

Decimation map in 2D for accelerating HMC

Nobuyuki Matsumoto

RIKEN BNL Research Center
(→ BU from Sep)



RIKEN BNL
Research Center



Brookhaven
National Laboratory

LATTICE 2023
Fermilab, 07.31.2023

Based on collaboration with
P. Boyle, R. Brower, T. Izubuchi, L. Jin, C. Jung, A. Tomiya

in preparation

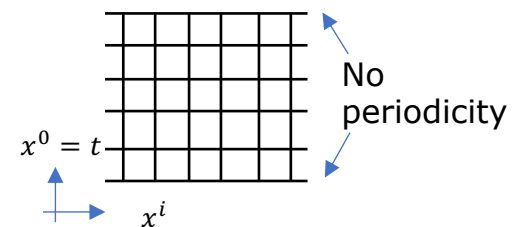


RIKEN's
Programs for
Junior Scientists
(SPDR)

Goal Overcome critical slowing down & topological freezing and perform efficient simulations on fine lattices.

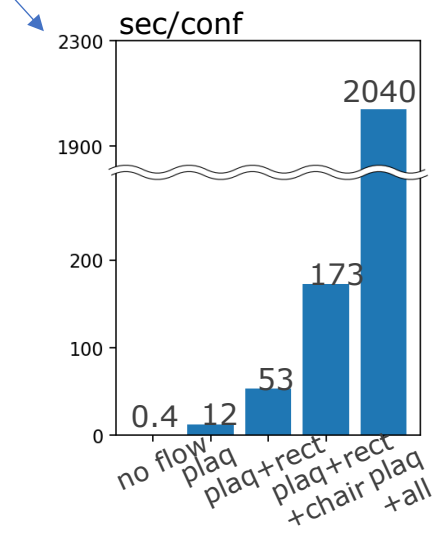
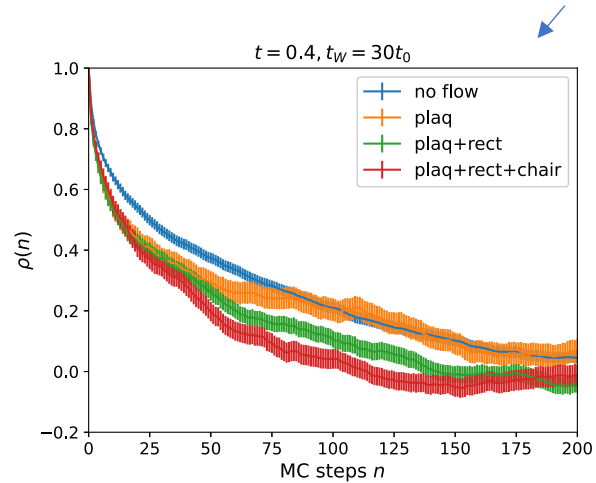
- Partial solution may be *open boundary condition*
Lüscher-Schaefer 11

But many statistical techniques assume translational invariance thus we want to keep the periodic boundary if possible.



Short timeline

- Proposal of *trivializing map* **Lüscher 09** (Name borrowed from **Nicolai 80**)
- LO approximation w/ CP^{N-1} **Engel-Schaefer 11**
"The reduction in the forces, ..., is compensated by the computational overhead"
- Wilson flow and its generalizations [4D SU(3) pure gauge]
Jin LATTICE 2021 poster **Boyle-Izubuchi-Jin-Jung-NM-Lehner-Tomiya LATTICE2022 [2212.11387]**



RIKEN HOKUSAI
 1 node
 (2 MPI x 40 OpenMP)

Tunneling rate can increase, but still overwhelmed by overhead.

➡ Step back and reconsider strategy using 2D U(1)

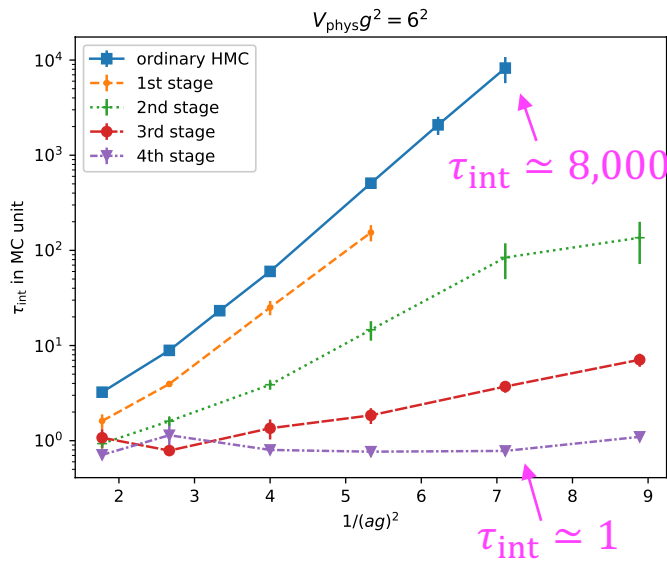
- Significant developments in normalizing flow/generative models

Albergo-Kanwar-Shanahan 19
Bacchio-Kessel-Schaefer-Vaitl 22

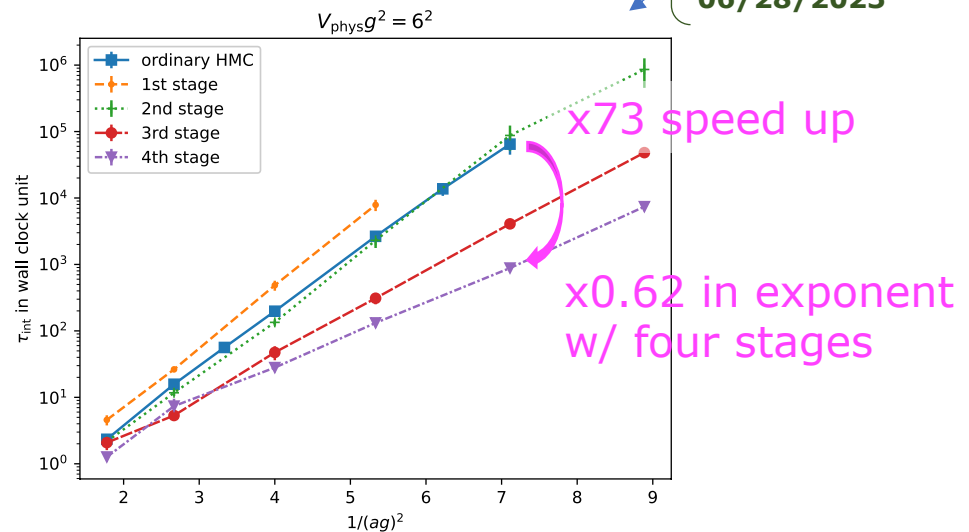
See Plenary of G. Kanwar this morning (Mon)

- Develop a variant of trivializing map: *"decimation map"*
Trivialization decomposed into several stages.
Each stage corresponds to integrating out local DOF
- Apply to 2D U(1) pure gauge with the guided MC (a generalization of HMC)
Horowitz 91

$\tau_{\text{int}}(Q)$ in MC unit



$\tau_{\text{int}}(Q)$ in wall-clock time

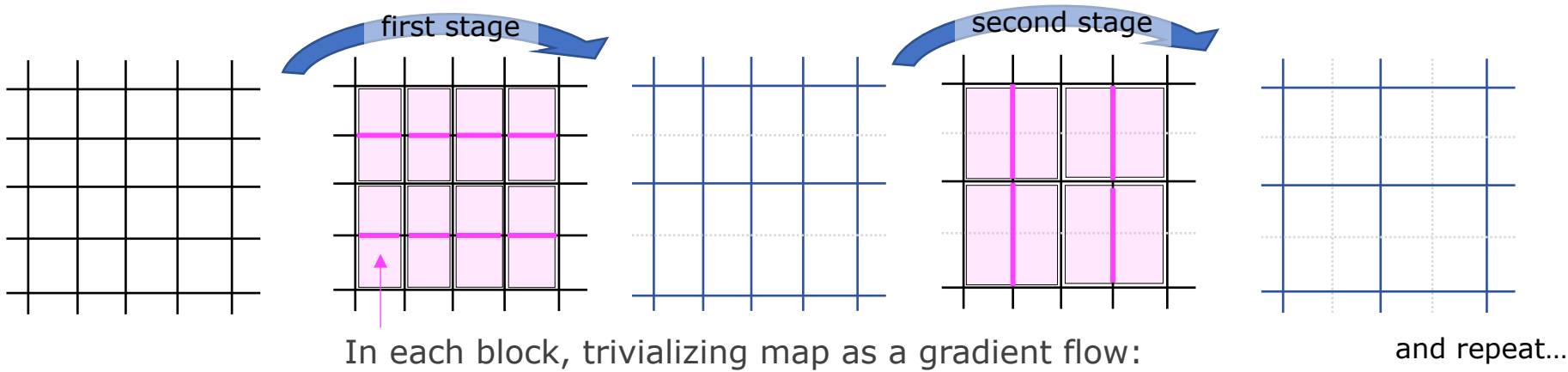


Updated from ECT* ML workshop 06/28/2023

- Algorithm is exact, scalings known and controlled.

Decimation map (1/2)

- Trivialization w/ several stages



$$\dot{U}_t U_t^{-1} = -\partial K_t \quad \left[K_t: \text{flow kernel} \right]$$

cf. Lüscher 09

- K_t can be obtained from solving the linear equation: $\mathcal{L}_t K_t = -S^{(\text{loc})}$. $\left[\mathcal{L}_t \equiv -\partial^2 - t \partial S^{(\text{loc})} \cdot \partial \right]$

We directly solve this with CGNE; only once as pre-calculation.

- Dividing the system into \mathcal{T} and \mathcal{R} ,
to be trivialized remainder

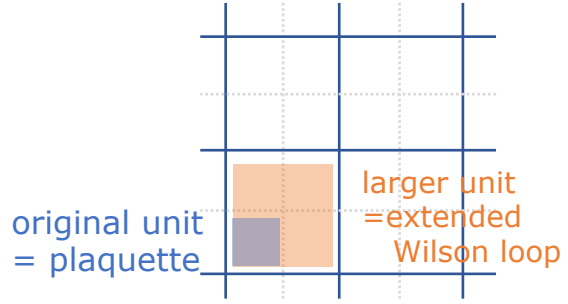
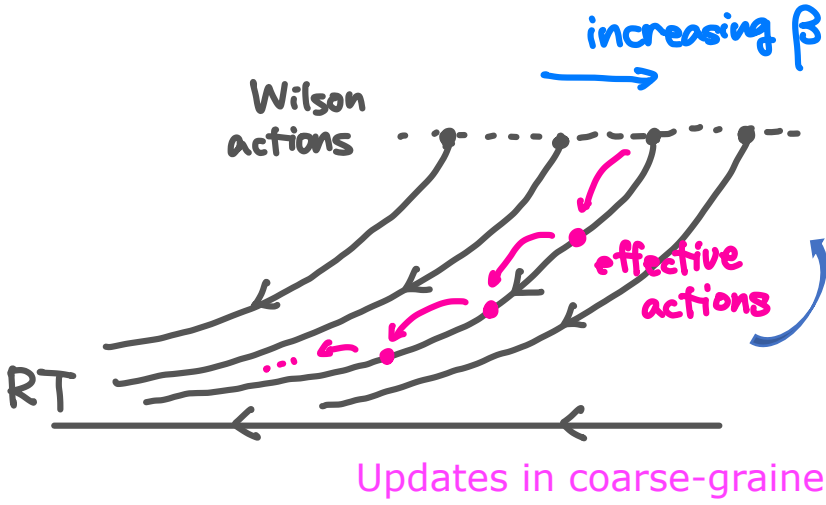
$$\begin{aligned}
 Z &= \int (d\mathcal{T})(d\mathcal{R}) e^{-S(\mathcal{T}, \mathcal{R})} &= & \int (d\mathcal{T})(d\mathcal{R}) e^{-S_{\text{eff}}(\mathcal{T}, \mathcal{R})} \\
 &\downarrow \text{integrate} && \downarrow \text{integrate} \\
 &\int (d\mathcal{R}) e^{-S_{\text{eff}}(\mathcal{R})} &= & \int (d\mathcal{R}) e^{-S_{\text{eff}}(\mathcal{R})}
 \end{aligned}$$

! =

must be the same!

- This map leaves the physics in large units exactly unchanged.
 - ∴ can be regarded as a coarse-graining map

cf. Wilson-Kogut 74, Kadanoff 75, Migdal 76
 see also U. Wenger Mon, R. Abbott Mon



Integrate in the decimated links when calculating observables.

- Acceleration expected because of
 - Increase of the trivialized region
 - Coarser action for the large-unit variables

Wilson action

$$S(U) \equiv -\beta \sum_x \cos \kappa_x \quad \left(\begin{array}{l} \kappa_x \equiv \frac{1}{i} \log(U_{x,0} U_{x+0,1} U_{x+1,0}^\dagger U_{x,1}^\dagger) : \text{plaquette angle} \\ \beta = \frac{1}{(ag)^2} \quad \quad \quad -\pi \leq \kappa_x < \pi \end{array} \right)$$

Solvable system, exact formulas from Fourier expansion.

see also D. Hoying Tue

Characteristic features

- Topological charge: $Q = \frac{-1}{2\pi} \sum_x \kappa_x \in \mathbb{Z}$ Simplest lattice gauge field theory with topology
- correlation length zero (ultralocal)
 deal with topological freezing rather than entire critical slowing down.
- topological susceptibility: $\chi_Q \sim \frac{g^2}{(2\pi)^2}$. \therefore typical instanton size $\sim (2\pi)^2/g^2$.

Guided-Hamiltonian Monte Carlo (a variant of HMC):

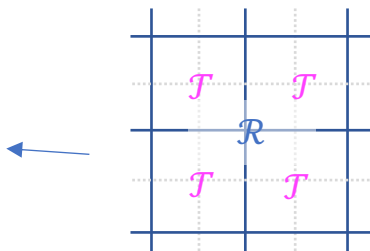
Horowitz 91

Duane et al. 87

- Replace the action in H by an approximate effective action (calculated in CG).
- Detailed balance still holds \because Liouville theorem and reversibility

Pros

- Simplifies the calculation of force
- Transparent separation of UV/IR
 \mathcal{T} \mathcal{R}

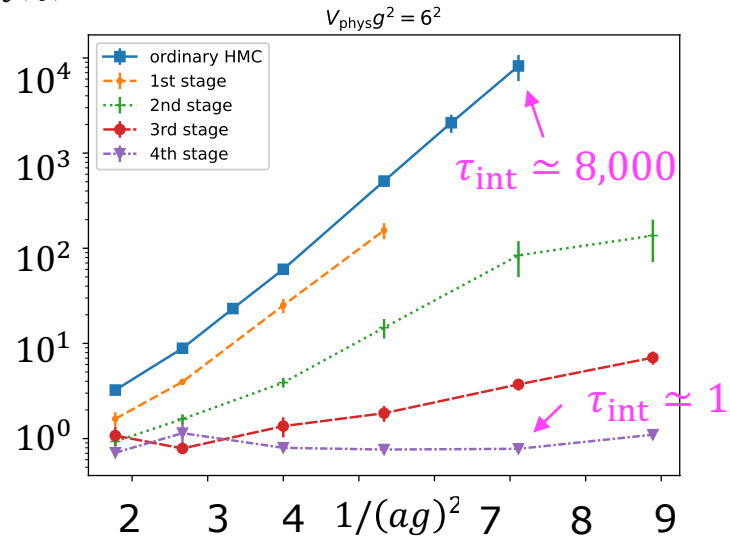


Cons

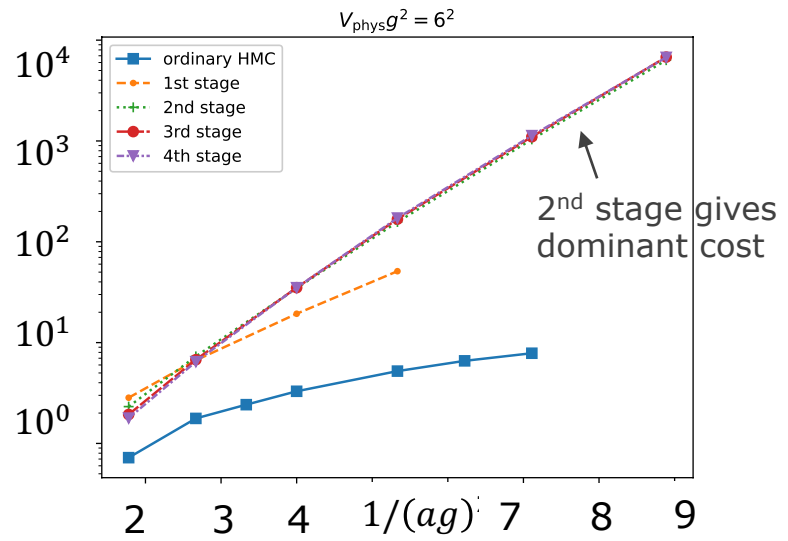
- Acceptance rate needs to be controlled additionally by the flow step size ϵ : $\delta H \propto \epsilon^n$
($n = 2$ below)
- no gauge fixing
- periodic boundary
- Flow equation solved with the midpoint integrator (\because net discretization error = $o(\epsilon^2)$ as mentioned)
- Bijective ensured by keeping ϵ in a bound. **cf. Lüscher 09**

Physical volume fixed to $\frac{6 \times 6}{g^2}$, ϵ scaled to keep acceptance rate ~ 0.9
~typical instanton size

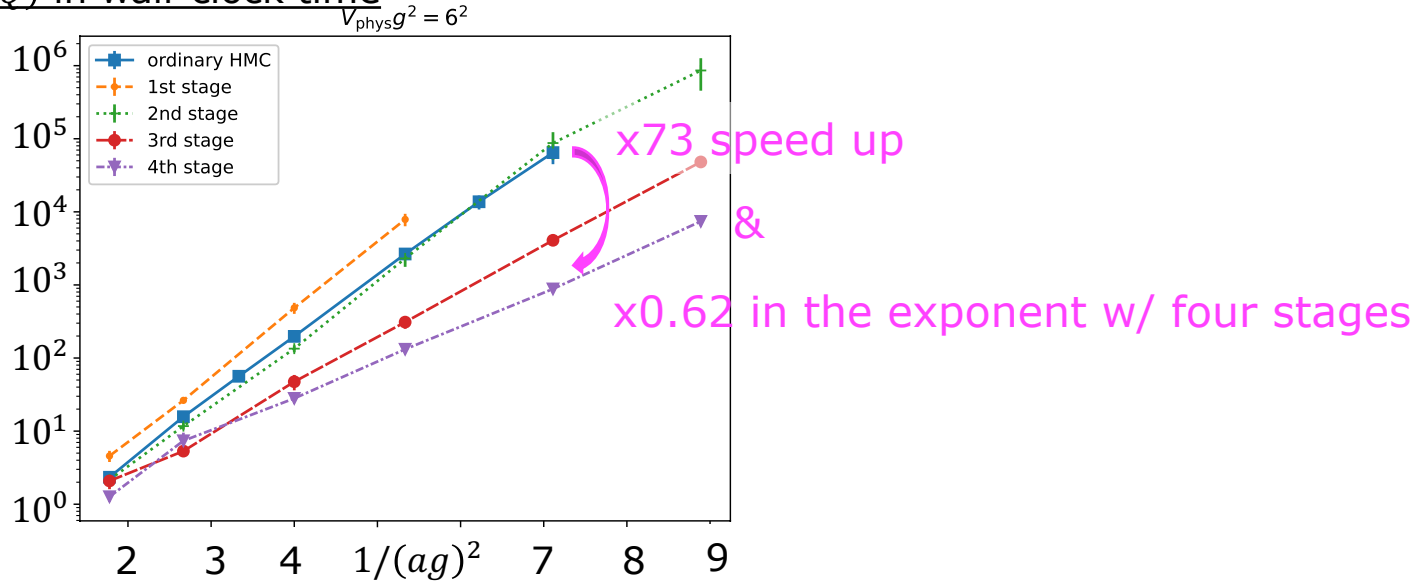
$\tau_{int}(Q)$ in MC unit



Execution time/conf (no parallelization)



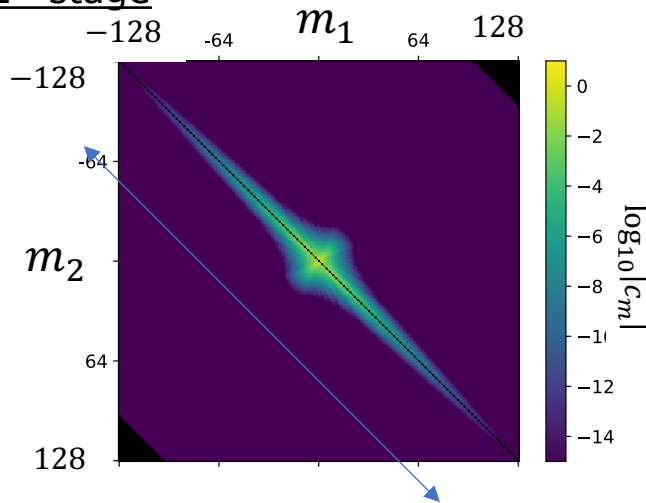
$\tau_{int}(Q)$ in wall-clock time



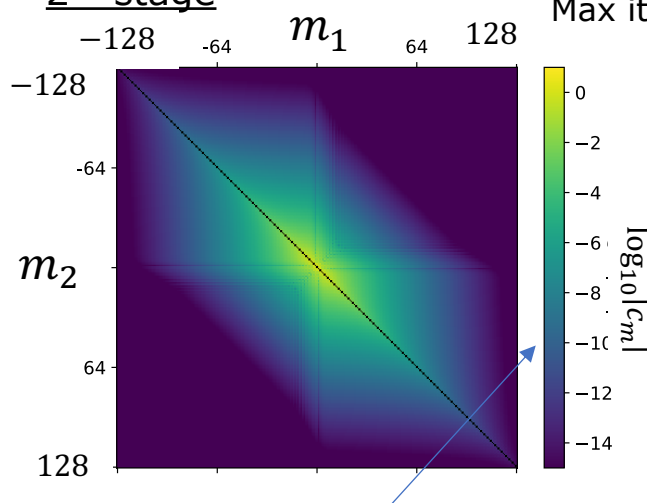
Kernel shapes @ $\beta=8.89$ (1/1)

Solves with about an hour using GPU.
Max iteration count: ~9000

1st stage

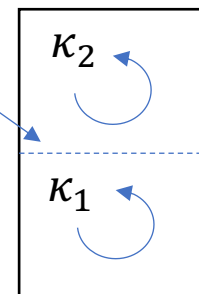


2nd stage



$$K_t = \sum_{\mathbf{m}} c_{t,\mathbf{m}} \cos(m_1 \kappa_1 + m_2 \kappa_2)$$

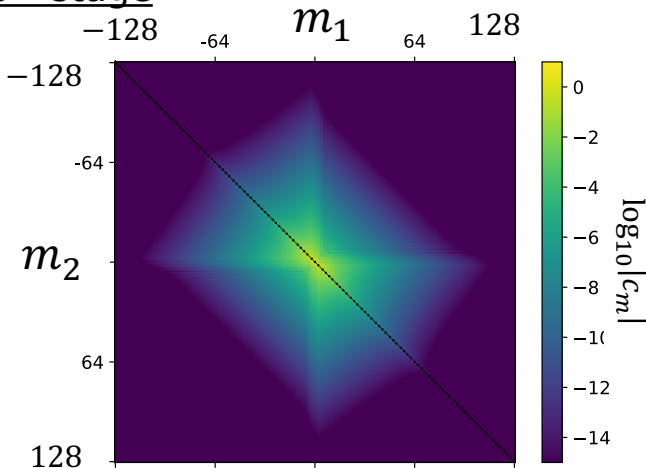
trivialized link



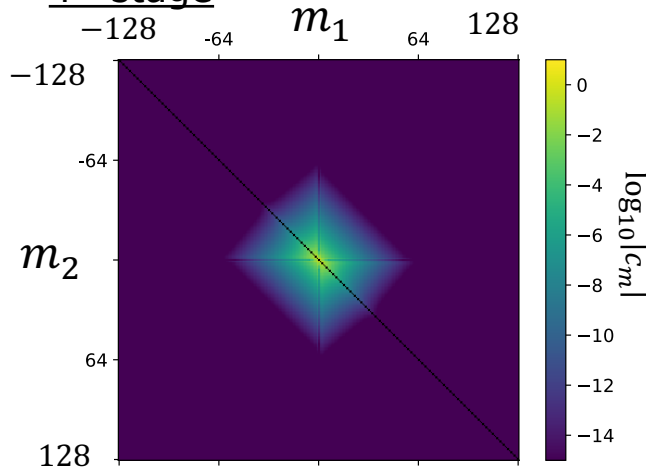
more spread out as $\beta \rightarrow \infty$

Due to higher-order loops in S_{eff}

3rd stage

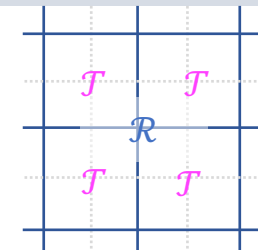


4th stage



#(relevant loops) shrinks
for higher stages=coarse actions

- 2nd stage is the most expensive simply because \exists many terms.
- From $1/\beta$ expansion of the kernel, suppression of large winding is seemingly power-law.

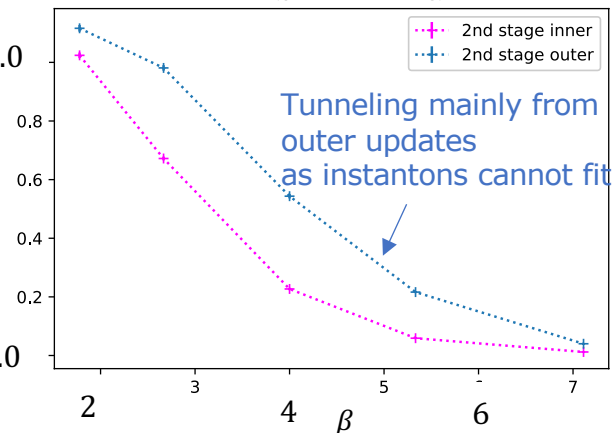


Analyzing the tunneling (1/1)

- By switching on/off the inner/outer updates alternatively, we investigate how tunneling is induced.

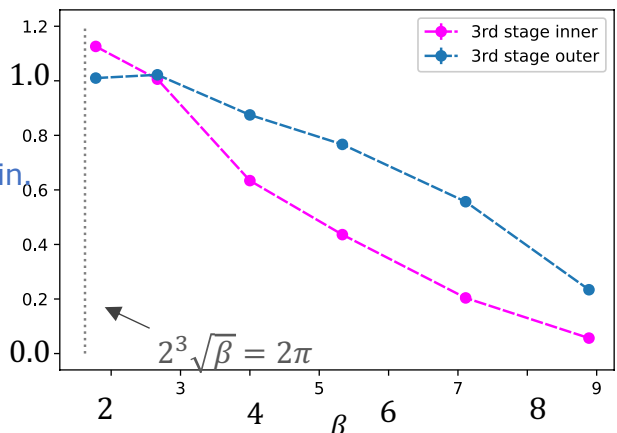
2nd stage

tunneling rate ($\equiv \langle |Q_n - Q_{n+1}| \rangle_{MC}$)



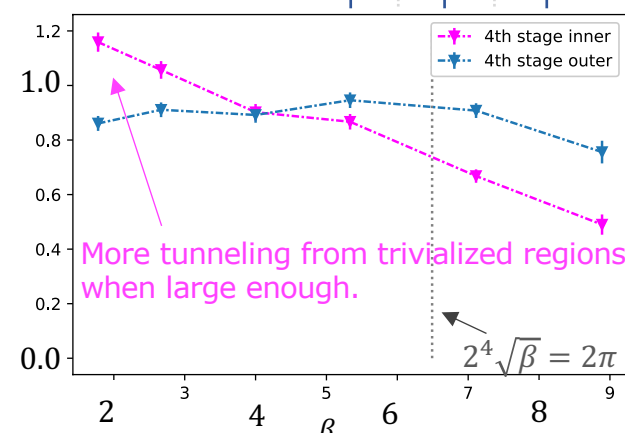
3rd stage

tunneling rate



4th stage

tunneling rate



- At the large β regime, the ratio of the tunneling rates (in MC units) can be understood from the probability:

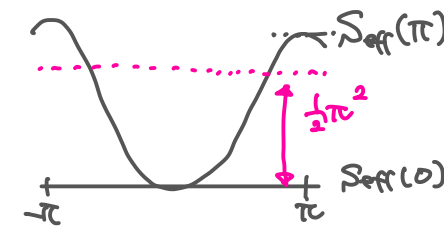
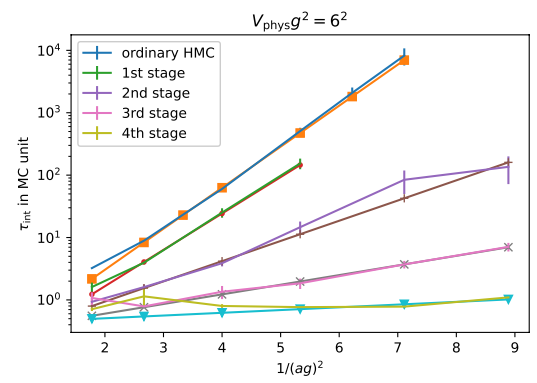
$$P\left(\frac{\pi^2}{2} > S_{\text{eff}}^{(\text{loc})}(\pi) - S_{\text{eff}}^{(\text{loc})}(0)\right), \pi \sim \mathcal{N}(0,1)$$

From fit

- 1: 0.886(52)
- 2: 0.493(40)
- 3: 0.235(14)
- 4: 0.067(42)

From prob

- 1: 0.856
- 2: 0.507
- 3: 0.242
- 4: 0.101

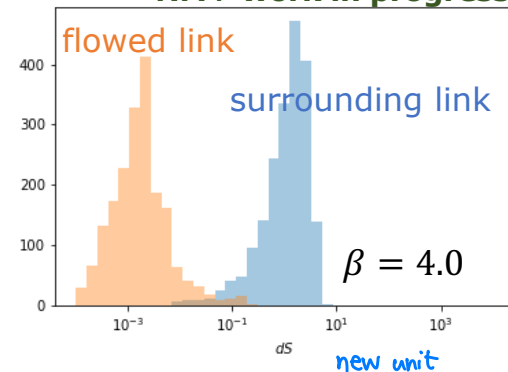


Aiming for large β , it seems more effective to get with field transformations a coarse-grained theory than enlarging completely trivialized regions.

Including fermion

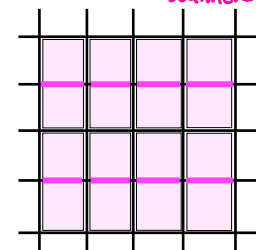
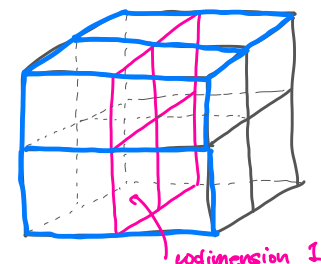
- We can expect exponential speed-up to remain when using exact S_{eff} in HMC (sometimes referred to as "FT-HMC")
see Lüscher 09
especially when fermion is present because of its additional cost.
cf. Engel-Schaefer 11, P.Boyle Plenary Tue
- However, effect of this peculiar smearing on fermion is nontrivial.
see also J. Finkenrath Mon

Forces (1st stage, exact S_{eff})
NM+ work in progress



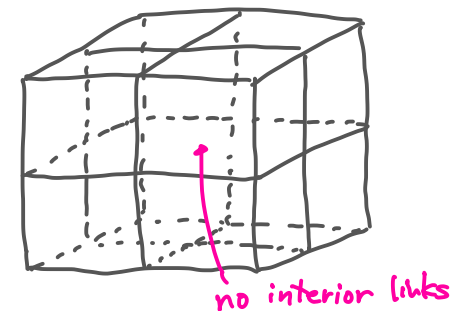
Higher-dimension

- trivializing map for codimension one surfaces → theory of 1x2 Wilson loops
∴ gauge invariance
2D Wilson is special that links on a line are independent;
much larger function space required for higher dimensions.
- Another possibility: **freeze links** surrounding local volume & trivialize interior.
→ "cage action", which can lower potential barriers for the remaining links.



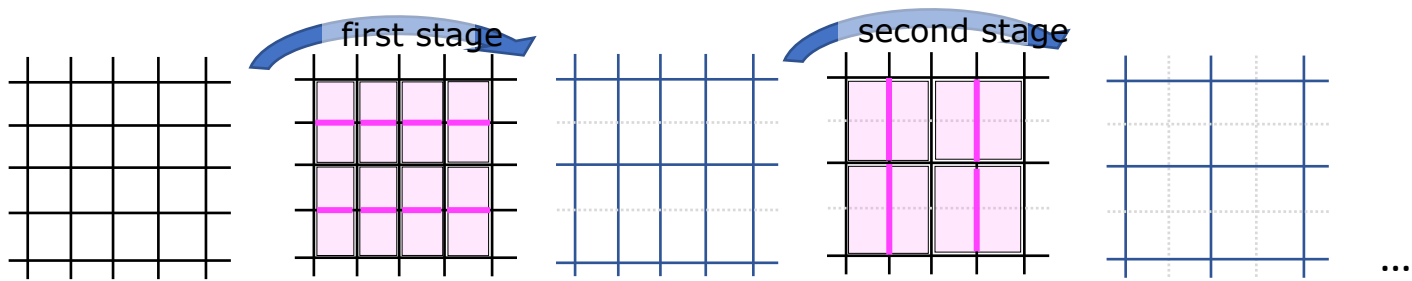
Non-Abelian groups

- We can work in the Wilson loop space to construct the function basis; labels may again be the winding numbers (though we need to take into account the noncommutativity and traces).
- Though there are Mandelstam constraints, **Mandelstam 79**
use of overcomplete basis seems possible when using CG.

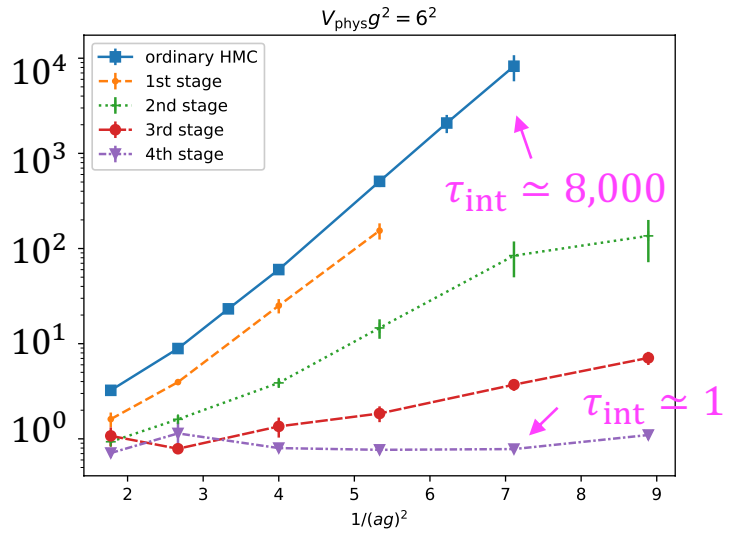


Summary (1/1)

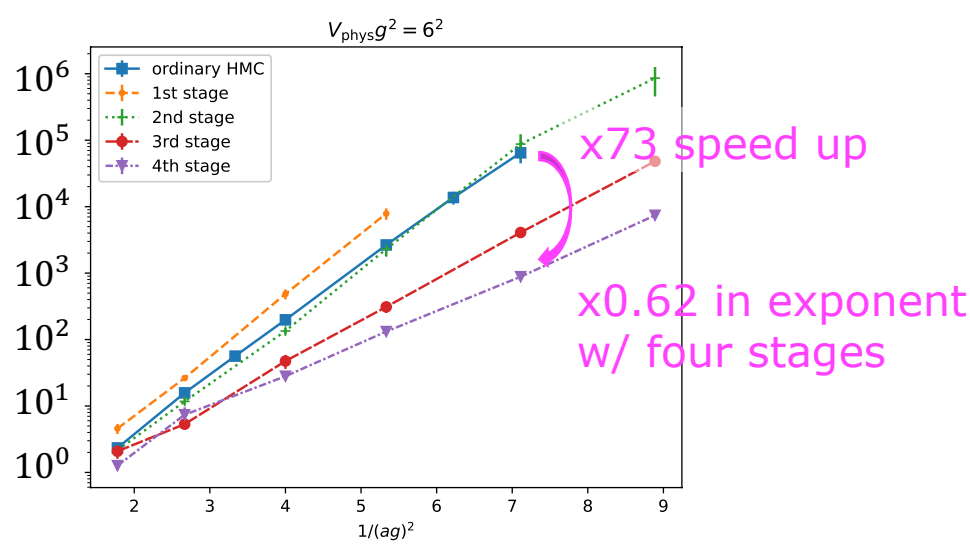
- We considered *decimation map* that can be regarded as a coarse-graining transformation.



$\tau_{int}(Q)$ in MC unit



$\tau_{int}(Q)$ in wall-clock time (no parallelization)



- It is true that the current investigation uses special features of 2D and U(1); however, we believe that having a method that works on this simplest model and has possible generalization directions will be a good starting point for developing algorithms for QCD.

Thank you.