

# Baryonic screening masses in high temperature QCD

**Pietro Rescigno**

University of Milano - Bicocca, INFN Milano - Bicocca

In collaboration with:

**Leonardo Giusti, Tim Harris, Davide Laudicina, Michele Pepe**

The 40<sup>th</sup> International Symposium on Lattice Field Theory

Fermilab, Batavia, U.S.A.

August 1<sup>st</sup> 2023



## Summary

- The **baryonic screening spectrum** is analysed at 12 temperatures between  $\sim 1$  GeV and  $\sim 160$  GeV
  - Simulations of **Lattice QCD** with  $N_f = 3$  Wilson quarks in the **chiral limit**
  - First continuum results in this temperature range, with **sub percent** precision
-

## Summary

- The **baryonic screening spectrum** is analysed at 12 temperatures between  $\sim 1$  GeV and  $\sim 160$  GeV
  - Simulations of **Lattice QCD** with  $N_f = 3$  Wilson quarks in the **chiral limit**
  - First continuum results in this temperature range, with **sub percent** precision
- 
- Study properties of **Quark-Gluon Plasma** from first principles
  - Probe nature of QCD in the **high temperature regime**

- o **Nucleon** interpolating operator:  $N(x) = \varepsilon^{abc} \left[ u_a^T(x) C \gamma_5 d_b(x) \right] d_c(x)$

- **Nucleon** interpolating operator:  $N(x) = \varepsilon^{abc} \left[ u_a^T(x) C \gamma_5 d_b(x) \right] d_c(x)$
- Thermal theory  $\implies$  Periodicity in  $x_0 \implies$  Fermionic **Matsubara frequencies**  
 $\omega_n = (2n + 1)\pi T$     Lowest frequency:  $\omega_0 = \pi T$

- **Nucleon** interpolating operator:  $N(x) = \varepsilon^{abc} \left[ u_a^T(x) C \gamma_5 d_b(x) \right] d_c(x)$
- Thermal theory  $\implies$  Periodicity in  $x_0 \implies$  Fermionic **Matsubara frequencies**  
 $\omega_n = (2n + 1)\pi T$     Lowest frequency:  $\omega_0 = \pi T$
- Correlator projected to given frequency and **definite  $x_3$  - parity**:

$$C_{N^\pm}(x_3) = \sum_{x_0, x_1, x_2} e^{i\omega_0 x_0} \langle \text{Tr} \{ N(x) \bar{N}(0) P_\pm \} \rangle, \quad P_\pm = \frac{\mathbb{1} \pm \gamma_3}{2}$$

- **Nucleon** interpolating operator:  $N(x) = \varepsilon^{abc} \left[ u_a^T(x) C \gamma_5 d_b(x) \right] d_c(x)$
- Thermal theory  $\implies$  Periodicity in  $x_0 \implies$  Fermionic **Matsubara frequencies**  
 $\omega_n = (2n + 1)\pi T$     Lowest frequency:  $\omega_0 = \pi T$
- Correlator projected to given frequency and **definite  $x_3$  - parity**:

$$C_{N^\pm}(x_3) = \sum_{x_0, x_1, x_2} e^{i\omega_0 x_0} \langle \text{Tr} \{ N(x) \bar{N}(0) P_\pm \} \rangle, \quad P_\pm = \frac{\mathbb{1} \pm \gamma_3}{2}$$

- **Screening mass** characterizes exponential fall-off (inverse correlation length)

$$C_{N^\pm}(x_3) \underset{x_3 \rightarrow \infty}{\sim} \exp\{-m_{N^\pm} x_3\} \left(1 + O\left(e^{-\Delta m_{N^\pm} x_3}\right)\right)$$

$$m_{N^\pm} = \lim_{x_3 \rightarrow \infty} \underbrace{-\frac{d}{dx_3} \ln [C_{N^\pm}(x_3)]}_{m_{N^\pm}^{\text{eff}}(x_3)} \quad \text{Free theory: } m_{N^\pm} = 3\pi T$$

- **Nucleon** interpolating operator:  $N(x) = \varepsilon^{abc} \left[ u_a^T(x) C \gamma_5 d_b(x) \right] d_c(x)$
- Thermal theory  $\implies$  Periodicity in  $x_0 \implies$  Fermionic **Matsubara frequencies**  
 $\omega_n = (2n + 1)\pi T$     Lowest frequency:  $\omega_0 = \pi T$

- Correlator projected to given frequency and **definite  $x_3$  - parity**:

$$C_{N^\pm}(x_3) = \sum_{x_0, x_1, x_2} e^{i\omega_0 x_0} \langle \text{Tr} \{ N(x) \bar{N}(0) P_\pm \} \rangle, \quad P_\pm = \frac{\mathbb{1} \pm \gamma_3}{2}$$

- **Screening mass** characterizes exponential fall-off (inverse correlation length)

$$C_{N^\pm}(x_3) \underset{x_3 \rightarrow \infty}{\sim} \exp\{-m_{N^\pm} x_3\} \left(1 + O\left(e^{-\Delta m_{N^\pm} x_3}\right)\right)$$

$$m_{N^\pm} = \lim_{x_3 \rightarrow \infty} \underbrace{-\frac{d}{dx_3} \ln [C_{N^\pm}(x_3)]}_{m_{N^\pm}^{\text{eff}}(x_3)} \quad \text{Free theory: } m_{N^\pm} = 3\pi T$$

- Opposite parity channels related by **chiral symmetry**:  $C_{N^+}(x_3) \xrightarrow[\text{Continuum}]{SU(2)_A \text{ W.I.}} -C_{N^-}(x_3)$

**Degeneracy of screening masses when chiral symmetry is restored**



$T$	$T[\text{GeV}]$	$L_0/a$	$n_{\text{mdu}}$	$n_{\text{skip}}$
$T_0$	164.6(5.6)	4	300	10
		6	390	10
$T_1$	82.3(2.8)	4	300	10
		6	310	10
		8	500	10
		10	500	10
$T_2$	51.4(1.7)	4	300	10
		6	320	10
		8	490	10
		10	500	10
$T_3$	32.8(1.0)	4	300	10
		6	340	10
		8	490	10
		10	500	10
$T_4$	20.63(63)	4	440	10
		6	310	10
		8	490	10
		10	500	10
$T_5$	12.77(37)	4	310	10
		6	310	10
		8	500	10
		10	500	10
$T_6$	8.03(22)	4	300	10
		6	320	10
		8	500	10
		10	500	10
$T_7$	4.91(13)	4	320	10
		6	310	10
		8	500	10
		10	500	10

$T$	$T[\text{GeV}]$	$L_0/a$	$n_{\text{mdu}}$	$n_{\text{skip}}$
$T_8$	3.040(78)	4	320	10
		6	300	10
		8	500	10
		10	500	10
$T_9$	2.833(68)	4	400	10
		6	390	10
		8	390	10
$T_{10}$	1.821(39)	4	410	10
		6	400	10
		8	400	10
$T_{11}$	1.167(23)	4	400	10
		6	390	10
		8	400	10

- 12 temperature values  $T_0, \dots, T_{11}$  between **1.167 GeV** and **164.6 GeV**
- 4 lattice spacings  $L_0/a = 4, 6, 8, 10$ ,  
 $L_i/a = 288 \Rightarrow TL \sim 20 - 50$   
Finite volume effects suppressed  
**exponentially** in  $g^2 TL$
- $N_f = 3$  flavors of  $O(a)$  improved massless Wilson fermions

$T$	$T[\text{GeV}]$	$L_0/a$	$n_{\text{mdu}}$	$n_{\text{skip}}$
$T_0$	164.6(5.6)	4	300	10
		6	390	10
$T_1$	82.3(2.8)	4	300	10
		6	310	10
		8	500	10
		10	500	10
$T_2$	51.4(1.7)	4	300	10
		6	320	10
		8	490	10
		10	500	10
$T_3$	32.8(1.0)	4	300	10
		6	340	10
		8	490	10
		10	500	10
$T_4$	20.63(63)	4	440	10
		6	310	10
		8	490	10
		10	500	10
$T_5$	12.77(37)	4	310	10
		6	310	10
		8	500	10
		10	500	10
$T_6$	8.03(22)	4	300	10
		6	320	10
		8	500	10
		10	500	10
$T_7$	4.91(13)	4	320	10
		6	310	10
		8	500	10
		10	500	10

$T$	$T[\text{GeV}]$	$L_0/a$	$n_{\text{mdu}}$	$n_{\text{skip}}$
$T_8$	3.040(78)	4	320	10
		6	300	10
		8	500	10
		10	500	10
$T_9$	2.833(68)	4	400	10
		6	390	10
		8	390	10
$T_{10}$	1.821(39)	4	410	10
		6	400	10
		8	400	10
$T_{11}$	1.167(23)	4	400	10
		6	390	10
		8	400	10

o Shifted boundary conditions [Giusti, Meyer 2011-13]

$$U_\mu(x_0 + L_0, \mathbf{x}) = U_\mu(x_0, \mathbf{x} - L_0 \hat{\xi})$$

$$\psi(x_0 + L_0, \mathbf{x}) = -\psi(x_0, \mathbf{x} - L_0 \hat{\xi})$$

$T$	$T[\text{GeV}]$	$L_0/a$	$n_{\text{mdu}}$	$n_{\text{skip}}$
$T_0$	164.6(5.6)	4	300	10
		6	390	10
$T_1$	82.3(2.8)	4	300	10
		6	310	10
		8	500	10
		10	500	10
$T_2$	51.4(1.7)	4	300	10
		6	320	10
		8	490	10
		10	500	10
$T_3$	32.8(1.0)	4	300	10
		6	340	10
		8	490	10
		10	500	10
$T_4$	20.63(63)	4	440	10
		6	310	10
		8	490	10
		10	500	10
$T_5$	12.77(37)	4	310	10
		6	310	10
		8	500	10
		10	500	10
$T_6$	8.03(22)	4	300	10
		6	320	10
		8	500	10
		10	500	10
$T_7$	4.91(13)	4	320	10
		6	310	10
		8	500	10
		10	500	10

$T$	$T[\text{GeV}]$	$L_0/a$	$n_{\text{mdu}}$	$n_{\text{skip}}$
$T_8$	3.040(78)	4	320	10
		6	300	10
		8	500	10
		10	500	10
$T_9$	2.833(68)	4	400	10
		6	390	10
		8	390	10
$T_{10}$	1.821(39)	4	410	10
		6	400	10
		8	400	10
$T_{11}$	1.167(23)	4	400	10
		6	390	10
		8	400	10

- o Shifted boundary conditions [Giusti, Meyer 2011-13]

$$U_\mu(x_0 + L_0, \mathbf{x}) = U_\mu(x_0, \mathbf{x} - L_0 \hat{\xi})$$

$$\psi(x_0 + L_0, \mathbf{x}) = -\psi(x_0, \mathbf{x} - L_0 \hat{\xi})$$

- o Periodicity in a tilted system of length

$$L_0 \sqrt{1 + \xi^2} \implies T = 1 / \left( L_0 \sqrt{1 + \xi^2} \right)$$

$$\xi = (1, 0, 0) \implies T = 1 / \left( \sqrt{2} L_0 \right)$$

$$\text{Phase factor: } \omega_0 x_0 \rightarrow \omega_0 (x_0 + x_1) / 2$$

$T$	$T[\text{GeV}]$	$L_0/a$	$n_{\text{mdu}}$	$n_{\text{skip}}$
$T_0$	164.6(5.6)	4	300	10
		6	390	10
$T_1$	82.3(2.8)	4	300	10
		6	310	10
		8	500	10
		10	500	10
$T_2$	51.4(1.7)	4	300	10
		6	320	10
		8	490	10
		10	500	10
$T_3$	32.8(1.0)	4	300	10
		6	340	10
		8	490	10
		10	500	10
$T_4$	20.63(63)	4	440	10
		6	310	10
		8	490	10
		10	500	10
$T_5$	12.77(37)	4	310	10
		6	310	10
		8	500	10
		10	500	10
$T_6$	8.03(22)	4	300	10
		6	320	10
		8	500	10
		10	500	10
$T_7$	4.91(13)	4	320	10
		6	310	10
		8	500	10
		10	500	10

$T$	$T[\text{GeV}]$	$L_0/a$	$n_{\text{mdu}}$	$n_{\text{skip}}$
$T_8$	3.040(78)	4	320	10
		6	300	10
		8	500	10
		10	500	10
$T_9$	2.833(68)	4	400	10
		6	390	10
		8	390	10
$T_{10}$	1.821(39)	4	410	10
		6	400	10
		8	400	10
$T_{11}$	1.167(23)	4	400	10
		6	390	10
		8	400	10

- Shifted boundary conditions [Giusti, Meyer 2011-13]

$$U_\mu(x_0 + L_0, \mathbf{x}) = U_\mu(x_0, \mathbf{x} - L_0 \hat{\xi})$$

$$\psi(x_0 + L_0, \mathbf{x}) = -\psi(x_0, \mathbf{x} - L_0 \hat{\xi})$$

- Periodicity in a tilted system of length

$$L_0 \sqrt{1 + \xi^2} \implies T = 1 / \left( L_0 \sqrt{1 + \xi^2} \right)$$

$$\xi = (1, 0, 0) \implies T = 1 / \left( \sqrt{2} L_0 \right)$$

Phase factor:  $\omega_0 x_0 \rightarrow \omega_0 (x_0 + x_1) / 2$

- Mesonic case [Dalla Brida et al. 2021]

## Renormalization & scale setting

- Hadronic scheme:  $M_{\text{Had}} = M_{\text{Had}}^{\text{Phys}}$  would require  $a \ll \frac{1}{T} \ll \frac{1}{M_{\text{Had}}} \ll L$

## Renormalization & scale setting

- **Hadronic scheme:**  $M_{\text{Had}} = M_{\text{Had}}^{\text{Phys}}$  would require  $a \ll \frac{1}{T} \ll \frac{1}{M_{\text{Had}}} \ll L$
- **Instead:** fix  $\bar{g}^2(g_0, a\mu) = \bar{g}^2(\mu)$  with  $a\mu \ll 1$ ,  $\mu \sim T$

## Renormalization & scale setting

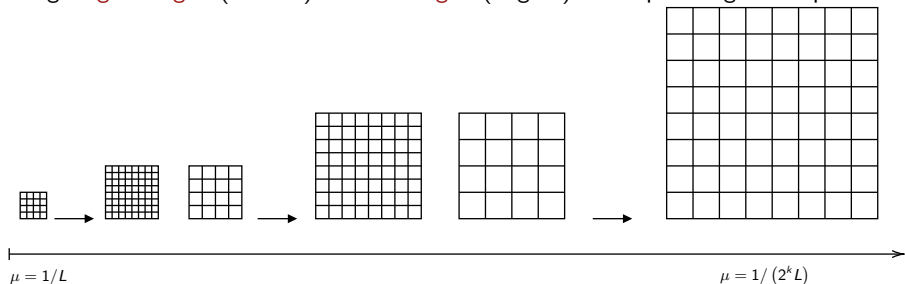
○ **Hadronic scheme:**  $M_{\text{Had}} = M_{\text{Had}}^{\text{Phys}}$  would require  $a \ll \frac{1}{T} \ll \frac{1}{M_{\text{Had}}} \ll L$

○ Instead: fix  $\bar{g}^2(g_0, a\mu) = \bar{g}^2(\mu)$  with  $a\mu \ll 1, \mu \sim T$

○ **Non-perturbative**  $\bar{g}^2(\mu)$   $\left\{ \begin{array}{l} T_0 - T_8: \text{ Schrödinger functional, } T = \frac{1}{\sqrt{2}L_0} = \frac{1}{\sqrt{2}L_{\text{SF}}} = \frac{\mu_{\text{SF}}}{\sqrt{2}} \\ T_9 - T_{11}: \text{ Gradient flow, } T = \frac{1}{\sqrt{2}L_0} = \frac{2}{\sqrt{2}L_{\text{GF}}} = \sqrt{2}\mu_{\text{GF}} \end{array} \right.$

[ALPHA collaboration, '93,2016-18]

Bridge **high energies** (small L) to **low energies** (large L) via step scaling techniques



## Renormalization & scale setting

○ **Hadronic scheme:**  $M_{\text{Had}} = M_{\text{Had}}^{\text{Phys}}$  would require  $a \ll \frac{1}{T} \ll \frac{1}{M_{\text{Had}}} \ll L$

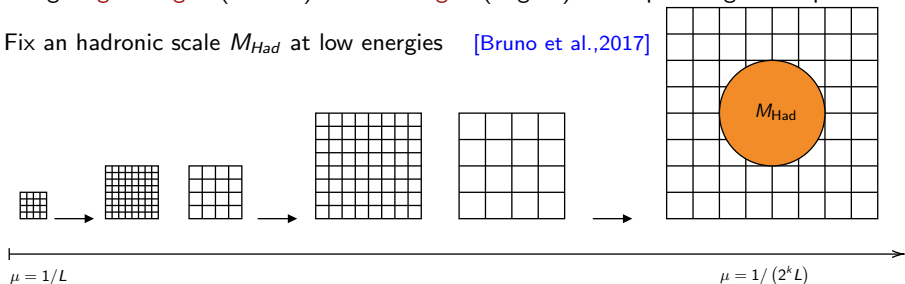
○ Instead: fix  $\bar{g}^2(g_0, a\mu) = \bar{g}^2(\mu)$  with  $a\mu \ll 1, \mu \sim T$

○ **Non-perturbative**  $\bar{g}^2(\mu)$   $\left\{ \begin{array}{l} T_0 - T_8: \text{ Schrödinger functional, } T = \frac{1}{\sqrt{2}L_0} = \frac{1}{\sqrt{2}L_{\text{SF}}} = \frac{\mu_{\text{SF}}}{\sqrt{2}} \\ T_9 - T_{11}: \text{ Gradient flow, } T = \frac{1}{\sqrt{2}L_0} = \frac{2}{\sqrt{2}L_{\text{GF}}} = \sqrt{2}\mu_{\text{GF}} \end{array} \right.$

[ALPHA collaboration, '93,2016-18]

Bridge **high energies** (small L) to **low energies** (large L) via step scaling techniques

○ Fix an hadronic scale  $M_{\text{Had}}$  at low energies [Bruno et al.,2017]





## Correlator and Noise-to-Signal ratio

- Correlator Noise-to-Signal ratio grows **exponentially** with distance **at zero temperature**  
[Parisi 1984, Lepage 1989]

## Correlator and Noise-to-Signal ratio

- Correlator Noise-to-Signal ratio grows **exponentially** with distance **at zero temperature**  
[Parisi 1984, Lepage 1989]
- In the **infinite temperature limit** there is **no exponential problem** for  $C_{N^\pm}(x_3)$

## Correlator and Noise-to-Signal ratio

- Correlator Noise-to-Signal ratio grows **exponentially** with distance **at zero temperature**  
[Parisi 1984, Lepage 1989]
- In the **infinite temperature limit** there is **no exponential problem** for  $C_{N\pm}(x_3)$
- At **high but finite** temperature **milder signal loss** is expected

$$\frac{\langle |C_{N\pm}(x_3)|^2 \rangle - \langle C_{N\pm}(x_3) \rangle^2}{\langle C_{N\pm}(x_3) \rangle^2} \sim \exp\{(2m_{N\pm} - 3m_{S/P})x_3\}$$

$\downarrow$   
 $3\pi T$

$T \rightarrow \infty$

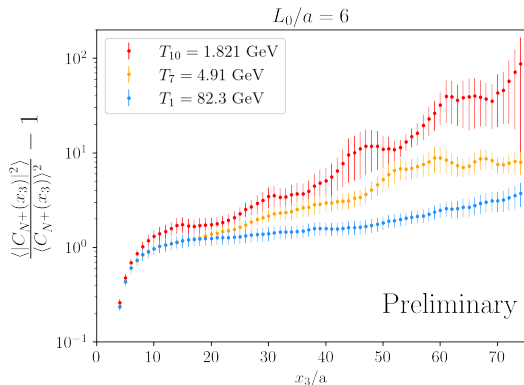
$\downarrow$   
 $2\pi T$

## Correlator and Noise-to-Signal ratio

- Correlator Noise-to-Signal ratio grows **exponentially** with distance **at zero temperature** [Parisi 1984, Lepage 1989]
- In the **infinite temperature limit** there is **no exponential problem** for  $C_{N\pm}(x_3)$
- At **high but finite temperature** milder **signal loss** is expected

$$\frac{\langle |C_{N\pm}(x_3)|^2 \rangle - \langle C_{N\pm}(x_3) \rangle^2}{\langle C_{N\pm}(x_3) \rangle^2} \sim \exp\{(2m_{N\pm} - 3m_{S/P})x_3\}$$

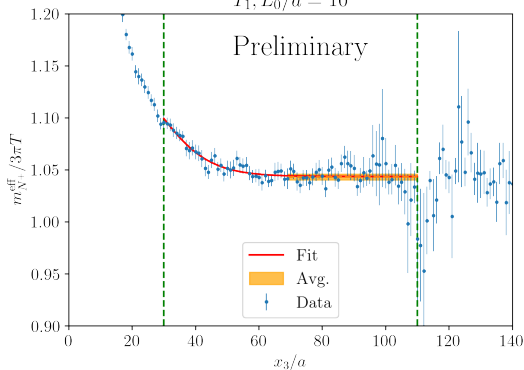
$$\begin{array}{ccc} \downarrow & T \rightarrow \infty & \downarrow \\ 3\pi T & & 2\pi T \end{array}$$



## Effective mass

$$T_1, L_0/a = 10$$

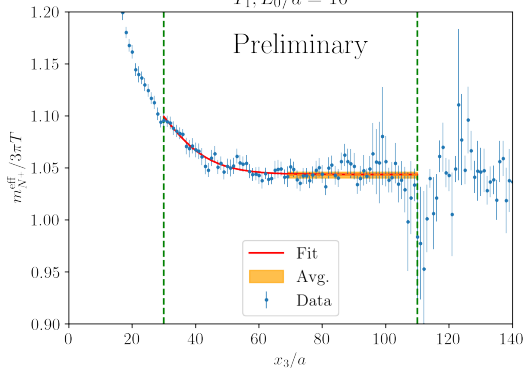
Preliminary



- Effective mass data is fitted to double exponential ansatz on correlator

## Effective mass

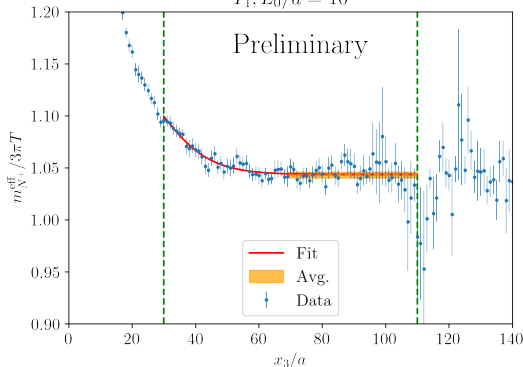
$$T_1, L_0/a = 10$$



- Effective mass data is fitted to double exponential ansatz on correlator
- Fit results used to determine plateau start

## Effective mass

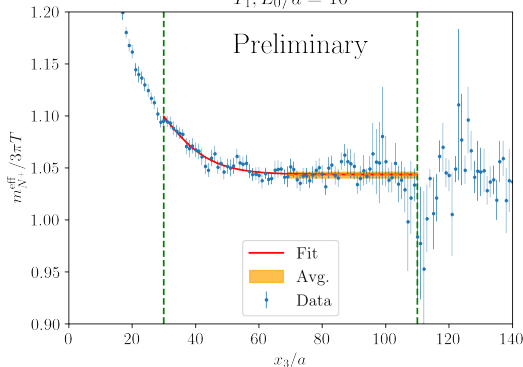
$$T_1, L_0/a = 10$$



- Effective mass data is fitted to double exponential ansatz on correlator
- Fit results used to determine plateau start
- High  $T$ : average in the plateau
- Low  $T$ : determine mass from fit

## Effective mass

$$T_1, L_0/a = 10$$



- Effective mass data is fitted to double exponential ansatz on **correlator**
- Fit results used to determine plateau start
- **High  $T$** : average in the plateau
- **Low  $T$** : determine mass from fit

## Tree Level Improvement

**Tree level screening mass** at finite lattice spacing computed analytically :  $m_{N^{\pm}}^{\text{Free}}(L_0/a)$

Accelerate continuum limit extrapolation by defining **tree level improved** screening mass:

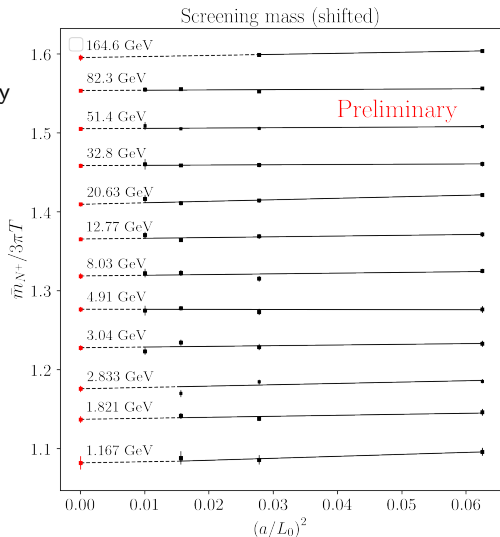
$$\bar{m}_{N^{\pm}}(L_0/a) \equiv m_{N^{\pm}}(L_0/a) - (m_{N^{\pm}}^{\text{Free}}(L_0/a) - 3\pi T)$$



## Continuum extrapolation

### Nucleon screening mass

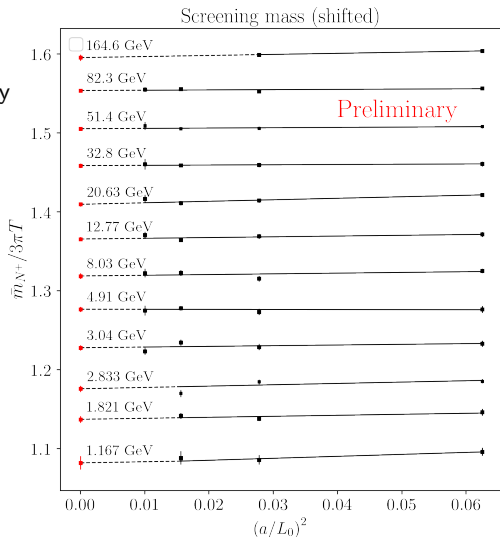
- $O(a)$  improvement: excellent fit quality for **constant** + **quadratic** polynomial



## Continuum extrapolation

### Nucleon screening mass

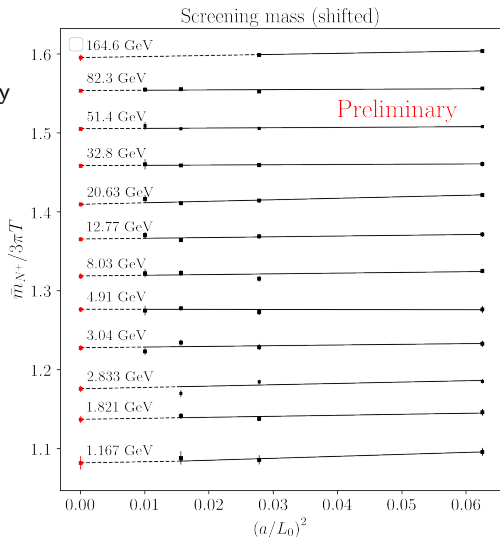
- $O(a)$  improvement: excellent fit quality for **constant** + **quadratic** polynomial
- TLI:  $(a/L_0)^2$  coefficients compatible with 0: small discretization effects



## Continuum extrapolation

### Nucleon screening mass

- $O(a)$  improvement: excellent fit quality for **constant** + **quadratic** polynomial
- TLI:  $(a/L_0)^2$  coefficients compatible with 0: small discretization effects
- Continuum extrapolated values obtained with **sub-percent** precision, in most cases **errors below 0.5%**



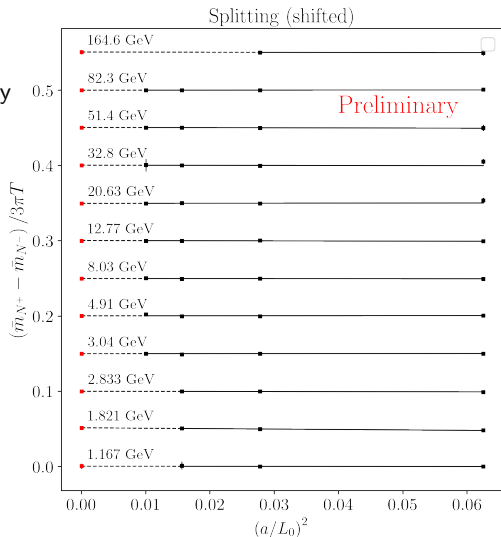
## Continuum extrapolation

### Nucleon screening mass

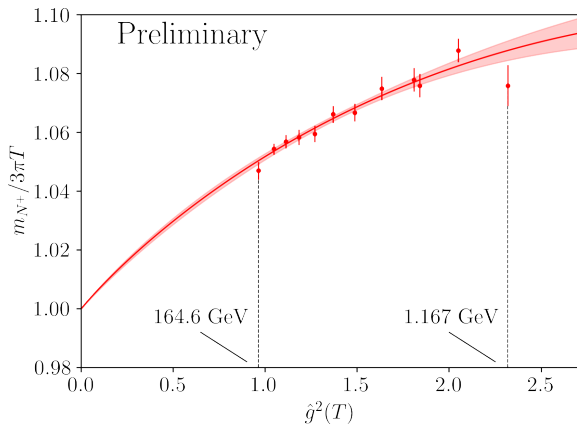
- $O(a)$  improvement: excellent fit quality for **constant** + **quadratic** polynomial
- TLI:  $(a/L_0)^2$  coefficients compatible with 0: small discretization effects
- Continuum extrapolated values obtained with **sub-percent** precision, in most cases **errors below 0.5%**

### Parity splitting

- $(m_{N^+} - m_{N^-})$  compatible with 0 for **all temperatures and lattice spacings**

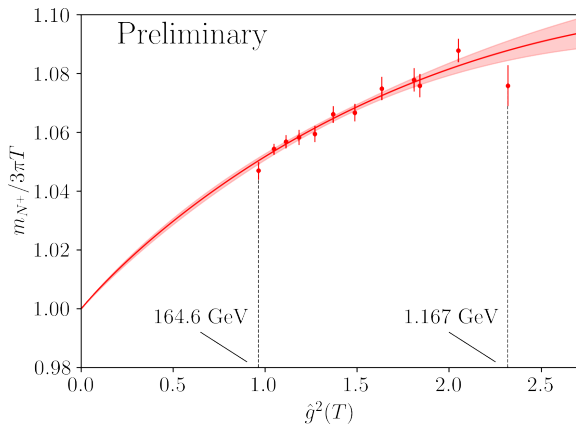


- $\frac{m_{N^+}}{3\pi T}$  displayed as a function of  $\hat{g}^2(T) \equiv \left( \frac{9}{8\pi^2} \ln \frac{2\pi T}{\Lambda_{\overline{\text{MS}}} } + \frac{4}{9\pi^2} \ln \left( 2 \ln \frac{2\pi T}{\Lambda_{\overline{\text{MS}}} } \right) \right)^{-1}$



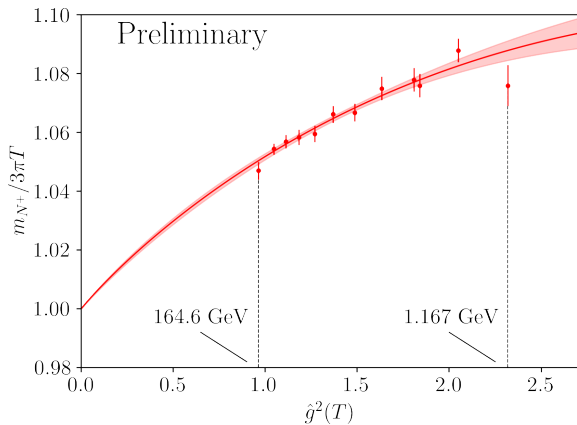
- $\frac{m_{N^+}}{3\pi T}$  displayed as a function of  $\hat{g}^2(T) \equiv \left( \frac{9}{8\pi^2} \ln \frac{2\pi T}{\Lambda_{\overline{\text{MS}}} } + \frac{4}{9\pi^2} \ln \left( 2 \ln \frac{2\pi T}{\Lambda_{\overline{\text{MS}}} } \right) \right)^{-1}$

- Data is fitted to a polynomial in  $\hat{g}$



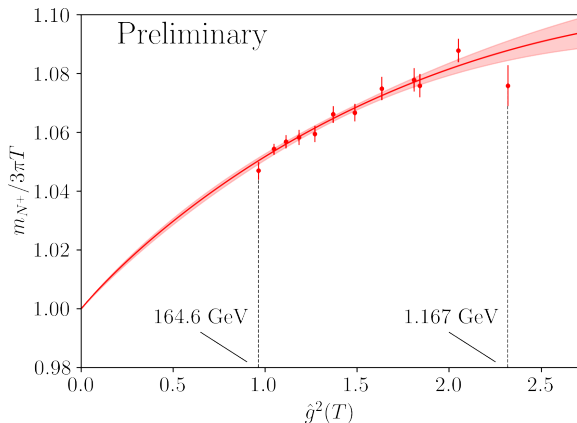
- $\frac{m_{N^+}}{3\pi T}$  displayed as a function of  $\hat{g}^2(T) \equiv \left( \frac{9}{8\pi^2} \ln \frac{2\pi T}{\Lambda_{\overline{\text{MS}}} } + \frac{4}{9\pi^2} \ln \left( 2 \ln \frac{2\pi T}{\Lambda_{\overline{\text{MS}}} } \right) \right)^{-1}$

- Data is fitted to a **polynomial** in  $\hat{g}$
- Free theory result: constant term set to 1  
 $\Rightarrow$  Single  $\hat{g}^2$  term **not enough** to reproduce behaviour



- $\frac{m_{N^+}}{3\pi T}$  displayed as a function of  $\hat{g}^2(T) \equiv \left( \frac{9}{8\pi^2} \ln \frac{2\pi T}{\Lambda_{\overline{\text{MS}}} } + \frac{4}{9\pi^2} \ln \left( 2 \ln \frac{2\pi T}{\Lambda_{\overline{\text{MS}}} } \right) \right)^{-1}$

- Data is fitted to a **polynomial** in  $\hat{g}$
- Free theory result: constant term set to 1  
 $\Rightarrow$  Single  $\hat{g}^2$  term **not enough** to reproduce behaviour
- Leaving constant term free (with curvature) yields fit results **compatible with 1**





## Conclusions

- Bulk of the nucleon screening mass is given by **free theory** contribution  $3\pi T$
- Interactions induce  $\sim 4\%$  to  $\sim 8\%$  deviations, clearly visible with **high statistical precision** of few per-mil
- **Chiral symmetry restoration**: opposite parity channels **degenerate within accuracy**
- Single  $\hat{g}^2$  term **unable to explain results** in this temperature range

## Outlook

- Conclude analysis and publish results
- Investigate screening masses of **non-static** mesonic fields in the same temperature range

---

**Thank you for your attention!**