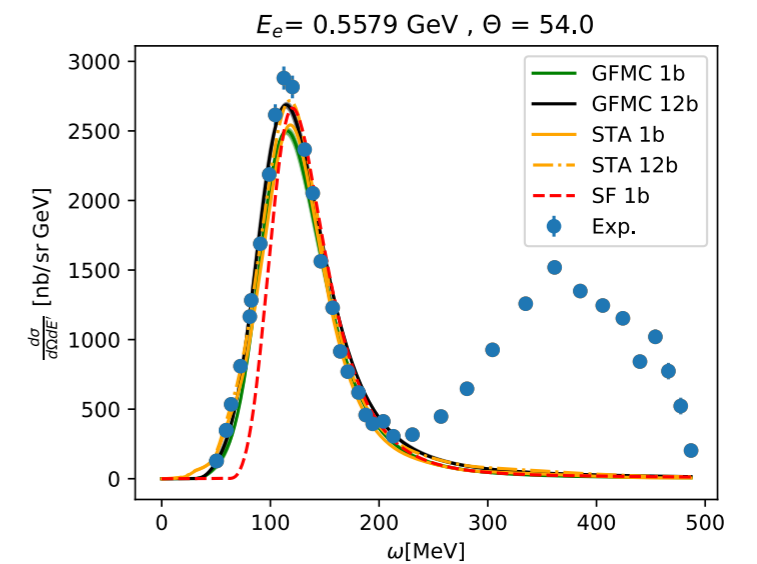
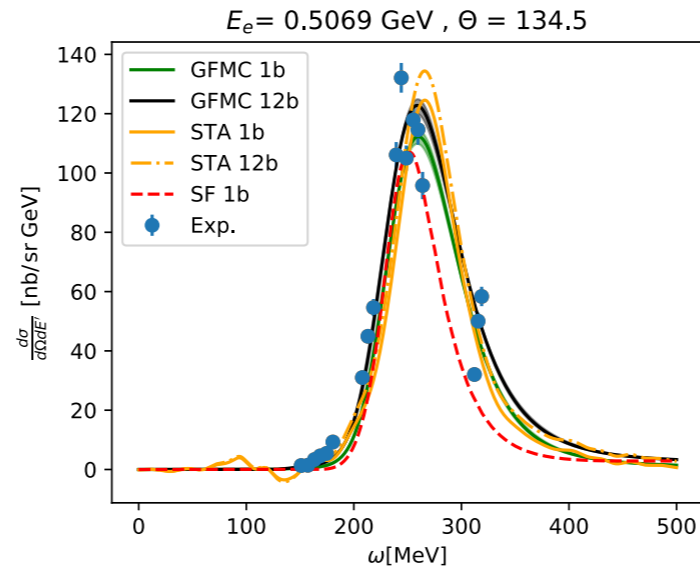
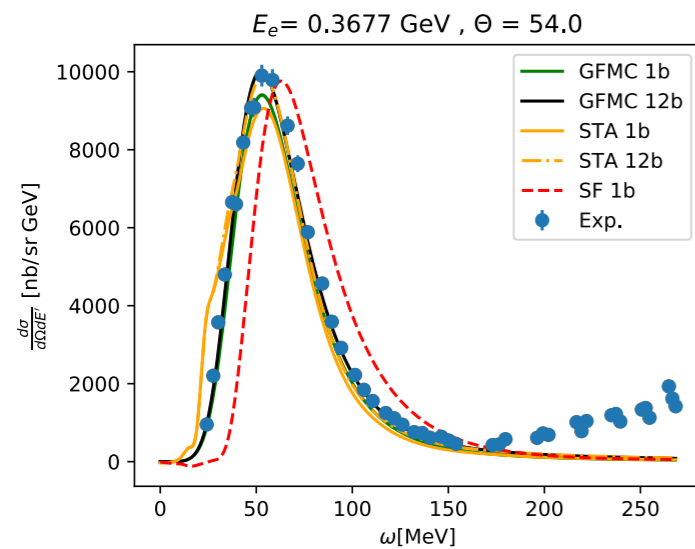
^3He



Cross sections



^3H



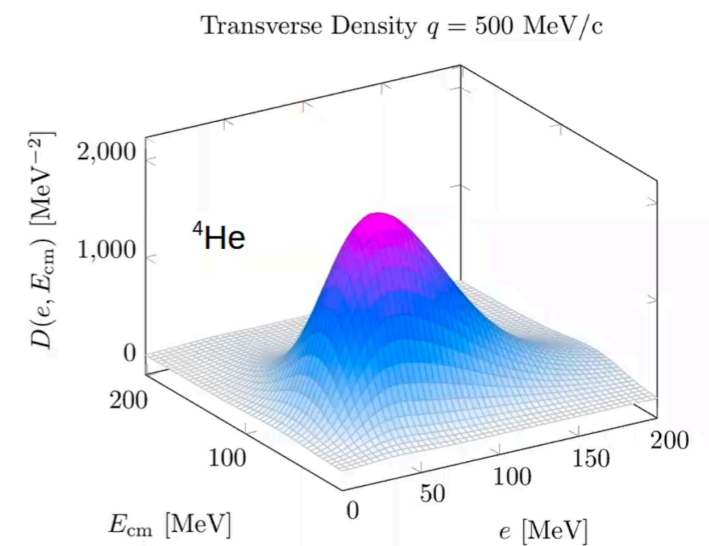
Heavier nuclei



Computational complexity of response functions and densities:

| | | | |
|---------------|-----------------------|-----------------|-------------------|
| Wave-function | | ${}^4\text{He}$ | ${}^{12}\text{C}$ |
| Spin | 2^A | 16 | 4096 |
| Isospin | $\frac{A!}{Z!(A-Z)!}$ | 6 | 924 |
| Pairs | $A(A-1)/2$ | 6 | 66 |

Response densities: E, e grid



$$R_{\alpha}(q, \omega) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{i(\omega + E_i)t} \langle \Psi_i | O_{\alpha}^{\dagger}(\mathbf{q}) e^{-iHt} O_{\alpha}(\mathbf{q}) | \Psi_i \rangle$$

Heavier nuclei

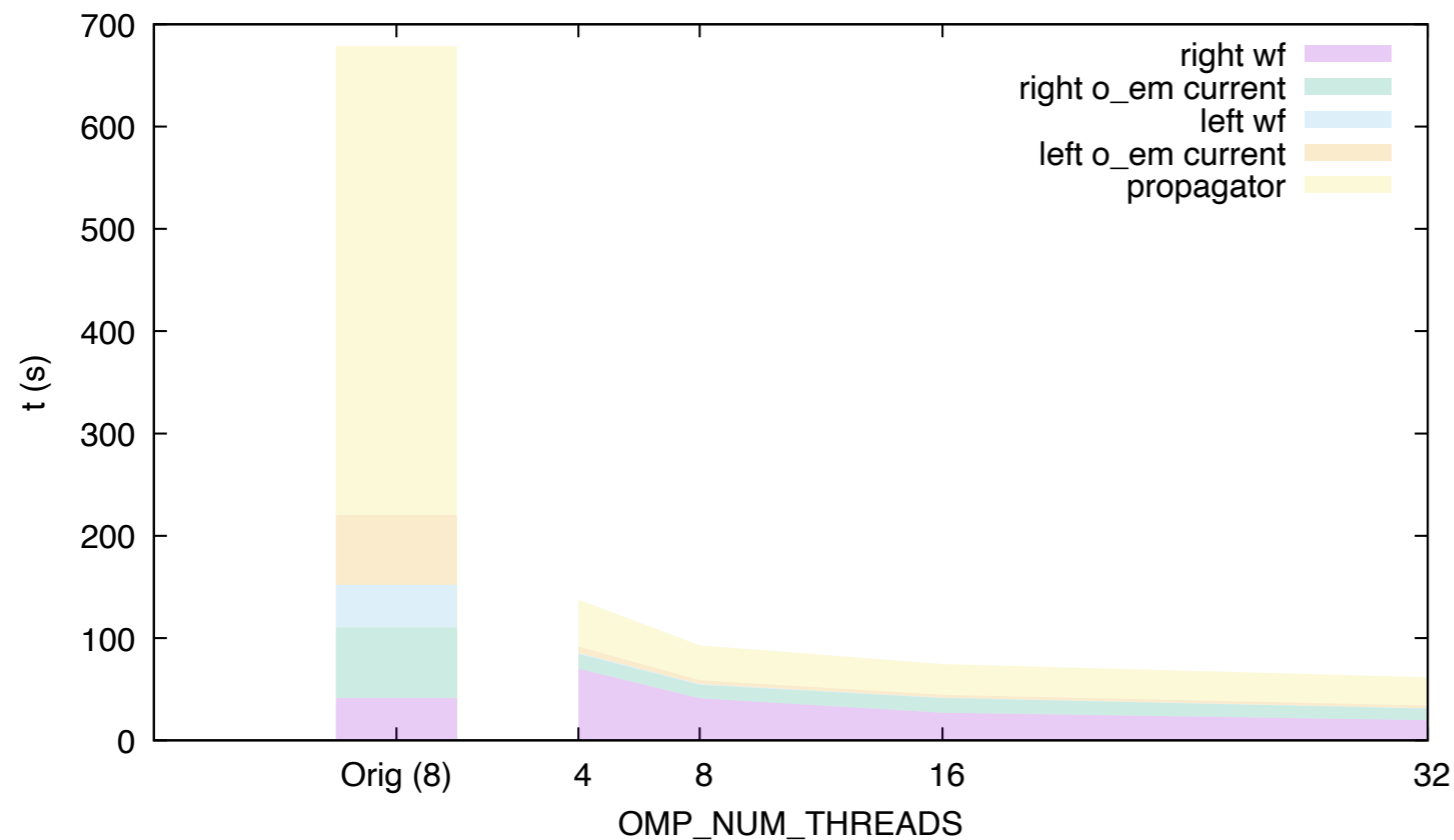


Optimization was necessary to tackle heavier nuclei

$$R_{\alpha}(q, \omega) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{i(\omega + E_i)t} \langle \Psi_i | O_{\alpha}^{\dagger}(\mathbf{q}) e^{-iHt} O_{\alpha}(\mathbf{q}) | \Psi_i \rangle$$

- Parallelization – MPI and OpenMP:
Variational Monte Carlo is almost perfectly parallelizable, but with increased system size memory becomes a constrain
- Refactoring of the code
- Computational algorithms and approximations:
variation of integration ranges (r , R) for struck nucleon pair

Heavier nuclei: ^{12}C



Optimization specific to ^{12}C was needed in order to perform full response densities calculations:

- **parallelization**
- **refactoring of the code**
- **reduction of memory usage**
- computational algorithms and approximations

$$R_{\alpha}(q, \omega) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{i(\omega + E_i)t} \langle \Psi_i | O_{\alpha}^{\dagger}(\mathbf{q}) e^{-iHt} O_{\alpha}(\mathbf{q}) | \Psi_i \rangle$$

Heavier nuclei



Comparison for 3H , 72k MC configurations, 40x40 points in r , R integration

Original:

- ~15k core hours
(LA et al arXiv:2108.10824)

After optimization:

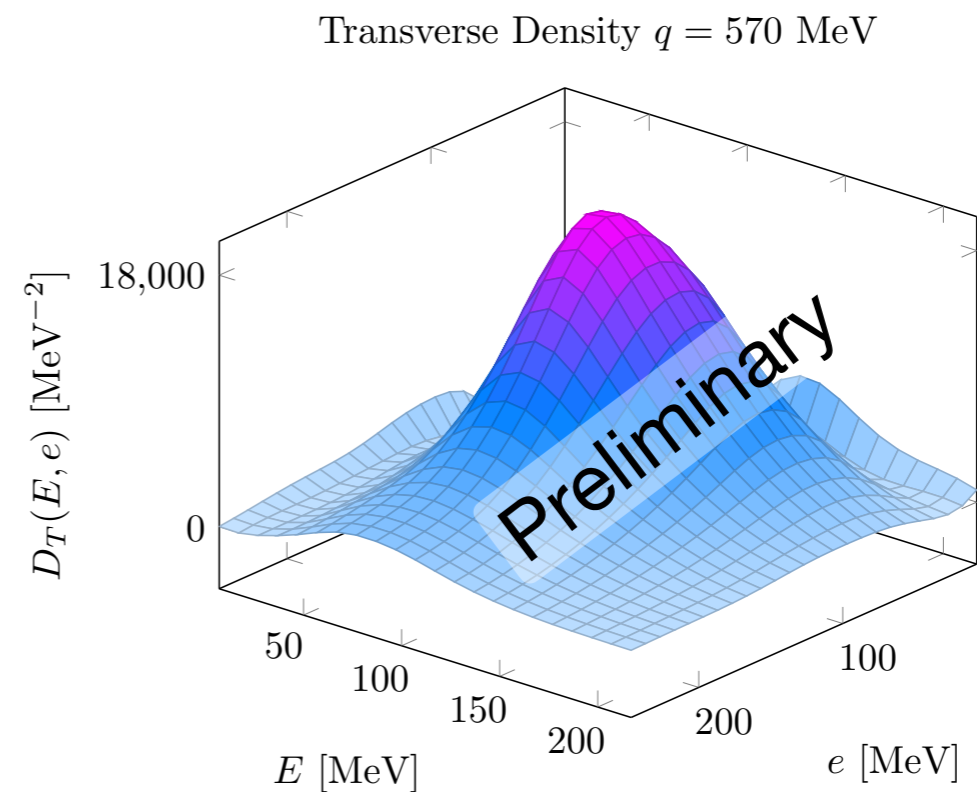
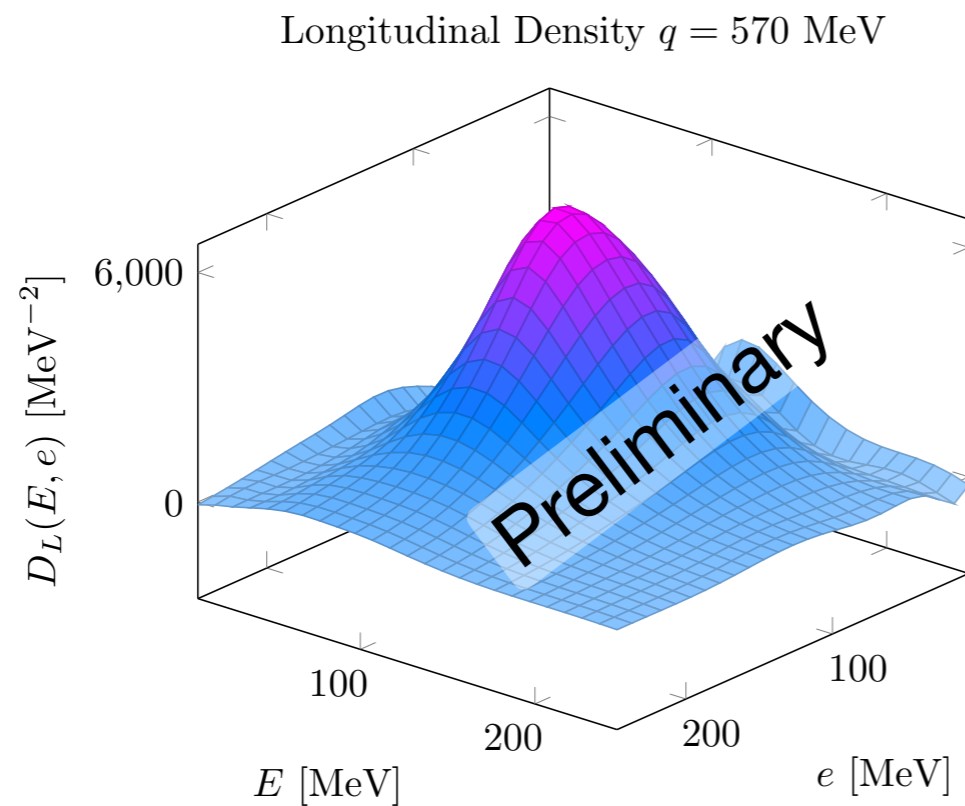
- ~1.8k core hours



Additional approximations reduce the computation time:

- r , R integration up to 8, 5 fm
- reduction of grid spacing in relative energy e

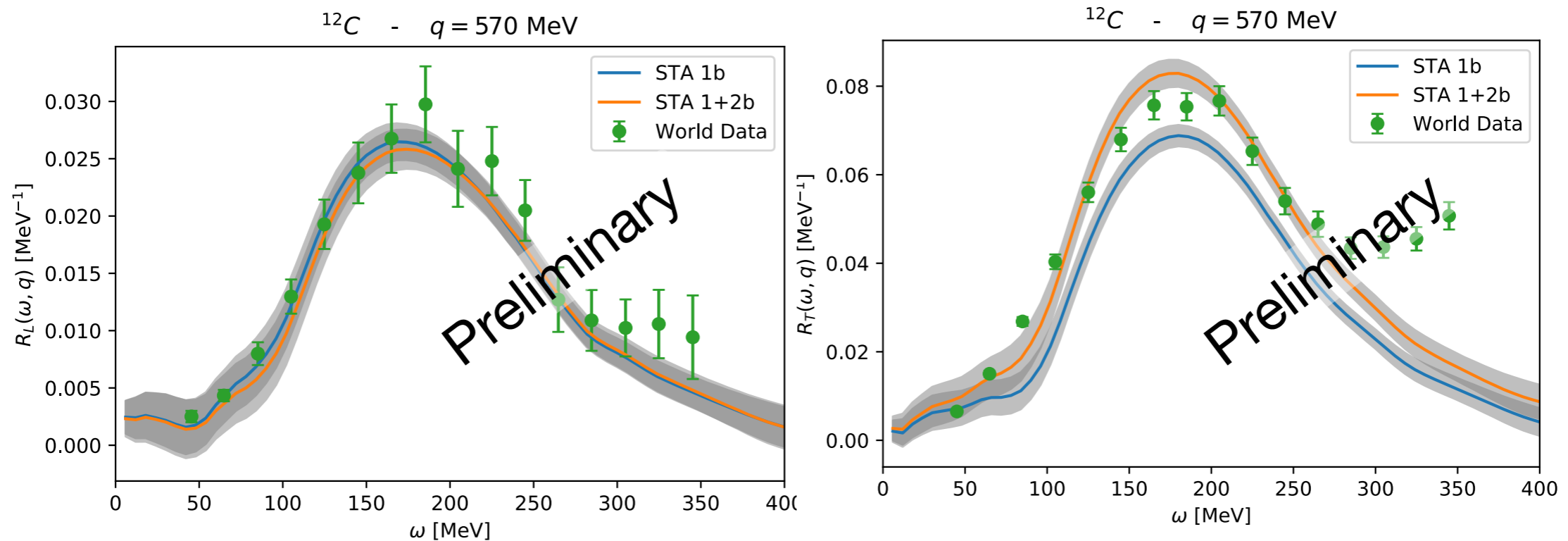
Response densities for ^{12}C



Longitudinal and transverse response densities in ^{12}C

$q = 570 \text{ MeV}$

Response functions for ^{12}C



Preliminary results for longitudinal and transverse response functions in ^{12}C

$q = 570 \text{ MeV}$

Current and future projects



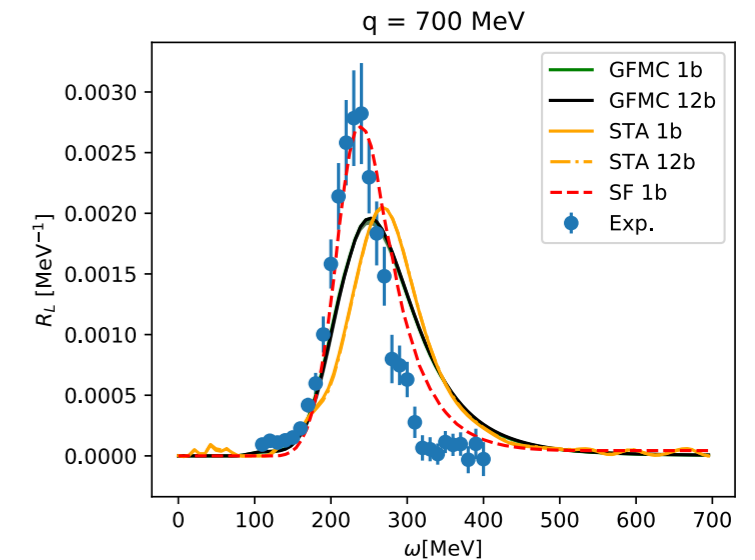
- Elastic contribution
- Relativistic corrections
- Collaboration with GENIE to include response **densities** in event generator

Relativistic corrections



Necessary to include relativistic correction at higher momentum q

- Include kinematics
- Full relativistic treatment of final states



$$E_{NR} \rightarrow E(\mathbf{p}, \mathbf{P}) = \sqrt{\left(\frac{\mathbf{P}}{2} - \mathbf{p}\right)^2 + m^2} + \sqrt{\left(\frac{\mathbf{P}}{2} + \mathbf{p}\right)^2 + m^2} - 2m = \sqrt{\frac{P^2}{4} - Pp \cos(\theta) + p^2 + m^2} + \sqrt{\frac{P^2}{4} + Pp \cos(\theta) + p^2 + m^2} - 2m$$

$$R_\alpha(\mathbf{q}, \omega) =$$

...

...

...

$$= \int \frac{d^3 K}{(2\pi)^3} \int \frac{d^3 k}{(2\pi)^3} \left[\sum_{M_i} \sum_{\alpha_{A-2}, \alpha} \int d^3 R_{A-2} \langle \Psi_i | O_L^\dagger(\mathbf{q}) | \mathbf{K} \mathbf{k} \alpha, \mathbf{R}_{A-2}, \alpha_{A-2} \rangle \langle \mathbf{K} \mathbf{k} \alpha, \mathbf{R}_{A-2}, \alpha_{A-2} | O_R(\mathbf{q}) | \Psi_i \rangle \right] \delta(\omega + E_i - E(\mathbf{k}, \mathbf{K}))$$

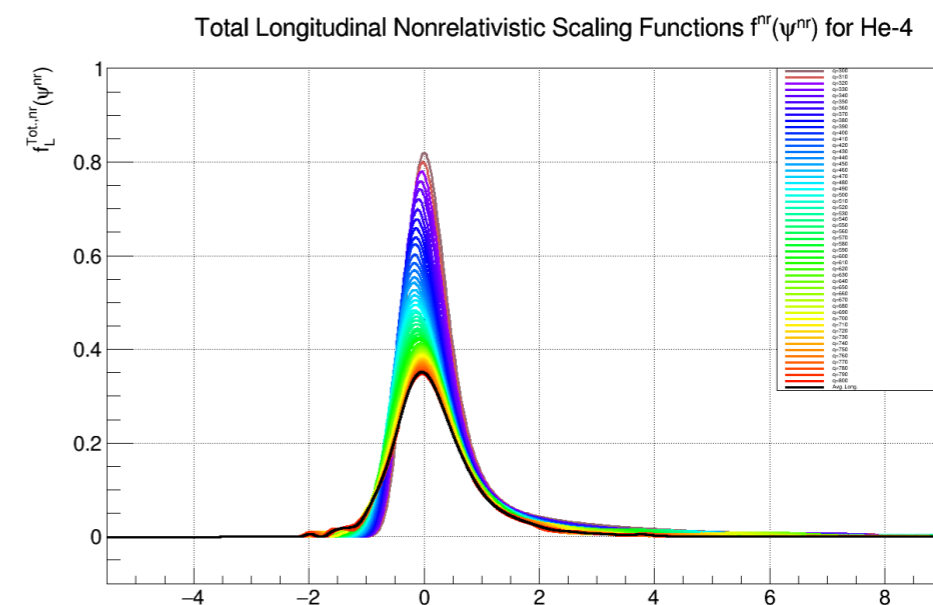
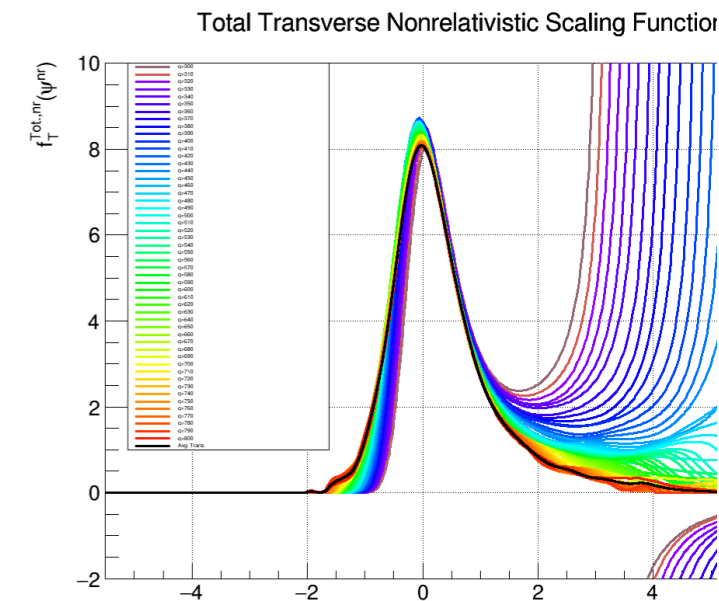
GENIE event generator



URA Visiting Scholars Program funding, collaboration with GENIE experts at Fermilab

Light systems:

- Responses and response densities w/o the elastic peak
- Finer grid and study of interpolation methods



GENIE event generator



URA Visiting Scholars Program funding, collaboration with GENIE experts at Fermilab

Light systems:

- Responses and response densities w/o the elastic peak
- Finer grid and study of interpolation methods
- Implement **densities** into GENIE:
 - currently responses have been implemented
(J. Barrow, S. Gardiner, S. Pastore, M. Betancourt, and J. Carlson. Phys. Rev. D 103, 052001)
 - at the moment GENIE has a way of “making up” that information rather than a direct theoretical prediction
 - test interpolation and morphing scheme
- More momenta for ^{12}C
- EW currents

Collaborators:

S. Pastore, M. Piarulli

J. Carlson, S. Gandolfi, A. Lovato, N. Rocco, R. Wiringa

S. Gardiner, J. Barrow, M. Betancourt

Thank you!

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