

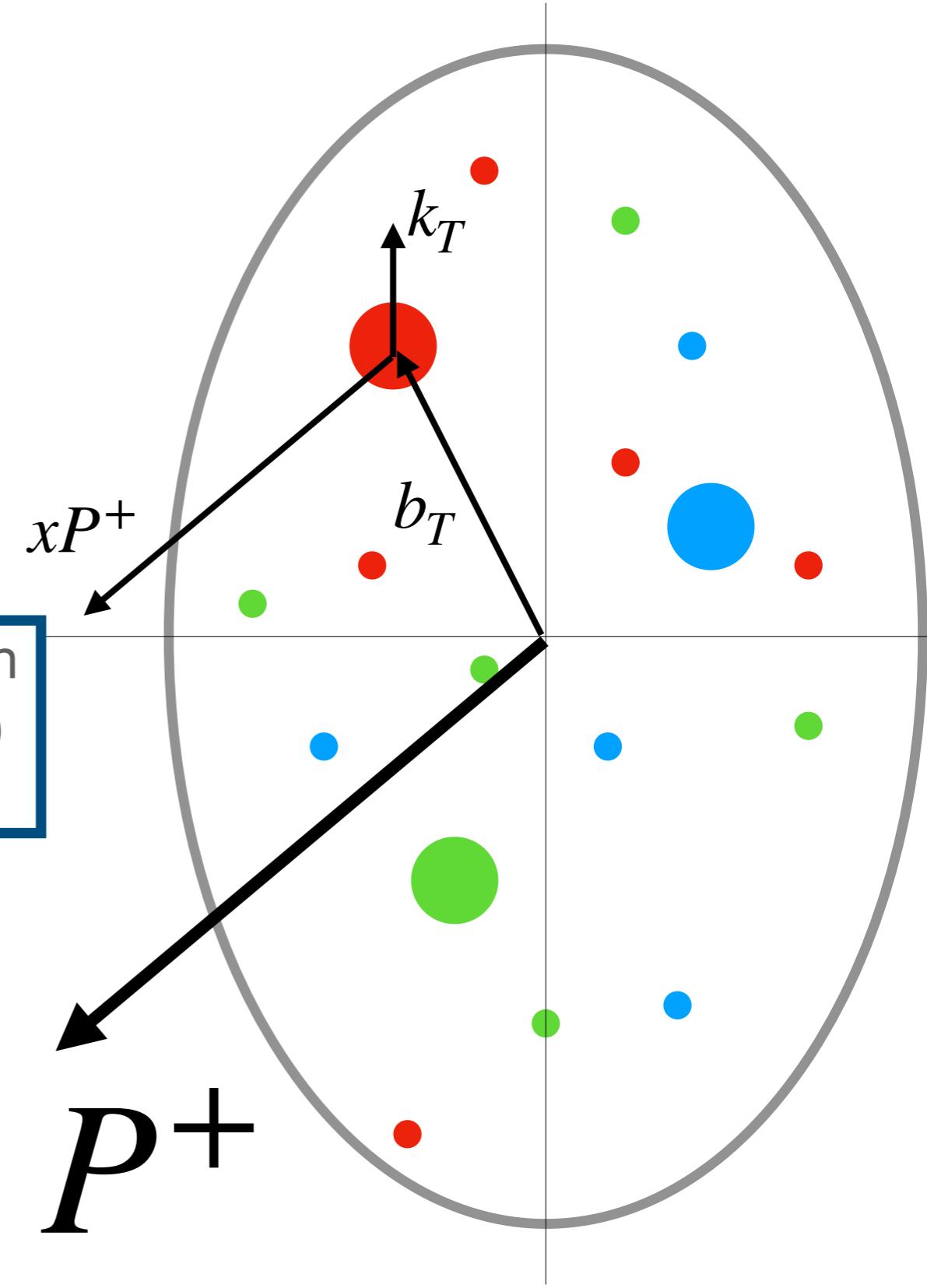
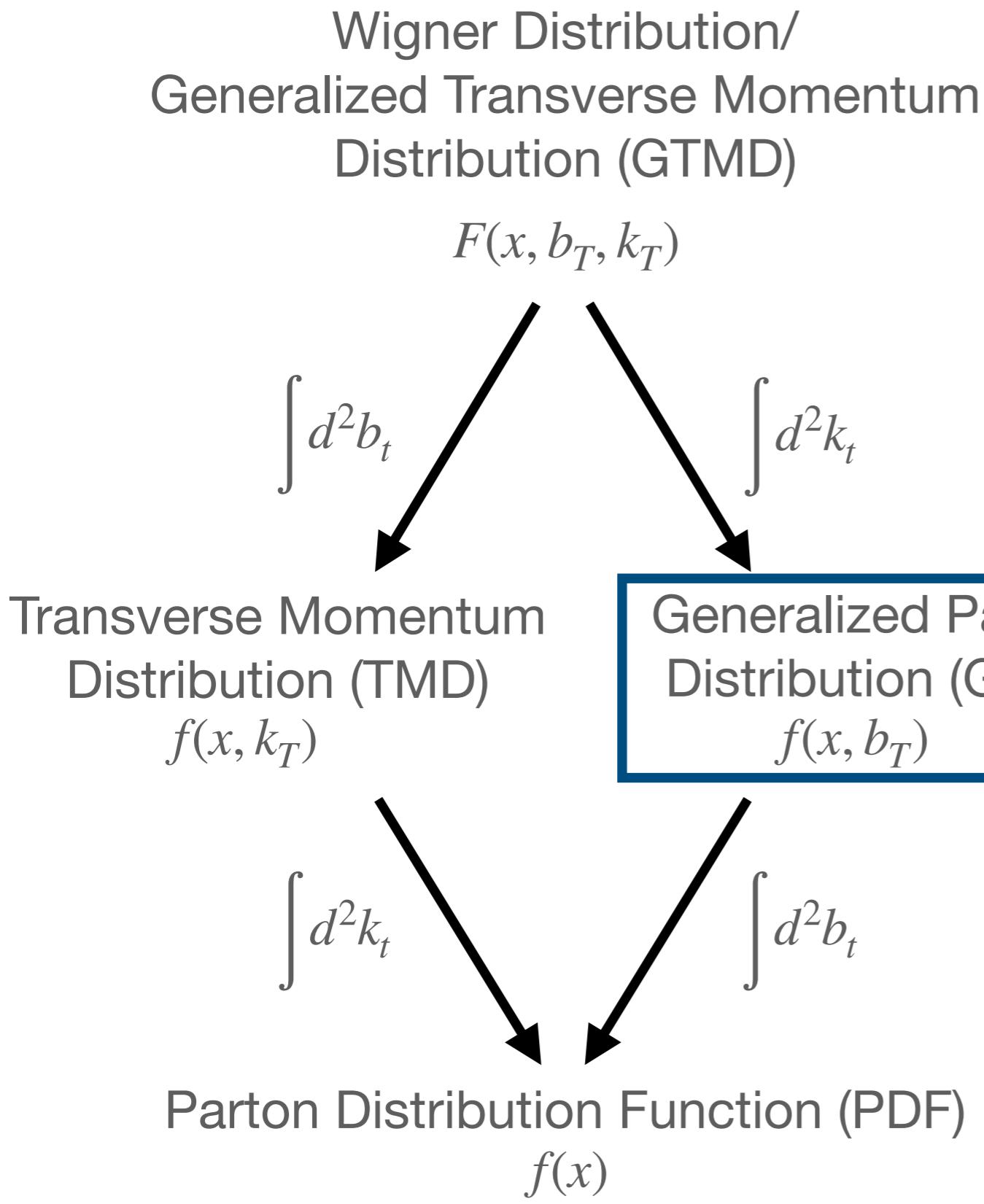
# Towards Unpolarized GPDs from Pseudo-Distributions

Joe Karpie (JLab) part of the HadStruc Collaboration



# Parton Structure

One for each flavor and spin combination



# Generalized Parton Distributions

- Generalized Ioffe time distributions

$$I^\mu(p', p, z = z^-, \mu^2) = \langle p' | \bar{q} \left( -\frac{z^-}{2} \right) \gamma^\mu W \left( -\frac{z^-}{2}, \frac{z^-}{2} \right) q \left( \frac{z^-}{2} \right) | p \rangle_{\mu^2}$$

Ioffe Time

$$\nu = \frac{p + p'}{2} \cdot z = P \cdot z$$

Momentum Transfer

squared

$$t = (p' - p)^2 = q^2$$

Skewness

$$\xi = \frac{q \cdot z}{P \cdot z}$$

- Generalized Parton Distributions

$$\sigma^{\mu a} = \sigma^{\mu\nu} a_\nu$$

$$\cdot \int \frac{d\nu}{2\pi} e^{i\nu x} I^\mu(\nu, t, \xi, \mu^2) = H(x, t, \xi) \bar{u}' \gamma^\mu u + E(x, t, \xi) \bar{u}' \frac{i\sigma^{\mu q}}{2m} u + \dots$$

X. Ji PRL 78 (1997) 610-613

Lorentz decomposition important!

# Double Distributions

- Generalized Ioffe time distributions

$$I^\mu(p', p, z = z^-, \mu^2) = \langle p' | \bar{q} \left( -\frac{z^-}{2} \right) \gamma^\mu W \left( -\frac{z^-}{2}, \frac{z^-}{2} \right) q \left( \frac{z^-}{2} \right) | p \rangle_{\mu^2}$$

Ioffe Time

$$\nu = \frac{p + p'}{2} \cdot z = P \cdot z$$

- Double Distributions

Momentum Transfer

squared

$$t = (p' - p)^2 = q^2$$

Skewness

$$\xi = \frac{q \cdot z}{P \cdot z}$$

$$\sigma^{ba} = b_\mu \sigma^{\mu\nu} a_\nu$$

$$z_\mu I^\mu(\nu, t, \xi, \mu^2) = \int d\beta d\alpha e^{i\nu(\beta + \xi\alpha)} \left[ f_q(\alpha, \beta, t) \bar{u}' z_\mu \gamma^\mu u + k_q(\alpha, \beta, t) \bar{u}' \frac{i\sigma^{zq}}{2m} u \right.$$

A.V. Radyushkin Phys Lett B380 (1996) 417-425

$$\left. + \delta(\beta) D(\alpha, t) \xi \nu \bar{u}' u \right]^2$$

M.V. Polyakov and C. Weiss PRD 60, 114017 (1999)

# Two Faced Distributions

## Radon Transform

GPDs:

$x, \xi$

$$f(x, \xi) = \int_{-1}^1 d\beta \int_{-1+|\beta|}^{1-|\beta|} d\alpha \delta(x - \beta - \xi\alpha) \tilde{f}(\alpha, \beta)$$

DDs:

$\beta, \alpha$

- Interpretation: average/change in parton momentum fraction
- Mellin moments give Form Factors and Angular Momentum decomposition
- Complex interrelation of variables
  - ERBL/DGLAP regions and polynomiality



Statue of Janus Bifrons  
(Wikipedia)

- Interpretation: Hybridize PDFs/DAs
- $\beta$  acts like PDF  $x$
- $\alpha$  acts like DA  $x$
- GPD evolution and polynomiality arise naturally from parameterized DD

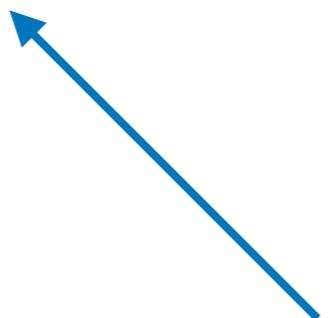
# Lorentz off the Lightcone (PDFs)

- Unpolarized Quark Ioffe time pseudo-distributions (pseudo-ITD)

$$M^\mu(p, z) = \langle p | \bar{q}(-\frac{z}{2}) \gamma^\mu W(-\frac{z}{2}, \frac{z}{2}) q(\frac{z}{2}) | p \rangle$$
$$= 2p^\mu \mathcal{M}(\nu, z^2) + 2z^\mu \mathcal{N}(\nu, z^2)$$

A. Radyushkin *PRD* 96 (2017) 3, 034025

Exist for  
lightcone at  $O(\alpha_s^0)$



# Lorentz off the Lightcone (PDFs)

- Unpolarized Quark Ioffe time pseudo-distributions (pseudo-ITD)

$$\begin{aligned} M^\mu(p, z) &= \langle p | \bar{q}\left(-\frac{z}{2}\right) \gamma^\mu W\left(-\frac{z}{2}, \frac{z}{2}\right) q\left(\frac{z}{2}\right) | p \rangle \\ &= 2p^\mu \mathcal{M}(\nu, z^2) + 2z^\mu \mathcal{N}(\nu, z^2) \end{aligned}$$

- Polarized Quark pseudo-ITD

C. Egerer et al (HadStruc) PRD 105 (2022) 3, 034507

R. Edwards et al (HadStruc) JHEP 03 (2023) 086

$$\begin{aligned} M^\mu(p, z) &= \langle p | \bar{q}\left(-\frac{z}{2}\right) \gamma^\mu \gamma^5 W\left(-\frac{z}{2}, \frac{z}{2}\right) q\left(\frac{z}{2}\right) | p \rangle & S^\mu = \bar{u} \gamma^\mu \gamma^5 u \\ &= 2mS^\mu \mathcal{M}(\nu, z^2) - 2im p^\mu z \cdot S \mathcal{N}(\nu, z^2) + 2m^3 z^\mu z \cdot S \mathcal{R}(\nu, z^2) \end{aligned}$$

Exist for

lightcone at  $O(\alpha_s^0)$

$$I^+(\nu, \mu^2) = 2mP^+ [\mathcal{M}(\nu, \mu^2) + i\nu \mathcal{N}(\nu, \mu^2)]$$

# Lorentz off the Lightcone (PDFs)

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$$M^\mu(p, z) = \langle p | \bar{q} \left( -\frac{z}{2} \right) \gamma^\mu W \left( -\frac{z}{2}, \frac{z}{2} \right) q \left( \frac{z}{2} \right) | p \rangle$$

$$= 2p^\mu \mathcal{M}(\nu, z^2) + 2z^\mu \mathcal{N}(\nu, z^2)$$

- Polarized Quark pseudo-ITD

$$M^\mu(p, z) = \langle p | \bar{q} \left( -\frac{z}{2} \right) \gamma^\mu \gamma^5 W \left( -\frac{z}{2}, \frac{z}{2} \right) q \left( \frac{z}{2} \right) | p \rangle \quad S^\mu = \bar{u} \gamma^\mu \gamma^5 u$$

$$= 2mS^\mu \mathcal{M}(\nu, z^2) - 2im p^\mu z \cdot S \mathcal{N}(\nu, z^2) + 2m^3 z^\mu z \cdot S \mathcal{R}(\nu, z^2)$$

- Helicity Gluon pseudo-ITD

$$M^{\mu\nu;\alpha\beta}(p, z) = \langle p | F^{\mu\nu} \left( -\frac{z}{2} \right) W \left( -\frac{z}{2}, \frac{z}{2} \right) \tilde{F}^{\alpha\beta} \left( \frac{z}{2} \right) | p \rangle$$

$$= \sum_{i=1}^{12} K_i^{\mu\nu;\alpha\beta}(p, z) \mathcal{M}_i(\nu, z^2)$$

**Only 2 exist for  
lightcone at  $O(\alpha_s^0)$**

I. Balitsky, W. Morris, A. Radyushkin JHEP 02 (2022) 193  
C. Egerer et al (HadStruc) PRD 106 (2022) 9, 094511

# Lorentz off the Lightcone (GPDs)

- Generalized pseudo-ITD

$$M^\mu(p', p, z) = \langle p' | \bar{q}\left(-\frac{z}{2}\right) \gamma^\mu W\left(-\frac{z}{2}, \frac{z}{2}\right) q\left(\frac{z}{2}\right) | p \rangle$$

Rest frame has  
same  
renormalization  
constant

$$\mathfrak{M}^\mu(p', p, z) = \sum_{i=1}^8 K^\mu(p', p, z) A_i(\nu, \xi, t)$$

$$\mathfrak{M}^\mu = \frac{M^\mu(p', p, z)}{M^4(p' = 0, p = 0, z)}$$

S. Bhattacharya et al PRD 106 (2022) 11, 114512

Gordon Identity means  
decomposition is not unique

- Leading twist amplitude for Double Distributions

$$A_1 \rightarrow f_q : \quad K_1 = \bar{u}' \gamma^\mu u$$

# Isolation of Amplitudes: SVD

$$\mathfrak{M}^\mu(p', p, z) = \sum_{i=1}^8 K^\mu(p', p, z) A_i(\nu, \xi, t)$$

Calculate for fixed  $p', p, z$  and vary initial/final spin and  $\mu$  to build matrix equation

$$\mathfrak{M} = KA$$

Pseudo-inverse solution  $\tilde{A}$  gives minimum of  $\chi^2 = (K\tilde{A} - \mathfrak{M})^2$

$$\tilde{A} = K^+ \mathfrak{M}$$

**Already used for Lattice Form Factor calculations**

# Inverse Problems: In all Parton Physics

- **Structure Functions**

$$F_2(x, Q^2) = \sum_i \int_x^1 d\xi C(\xi, \frac{\mu^2}{Q^2}) q\left(\frac{x}{\xi}, \mu^2\right)$$

- **Local Matrix elements / HOPE / OPE-without-OPE**

- **LaMET (on the lattice)**

$$M(p_z, z) = \int_{-\infty}^{\infty} dy e^{iy p_z z} \tilde{q}(y, p_z^2)$$

$$a_n(\mu^2) = \int_{-1}^1 dx x^{n-1} q(x, \mu^2)$$

- **Hadronic Tensor**

- **pseudo-Distributions / Good Lattice Cross Sections**

$$\mathfrak{M}(\nu, z^2) = \int_{-1}^1 dx C(x\nu, \mu^2 z^2) q(x, \mu^2)$$

$$\tilde{W}_{\mu\nu}(\tau) = \int d\nu e^{-\nu\tau} W_{\mu\nu}(\nu)$$

# Inverse Problems: Most probable solution

JK, K. Orginos, A. Rothkopf, S. Zafeiropoulos JHEP 04 (2019) 057

- Exists because information is lost. Must add new “prior information”
- Find the most probable distribution  $\langle q(x) \rangle = \int D[q]q(x)P[q | \mathcal{M}, I]$
- Posterior Prob by Bayes’s Theorem  $P[q | \mathcal{M}, I] \propto \exp[-\frac{\chi^2}{2}]P[q | I]$
- Choice of prior introduce biases
  - Bayesian Model averaging to amortize the biases
  - Allow for rigorous comparison of cuts in  $\chi^2$

W. Jay and E. Neil PRD 103, 114502 (2021)

# Inverse Problems: Most probable solution

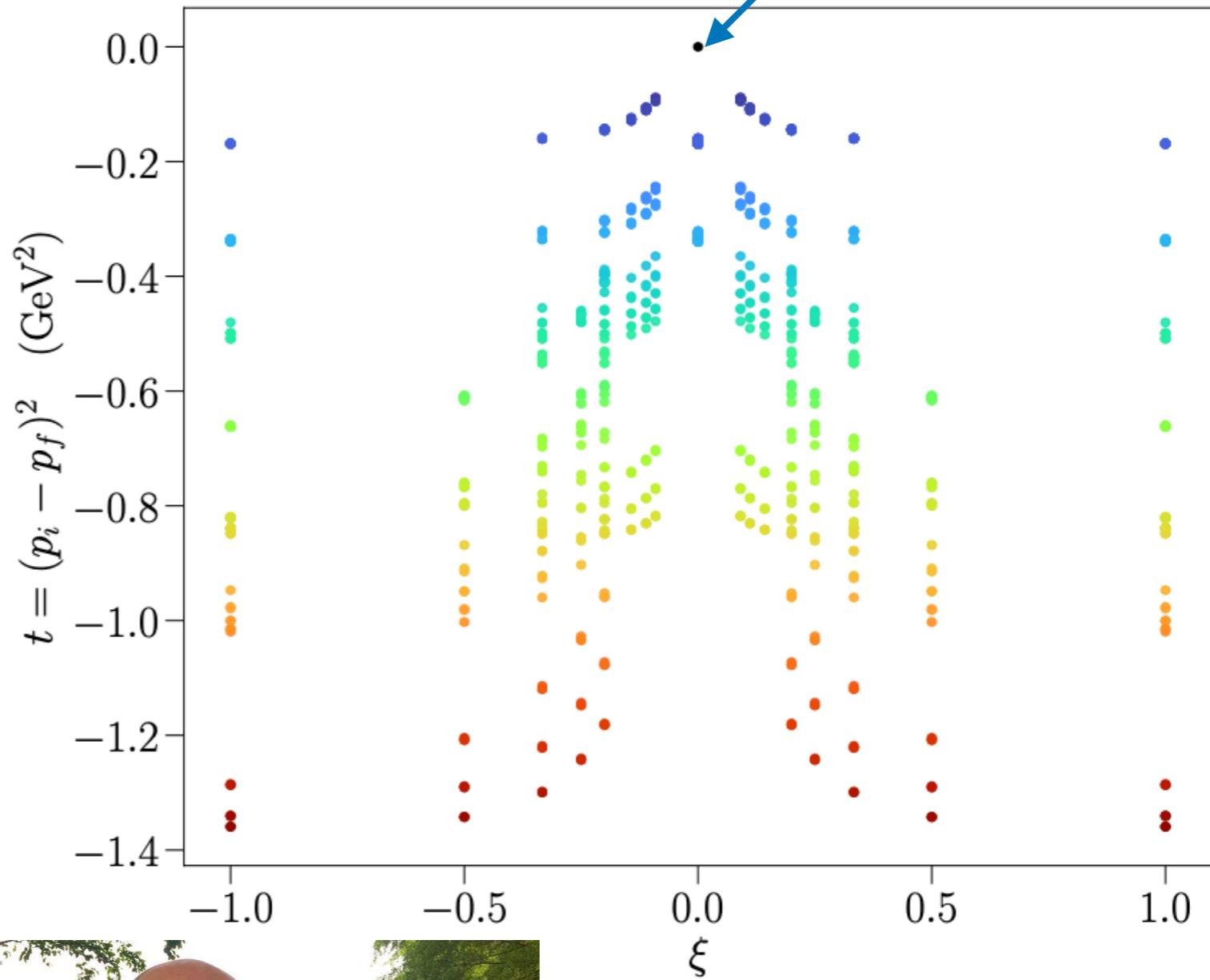
JK, K. Orginos, A. Rothkopf, S. Zafeiropoulos JHEP 04 (2019) 057

- Exists because information is lost. Must add new “prior information”

# Lattice Data

- Used Distillation to improve signal and allow easy substitution of different  $p$
- Summation method for controlling excited states
- Calculated 270  $(p', p)$  combinations
- Separations  $z \in [-8,8]\hat{3}$
- Different  $\gamma^\mu$  and spins to perform SVD
  - $\mu = 1, 2, 4$  and 2 spin combinations
- O(25k) matrix elements

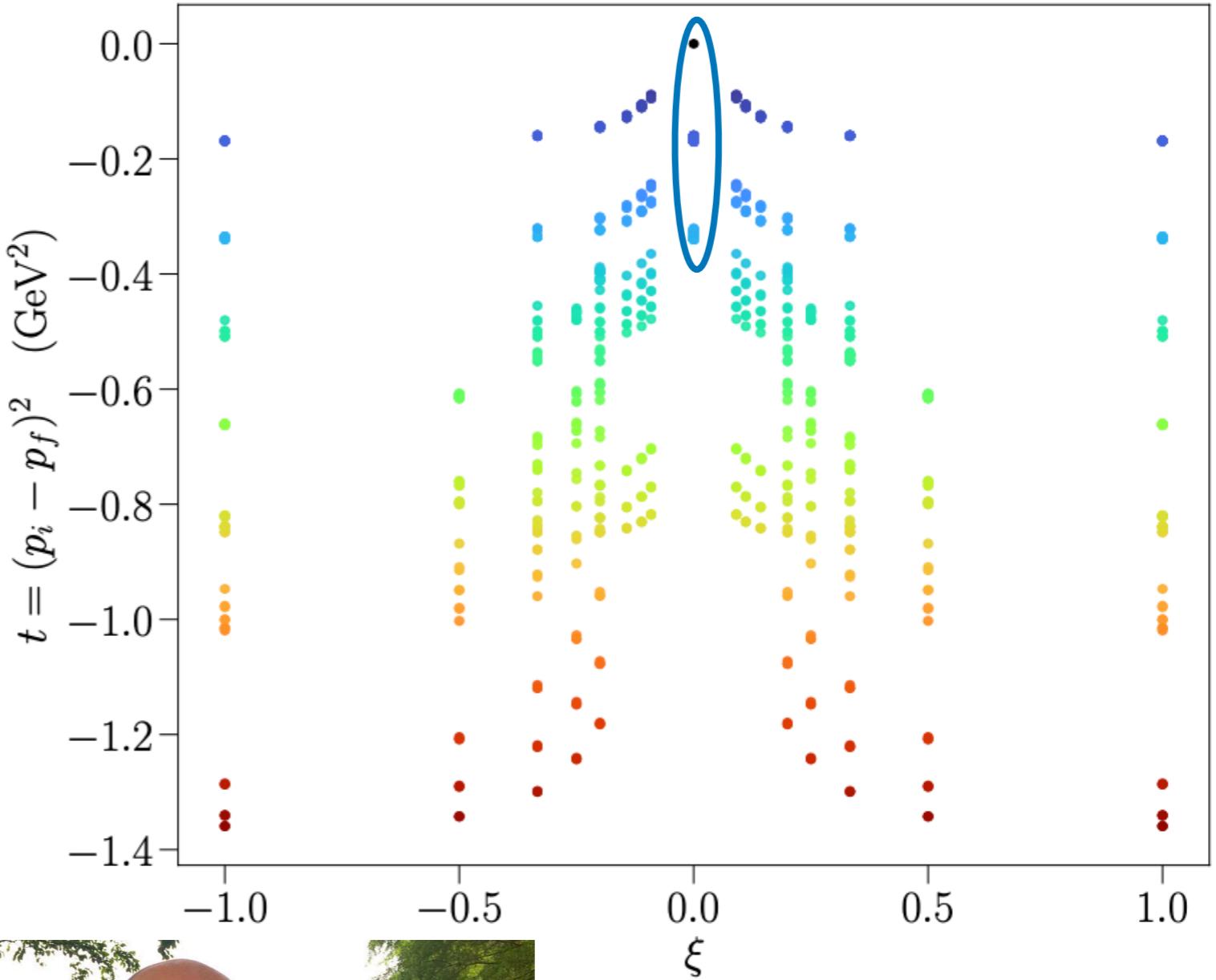
**PDF Calculation with  
72 matrix elements**



$$\begin{aligned} a &= 0.094(1) \text{ fm} \\ m_\pi &= 358(3) \text{ MeV} \\ L^3 \times T &= 32^3 \times 64 \end{aligned}$$

# Lattice Data

- Used Distillation to improve signal and allow easy substitution of different  $p$
- Summation method for controlling excited states
- Calculated 270  $(p', p)$  combinations
- Separations  $z \in [-8,8]\hat{3}$
- Different  $\gamma^\mu$  and spins to perform SVD
  - $\mu = 1, 2, 4$  and 2 spin combinations
- O(30k) matrix elements



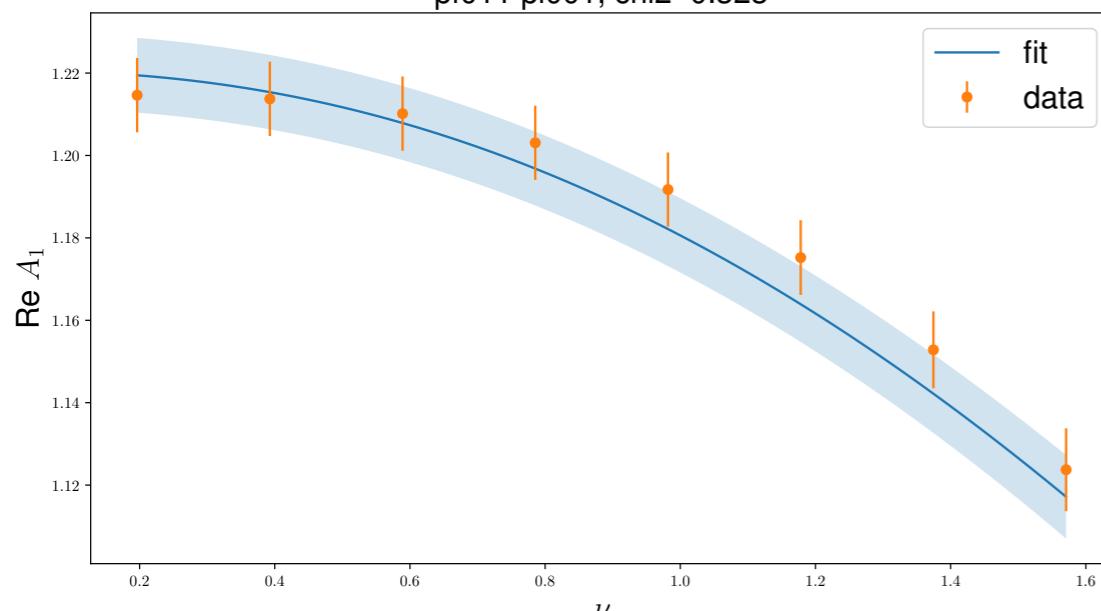
$$a = 0.094(1) \text{ fm}$$
$$m_\pi = 358(3) \text{ MeV}$$
$$L^3 \times T = 32^3 \times 64$$

$f_q$  at  $\xi = 0$

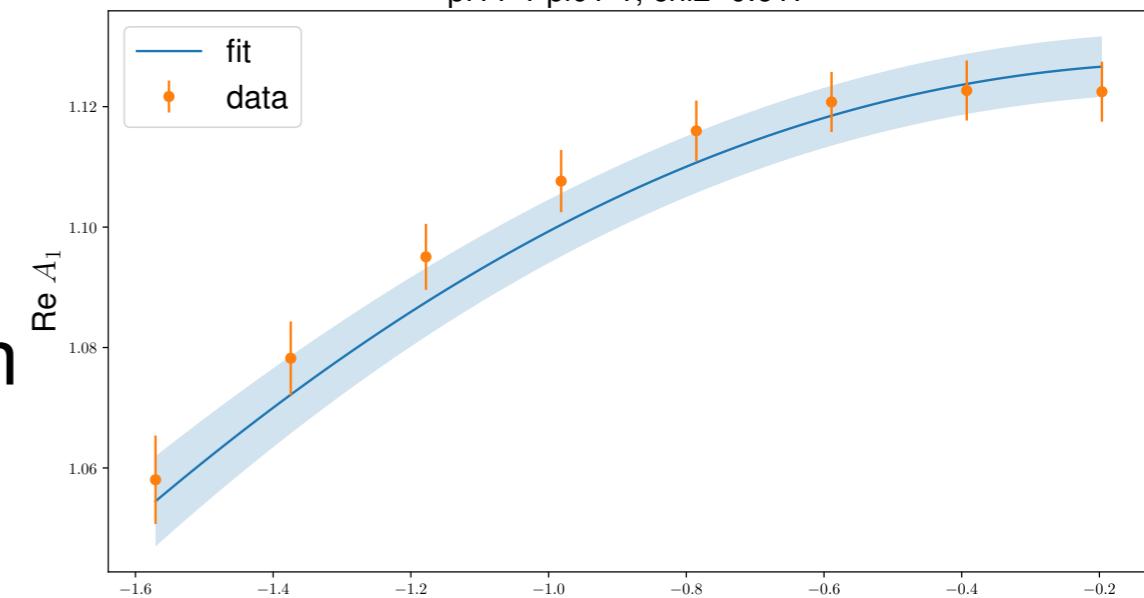
- Parametrize the DD with simplest form

$$A_1(\nu, t, \xi, z^2) = N \int d\beta e^{i\nu\beta} \beta^a (1 - \beta)^b$$

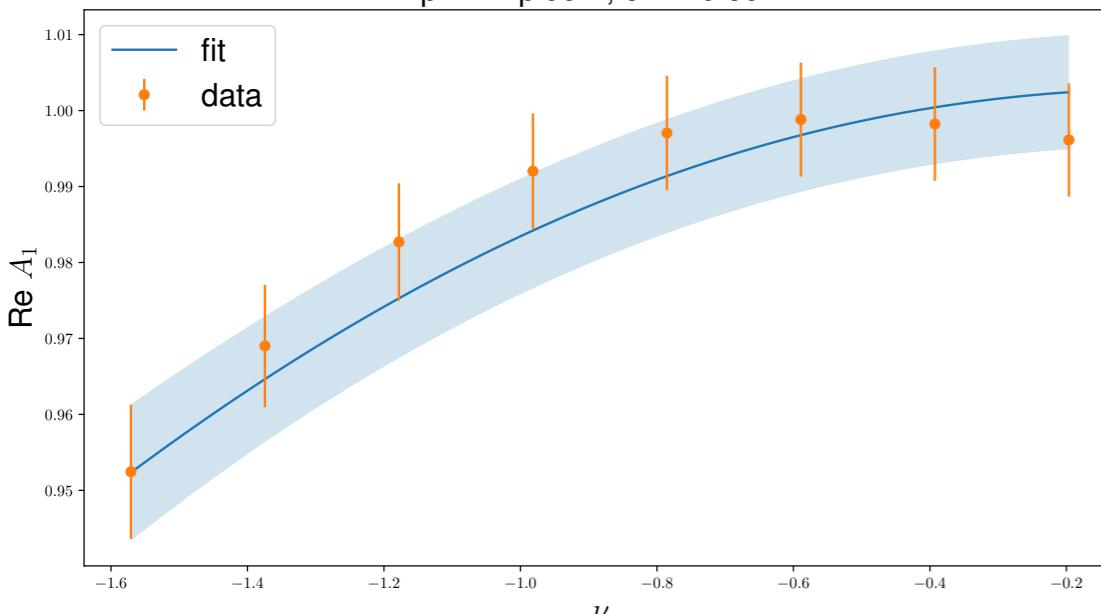
pf011 pi001, chi2=0.323



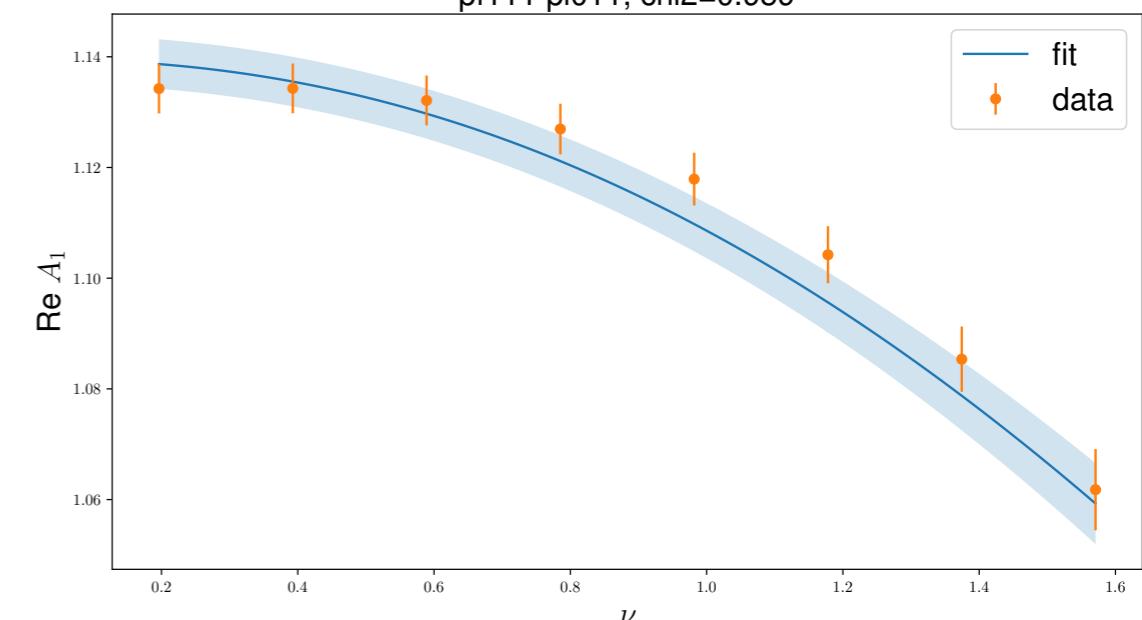
pf11-1 pi01-1, chi2=0.617



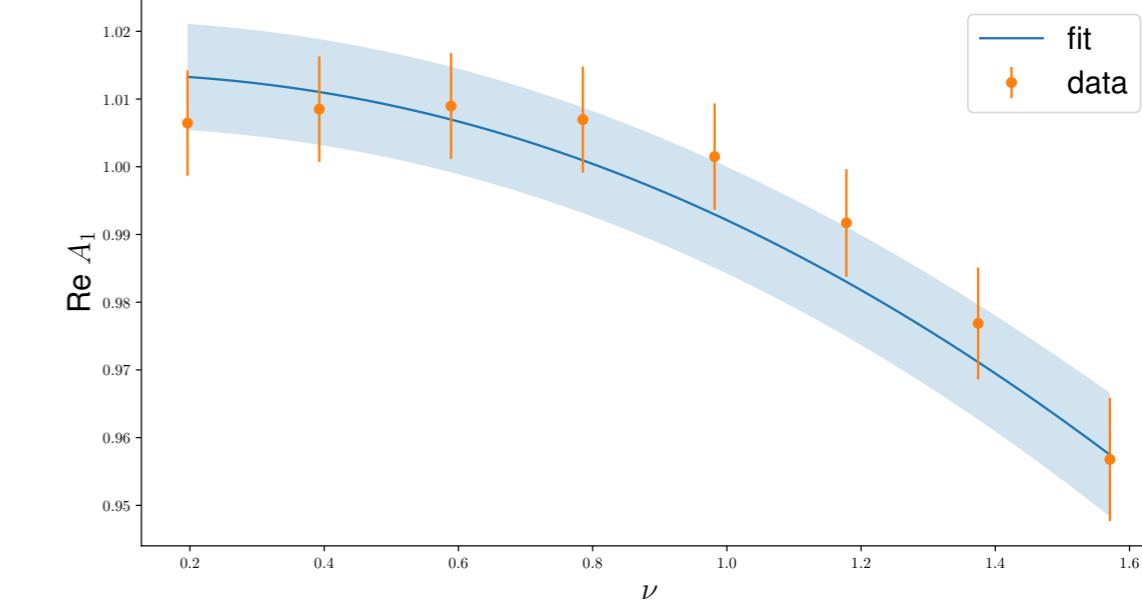
pf11-1 pi00-1, chi2=0.592



pf111 pi011, chi2=0.959



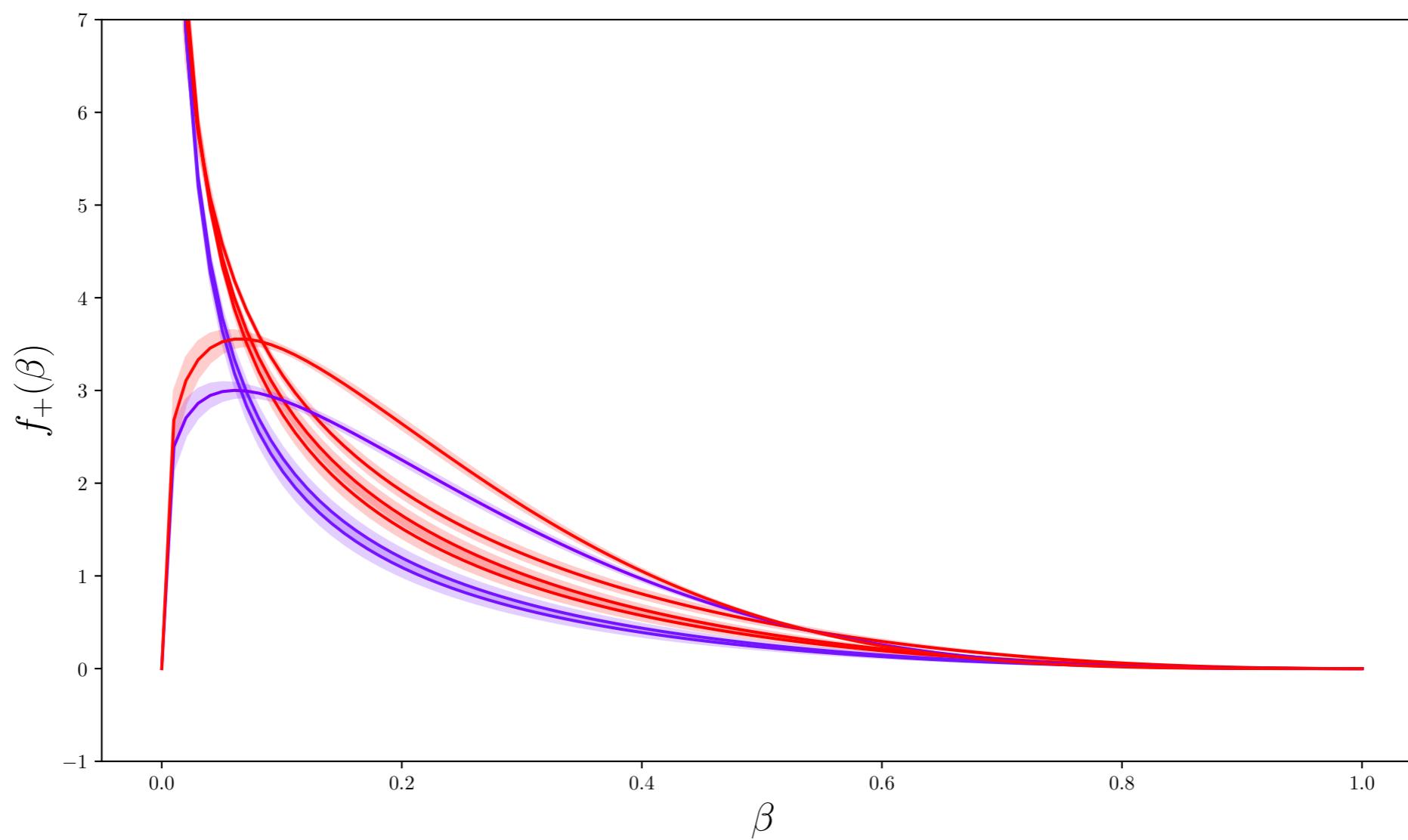
pf111 pi001, chi2=0.703



$$f_q \text{ at } \xi = 0$$

- Parametrize the DD with simplest form

$$A_1(\nu, t, \xi, z^2) = N \int d\beta e^{i\nu\beta} \beta^a (1 - \beta)^b$$



# Future work and Conclusion

- Double Distributions are a natural choice to identify “lightcone” combination of amplitudes for lattice GPDs
- Also a good parameterization of lattice GPDs
- Using Distillation gives ready access to a wide range of momenta
- A scan of all 3 variables, with hundreds of matrix elements, are necessary to constrain the GPD
- Bayesian Model averaging will be necessary to tamper the biases of the inverse problem

**Thank you and the organizers!**