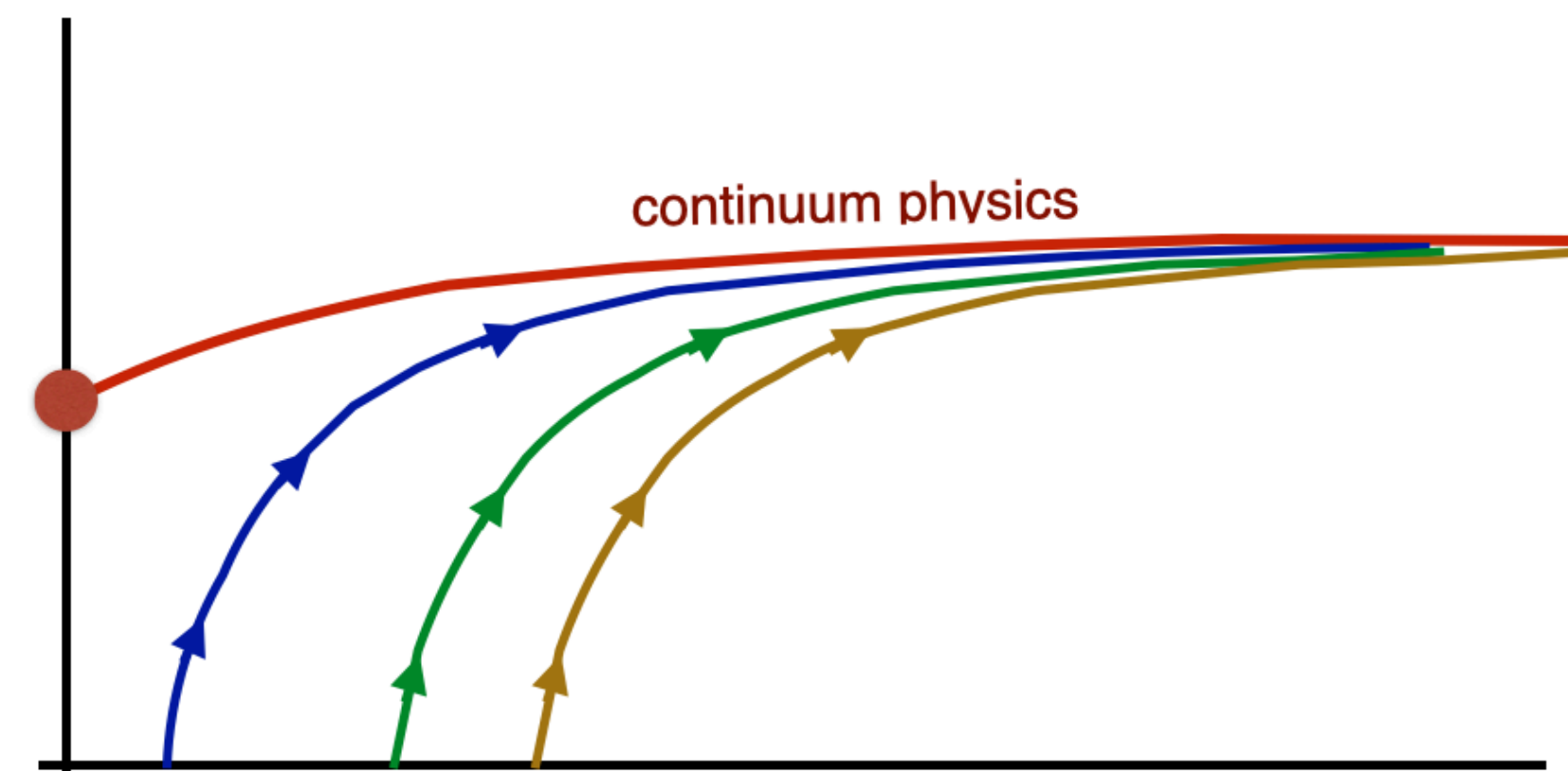


The conformal fixed point of the $SU(3)$ with 10 fundamental flavors

Anna Hasenfratz
University of Colorado Boulder

July 31, 2023

Based on recent work by
A.H., E.Neil, Y.Shamir,
B.Svetitsky, O.Witzel,
arXiv:2306.07236



Gradient flow vs continuous RG transformations

A. Carosso, AH, E. Neil,
PRL 121,201601 (2018)

GF can be *interpreted* as **continuous real space RG** with $\mu \propto 1/\sqrt{8t}$

- for *local* operators
- in the infinite volume limit

$$\text{Ex-1: } g_{GF}^2 = \mathcal{N} t^2 \langle E(t) \rangle \implies \beta(g^2)$$

$$\text{Ex-2: } \mathcal{O} = \bar{\psi}(x)\Gamma\psi(x) \implies \gamma_{\mathcal{O}}(g^2)$$

Continuous β function

$$g_{GF}^2 = \mathcal{N} t^2 \langle E(t) \rangle \implies \beta_{GF}(a; g_{GF}^2) = -t \frac{dg_{GF}^2(a; t)}{dt}$$

- estimate flowed energy density $\langle E(t) \rangle$: W, C or S
- $\beta(g^2)$ defined in $am_f = 0$ chiral limit : simulations are at κ_{cr}

- Fodor et al, EPJ Web Conf. 175, 08027 (2018)
- AH, O. Witzel
Phys. Rev. D 101 (2020) 3, 034514

Analysis:

- (1) infinite volume limit : extrapolate in $(a/L)^4 \rightarrow 0$ at fixed bare coupling and t/a^2
- (2) interpolate $\beta(g^2; t)$ vs $g^2(t)$, all t/a^2
- (3) continuum limit : $t/a^2 \rightarrow \infty$ while keeping g_{GF}^2 fixed

Same approach as $N_f = 0$

- AH, C.T. Peterson, J. VanSickle, O. Witzel,
Phys. Rev. D 108 (2023) 1, 014502

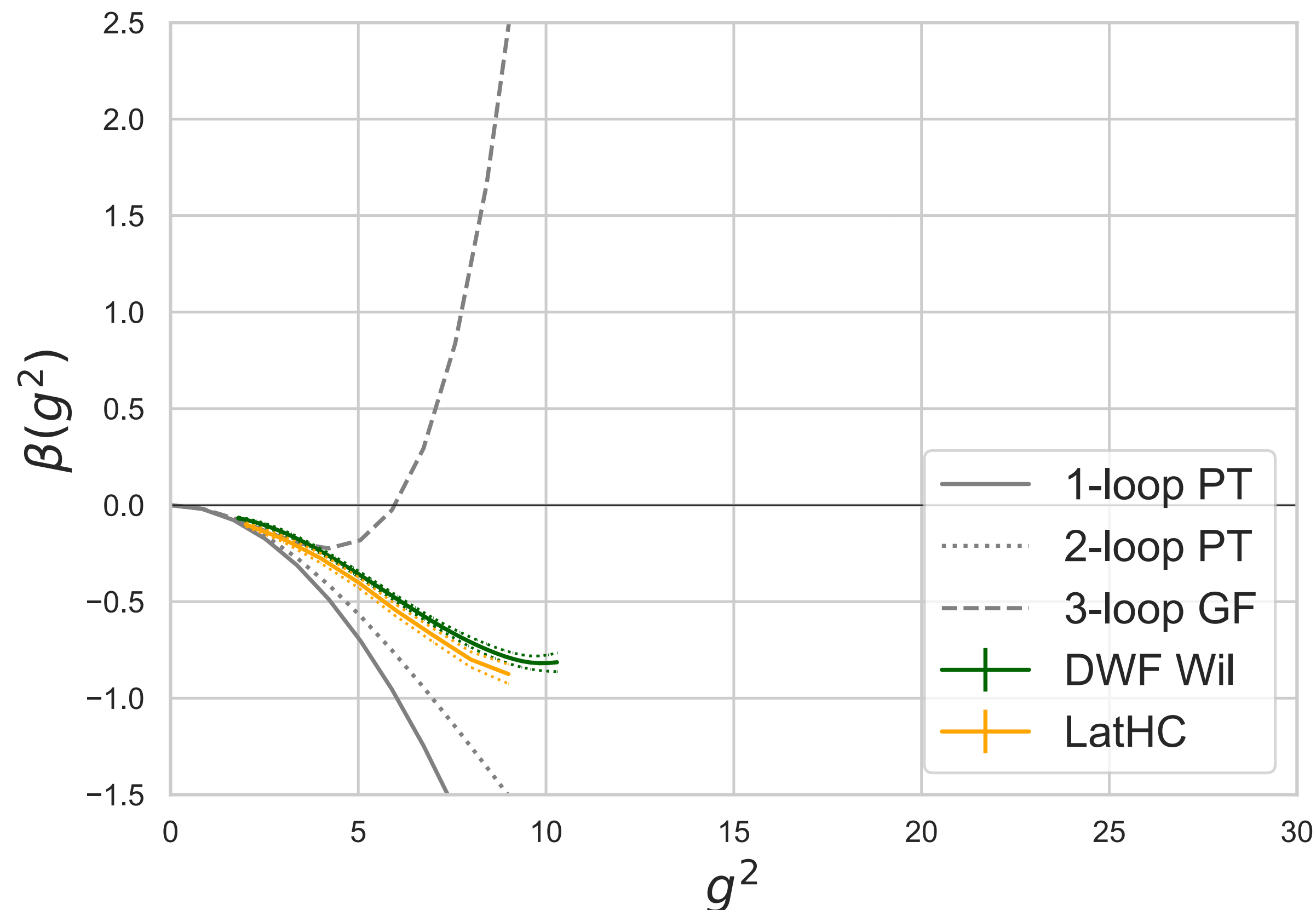
SU(3) gauge, $N_f = 10$ flavors

This system is close to the conformal sill

Prior simulations were limited by strong lattice artifacts

- ➔ restricted to $g^2 \lesssim 10$
- ➔ no prediction of IR dynamics

- AH, C. Rebbi, O. Witzel,
Phys.Rev.D 101 (2020) 11, 114508
- O. Witzel, this conference
- J. Kuti, Z. Fodor, K. Hollad, C-H.
Wong arXiv: [2203.15847](https://arxiv.org/abs/2203.15847)



Continuous β function

- staggered fermions (LatHC)
- Mobius DW

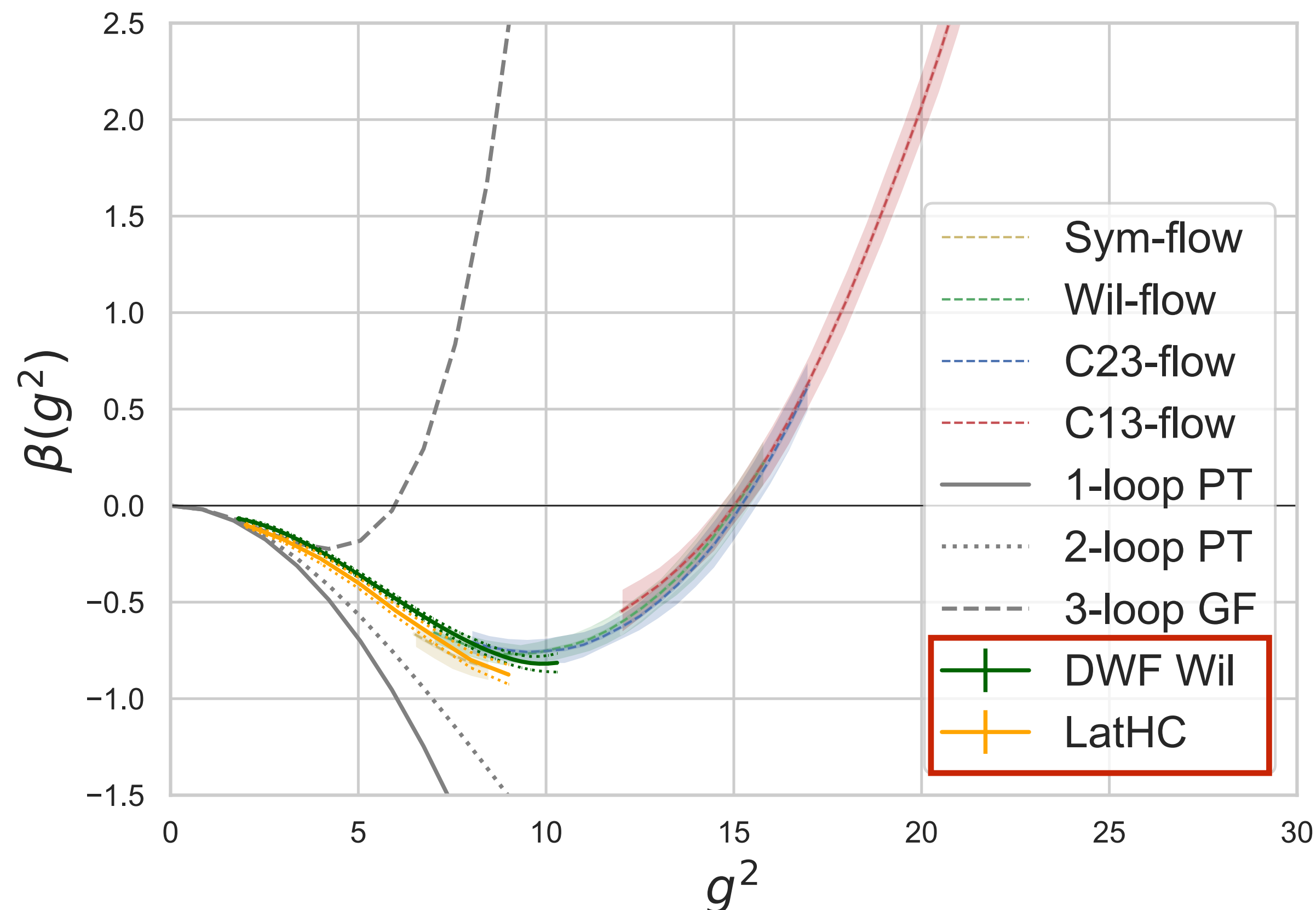
give consistent results

SU(3) gauge, $N_f = 10$ flavors

Our work with Wilson fermions reaches $g^2 \gtrsim 25$

- ➔ reveals IRFP at $g_*^2 = 15.0(5)$
- ➔ consistent with prior results

- A.H., E.Neil, Y.Shamir,
B.Svetitsky, O.Witzel,
arXiv:2306.07236



We utilize several improvements[#]

- PV improved action
- different flows
- bootstrap error analysis

[#]See applications in:

- A.H., E.Neil, Y.Shamir,
B.Svetitsky, O.Witzel,
Phys.Rev.D 107 (2023) 11, 114504
- Y. Shamir, this conference

Improved action: heavy PV bosons

AH, Y. Shamir, B. Svetitsky,
PRD104, 074509 (2021)

Many flavor systems suffer from *bulk phase transitions* in strong coupling

- Fermions generate a *positive* effective gauge action (hopping expansion)

$$S_{eff}^{(f)} = \frac{N_s}{(2am_f)^4} \sum_p \text{ReTr} V_{\square} + c \frac{N_s}{(2am_f)^6} \sum_{6link} \text{ReTr} V_{6-link} \dots$$

Bare gauge coupling g_0^2 *increases* to compensate:

➔ rough gauge configurations, large cutoff effects

Improved action: heavy PV bosons

AH, Y. Shamir, B. Svetitsky,
PRD104, 074509 (2021)

Compensate with **heavy Pauli-Villars bosons**

- same interaction as fermions but with *bosonic statistics* : $S_{eff}^{(PV)} = - S_{eff}^{(f)}$
- g_0^2 decreases, cutoff effects decrease
- keep $am_{PV} \sim \mathcal{O}(1)$ fixed: in the IR ($a \rightarrow 0$) the PV bosons decouple
 - e.g.: $N_f = 10, \beta = 6.3, \kappa_f = 0.12677$: $am_{AWI} \lesssim 10^{-4}$, $am_{PS} \propto 1/L$
 - $\kappa_{PV} = 0.1$ ($am_{PV} = 1$) : $am_{AWI} = 0.73$, $am_{PS} = 1.80$

Improved action: heavy PV bosons

AH, Y. Shamir, B. Svetitsky,
PRD104, 074509 (2021)

Add many PV bosons to reduce the lattice fluctuations:

- ➔ extends accessible renormalized coupling range
- ➔ only a local effective gauge action - IR properties are unchanged

Different flows / RG transformation

Vary the action for gradient flow (improved RG):

$$S_{flow} = c_p S_p + c_r S_{rec} \quad , \quad c_p + 8c_r = 1$$

$$S : \quad c_p = 5/3$$

$$W : \quad c_p = 1$$

$$C23 : \quad c_p = 2/3$$

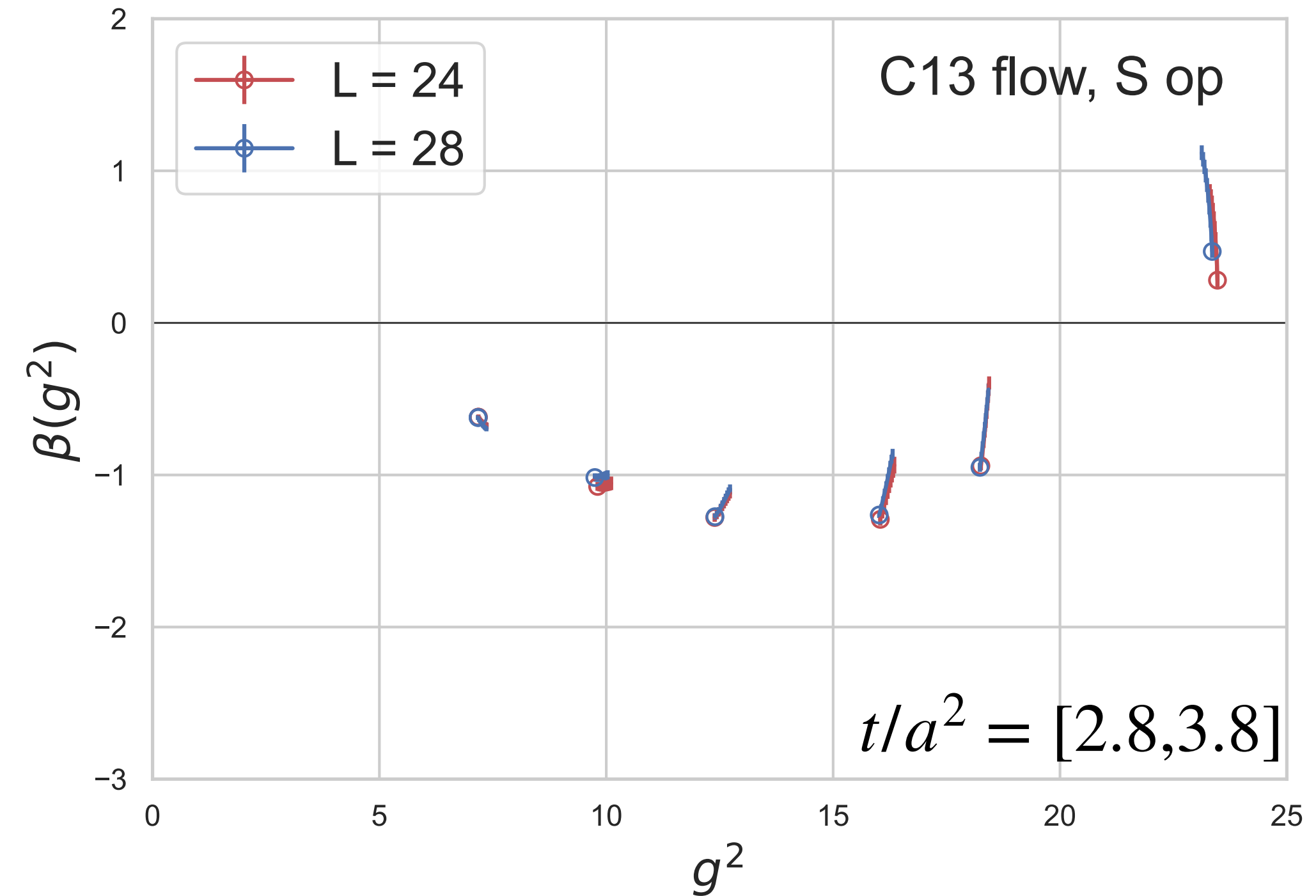
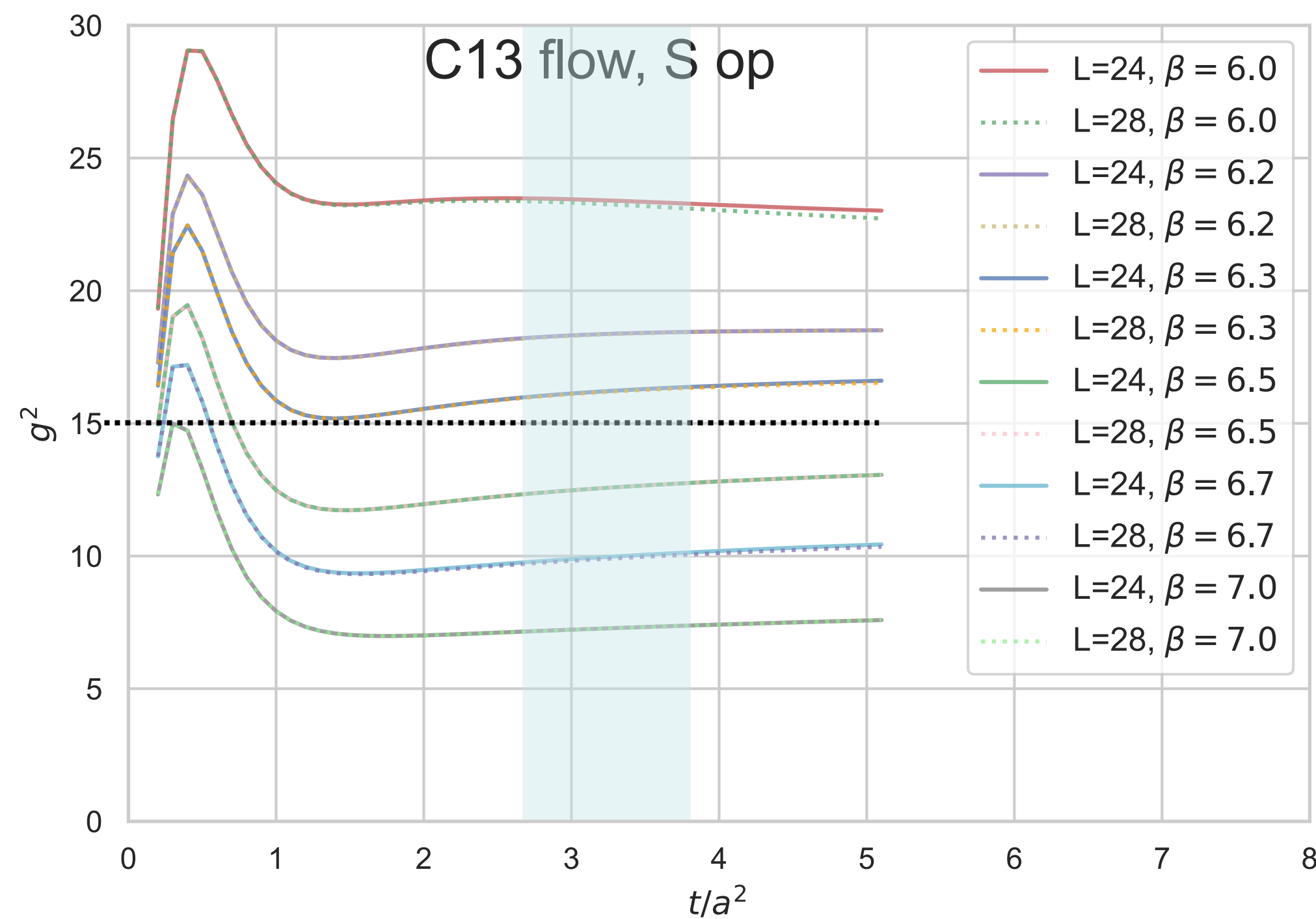
$$C13 : \quad c_p = 1/3$$

All flows should give the same continuum limit
 g^2 range is increasing as c_p decreases

(1) Infinite volume extrapolation

Volume dependence at leading order $\propto (t^2/L^4)$ (Wilsonian RG in finite volume)

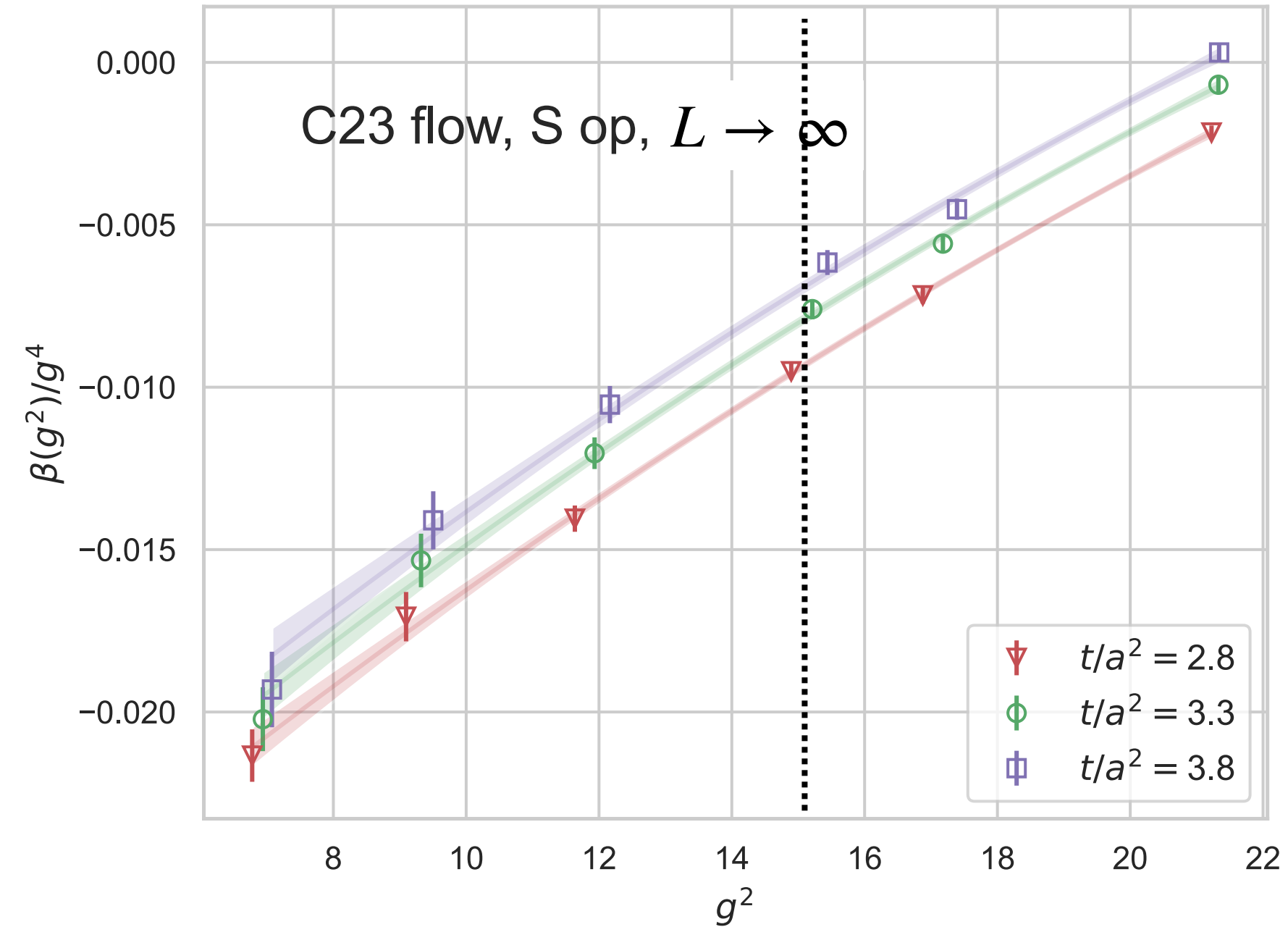
- we use $t/a^2 = [2.8, 3.8]$, $L = 24, 28$: $t^2/L^4 < 6 \times 10^{-5}$
- take $(a/L)^4 \rightarrow 0$ for g^2 , $\beta(g^2)$ for *each* bare coupling and flow time



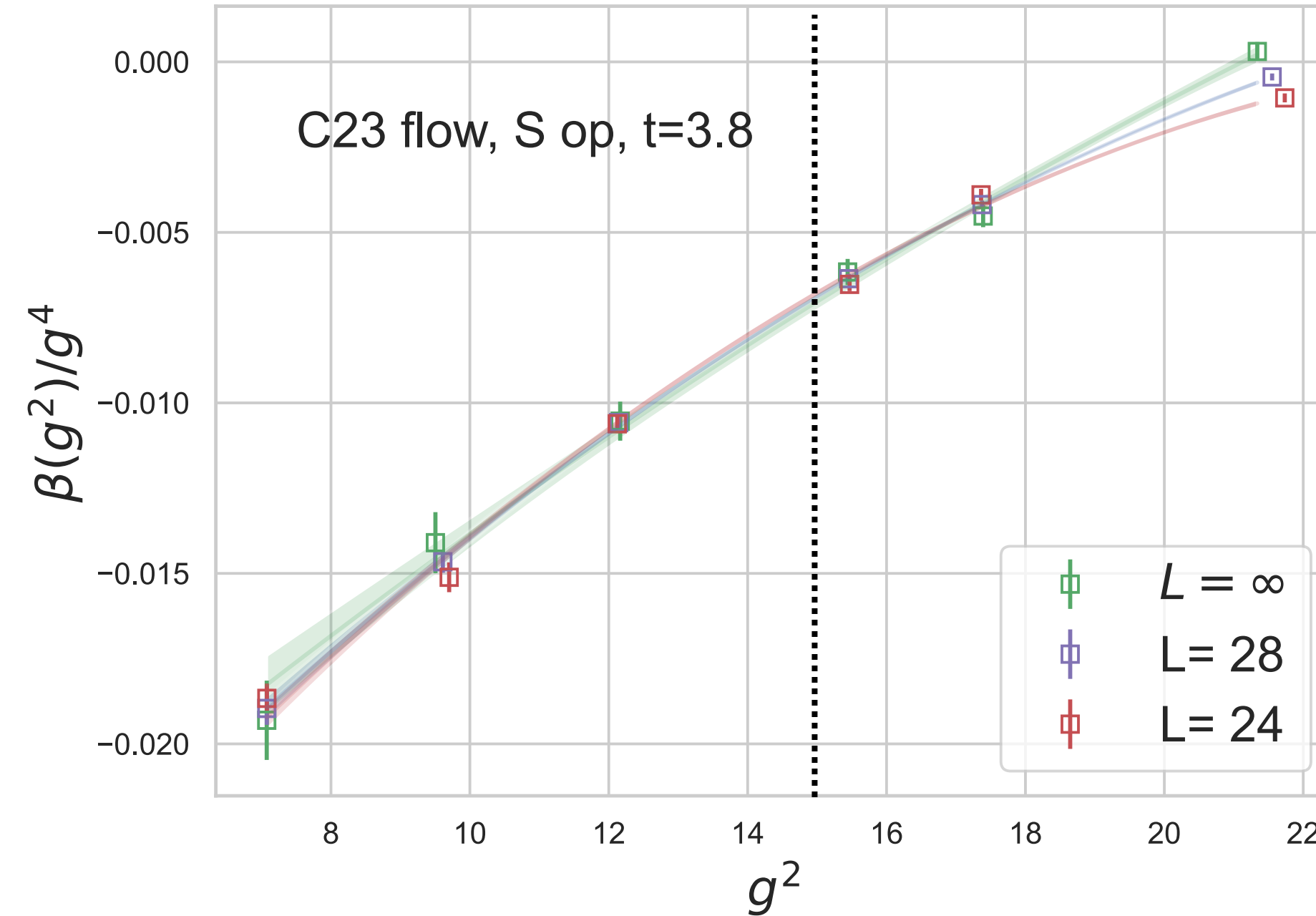
..... IRFP at $g^2 \approx 15$

(2) Interpolation

We need the value of $\beta(g^2)$ for every g^2 and t/a^2 : polynomial interpolation



Compare $L=24, 28$ and ∞

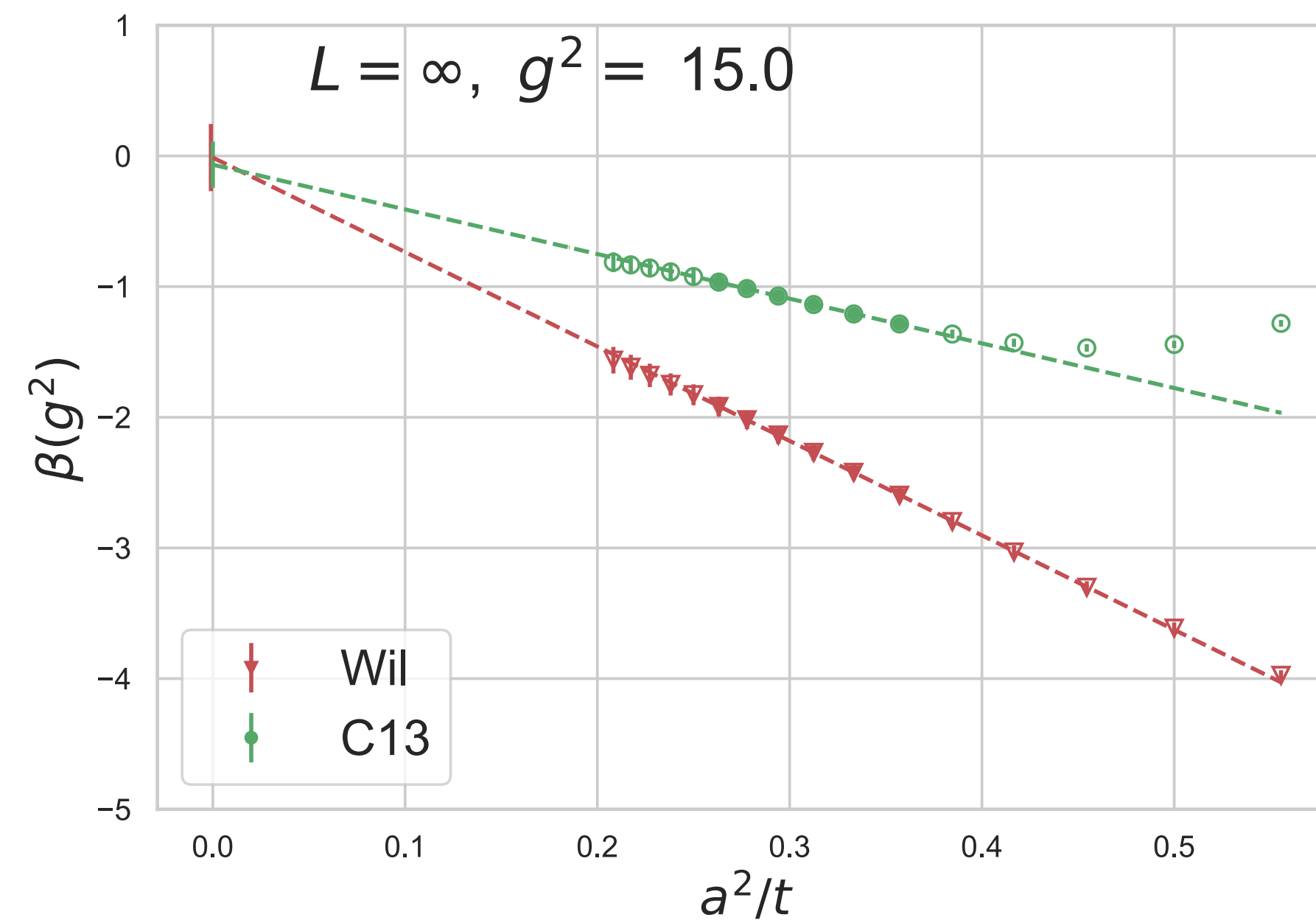
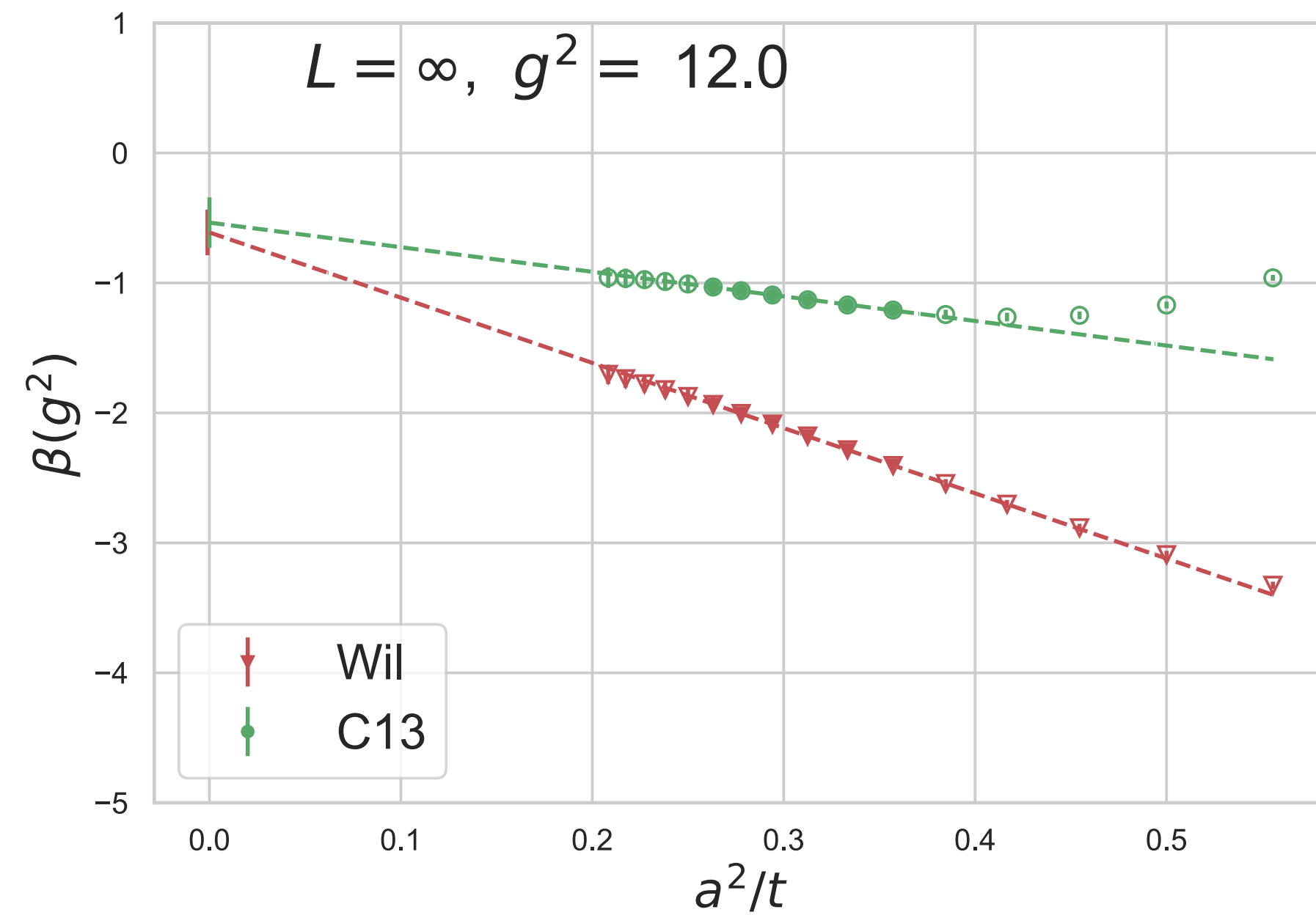


..... IRFP at $g^2 \approx 15$

(3) Continuum limit

Extrapolate linearly in $a^2/t \rightarrow 0$

$L = \infty$ continuum limit, S op

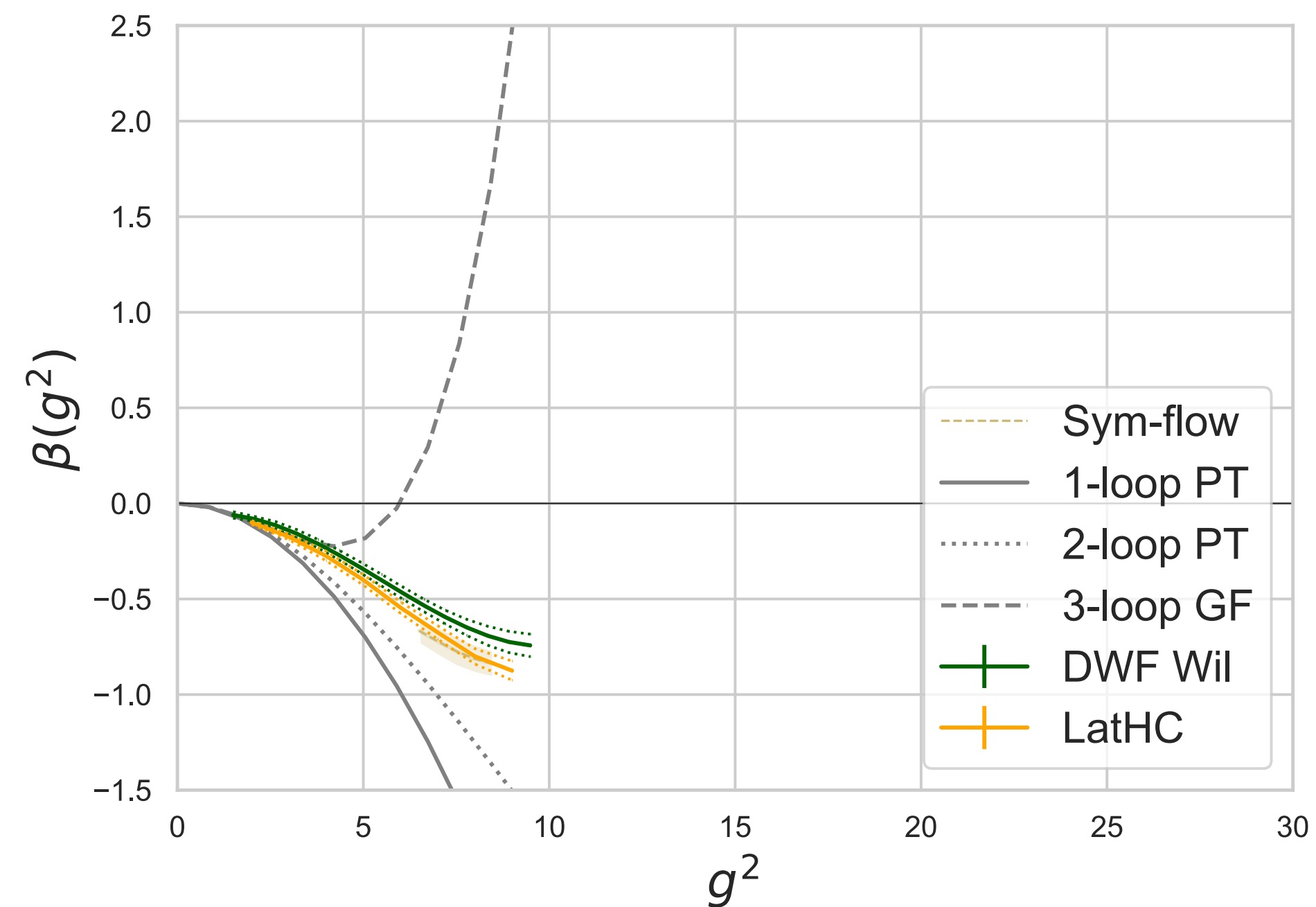


..... IRFP at $g^2 \approx 15$

Continuum limit

Require consistency between different operators :

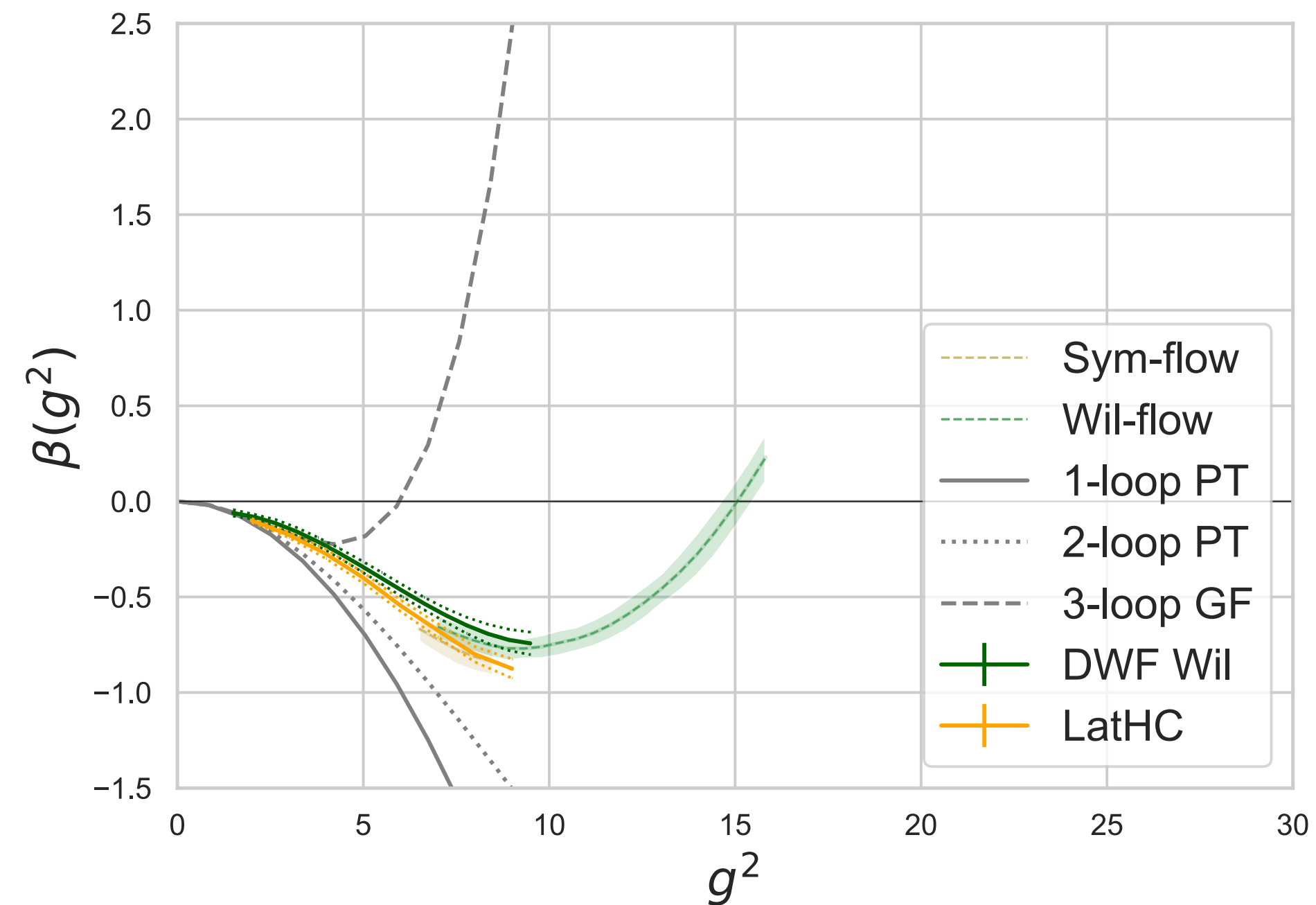
- W agrees with S within $1\ 1\sigma$
 - C and W agree with S within 1σ
- } Constrains g^2 range (details in extra slides)



Continuum limit

Require consistency between different operators :

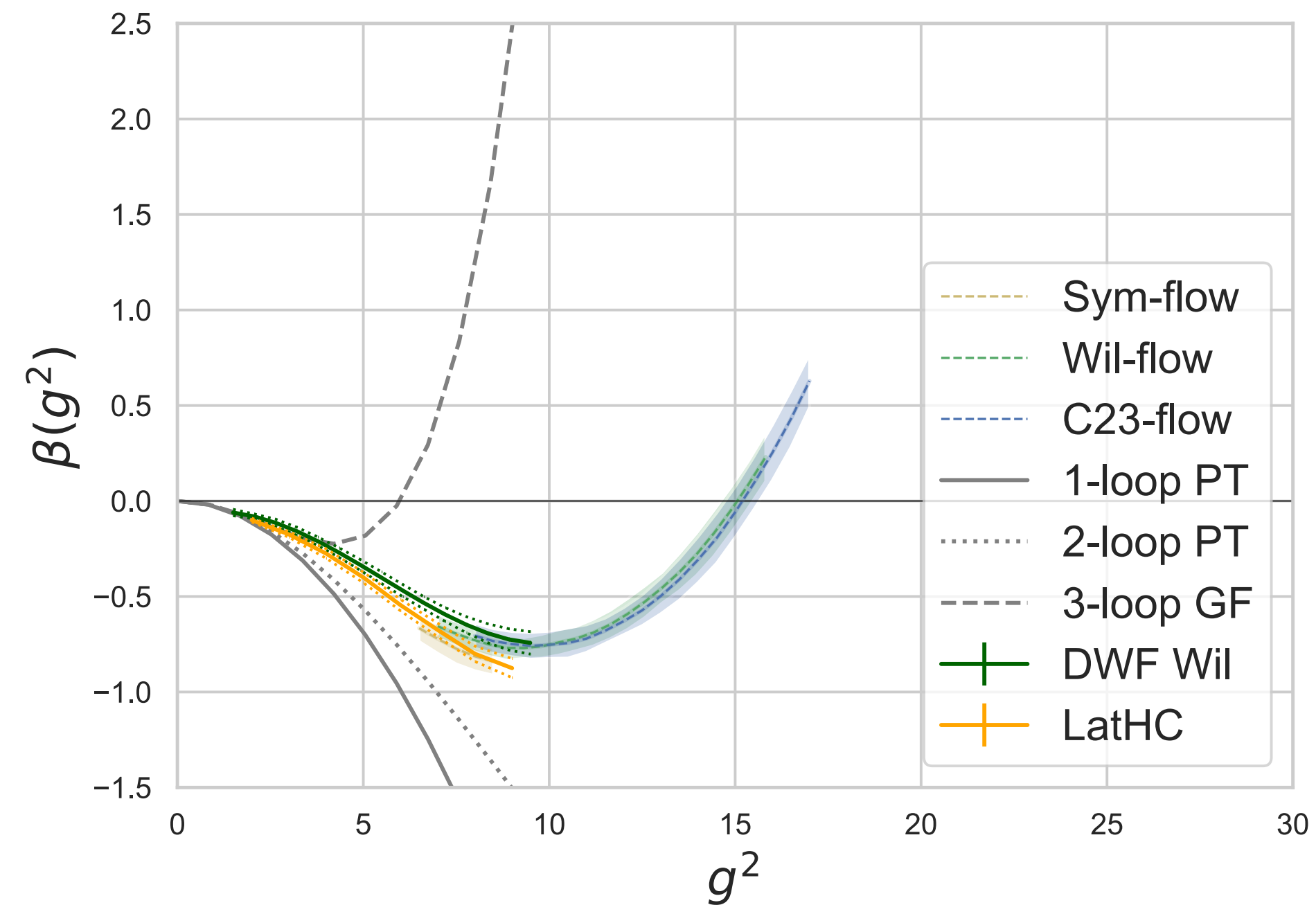
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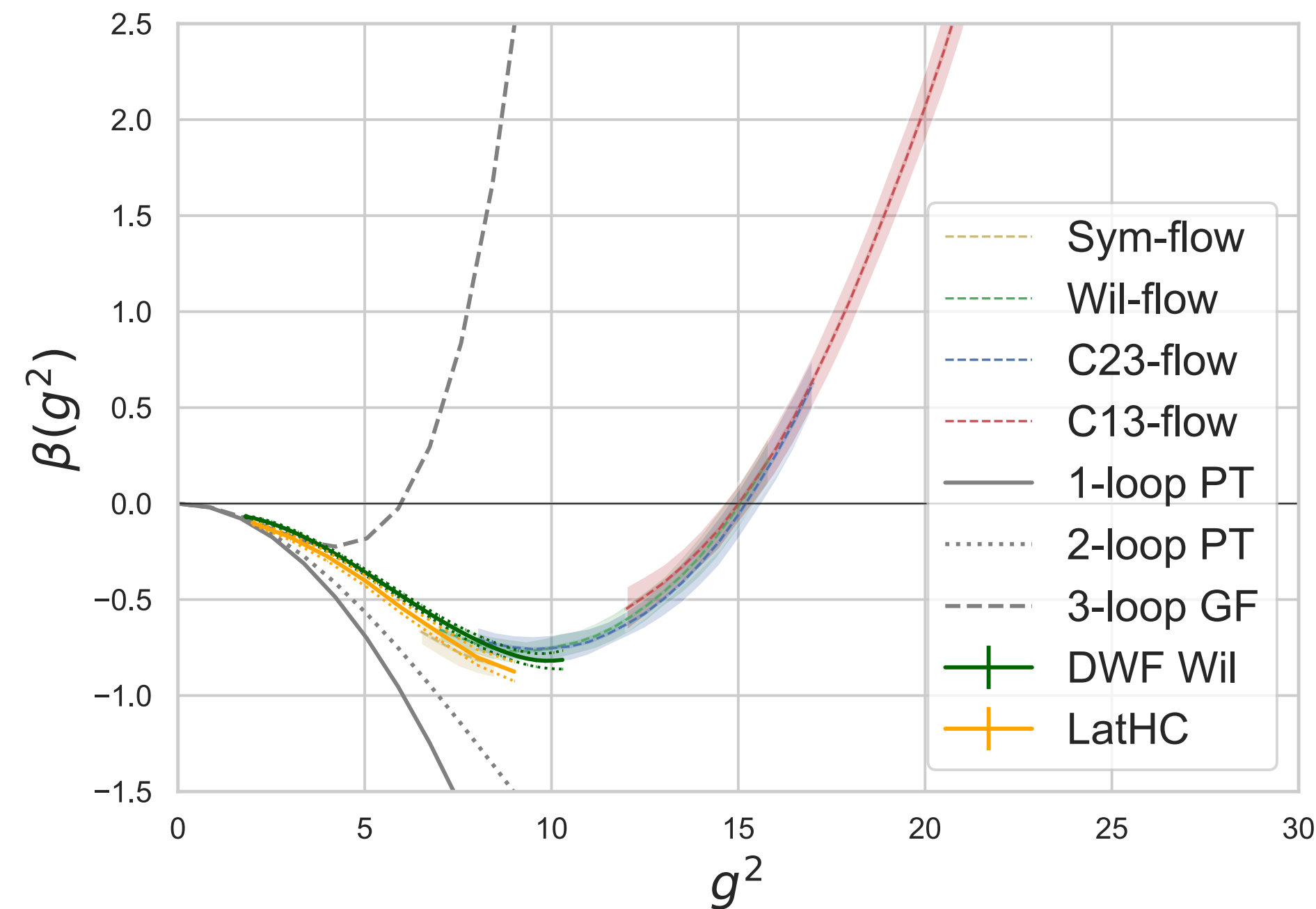


Continuum limit

Require consistency between different operators :

- W agrees with S within 1σ (bootstrap)
- Or: C and W agree with S within 1σ (bootstrap)

Constrains g^2 range



All flows are consistent
IRFP at $g^2 \simeq 15$

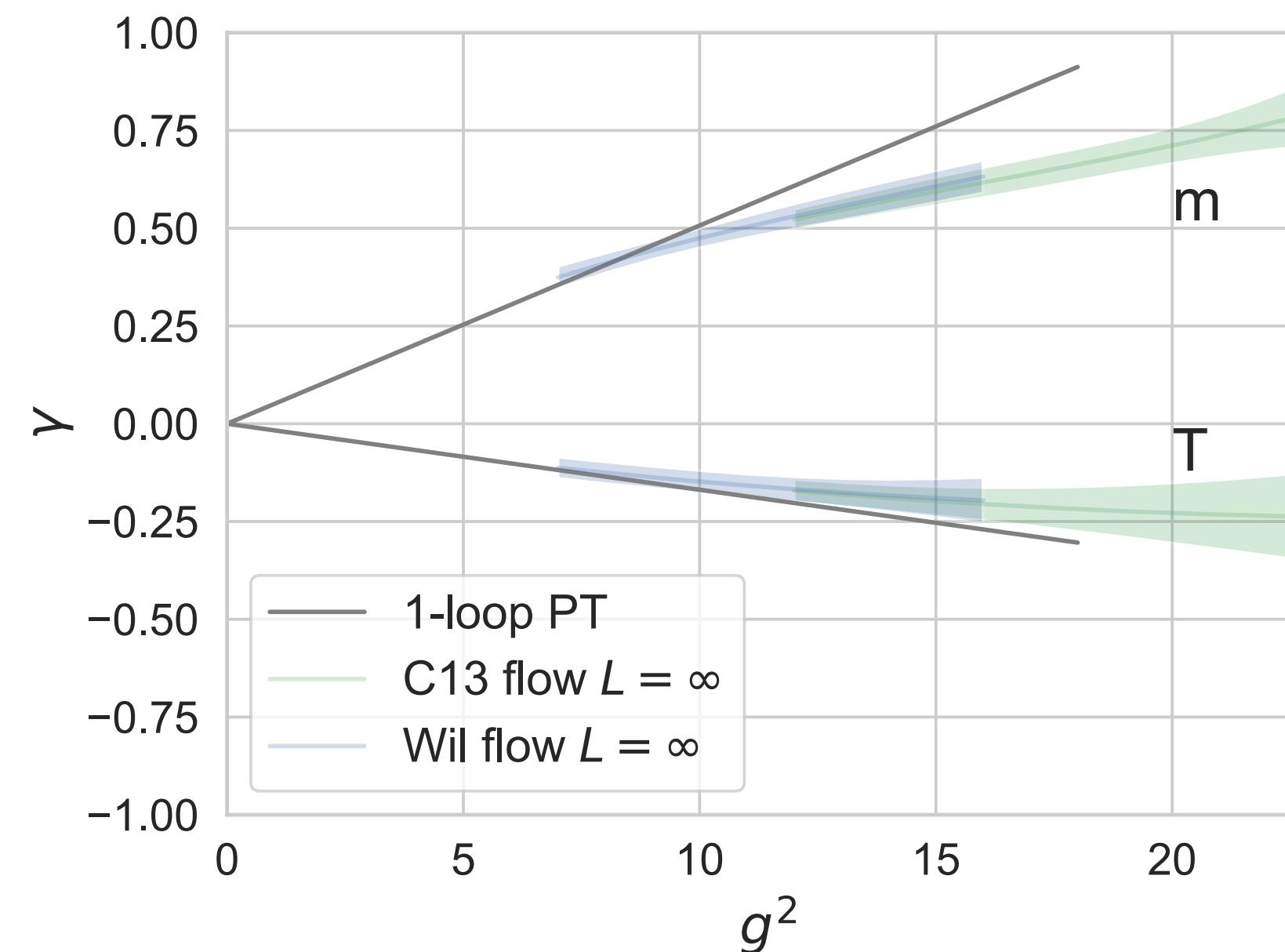
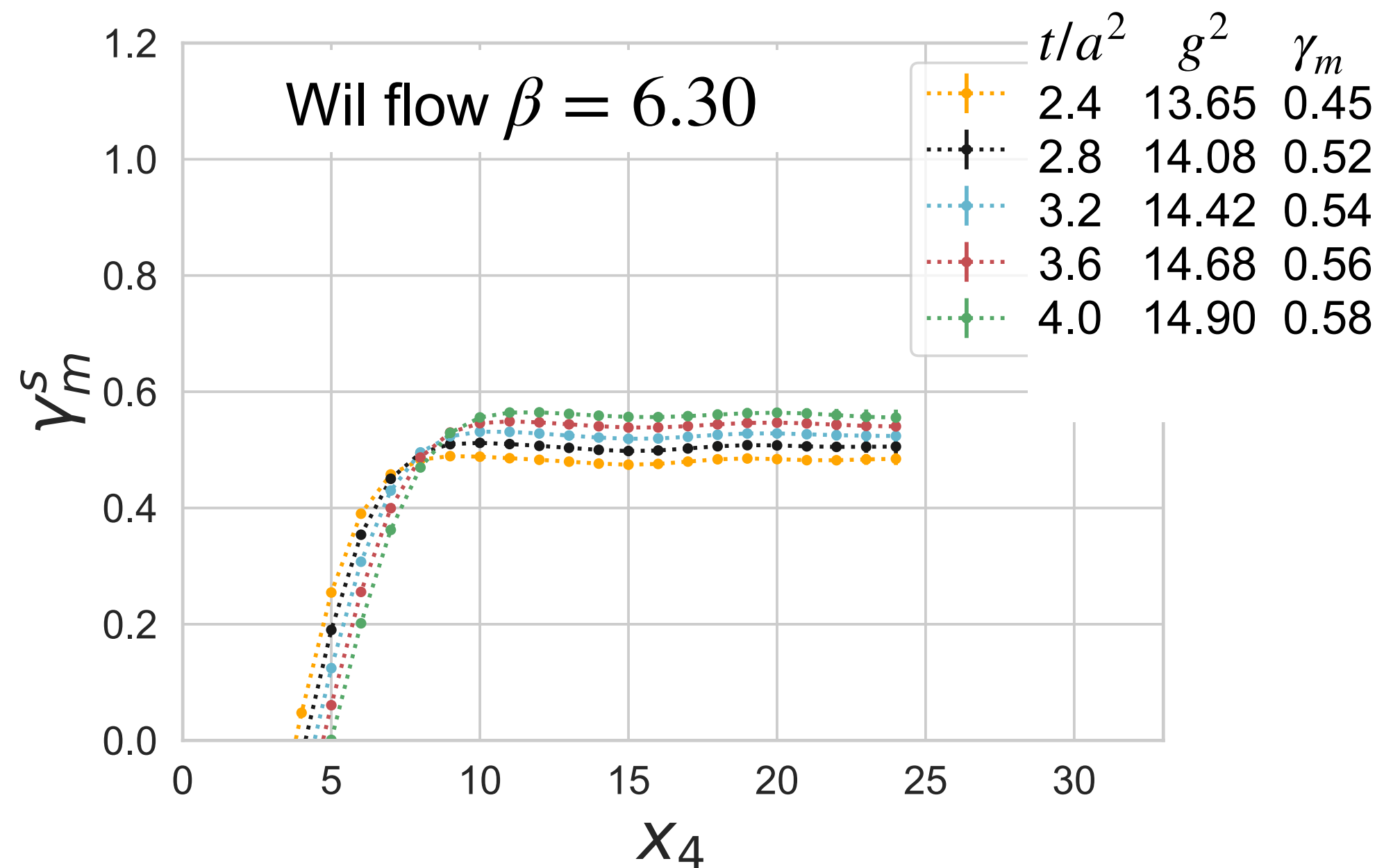
Anomalous dimension

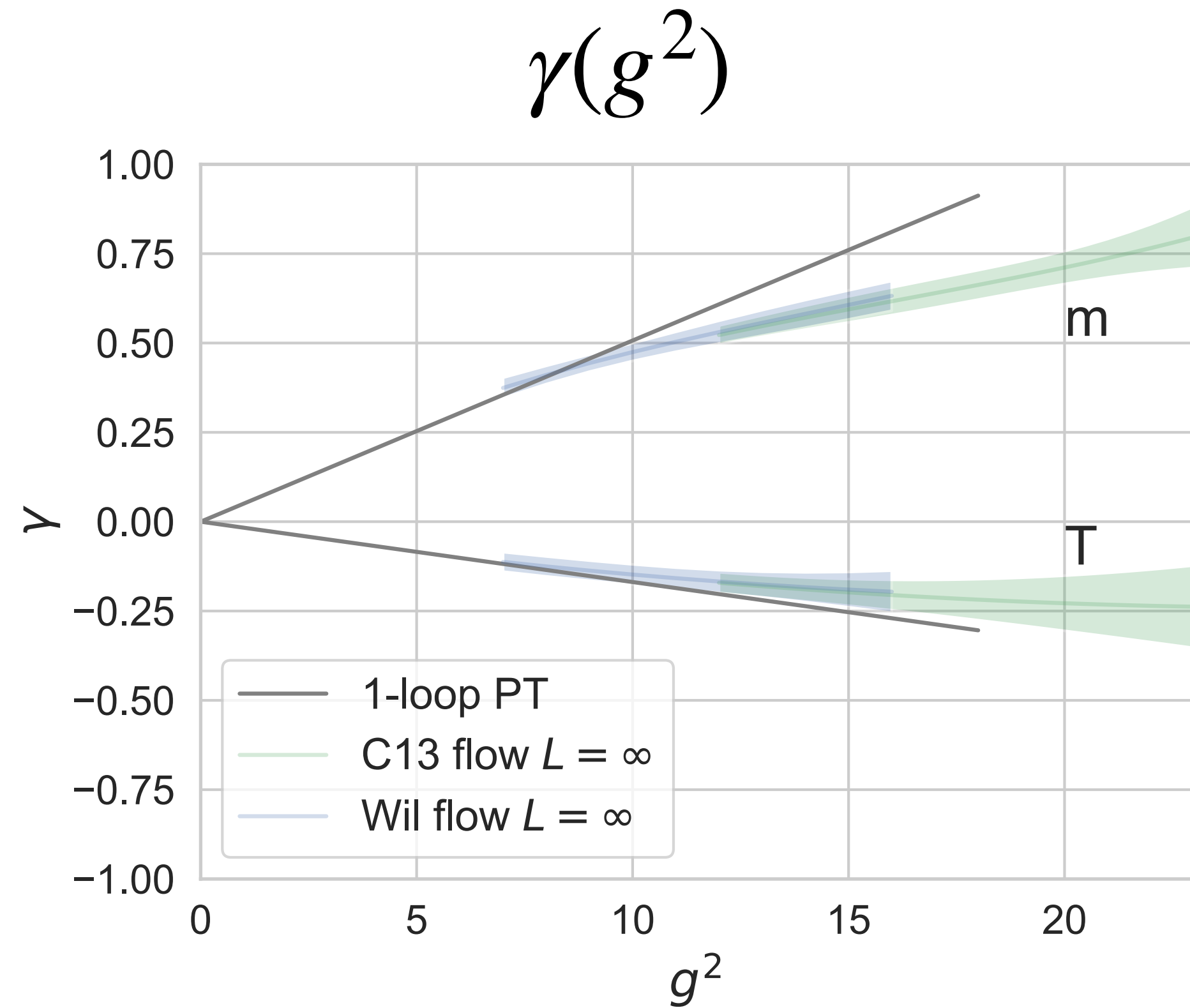
$$\mathcal{O} = \bar{\psi}(x)\Gamma\psi(x) \quad \text{or} \quad G_{\mathcal{O}}(x_4, t) = \langle \mathcal{O}(\bar{p} = 0, x_4; \mathbf{t}) \mathcal{O}(\bar{p} = 0, 0; \mathbf{t} = \mathbf{0}) \rangle$$

- remove η_{ψ} (Z_{ψ}) by dividing with the vector operator :

$$\mathcal{R}_{\mathcal{O}}(x_4, t) = \frac{G_{\mathcal{O}}(x_4, t)}{G_{\mathcal{V}}(x_4, t)} \quad \Rightarrow \quad \gamma_{\mathcal{O}}(a; t) = t \frac{d \log \mathcal{R}_{\mathcal{O}}(a; t)}{dt}, \quad \text{no dependence on } x_4 \gg \sqrt{8t}$$

- combine $\gamma_{\mathcal{O}}(t)$ with $g_{GF}^2(t)$ to obtain $\gamma_{\mathcal{O}}(g^2)$
- continuum limit : $a^2/t \rightarrow 0$ at fixed g^2

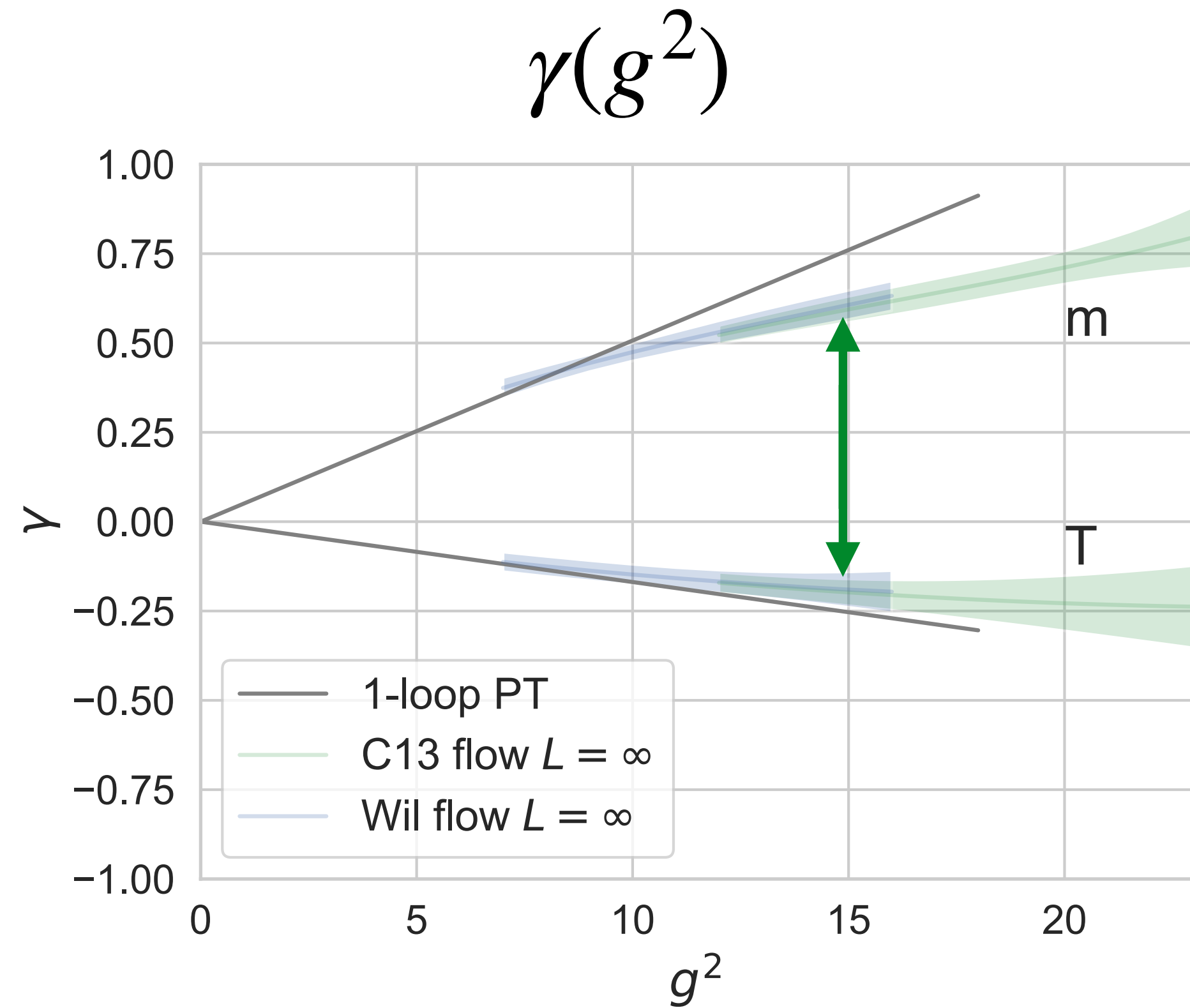




Anomalous dimension $\gamma_m^* \simeq 0.60$ (at $g_{IRFP}^2 = 15.0(5)$)
(not even close to the conformal sill)

γ_m is consistent with prior hyperscaling determination
at $g^2 \simeq 10$, $\gamma_m = 0.46$

LSD coll., *Phys.Rev.D* 103 (2021) 1, 014504



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LSD coll., *Phys.Rev.D* 103 (2021) 1, 014504

Summary:

With new methods we cover strongly coupled regime that *was not accessible before*

- Heavy PV bosons remove cutoff effects → stronger renormalized gauge couplings
- Use of different flows further extends the reach in g^2
- Different flows at same g^2 rely on different bare couplings \implies consistency check!
- Bootstrap analysis give reliable errors and consistency checks

Our results show that SU(3) with 10 flavors is IR conformal

- $g_*^2 = 15.0(5)$
- $\gamma_m^* \simeq 0.6$

EXTRA SLIDES

Simulation details

Plaquette action and Wilson fermions with:

- Clover term
- nHYP smeared links
- nHYP dislocation suppressing gauge term
- PV bosons : 3 per fermions, $\kappa_{PV} = 0.1$ ($am_{PV} \gtrsim 1.0$)
- well behaved effective gauge action

Different flows / RG transformation

Vary the action for gradient flow (improved RG):

$$S_{flow} = c_p S_p + c_r S_{rec} , \quad c_p + 8c_r = 1$$

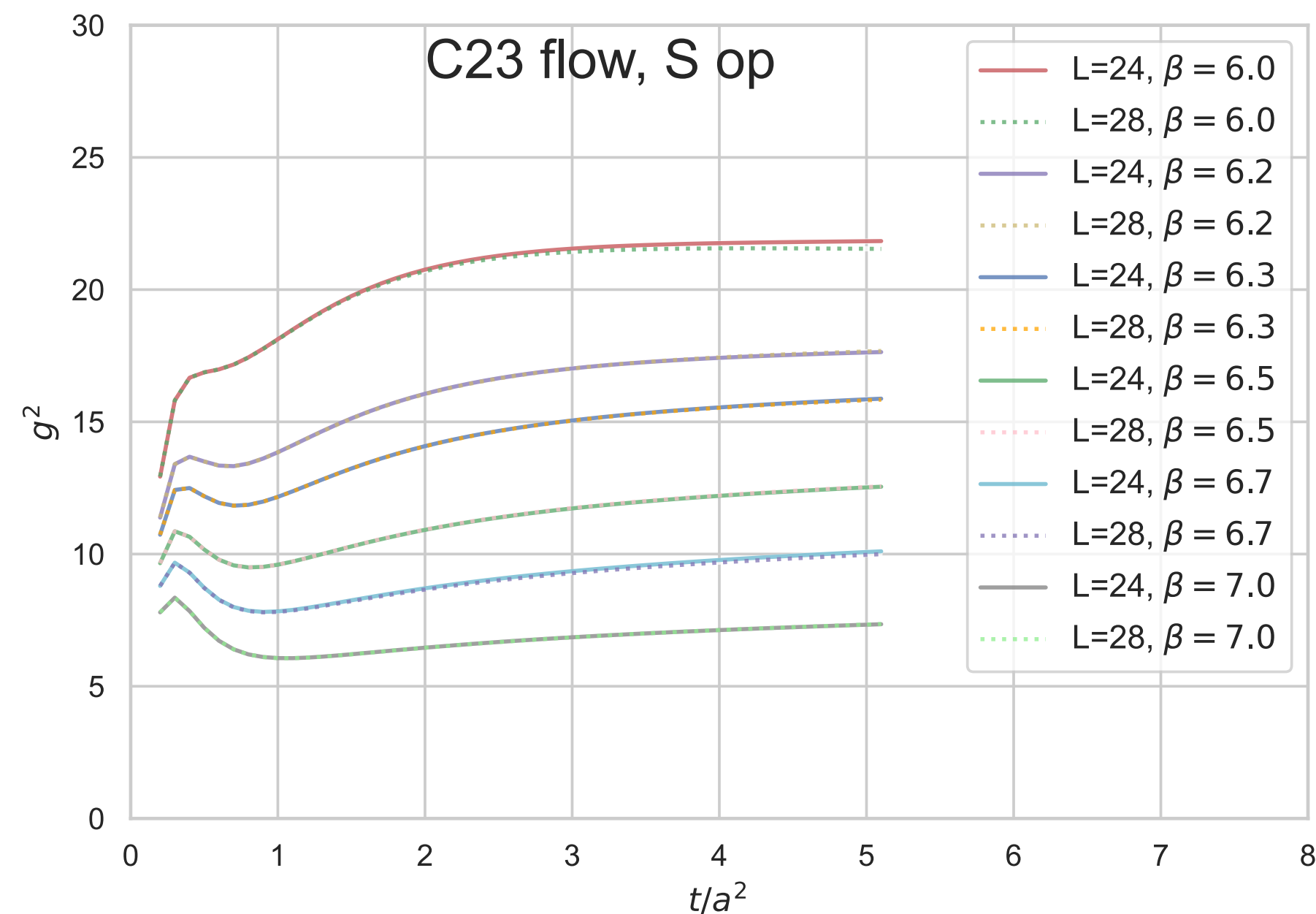
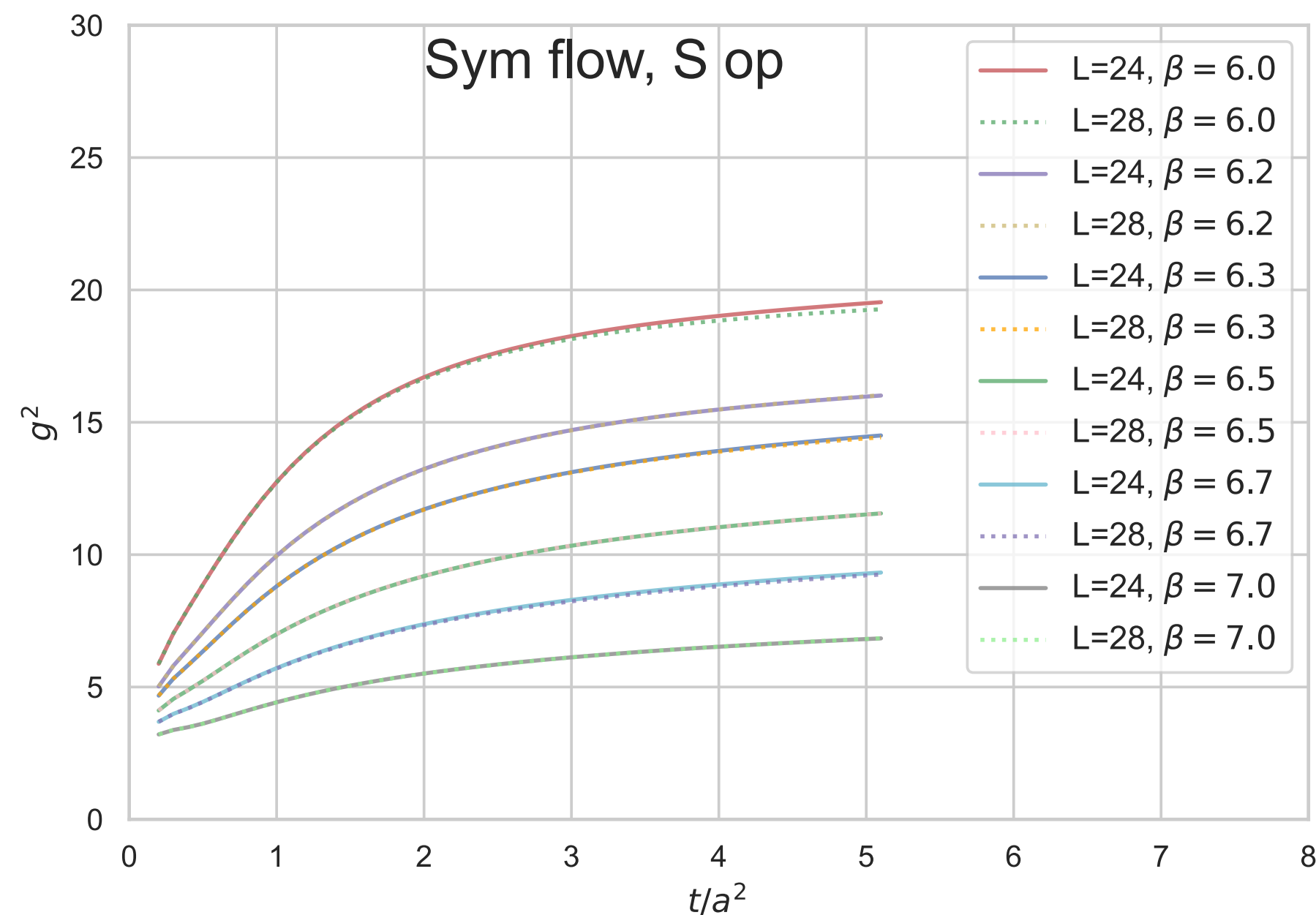
S : $c_p = 5/3$

W: $c_p = 1$

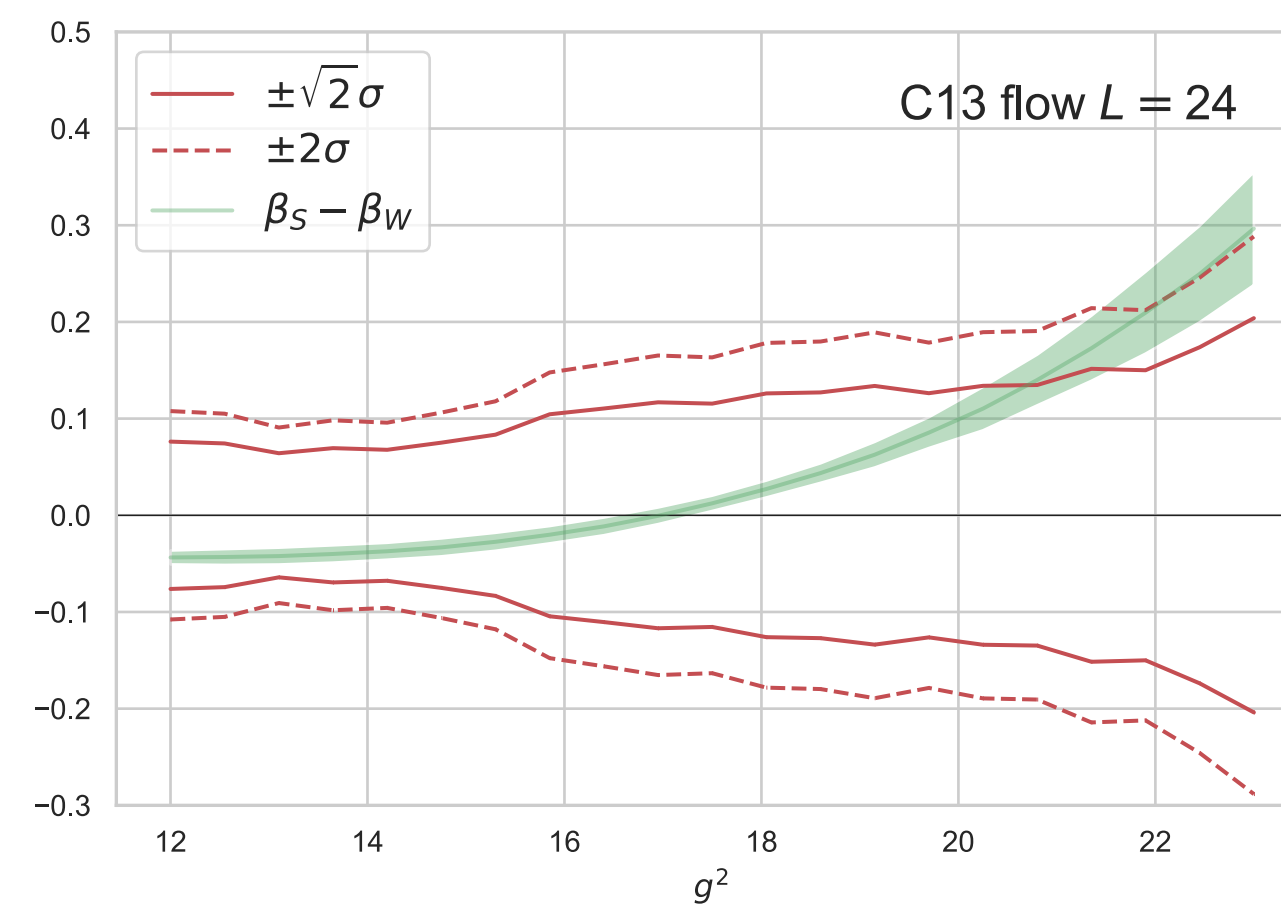
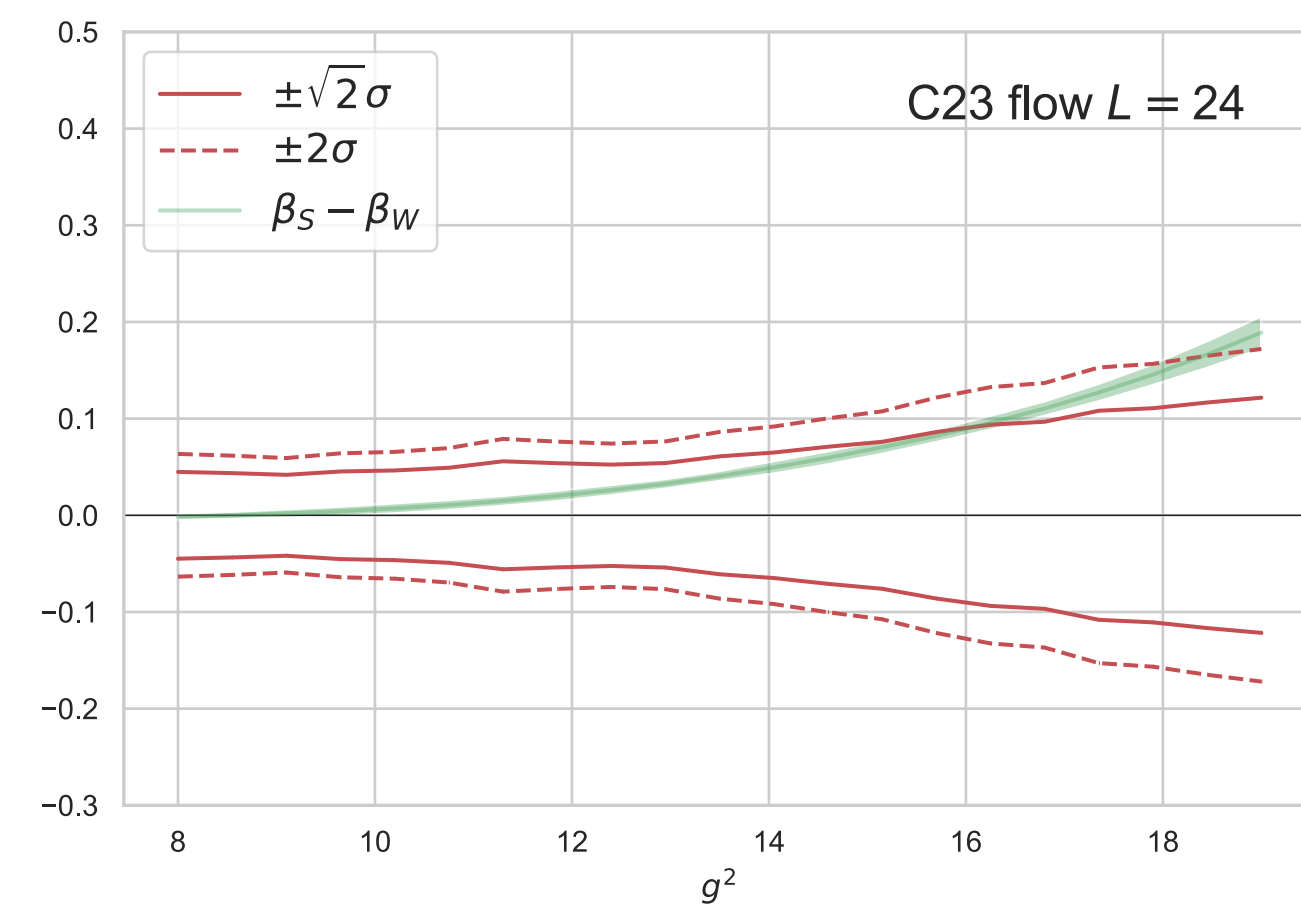
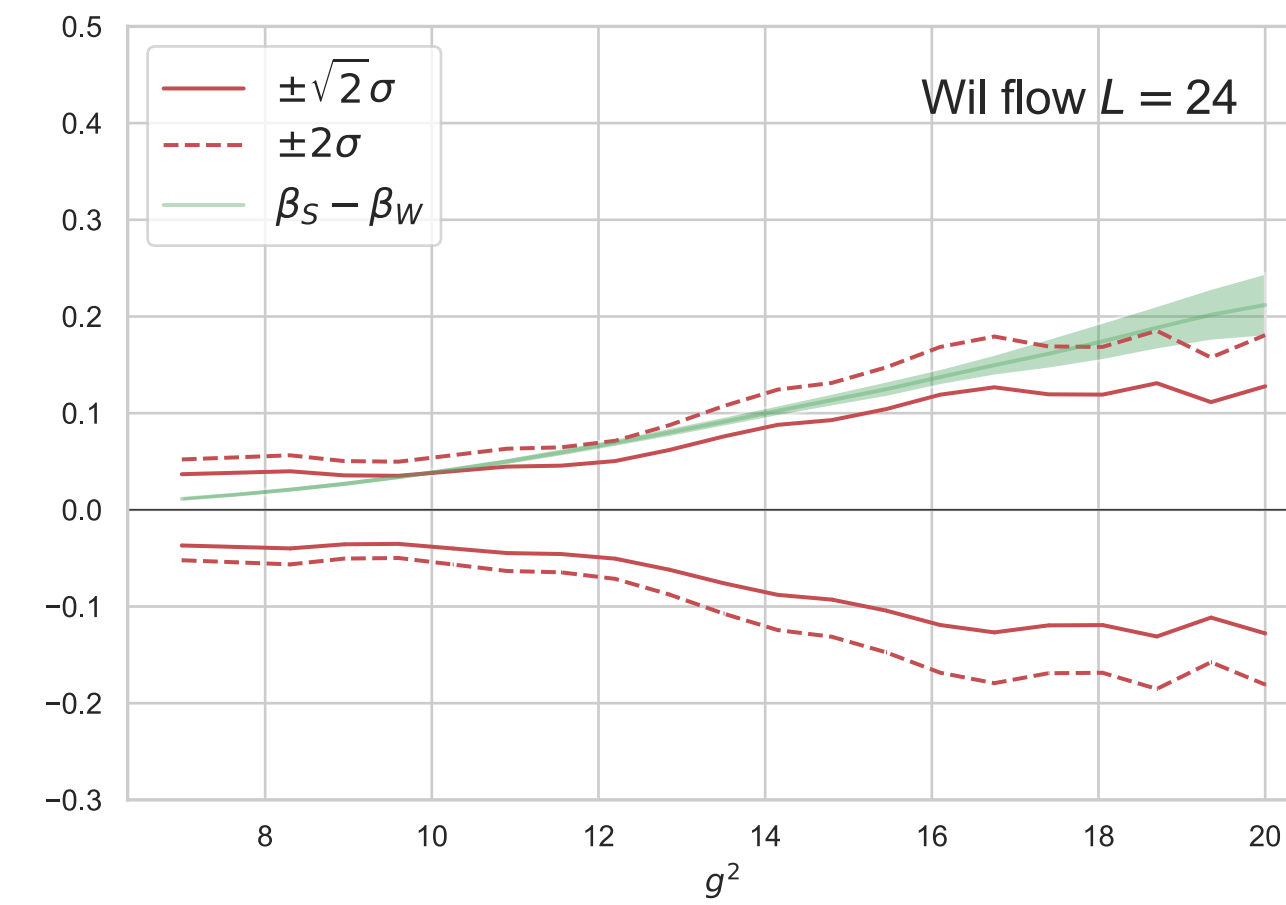
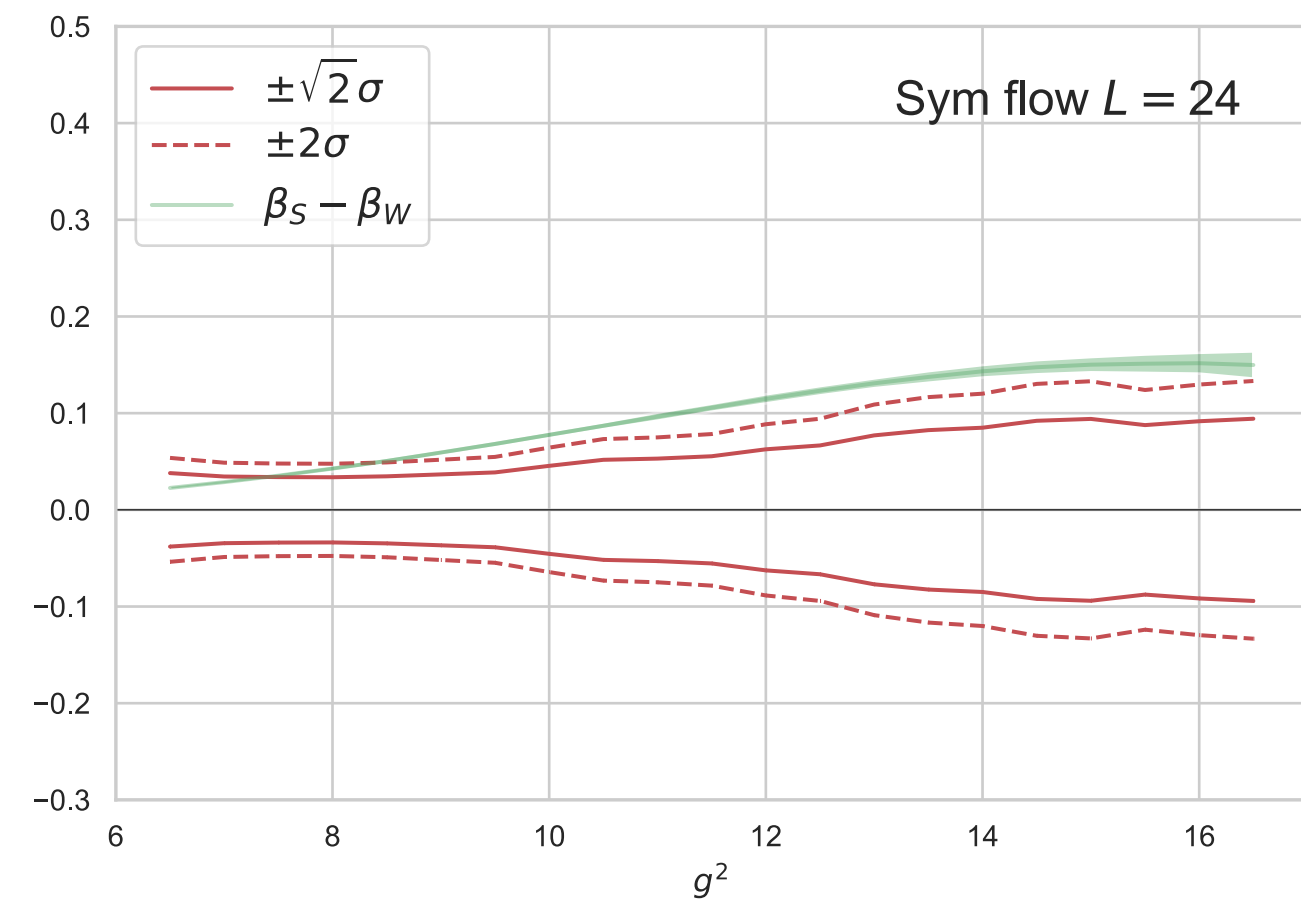
C23: $c_p = 2/3$

C13: $c_p = 1/3$

All flows should give the same continuum limit
 g^2 range is increasing as c_p decreases



Consistency of flows



Continuum limit, $L \rightarrow \infty$

