



Achilles: A Modern Theorist-Driven Event Generator

Joshua Isaacson

In Collaboration with: S. Höche, W. Jay, D. Lopez Gutierrez, A. Lovato, P.A.N. Machado, N. Rocco

Based on: 2007.15570, 2110.15319, 2205.06378

Cross Theory Generator Working Group Meeting

7 June 2022



Achilles: A CHIcago Land Lepton Event Simulator

Project Goals:

- Theory driven
- Develop modular neutrino event generator
- Provide means for easy extension by end users
- Provide automated BSM calculations for neutrino experiments
- Evaluate theory uncertainties

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Version: Authors:	1.0.0 Joshua Pedro #	Isaacso A. Macha	n, Williar do, Noemi	n Jay, Al Rocco			



ion

Cascade

Why a new generator?

Oscillation Measurements

•
$$P_{\nu_{\mu} \to \nu_{e}} \sim \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 L}{4E}\right) \to \Phi_e(E,L)/\Phi_{\mu}(E,0)$$

- Only measure events and not fluxes directly: $N(E_{reco}, L) \propto \sum_{i} \Phi(E, L) \sigma_i(E) f_{\sigma_i}(E, E_{reco})$
- Fit oscillation parameters by taking ratio of number of events in ${\cal E}_{reco}$ bins
- Cross sections do not exactly cancel in ratio, thus they are crucial
- Requires fully differential predictions

Other Measurements

- DUNE and HyperK near detectors are general purpose
- Leverage them for BSM searches
- Requires both SM and BSM fully differential predictions

Introduction

Hard Interactic

Cascad

Why a new generator? $(e4\nu)$



• State of the art is ok for inclusive

• Exclusive results need significant work

[Nature 599, 565–570 (2021)]



Why a new generator? (NOvA)

Source of Uncertainty	$\nu_e \text{ signal } (\%)$	Total beam background (%)
Cross-section and FSI	7.7	8.6
Normalization	3.5	3.4
Calibration	3.2	4.3
Detector response	0.67	2.8
Neutrino flux	0.63	0.43
ν_e extrapolation	0.36	1.2
Total systematic uncertainty	9.2	11
Statistical uncertainty	15	22
Total uncertainty	18	25

Table 1: Effect of 1 or variations of the systematic uncertainties in the total v_e signal and background predictions in the NOvA experiment². The systematic uncertainties are from the latest NOvA results with $8.85x10^{20}$ protons on target.

- Cross section uncertainty one of dominant uncertainties
- Unclear if correctly fully estimated
- NOvA systematics and statistical uncertainty equal
- DUNE and HyperK will have significantly more events



Separating Primary Interaction and Cascade

General Lepton-Nucleus Scattering Cross Section

$$\mathrm{d}\sigma = \left(\frac{1}{|v_A - v_\ell|} \frac{1}{4E_A^{\mathrm{in}} E_\ell^{\mathrm{in}}}\right) |\mathcal{M}|^2 \prod_f \frac{\mathrm{d}^3 p_f}{(2\pi)^3} (2\pi)^4 \delta^4 \left(k_A + k_\ell - \sum_f p_f\right)$$



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Matrix Element Schematically

$$\left|\mathcal{M}\left(\{k\} \to \{p\}\right)\right|^2 = \left|\sum_{p'} \mathcal{V}(\{k\} \to \{p'\}) \times \mathcal{P}(\{p'\} \to \{p\})\right|^2$$



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Primary interaction —



Separating Primary Interaction and Cascade

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- Primary interaction
- Evolution out of nucleus _____



Separating Primary Interaction and Cascade

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Matrix Element Schematically Approximation

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- Primary interaction
- Evolution out of nucleus
- Approximate as incoherent product of primary interaction and cascade

Hard Interaction

Cascade

Factorization

- For Quasielastic scattering, factorize primary interaction as: $|\Psi_f\rangle = |p\rangle \otimes |\Psi_f^{A-1}\rangle$
- Initial state given via spectral function (probability distribution of removing a "hole" nucleon):

$$S_h(\mathbf{k}_h, E') = \sum_{f_{A-1}} |\langle \Psi_0 | k \rangle \otimes |\Psi_f^{A-1}\rangle|^2 \delta(E' + E_0^A - E_f^{A-1})$$

- Here we use spectral function obtained from correlated basis function theory [Phys. A 579, 493 (1994)]
- Spectral function normalized as:

$$\int \frac{\mathrm{d}k_h}{(2\pi)^3} \mathrm{d}E' S_h(\mathbf{k}_h, E') = \begin{cases} Z, & h = \mathrm{p}, \\ A - Z, & h = \mathrm{n}. \end{cases}$$



[Rev. Mod. Phys. 80, 189 (2008)]

J



BSM Motivation: MiniBooNE and MicroBooNE



[arXiv:2110.14054]

- MiniBooNE sees excess of events
- MicroBooNE does not see excess of single electron events
- Excess can be from multiple lepton final states
- Event generators can not simulate these processes





Using Currents

Using tensors:

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = \sum_{i,j} L^{(ij)}_{\mu\nu} W^{(ij)\mu\nu} = L^{(\gamma\gamma)}_{\mu\nu} W^{(\gamma\gamma)\mu\nu} + L^{(\gamma Z)}_{\mu\nu} W^{(\gamma Z)\mu\nu} + L^{(Z\gamma)}_{\mu\nu} W^{(Z\gamma)\mu\nu} + L^{(ZZ)}_{\mu\nu} W^{(ZZ)\mu\nu} + \cdots$$

Using Currents:

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = \left|\sum_i L^{(i)}_{\mu} W^{(i)\mu}\right|^2$$

Interferences handled automatically using currents

Interface to tensors provided for nuclear calculations that **must** be expressed using tensors.



Hard Interaction

Cascad

Handling Form Factors

Nuclear one-body current operators:

$$\begin{aligned} \mathcal{J}^{\mu} &= \left(\mathcal{J}^{\mu}_{V} + \mathcal{J}^{\mu}_{A}\right) \\ \mathcal{J}^{\mu}_{V} &= \gamma^{\mu} \mathcal{F}^{a}_{1} + i \sigma^{\mu\nu} q_{\nu} \frac{\mathcal{F}^{a}_{2}}{2M} \\ \mathcal{J}^{\mu}_{A} &= -\gamma^{\mu} \gamma_{5} \mathcal{F}^{a}_{A} - q^{\mu} \gamma_{5} \frac{\mathcal{F}^{a}_{P}}{M} \end{aligned}$$

Coherent Form Factors (spin-0 nucleus):

$$\mathcal{J}^{\mu} = (p_{\mathsf{in}} + p_{\mathsf{out}})^{\mu} \mathcal{F}_{\mathsf{coh}}$$

Standard Model Form Factors:

$$\begin{aligned} \mathcal{F}_{i}^{\gamma(p,n)} &= F_{i}^{p,n}, \qquad \mathcal{F}_{A}^{\gamma} = 0\\ \mathcal{F}_{i}^{W(p,n)} &= F_{i}^{p} - F_{i}^{n}, \qquad \mathcal{F}_{A}^{W} = F_{A}\\ \mathcal{F}_{i}^{Z(p)} &= \left(\frac{1}{2} - 2\sin^{2}\theta_{W}\right)F_{i}^{p} - \frac{1}{2}F_{i}^{n},\\ \mathcal{F}_{i}^{Z(n)} &= \left(\frac{1}{2} - 2\sin^{2}\theta_{W}\right)F_{i}^{n} - \frac{1}{2}F_{i}^{p}\\ \mathcal{F}_{A}^{Z(p)} &= \frac{1}{2}F_{A}, \qquad \mathcal{F}_{A}^{Z(n)} = -\frac{1}{2}F_{A}\end{aligned}$$

Hard Interaction

Cascad

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Straight-forward to extend to BSM if CVC is valid

Conclusions

Recursive Matrix Element Generation

$$\mathcal{J}_{\alpha}(\pi) = P_{\alpha}(\pi) \sum_{\mathcal{V}_{\alpha}^{\alpha_{1},\alpha_{2}}} \sum_{\mathcal{P}_{2}(\pi)} \mathcal{S}(\pi_{1},\pi_{2}) V_{\alpha}^{\alpha_{1},\alpha_{2}}(\pi_{1},\pi_{2}) \mathcal{J}_{\alpha_{1}}(\pi_{1}) \mathcal{J}_{\alpha_{2}}(\pi_{2})$$

$$L^{(i)}_{\mu\nu}(1,\ldots,m) = \mathcal{J}^{(i)}_{\mu}(1,\ldots,m)$$
$$L^{(i,j)}_{\mu\nu}(1,\ldots,m) = \mathcal{J}^{(i)}_{\mu}(1,\ldots,m) \mathcal{J}^{(j)\dagger}_{\nu}(1,\ldots,m)$$

Berends-Giele Recursion

- Reuse parts of calculation
- Most efficient for high multiplicity
- Reduces computation from $\mathcal{O}\left(n!\right)$ to $\mathcal{O}\left(3^{n}\right)$

[Nucl. Phys. B306(1988), 759]





Cascac

Phase Space Generation

$$d\Phi_n(a,b;1,\ldots,n) = \delta^{(4)} \left(p_a + p_b - \sum_{i=1}^n p_i \right) \left[\prod_{i=1}^n \frac{d^4 p_i}{(2\pi)^3} \delta\left(p_i^2 - m_i^2 \right) \Theta\left(p_{i_0} \right) \right]$$

The above phase space definition does not contain the handling of initial states.

Algorithms for n-body phase space generation

- RAMBO [Comput. Phys. Commun. 40(1986) 359]
- Multi-channel techniques [hep-ph/9405257]
 - Recursive Phase Space [arXiv:0808.3674]



Introduction	Hard Interaction	Cascade	Conclusions	
Results				
Processes Consid	lered:	Parameters:		
 Electron-Ca 	arbon Scattering	• Only quasielastic scatte	ering (coherent for	
Neutrino-Ca	arbon Scattering	HNL) is included and n	io FSI	
• Neutrino Tridents		 EM Form Factors: Kelly [PRC 70, 068202 (2004)] Coherent Form Factor: Lovato [1305.6959] 		
 Dirac/Majorana Heavy Neutral Lepton [1807.09877] 				
Experimental Setup:		• Axial Form Factor:		
• Target Nuc	leus: Carbon (Argon for HNL)	• Dipole • $M_A = 1.0$ GeV		
Electron: 9	61 MeV and 1299 MeV	• $g_A = 1.2694$		
• Neutrino: 1000 MeV		• $\alpha = 1/137$		
 Validating beam fluxes 		• $G_F = 1.16637 \times 10^{-5}$		
NOTE: All processes are fully differential		• $M_Z = 91.1876 \text{ GeV}$		

Introduction

Hard Interaction

Cascade

Electron Scattering



Neutrino Total Cross Section





Neutrino Differential Cross Section





Heavy Neutral Lepton



Parameters:

- $m_{N'} = 420 \text{ MeV}$
- $m_{Z'} = 30 \text{ MeV}$

•
$$\alpha_D = 0.25$$

•
$$\alpha \epsilon^2 = 2 \times 10^{-10}$$

•
$$|U_{42}^{\mu}| = 9 \times 10^{-7}$$

- Widths of N^\prime and Z^\prime automatically calculated based on input parameters
- Handles both Dirac and Majorana fermions
- Results are flux-averaged over the MiniBooNE / MicroBooNE neutrino flux

Heavy Neutral Lepton



Heavy Neutral Lepton

- No cuts applied yet
- Typical opening angle around 5-6 degrees
- Working on scanning parameter space





- Need to include background to compare to MiniBooNE data
- Simulate possible MicroBooNE limits



Conclusions

MicroBooNE Simulation



Image generated by the MicroBooNE collaboration using Achilles

- Working on implementing into MicroBooNE Pipeline
- Developing interface to LArSoft



Final State Interactions

Modify Primary Interaction:

- Captures rate change from FSI
- Loses all information about hadronic final state
- Primarily done using folding functions

Intranuclear Cascade:

- Unitary process (*i.e.* no rate change)
- Contains information about hadronic final state
- Primarily done via Monte Carlo methods

Note: Both approaches attempt to capture effects from nuclear potential. Therefore, can only use one or the other to avoid double counting effects.



Hard Interaction

Cascade

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Algorithm Overview





Propagation with Potential

Initial Momentum: 250 MeV



• Blue: Non-relativistic potential ($E = \sqrt{p^2 + m^2} + V$)

[Phys. Rev. C. 38, 2967]

- Propagation using symplectic integrator for non-separable Hamiltonians [1609.02212]
- Energy is conserved to a high degree of precision
- Extremely stable

• Red: Relativistic potential ($E = \sqrt{p^2 + (m+S)^2} + V$)

[Phys. Rev. C. 80, 034605]



Hard Interaction

Cascade

Propagation with Potential



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CLAS/e4v Comparison

CLAS/e4v cuts:

- Select $1p0\pi$ events
- Protons: $p_p > 300$ MeV, $12^{\circ} < \theta_p$.
- electrons: $E_e > 0.4, 0.55, 1.1 \text{ GeV}$, $\theta_e^i > \theta_0^i + \frac{\theta_1^i}{p_e[\text{GeV}]}, \ \theta_0^i = 17^\circ, 16^\circ, 13.5^\circ$, $\theta_1^i = 7^\circ, 10.5^\circ, 15^\circ \text{ for}$ $E_{\text{beam}} = 1.159, 2.257, 4.453 \text{ GeV}$ respectively.

Simulation details:

- ACHILLES only has Quasielastic channel so far
- Events are reweighted by Q^4/GeV^4 (as done in the analysis)







CLAS/e4v Comparison



• Mimics Cherenkov detectors

• Mimics LArTPC detectors



CLAS/e4v Comparison



$$\mathbf{p}^T = \mathbf{p}_e^T + \mathbf{p}_p^T$$



CLAS/e4v Comparison



$$\mathbf{p}^T = \mathbf{p}_e^T + \mathbf{p}_p^T$$



CLAS/e4v Comparison



$$\mathbf{p}^T = \mathbf{p}_e^T + \mathbf{p}_p^T$$



New Observables

New Observables:

• Momentum of 1st proton





New Observables

New Observables:

- Momentum of 1st proton
- Momentum of 2nd proton



New Observables

New Observables:

- Momentum of 1st proton
- Momentum of 2nd proton
- Momentum of 3rd proton





New Observables

New Observables:

- Momentum of 1st proton
- Momentum of 2nd proton
- Momentum of 3rd proton
- Reconstructed beam direction:

$$\cos \theta_{\rm rec} \equiv \frac{\hat{\mathbf{k}}_e \cdot \mathbf{p}_{\rm out}}{|\mathbf{p}_{\rm out}|}$$





Conclusions

Current Status:

- DUNE and HK will require precision neutrino event generators
- ACHILLES aims to be a modular theory driven generator to address these needs
- BSM important for the current and next generation neutrino experiments
- Robust BSM program requires automating theory calculations
- Comparison of cascade results with CLAS/e4v experiment

Future Steps:

- Implement QED showers to handle radiative corrections
- Interface with LArSoft
- Implement MEC, Resonance, and DIS processes
- Continue to improve cascade modeling



FeynRules

- Mathematica Program
- Takes model file and Lagrangian as input
- Calculates the Feynman rules
- Outputs in Universal FeynRules Output (UFO) format



[arXiv:0806.4194, arXiv:1310.1921]

Universal FeynRules Output (UFO)

Example QED ($e^+e^-\gamma$ Vertex):

- Python output files
- Contains model-independent files and model-dependent files
- Contains all information to calculate any tree level matrix element
- Has parameter file to adjust model parameters to scan allowed regions

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \bar{\psi} \left(i D^{\mu} \gamma_{\mu} - m \right) \psi$$

$$V_{e^+e^-\gamma} = ie\gamma^\mu = \gamma \checkmark$$

[arXiv:1108.2040]

Universal FeynRules Output (UFO) Example for photon-electron vertex

```
e minus = Particle(pdg code=11, name='e-', antiname='e+',
                      spin=2, color=1, mass=Param.ZERO,
                      width=Param.ZERO, texname='e-',
                      antitexname='e+', charge=-1,
                      GhostNumber=0, LeptonNumber=1,
                      Y=0)
V 77 = Vertex(name='V 77')
              particles=[ P.e plus , P.e minus , P.a ],
              color=[ '1' ], lorentz=[ L.FFV1 ],
              couplings = \{(0,0): C, GC \}
FFV1 = Lorentz(name='FFV1', spins=[ 2, 2, 3 ],
               structure = 'Gamma(3,2,1)')
GC_3 = Coupling(name='GC_3', value='-(ee*complex(0,1))'.
                order={'QED':1})
```

Tree Level Matrix Element Generators

- ALPGEN [arXiv:hep-ph/0206293]
- AMEGIC [arXiv:hep-ph/0109036]
- COMIX [arXiv:0808.3674]
- CALCHEP [arXiv:1207.6082]
- HERWIG [arXiv:0803.0883]
- MADGRAPH [arXiv:1405.0301]
- WHIZARD [arXiv:0708.4233]
- etc.





[arXiv:1702.05725]

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[arXiv:1702.05725]

 $\mathrm{d}\sigma \propto \mathrm{d}\Phi_2(a,b;1,2) \ \mathrm{d}^4 p_a \ \mathrm{d}^3 p_b$



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• Phase space:
$$d\Phi_2(a,b;1,2) = \frac{\lambda(s_{ab},s_1,s_2)}{16\pi^2 2s_{ab}} d\cos\theta_1 d\phi_1$$



$$\mathrm{d}\sigma \propto \ \mathrm{d}\Phi_2(a,b;1,2) \ \mathrm{d}^4p_a \ \mathrm{d}^3p_b$$

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• Initial nucleon:
$$d^4p_a = |\vec{p}_a|^2 dp_a dE_r d\cos\theta_a d\phi_a \sqrt{1-1}$$

$$\mathrm{d}\sigma \propto \mathrm{d}\Phi_2(a,b;1,2) \ \mathrm{d}^4 p_a \ \mathrm{d}^3 p_b$$

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• Initial nucleon:
$$d^4p_a = |\vec{p_a}|^2 dp_a dE_r d\cos\theta_a d\phi_a$$

• Initial lepton (Here only monochromatic): $d^3p_b = \delta^3(p_b - p_{beam})d^3p_b$

 $\mathrm{d}\sigma \propto \mathrm{d}\Phi_2(a,b;1,2) \,\mathrm{d}^4 p_a \,\mathrm{d}^3 p_b$

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• Initial lepton (Here only monochromatic): $d^3p_b = \delta^3(p_b - p_{beam})d^3p_b$

Quasielastic Delta Function: $\delta(E_b - E_1 - E_r + m - E_2)$ Phase Space Delta Function: $\delta(E_a + E_b - E_1 - E_2)$ Define initial nucleon energy as $E_a = m - E_r$. Allows use of phase space tools developed at LHC.

Multi-channel Integration



- Both diagrams contribute to cross section
- They have different pole structures
- Need method to sample these structures efficiently (i.e. $|A_1 + A_2|^2$)

Multi-channel Integration and VEGAS

Multi-channel Integration

- Generate PS efficiently for $|\mathcal{A}_1|^2$ or $|\mathcal{A}_2|^2$
- Do not know how to efficiently sample $2Re(\mathcal{A}_1\mathcal{A}_2^\dagger)$
- Define channels: C_1 and C_2
- Generate events according to distributions g_i for channel i

$$\int d\vec{x} f(\vec{x}) = \sum_{i} \alpha_{i} \int d\vec{x} g_{i}(\vec{x}) \frac{f(\vec{x})}{g_{i}(\vec{x})}$$

• Optimize α_i to minimize variance

VEGAS

- Adaptive importance sampling
- Use this to get interference terms more accurately



Phase space can be decomposed as:

$$\mathrm{d}\Phi_n(a,b;1,\ldots,n) = \mathrm{d}\Phi_{n-m+1}(a,b;m+1,\ldots,n)\frac{\mathrm{d}s_\pi}{2\pi}\mathrm{d}\Phi_m(\pi;1,\ldots,m)$$

Iterate until only $1 \rightarrow 2$ phase spaces remain. Basic building blocks:

$$S_{\pi}^{\rho,\pi\setminus\rho} = \frac{\lambda(s_{\pi}, s_{\rho}, s_{\pi\setminus\rho})}{16\pi^2 2 s_{\pi}} \operatorname{d}\cos\theta_{\rho} \operatorname{d}\phi_{\rho}$$
$$T_{\alpha,b}^{\pi,\overline{\alpha}b\overline{n}} = \frac{\lambda(s_{\alpha b}, s_{\pi}, s_{\overline{\alpha}b\overline{n}})}{16\pi^2 2 s_{\alpha b}} \operatorname{d}\cos\theta_{\pi} \operatorname{d}\phi_{\pi}$$

Momentum conservation: $(2\pi)^4 d^4 p_{\overline{\alpha b}} \delta^{(4)}(p_{\alpha} + p_b - p_{\overline{\alpha b}})$

Neutrino Tridents







J. Isaacson

Symplectic Integration for non-separable Hamiltonians

• Create copy of Hamiltonian:

$$\overline{H}(q, p, x, y) \equiv H_A(q, y) + H_B(x, p) + \omega H_C(q, p, x, y)$$
$$H_C(q, p, x, y) = |q - x|^2 / 2 + |p - y|^2 / 2$$

• Time step for each Hamiltonian:

$$\begin{split} \phi_{H_A}^{\delta} &: \begin{bmatrix} q \\ p \\ x \\ y \end{bmatrix} \to \begin{bmatrix} q \\ p - \delta \partial_q H(q, y) \\ x + \delta \partial_y H(q, y) \\ y \end{bmatrix}, \phi_{H_B}^{\delta} &: \begin{bmatrix} q \\ p \\ x \\ y \end{bmatrix} \to \begin{bmatrix} q + \delta \partial_p H(x, p) \\ p \\ x \\ y - \delta \partial_x H(x, p) \end{bmatrix}, \\ \phi_{H_C}^{\delta} &: \begin{bmatrix} q \\ p \\ x \\ y \end{bmatrix} \to \frac{1}{2} \begin{bmatrix} \begin{pmatrix} q + x \\ p + y \end{pmatrix} + R(\delta) \begin{pmatrix} q - x \\ p - y \end{pmatrix} \\ \begin{pmatrix} q + x \\ p + y \end{pmatrix} - R(\delta) \begin{pmatrix} q - x \\ p - y \end{pmatrix} \end{bmatrix}, \end{split}$$

• Full second order time step:

$$\phi_2^{\delta} = \phi_{H_A}^{\delta/2} \circ \phi_{H_B}^{\delta/2} \circ \phi_{\omega H_C}^{\delta} \circ \phi_{H_B}^{\delta/2} \circ \phi_{H_A}^{\delta/2}.$$