

QUANTUM SIMULATION OF FINITE TEMPERATURE SCHWINGER MODEL VIA QUANTUM IMAGINARY TIME EVOLUTION

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Finite-T QFT on quantum computers (QC)

How to simulate finite-T QFTs on QC?

Finite T \sim Imaginary time evolution (non-unitary operation) \Rightarrow Difficult on QC

Our work:

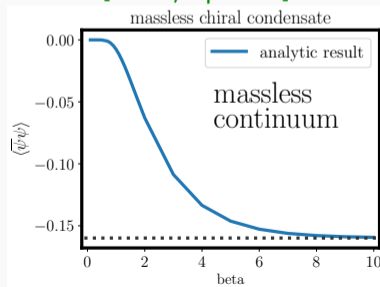
Simulate finite-T Schwinger model based on a quantum-classical hybrid algorithm

- Hamiltonian:

$$H = \frac{1}{4a} \sum_{n=1}^{N-1} [X_n X_{n+1} + Y_n Y_{n+1}] + \frac{m}{2} \sum_{n=1}^N (-1)^n Z_n + \frac{ag^2}{2} \sum_{n=1}^{N-1} \left[\sum_{i=1}^{n-1} \frac{Z_i + (-1)^i}{2} + \frac{\theta}{2\pi} \right]^2$$

- $\langle \bar{\psi} \psi \rangle \approx 0$ at low β ($\beta = 1/T$)
 $\langle \bar{\psi} \psi \rangle \neq 0$ at high β

[Sachs,Wipf 2010]



Introduction

Thermal states are mixed states:

$$\rho_\beta = Z^{-1} e^{-\beta \hat{H}}, \quad \langle \mathcal{O} \rangle_\beta^{\text{ens}} = Z^{-1} \text{Tr}(\mathcal{O} \rho_\beta), \quad (Z = \text{Tr} \rho_\beta)$$

⇒ Mixed state: Difficult to prepare on QC

Thermal Pure Quantum (TPQ) state [Sugiura, Shimizu 2011] [Sugiura, Shimizu 2013]

- Typical pure state in thermal system
- Able to calculate local thermodynamic quantities
- We use canonical TPQ:

$$|\beta, N\rangle \equiv e^{-\frac{\beta}{2} \hat{H}} |\psi_R\rangle \quad (|\psi_R\rangle : \text{random state})$$

- Average of initial random state

$$\overline{\left(\langle \hat{A} \rangle_{\beta, N}^{\text{TPQ}} - \langle \hat{A} \rangle_{\beta, N}^{\text{ens}} \right)^2} \rightarrow 0 \quad \text{w/ Thermodynamic limit} \quad N \rightarrow \infty$$

Preparing TPQ state on QC

Preparation of canonical TPQ state $|\beta, N\rangle \equiv \exp\left(-\frac{\beta}{2}\hat{H}\right)|\psi_R\rangle$ on QC

⇒ **Two tasks**

1. Preparation of random state $|\psi_R\rangle$
...Random state (unitary t-design) based on random circuits [Hunter-Jones 2019]
2. Implementation of non-unitary operation $\exp\left(-\frac{\beta}{2}\hat{H}\right)$
 - Variational method [McArdle+ 2018]
 - Probabilistic method [Liu+ 2020]
 - Quantum Imaginary Time Evolution (QITE) [Motta+ 2020]

Algorithm

QITE algorithm: Calculation step

1. Trotterization $e^{-\frac{\beta}{2}\hat{H}} \simeq (e^{-\Delta\beta\hat{h}_1}e^{-\Delta\beta\hat{h}_2} \dots)^{n_{step}}$
2. Substitute local non-unitary operation with larger unitary operation

$$Ce^{-\Delta\beta\hat{h}_1}|\psi\rangle \simeq e^{-i\Delta\beta\hat{A}}|\psi\rangle$$

3. Expand \hat{A} with Pauli strings and parameter \mathbf{a}

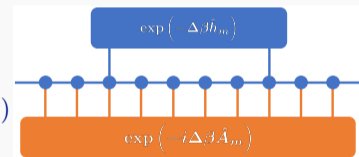
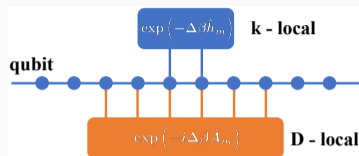
$$\hat{A}(\mathbf{a}) = \sum_{i_1 \dots i_k} a_{i_1 \dots i_k} \hat{\tau}_{i_1} \dots \hat{\tau}_{i_k} \equiv \sum a_I \hat{\sigma}_I$$

Pauli matrices $(\hat{\sigma}_I \in \{I, X, Y, Z\}^{\otimes D})$

4. Determine the optimal parameter $\mathbf{a} = \{a_I\}$
Then perform real time evolution

✓ local Hamiltonian (e.g.) Heisenberg model [Motta+ 2020], NJL model [Czajka+ 2022], Z_2 gauge theory (TPQ) [Davoudi+ 2022]

⇒ **How about Schwinger model, which includes non-local terms?**



Difficulty of QITE with non-local Hamiltonian

- For each term $e^{-\Delta\beta\hat{h}_m}$, we approximate $|\Phi\rangle_T$ with $|\Phi\rangle_A$:

$$\text{target state } |\Phi\rangle_T = C e^{-\Delta\beta\hat{h}_m} |\psi\rangle \quad \left(C = \langle\psi|e^{-2\Delta\beta\hat{h}_m}|\psi\rangle^{-1/2} \right)$$

$$\text{approximate state } |\Phi\rangle_A = e^{-i\Delta\beta\hat{A}(\mathbf{a})} |\psi\rangle$$

$$\hat{A}(\mathbf{a}) = \sum a_I \hat{\sigma}_I \quad \left(\hat{\sigma}_I \in \{I, X, Y, Z\}^{\otimes D} \right)$$

- Find the optimal parameter \mathbf{a} to minimize $\| |\Phi\rangle_T - |\Phi\rangle_A \|$

$$(\mathbf{S} + \mathbf{S}^T) \mathbf{a} = -\mathbf{b}$$

$$\text{w/ } S_{IJ} = \langle\psi|\hat{\sigma}_I\hat{\sigma}_J|\psi\rangle, \quad b_I = -2c^{-1/2} \text{Im}\langle\psi|\hat{\sigma}_I\hat{h}_m|\psi\rangle$$

- Measure $\langle\psi|**|\psi\rangle$ on QC

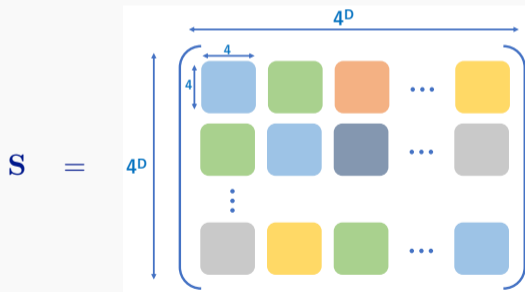
Solve the optimization problem by CG algorithm on CC

- Need to calculate $4^D \times 4^D$ matrix
non-local Hamiltonian \rightarrow large D

Practical improvements of QITE (Our idea)

To reduce cost and memory, we focus on the structure of S

1. Decompose S into 4×4 matrix
2. Independent elements = only first row (up to phase)



⇒ Result: cost and memory $O(16^D) \rightarrow O(4^D)$

- Hamiltonian
(Staggered fermion, Jordan-Wigner transformation, Gauss's law, open b.c.)

$$H = \frac{1}{4a} \sum_{n=1}^{N-1} [X_n X_{n+1} + Y_n Y_{n+1}] + \frac{m}{2} \sum_{n=1}^N (-1)^n Z_n + \frac{ag^2}{2} \sum_{n=1}^{N-1} \left[\sum_{i=1}^{n-1} \frac{Z_i + (-1)^i}{2} + \frac{\theta}{2\pi} \right]^2$$

- Parameters

$$N = 4 - 12, \quad a = 0.8, \quad g = 1.0,$$

$m = 0.00$: Exact solvable \rightarrow feasibility test

$m = 0.15$ with θ : prediction

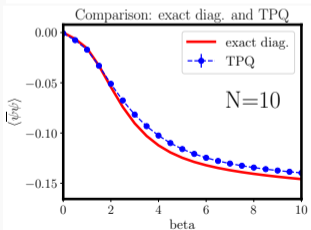
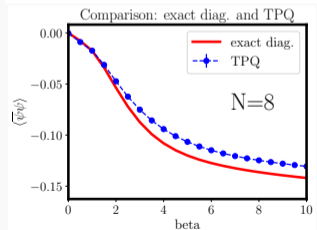
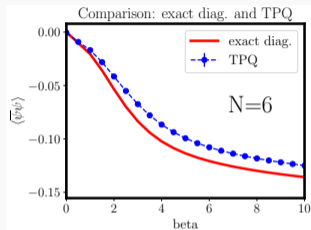
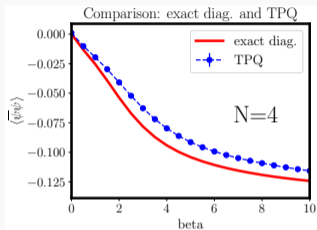
averaged over 100 – 1200 initial random states

- Methods:

1. Exact diagonalization (exact at finite N)
2. TPQ (classical algorithm)
3. QITE (quantum algorithm, statevector simulation)

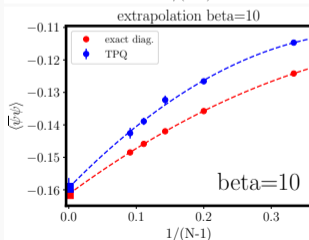
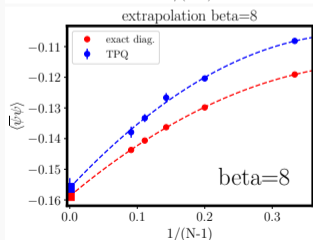
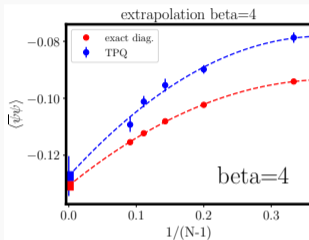
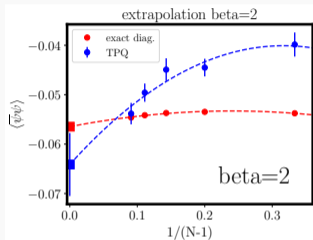
Results

Comparison: TPQ vs Exact diagonalization (massless case)



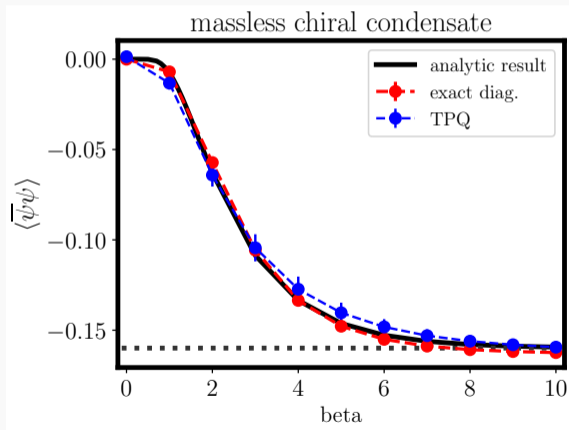
- At finite N , TPQ \neq exact diag.
- At $\beta \simeq 0$ or ∞ , TPQ = exact diag.

Extrapolation toward thermodynamic limit (massless case)



- $1/(N - 1)$ Extrapolation by quadratic function
- Extrapolated results of TPQ are consistent with one of exact diag. at $N \rightarrow \infty$ (within the statistical errors)

Thermodynamic limit of chiral condensate (massless case)



- Analytic result of massless chiral condensate at continuum:
[Sachs,Wipf 2010]

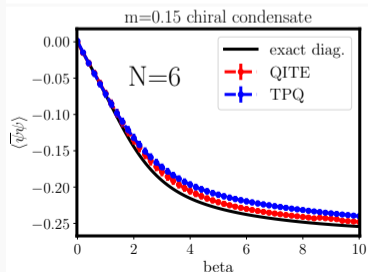
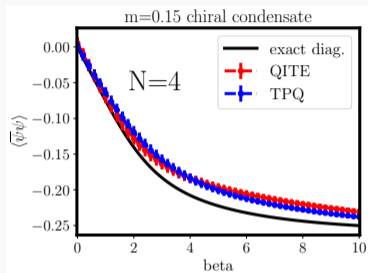
$$\langle \bar{\psi}\psi \rangle = -\frac{m_\gamma}{2\pi} e^\gamma e^{2I(\beta m_\gamma)}$$
$$I(x) = \int_0^\infty \frac{1}{1 - e^{x \cosh(t)}} dt$$

- At zero-T [Gross+ 1996],

$$\langle \bar{\psi}\psi \rangle = -e^{\frac{\exp(\gamma)}{2\pi^{3/2}}} \simeq -0.1599$$

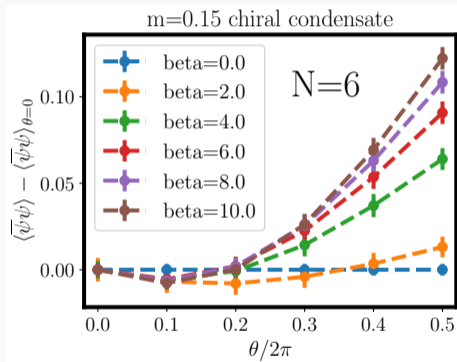
TPQ works well!

Comparison: QITE vs TPQ (massive case)



- Massive calculation ($m = 0.15$)
- Perform QITE algorithm with $\Delta\beta = 0.1$
- We found the difference between TPQ and QITE scales $O(\Delta\beta^2)$
- Systematic error from $\Delta\beta$ is under control

θ dependence of chiral condensate (massive case)

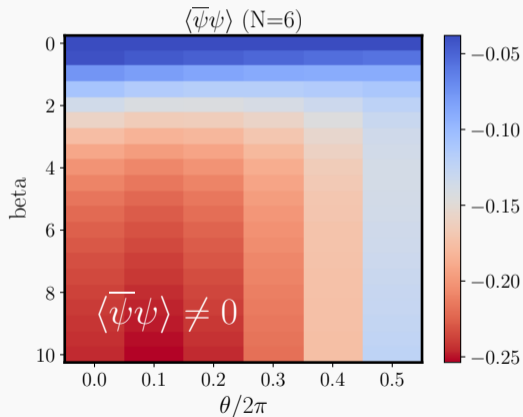


- QITE result
- Vertical axis denotes

$$\langle \bar{\psi}\psi \rangle_{\beta,\theta} - \langle \bar{\psi}\psi \rangle_{\beta,\theta=0}$$

- Low/High beta
...small/large θ -dependence

θ dependence of chiral condensate (massive case)



- $\langle \bar{\psi}\psi \rangle \neq 0$ at high β and small θ
 - In large θ , chiral symmetry seems to be restored
 - Taking thermodynamic limit: ongoing work
- ⇒ **TPQ+QITE works at massive and nonzero θ region !**

Summary

- We investigated the Schwinger model at finite T with TPQ and QITE
- Non-locality of the Hamiltonian requires large D , but we improved the QITE method and successfully handled large enough number of site
- TPQ result is consistent with T -dependence of chiral condensate at thermodynamic limit (massless case)
- Our method can be extended to massive and nonzero θ region
- Our calculations have been done classically, but **completely implementable on quantum circuits**

Discussion

- Application to other systems (higher dim., non-Abelian, finite density,...)
- Non-locality: Further improvements?
- Random state preparation
- Imaginary time evolution with local hamiltonian and Gauss' law projection

Thank you!

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TPQ (Thermal pure quantum) state

TPQ state ... Pure state that approximates the thermal state

- Definition:

$$P \left(\left| \langle \psi | \hat{A} | \psi \rangle_{TPQ} - \langle \hat{A} \rangle_{E,N}^{\text{eq}} \right| \geq \epsilon \right) \leq \eta_\epsilon(N)$$

$\langle \cdot \rangle_{E,N}^{\text{eq}}$: ensemble average

$\eta_\epsilon(N)$: function s.t. vanishes as $N \rightarrow \infty$

- Canonical TPQ state[Sugiura,Shimizu 2013]:

$$|\beta, N\rangle \equiv e^{-\frac{\beta}{2}\hat{H}} |\psi_R\rangle \quad (|\psi_R\rangle : \text{random state})$$

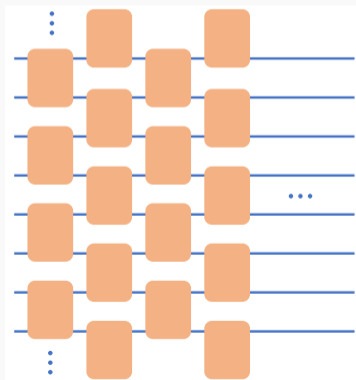
- Grand Canonical TPQ state[Sugiura 2014]:

$$|\beta, \mu, N\rangle \equiv e^{-\frac{\beta}{2}(\hat{H} - \mu\hat{N})} |\psi_R\rangle \quad (\hat{N} : \text{number operator})$$

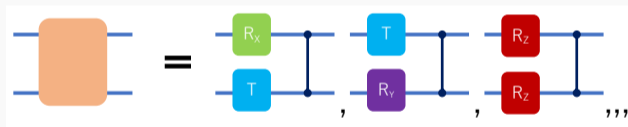
Random state on QC

Random state (unitary t-design) based on random circuits [Jones 2019]

...built from staggered layers of 2-site unitaries



- Each gate consists of single qubit gates randomly chosen from a universal gate set, and control-Z gate:



- Universal gate set (e.g.) :



w/ angle = $\pi/2$

QITE - practical improvements

- Utilizing the efficient Pauli measurement

“Pauli word” W : $W = W_1 W_2 \dots W_N \in \{X, Y, Z\}^N$

(e.g.) 10-sites

$$\begin{aligned} W &= XYZXYZXYZX \\ &= \text{XYZ XYZ XYZ } X \rightarrow \langle X_1 Y_2 Z_3 \rangle, \langle X_4 Y_5 Z_6 \rangle, \langle X_7 Y_8 Z_9 \rangle \end{aligned}$$

Improvement

Memory: $O(16^D) \rightarrow O(4^D)$

Calculation cost: $O(16^D) \rightarrow O(4^D)$ ($\rightarrow O(3^D)$)