Exploring the Higgs Sector (Beyond the SM)

S. Dawson Fermilab, 2012 Lecture 4

Many, many possibilities

Reprise on spin



Is the Standard Model Self-Consistent?

- M_h is a free parameter in the Standard Model
- Can we derive limits on the basis of consistency?
- Consider a scalar potential:

$$V = \frac{M_h^2}{2}h^2 + \frac{\lambda}{4}h^4$$

- This is potential at electroweak scale
- Parameters evolve with energy in a calculable way

Standard Model: $\lambda = M_h^2/(2v^2) = .13$ for $M_h = 125$ GeV

Perturbative Regime

Does Spontaneous Symmetry Breaking Happen?

• $\lambda(\Lambda) > 0$ gives lower bound on M_h

 $M_h^2 > \frac{3v^2}{2\pi^2} g_t^2 \log\left(\frac{\Lambda^2}{v^2}\right)$

• For any given M_h , there is an upper bound on Λ

[DeGrassi et al, 1205.6497]



Don't want λ to be infinite

- Point where $\lambda \rightarrow \infty$ called Landau pole
- Without λh^4 interactions, theory is non-interacting
- Require quartic coupling be finite gives upper bound

$$\frac{1}{\lambda(\Lambda)} > 0 \qquad \qquad M_h^2 < \frac{32\pi^2 v^2}{9\log\left(\frac{\Lambda^2}{v^2}\right)}$$

Maybe vacuum can be unstable if lifetime is longer than lifetime of the universe

[Ellis,0906.0954]



Does WW Scattering Make sense?

• Without a Higgs



$$\sigma \sim rac{s}{v^4}$$

Not a sensible theory!

• With a Higgs



 σ finite at high energy

Higgs plays a special role

• Consider $2 \rightarrow 2$ particle elastic scattering

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2 s} \left| A \right|^2$$

• Partial wave decomposition of amplitude

$$A = 16\pi \sum_{l=0}^{\infty} (2l+1)P_l(\cos\theta)a_l$$

• a_l are the spin *l* partial waves

Optical theorem requires: $|\operatorname{Re}(a_l)| \le \frac{1}{2}$

UNITARITY CONSTRAINT

Use Unitarity to Bound Higgs $W^+W^- \rightarrow W^+W^-$

• High energy limit:

$$a_0^0 \rightarrow -\frac{M_h^2}{8\pi v^2}$$

• Heavy Higgs limit:

$$a_0^0 \rightarrow -\frac{s}{32\pi v^2} \implies E_c \sim 1.7 \text{ TeV}$$

 $\rightarrow \text{ New physics at the TeV scale}$

Light Higgs Boson makes Standard Model consistent

Problems with the Higgs Mechanism

- We often say that the Standard Model cannot be the entire story because of the quadratic divergences of the Higgs Boson mass
- But don't we just renormalize the mass using dimensional regularization ????
- Should experimentalists care about this?

Motivation for physics beyond the Standard Model

Renormalized Fermion Mass

• Renormalization of fermion mass:

$$\delta m_F = -\frac{3\lambda_F^2 m_F}{32\pi^2} \log\left(\frac{\Lambda^2}{m_F^2}\right) + \dots$$

Fermion mass renormalization is logarithmic

 $\delta m_F \approx m_F$

- $m_F \rightarrow 0$ increases the symmetry of the theory
- Yukawa coupling (proportional to mass) breaks symmetry and so corrections ≈ m_F

Scalars are very different

$$\delta M_h^2 = -\frac{\lambda_F^2 \Lambda^2}{8\pi^2} + \left(m_s^2 - m_F^2\right) \log\left(\frac{\Lambda}{m_F}\right) + \dots$$

M_h has quadratic sensitivity to high mass scales

 $M_h \rightarrow 0$ doesn't increase symmetry of theory

Nothing protects Higgs mass from large corrections

What's the problem?

 Compute M_h in dimensional regularization and absorb infinities into definition of M_h

$$M_h^2 = M_{h0}^2 + \frac{1}{\varepsilon}(...)$$

- Perfectly valid approach
- Except we know there is a high scale

SUSY....Our favorite model

- Quadratic sensitivity to high scale physics cancelled automatically if SUSY particles at TeV scale
- Cancellation result of *supersymmetry*, so happens at every order



Stop mass should be TeV scale*

* Not an exact statement

Supersymmetric Models as Alternative to Standard Model

Many New Particles:

- Spin $\frac{1}{2}$ quarks \Rightarrow spin 0 squarks
- Spin $\frac{1}{2}$ leptons \Rightarrow spin 0 sleptons
- Spin 1 gauge bosons \Rightarrow spin $\frac{1}{2}$ gauginos
- Spin 0 Higgs \Rightarrow spin $\frac{1}{2}$ Higgsino

Unbroken supersymmetry \Rightarrow degenerate masses of partners

SUSY must be a broken symmetry

Supersymmetric Theories

- Predict many new undiscovered particles (>29!)
- Very predictive models
 - Can calculate particle interactions in terms of a few parameters
 - Solve naturalness problem of Standard Model
- Any Supersymmetric particle eventually decays to the lightest supersymmetric particle (LSP) which is stable and neutral (assuming R parity)
 - Dark Matter Candidate

Supersymmetry (MSSM version)

 Good agreement with electroweak measurements if SUSY masses are 1-2 TeV



Fermion Masses

• In Standard Model, M_u from $\Phi_c=i\sigma_2\Phi^*$

$$L_{SM} = -\lambda_u \overline{Q}_L \Phi_c u_R + hc \qquad \Phi_c = \begin{pmatrix} \overline{\phi}^0 \\ -\phi^- \end{pmatrix} \qquad \lambda_u = -\frac{M_u \sqrt{2}}{v_{SM}}$$

- SUSY models don't allow Φ_c interactions
- Supersymmetric models always have at least *two Higgs doublets* with opposite hypercharge in order to give mass to up and down quarks

Two Higgs Doublet Models

- 8 degrees of freedom
- 3 form W_L^{\pm} , Z_L
- 5 physical Higgs bosons
 - h^0, H^0, A^0, H^{\pm}

$$H_{!} = \begin{pmatrix} \phi_{1}^{0*} \\ -\phi_{1}^{-} \end{pmatrix}$$

Gives up quark mass

$$H_2 = \begin{pmatrix} \phi_2^+ \\ \phi_2^0 \end{pmatrix}$$

Gives down quark mass

General 2 Higgs Doublet Model

• 6 free parameters, plus a phase

$$V(H_{1}, H_{2}) = \lambda_{1} (H_{1}^{+}H_{1} - v_{1}^{2})^{2} + \lambda_{2} (H_{2}^{+}H_{2} - v_{2}^{2})^{2} + \lambda_{3} [(H_{1}^{+}H_{1} - v_{1}^{2}) + (H_{2}^{+}H_{2} - v_{2}^{2})]^{2} + \lambda_{4} [(H_{1}^{+}H_{1})(H_{2}^{+}H_{2}) - (H_{1}^{+}H_{2})(H_{2}^{+}H_{1})] + \lambda_{5} [Re(H_{1}^{+}H_{2}) - v_{1}v_{2}\cos\xi]^{2} + \lambda_{6} [Im(H_{1}^{+}H_{2}) - v_{1}v_{2}\sin\xi]^{2}$$

- W and Z masses just like in Standard Model $M_W^2 = \frac{g^2(v_1^2 + v_2^2)}{2}$
- *ρ* parameter:

$$\rho = \frac{M_W}{M_Z \cos \theta_W} = 1$$

 ρ =1 for any number of Higgs doublets or singlets

Higgs Potential Restricted in SUSY Models

• Two Higgs doublets with opposite hypercharge

$$H_2 = \begin{pmatrix} \phi_2^+ \\ \phi_2^0 \end{pmatrix} \qquad \qquad H_1 = \begin{pmatrix} \varphi_1^{0^*} \\ -\varphi_1^- \end{pmatrix}$$

- Couplings fixed by supersymmetry
- H_1 has Y=-1/2, H_2 has Y=+1/2 [$Q_{em}=(T_3+Y)/2$]

$$V = \left|\mu\right|^{2} \left(\left|H_{1}\right|^{2} + \left|H_{2}\right|^{2}\right) + \frac{g^{2} + {g'}^{2}}{8} \left(\left|H_{2}\right|^{2} - \left|H_{1}\right|^{2}\right)^{2} + \frac{g^{2}}{2} \left|H_{1}^{*} \cdot H_{2}\right|^{2}$$

V is positive definite: minimum at $\langle V \rangle = \langle H_1 \rangle = \langle H_2 \rangle = 0$

No electroweak symmetry breaking!

The MSSM Philosophy

- Add all soft (dimension 3 or less) terms allowed
 They don't introduce quadratic Λ² contributions
- Potential has 3 free parameters (1 of which is fixed by M_W)

$$V = (m_1^2 + |\mu|^2)H_1H_1^{+} + (m_2^2 + |\mu|^2)H_2H_2^{+} - m_{12}^2(\varepsilon_{ab}H_1^aH_2^b + h.c.) + (\frac{g'^2 + g^2}{8})(H_1H_1^{+} - H_2H_2^{+})^2 + \frac{g^2}{2}|H_1H_2^{+}|^2$$

Gauge Couplings

 If m₁₂=0, potential is positive definite and no symmetry breaking

EWSB and SUSY Models

• Electroweak symmetry broken by vevs

$$\langle H_1 \rangle = \begin{pmatrix} v_1 \\ 0 \end{pmatrix} \qquad \langle H_2 \rangle = \begin{pmatrix} 0 \\ v_2 \end{pmatrix}$$

- 5 Physical Higgs bosons, h⁰, H⁰, H[±], A⁰
- W gets mass, $M_W^2 = g^2 (v_1^2 + v_2^2)/2$
- 2 free parameters, typically pick M_A , tan $\beta = v_2/v_1$
- Predict M_h, M_H, M_{H±}

$$M_{A}^{2} = m_{12}^{2} (\tan \beta + \cot \beta)$$
$$M_{H^{\pm}}^{2} = M_{A}^{2} + M_{W}^{2}$$

Neutral Higgs Masses

$$M_{h,H}^{2} = \frac{1}{2} \left[M_{A}^{2} + M_{Z}^{2} \pm \sqrt{\left(M_{A}^{2} + M_{Z}^{2}\right)^{2} - 4M_{Z}^{2}M_{A}^{2}\cos^{2}2\beta} \right]$$

- $M_h < M_Z \cos 2\beta$
- Theory implies light Higgs boson!

Theoretical Upper Bound on M_h

At tree level, $M_h < M_Z$

- Large corrections $O(G_Fm_t^2)$
 - Predominantly from stop squark loop

$$M_{h}^{2} \le M_{Z}^{2} \cos^{2} 2\beta + \frac{3G_{F}m_{t}^{4}}{\sqrt{2}\pi^{2} \sin^{2} \beta} \ln \left[\frac{\widetilde{m}_{t}^{2}}{m_{t}^{2}}\right]$$

• Stop mass should be TeV scale for naturalness

MSSM predicts a light Higgs boson!!!

Theoretical Upper Bound on M_h



- Mt⁴ enhancement
- Logarithmic dependence on stop mass

Now that we've found $M_h = 125 \text{ GeV}...$

- Maybe very heavy stops?
- Parameter space becomes quite restricted



SUSY is being Squeezed

 Need large stop masses or large SUSY breaking trilinear couplings to get M_h=125 GeV



[Pomarol]

Higgs Masses in MSSM

 $M_{H^{\pm}}^2 = M_A^2 + M_W^2$



Large M_A: Degenerate A, H, H[±] and light h

[Spira]

Find Higgs Couplings

• Higgs-fermion couplings:

$$L = -\frac{gm_d}{2M_W \cos\beta} \overline{d}d(H\cos\alpha - h\sin\alpha) + \frac{igm_d \tan\beta}{2M_W} \overline{d}\gamma_5 dA$$
$$-\frac{gm_u}{2M_W \sin\beta} \overline{u}u(H\sin\alpha + h\cos\alpha) + \frac{igm_d \cot\beta}{2M_W} \overline{u}\gamma_5 uA$$

- Couplings given in terms of $\alpha,\,\beta$
- Can be very different from Standard Model
- No new free parameters

 $\boldsymbol{\alpha}$ is angle that diagonalizes neutral Higgs mixing matrix

Higgs Couplings Different from SM

Lightest Neutral Higgs, h



Higgs Couplings in SUSY

Heavier Neutral Higgs, H



Higgs Couplings in MSSM

- At tree level, 2 free parameters (M_A , tan β)
 - At one loop tri-linear mixing terms and stop mass
- Lightest neutral Higgs

 $-M_A \rightarrow \infty$, h couplings look SM-like

- Heavier neutral Higgs, H⁰,A⁰,H[±]
 - $M_A \rightarrow \infty$, masses degenerate
 - Couplings to charge -1/3 enhanced for large tan β

Gauge Boson Couplings to Higgs

- $g_{hVV}^2 + g_{HVV}^2 = g_{hVV}^2(SM)$
- Vector boson fusion and Wh production always suppressed in MSSM

$$\frac{g_{hVV}}{g_{h,smVV}} = \sin(\beta - \alpha)$$
$$\frac{g_{HVV}}{g_{h,smVV}} = \cos(\beta - \alpha)$$



Normalized to SM couplings

Tools

- Calculate SM and MSSM Higgs branching ratios:
 - HDECAY
 - http://people.web.psi.ch/spira/hdecay/
- Calculate MSSM Higgs masses and Higgs branching ratios:
 - FEYNHIGGS
 - http://www.feynhiggs.de/

Both of these programs are very easy to use!

Limits from LEP



Limits on SUSY Higgs from LEP



Higgs Decays Changed at Large tan β

• MSSM: At large tan $\beta,$ rates to bb and $\tau^{*}\tau^{\scriptscriptstyle -}$ large



Rate to bb and $\tau^+\tau^-$ almost constant in MSSM for H, A

Large tanβ Changes Relative Importance of Production Modes



 $\tan\beta \ge 7$, b production mode larger than gg

[Kilgore]

Large Enhancements for large tan β



Light h (gluon fusion + bbh)



 $\sigma_{SM}(M_h=125 \text{ GeV}) = 19 \text{ pb}$



Light Higgs Branching ratios in MSSM



tan β =50, M_A=120 GeV, M_{stop}=1 TeV \rightarrow M_h =120 GeV

ATLAS Exclusion



Can also look for charged H⁺

- $pp \rightarrow tbH^+$ is largest production mode
- $t \rightarrow bH^+$ if kinematically allowed



Charged Higgs Limits

Mostly from B decays



Many Possibilities

- Generic 2 Higgs doublet model (not SUSY)
 - Similar decoupling features as MSSM
- More complicated SUSY model
 - NMSSM is SUSY model with additional singlet
- Extra Higgs singlets which mix with Higgs but don't couple to anything else
 - Reduces all Higgs couplings

Finale

- It's going to be a fun year!
- With 20-30 fb⁻¹ we will be firmly in the era of precision Higgs physics

