

Hadronic Structure, Conformal Maps, & Analytic Continuation

Thomas Bergamaschi

William I. Jay

Patrick R. Oare



Based on [arXiv:2305.16190](https://arxiv.org/abs/2305.16190)



Lattice 2023

Fermilab

7/31/2023—8/4/2023

Hadronic Structure

The role of spectral functions

$$e^+e^- \rightarrow \text{hadrons}$$

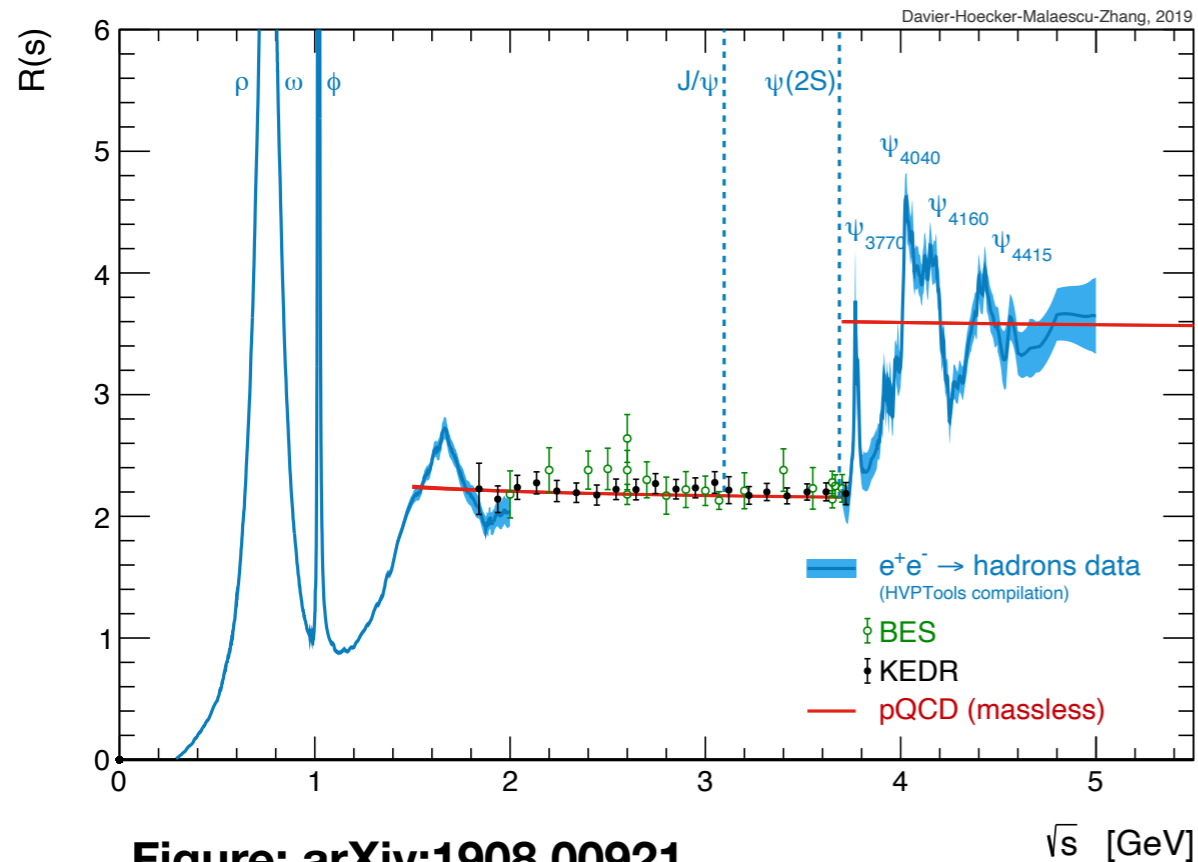
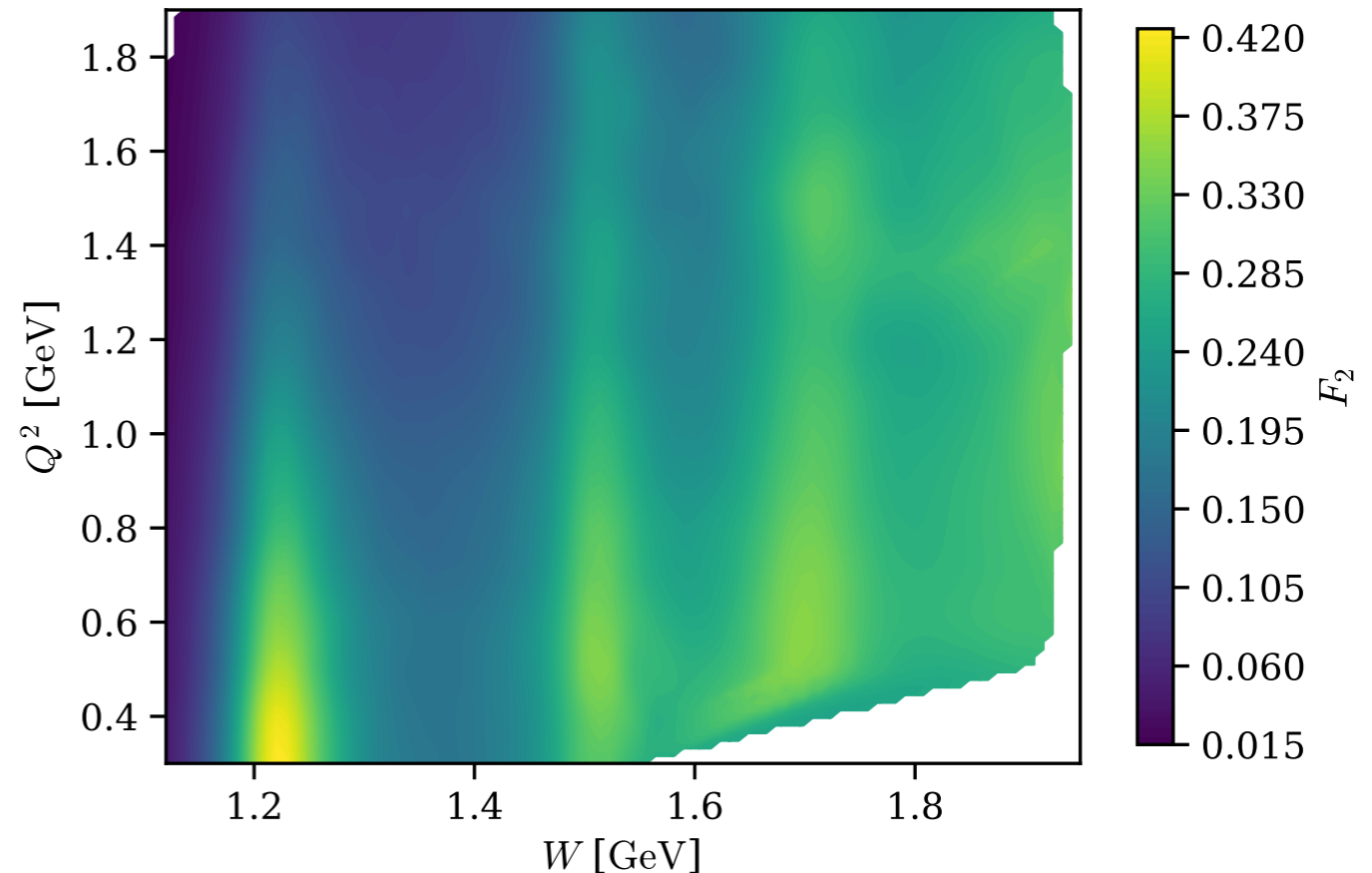


Figure: arXiv:1908.00921

Davier, Hoecker, Malaescu, Zhang

$$e^-p^+ \rightarrow e^-X$$



$$\rho_{\mu\nu}(q) = \frac{1}{2\pi} \int d^4x e^{iq \cdot x} \langle \emptyset | [j_\mu^{\text{EM}}(x), j_\nu^{\text{EM}}(0)] | \emptyset \rangle$$

$$\rho_{\mu\nu}(q) = (q_\mu q_\nu - q^2 g_{\mu\nu}) \rho(q^2)$$

$$W_{\mu\nu}(p, q) = \int \frac{d^4x}{4\pi} e^{iq \cdot x} \langle p | [j_\mu^{\text{EM}}(x), j_\nu^{\text{EM}}(0)] | p \rangle$$

$$W_{\mu\nu} = F_1 \times (\text{Lorentz projectors}) \\ + F_2 \times (\text{Lorentz projectors})$$

Hadronic Structure

The role of spectral functions

- Lattice QCD calculations occur in Euclidean time:

$$G(\tau) = \sum_n \left| \langle 0 | \mathcal{O} | n \rangle \right|^2 \left(e^{-E_n \tau} + e^{-E_n(\beta - \tau)} \right)$$

- In frequency space (take $a \ll 1$):

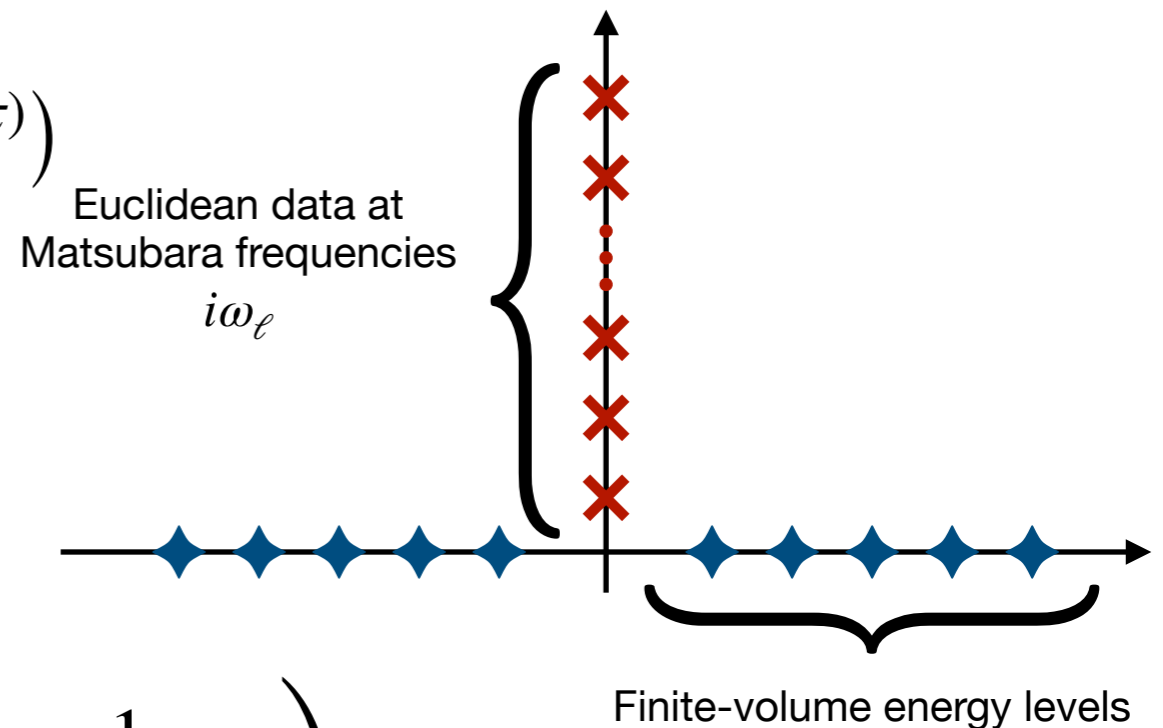
$$G(i\omega_\ell) = \int d\tau e^{i\omega_\ell \tau} G(\tau)$$

$$= \sum_n \left| \langle 0 | \mathcal{O} | n \rangle \right|^2 \left(\frac{1}{E_n + i\omega_\ell} + \frac{1}{E_n - i\omega_\ell} \right)$$

Spectral weight \iff Residue of pole(s)

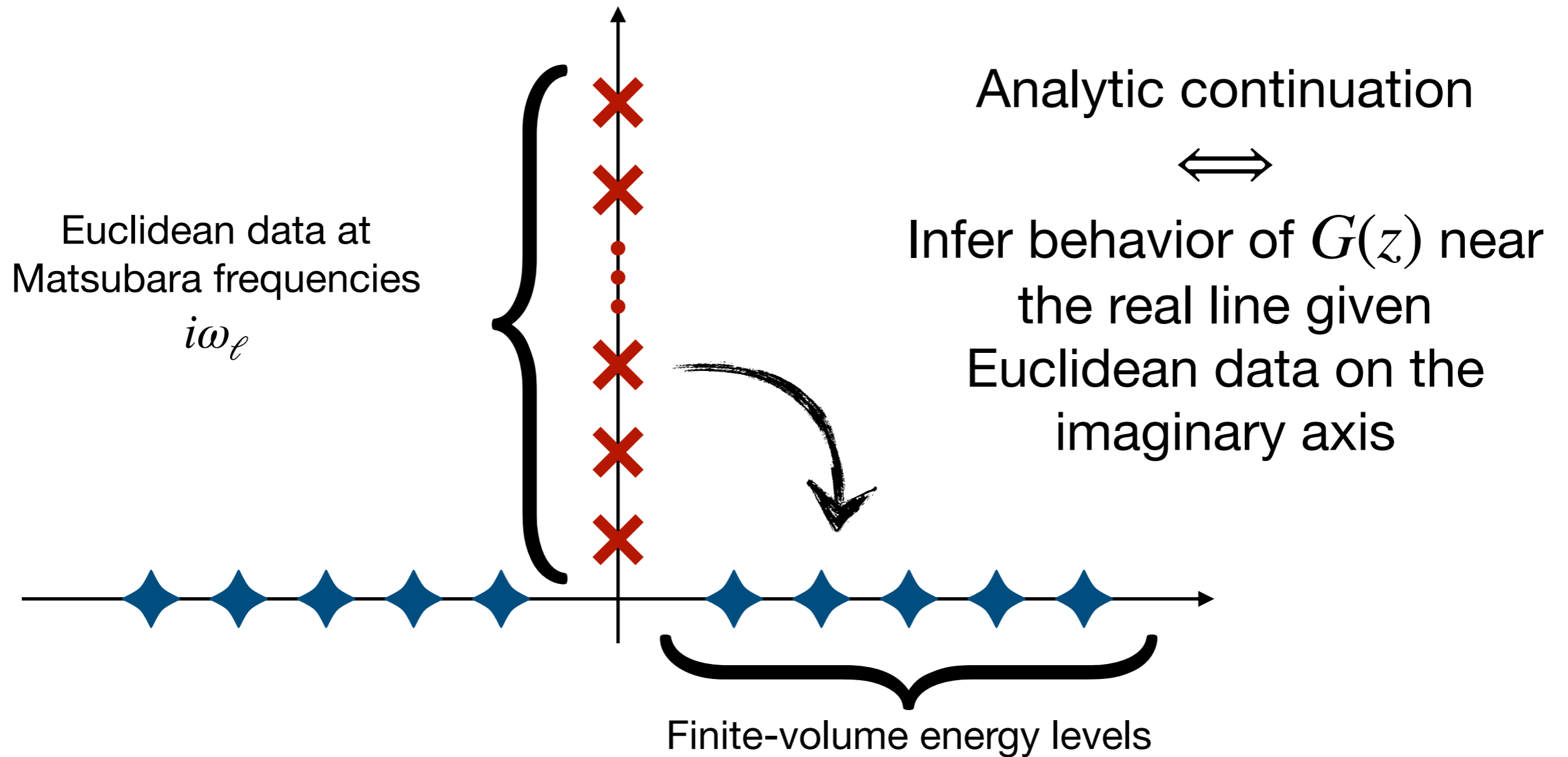
- General relation: $\rho(\omega) = \frac{1}{\pi} \text{Im} G(\omega)$

$$\frac{1}{\pi} \text{Im} \left(\frac{1}{x + i\epsilon} \right) \longrightarrow \delta(x)$$



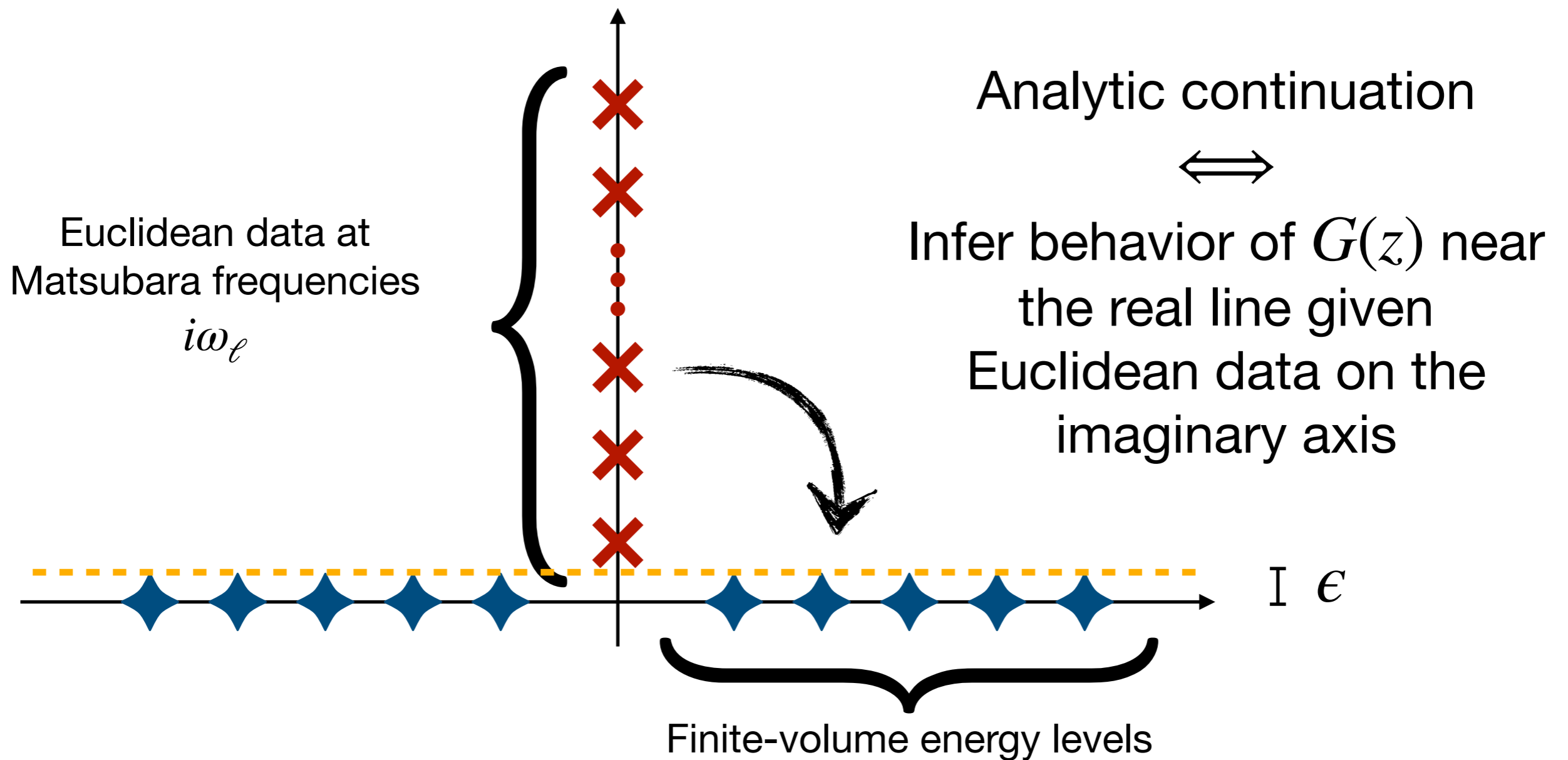
Analytic Continuation

Problem description



Analytic Continuation

Problem description



$\rho^\epsilon(\omega) \equiv \frac{1}{\pi} \text{Im} G(\omega + i\epsilon)$ can be viewed as smeared spectral function in the spirit of Hansen, Meyer, and Robaina [arXiv:1704.08993].

Analytic Continuation

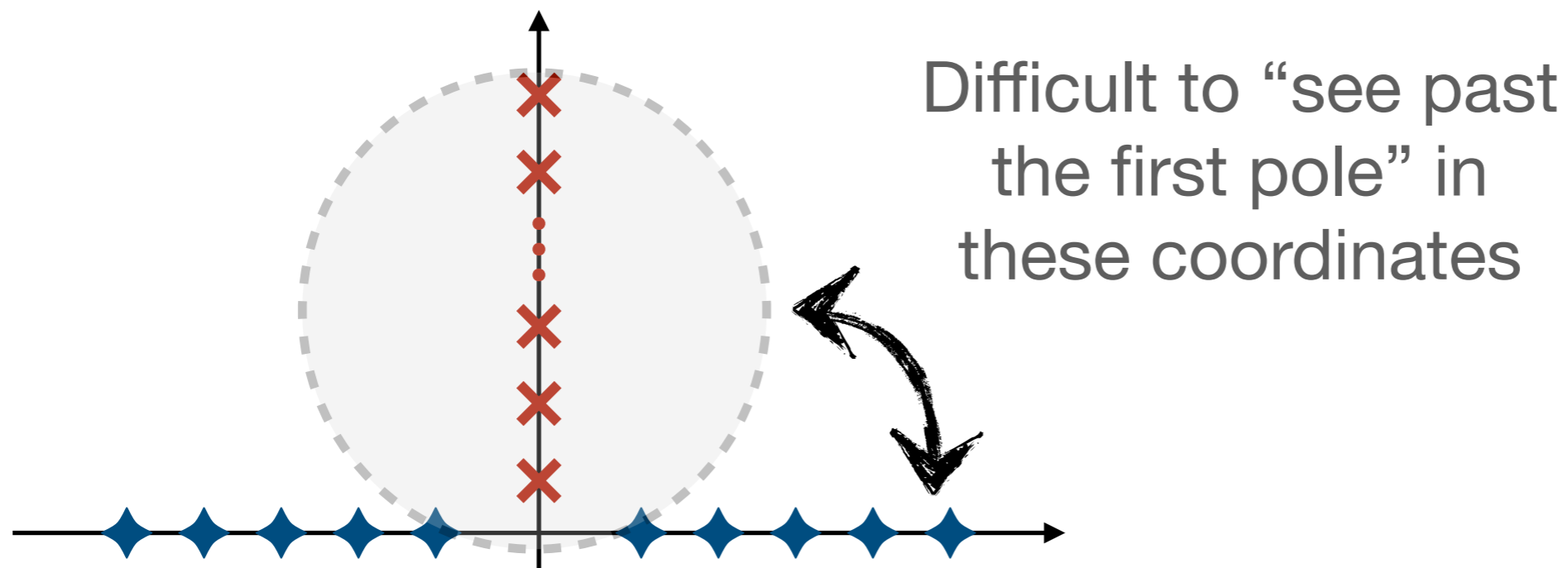
Challenges and hopes

- The received wisdom:
 - **“Excited states decay away exponentially.”** Analytic continuation *to the real line* amounts to an inverse Laplace transform. This problem is ill-defined.
 - **“Little can be said from a finite set of points.”** Complex functions are uniquely defined by analytic continuation given data for the function on a *full open set*, e.g., an open interval on the imaginary axis.
- Causes for cautious optimism:
 - **“Causal Green functions are remarkably rigid.”** Causality dictates that Green functions in QFT must be analytic in the upper half plane, with poles and branch cuts appearing only on the real line.
 - **Carlson’s theorem.** Roughly, *complex functions which do not grow quickly at infinity are uniquely specified by their value on the integers.*

Analytic Continuation

Conformal maps

- Recall: analytic functions are defined by convergent power series in an open set around each nonsingular point
- Radius of convergence is determined by the location of the nearest pole

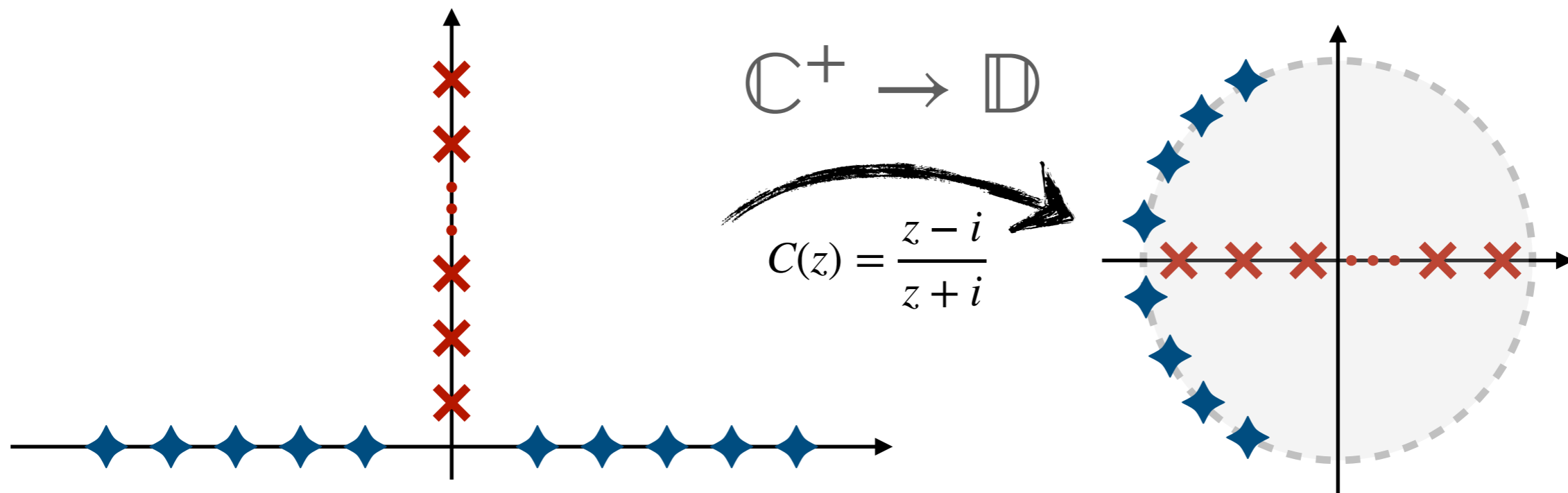


Analytic Continuation

Conformal maps

- Recall: analytic functions are defined by convergent power series in an open set around each nonsingular point
- Radius of convergence is determined by the location of the nearest pole

So change coordinates!



Analytic Continuation

The technical problem

- Recall: analytic functions are defined by convergent power series in an open set around each nonsingular point
- Radius of convergence is determined by the location of the nearest pole
- The Cayley transform maps the problem to the unit disk.

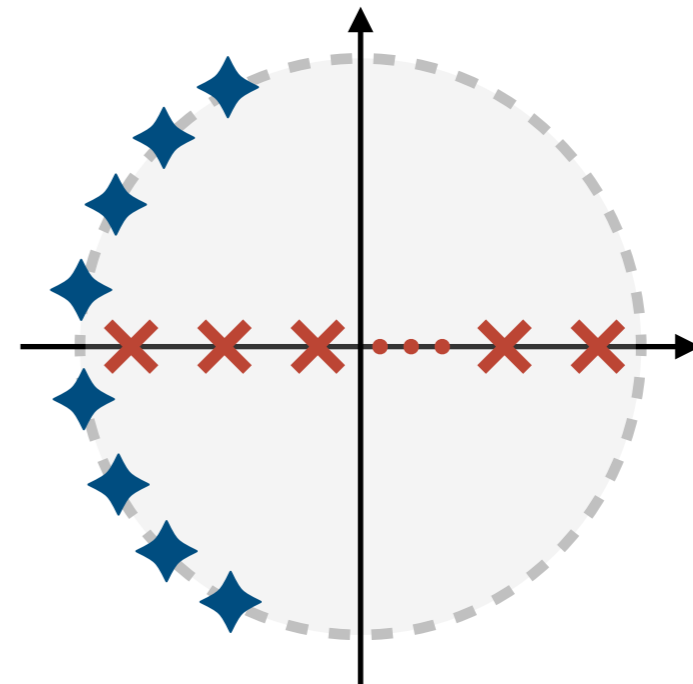
- **Given Euclidean data**

$$\{i\omega_\ell\} \rightarrow \zeta_\ell \in \mathbb{D},$$

$$\{G(i\omega_\ell)\} \mapsto w_\ell \in \mathbb{D},$$

construct an analytic function $f(\zeta)$

on the disk such that $f(\zeta_\ell) = w_\ell$.



Analytic Continuation

Nevanlinna's Theorem

R. Nevanlinna

Ann. Acad. Sci. Fenn. Ser. A 13 (1919)

Ann. Acad. Sci. Fenn. Ser. A 32 (1929)

- **Theorem (Nevanlinna, 1919/1929):**

- Any solution to the interpolation problem with N points can be written in the form

$$f(\zeta) = \frac{P_N(\zeta)f_N(\zeta) + Q_N(\zeta)}{R_N(\zeta)f_N(\zeta) + S_N(\zeta)}$$

where the coefficient functions P_N , Q_N , R_N , S_N are calculable using an inductive formula in terms of the input data $\{\zeta_\ell\}$ and $\{w_\ell\}$ and an arbitrary analytic function $f_N(\zeta) : \mathbb{D} \rightarrow \mathbb{D}$.

Derivation: See our preprint [arXiv:2305.16190], which follows modern treatment by mathematician Nicolau [<https://mat.uab.cat/~artur/data/nevanlinna-pick.pdf>]

- $P_N, Q_N, R_N, S_N \iff$ “Nevanlinna coefficients”
- Arbitrary function $f_N(\zeta) \iff$ “Freedom to specify further Euclidean data to constrain the interpolating function”
- Applicability to field-theory problems first recognized by condensed-matter theorists Fey, Yeh, and Gull [arXiv:2010.04572]

Analytic Continuation

The full space of solutions

- Question: For fixed N and ζ , **what are the possible values that an interpolating function $f(\zeta)$ can take**, by varying possible values of the arbitrary function $f_N(\zeta) \in \mathbb{D}$?
- The size of this set \iff ambiguity in the analytic continuation
- Remarkably, this set can be parameterized explicitly for each, N and $\zeta \in \overline{\mathbb{D}}$

Analytic Continuation

The full space of solutions

- Question: For fixed N and ζ , **what are the possible values that an interpolating function $f(\zeta)$ can take**, by varying possible values of the arbitrary function $f_N(\zeta) \in \mathbb{D}$?
- Answer: The space of possible values is given by the *Wertevorrat* $\Delta_N(\zeta)$, which is the disk of radius $r_N(\zeta)$ and centered at $c_N(\zeta)$.

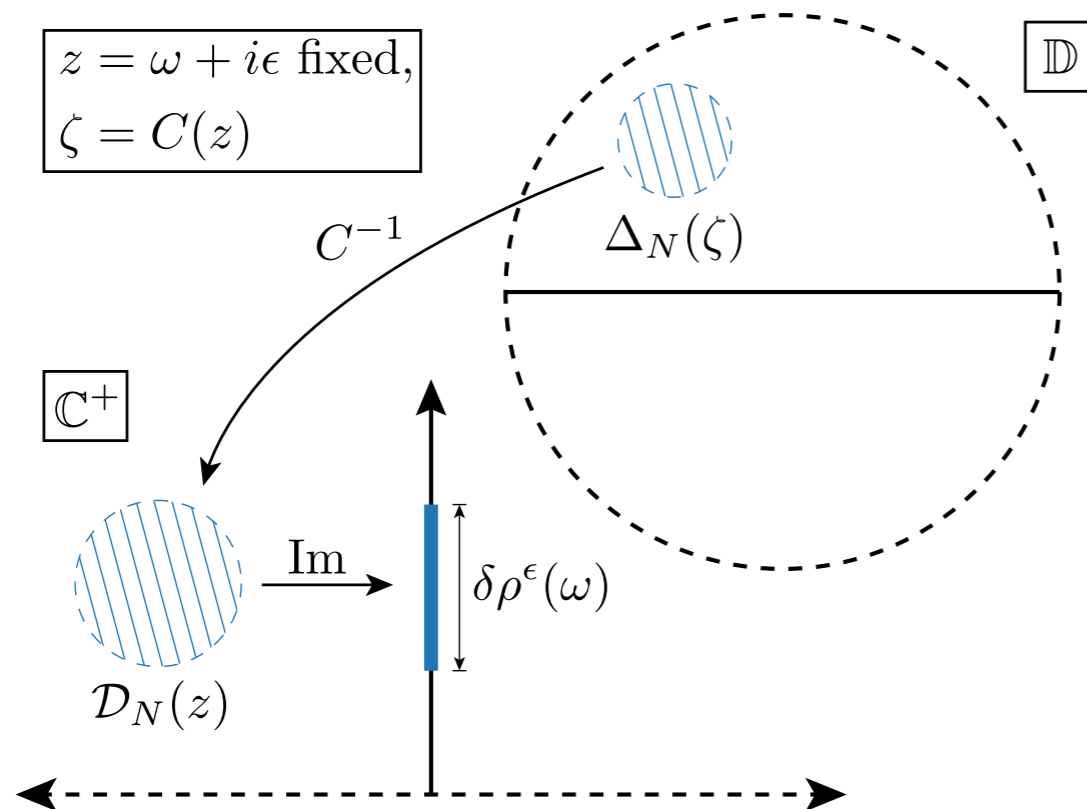
$$c_N = \frac{P_N \overline{(-R_N/S_N)} + Q_N}{R_N \overline{(-R_N/S_N)} + S_N} \quad r_N = \frac{|P_N S_N - Q_N R_N|}{|S_N|^2 - |R_N|^2}$$

- The Wertevorrat $\Delta_N(\zeta)$ rigorously contains the full infinite family of all possible analytic continuations at each point $\zeta \in \mathbb{D}$.

Analytic Continuation

The Wertevorrat and rigorous bounds on $\rho^\epsilon(\omega)$

- Finally we need to map the Wertevorrat back to the upper half plane. Use the inverse Cayley transform $z = C^{-1}(\zeta)$.



$$\rho^\epsilon(\omega) = \frac{1}{\pi} \text{Im } G(\omega + i\epsilon)$$

$$\delta\rho^\epsilon(\omega) = \frac{1}{\pi} \left[\max \text{Im } \partial D_N(\omega + i\epsilon) - \min \text{Im } \partial D_N(\omega + i\epsilon) \right]$$

Analytic Continuation

The Algorithm

1. Start with a Euclidean correlation function $G(t)$
2. Evaluate the Fourier coefficients to obtain $G(i\omega_\ell)$
3. Map the Euclidean data to the unit disk
4. Solve the interpolation problem
 - Evaluate the Nevanlinna coefficients
 - Compute the Wertevorrat
5. Map the Wertevorrat back to the upper half-plane
6. For each point $\omega + i\epsilon$, evaluate the space of possible smeared spectral densities $\delta\rho^\epsilon(\omega)$

Numerical Example

The R-ratio – reconstructing a parameterization

- Bernecker and Meyer [arXiv:1107.4388] give a useful parameterization of R-ratio data
- This parameterization can serve as input for a spectral reconstruction

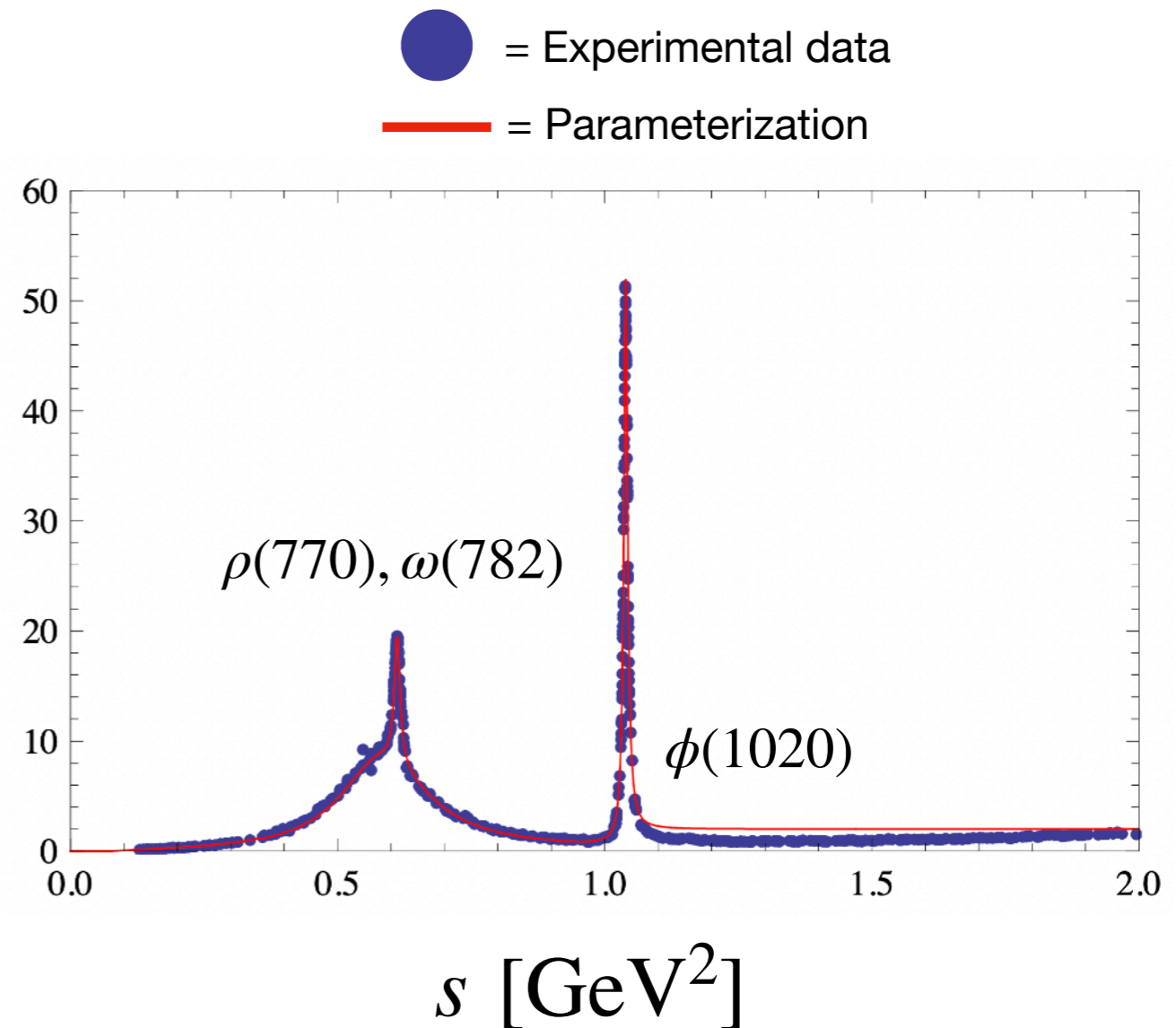
- Can easily convert:

$$R(s) \iff \rho(\omega) \iff G(i\omega_\ell)$$

Formula from
beginning of talk

“Laplace transform”

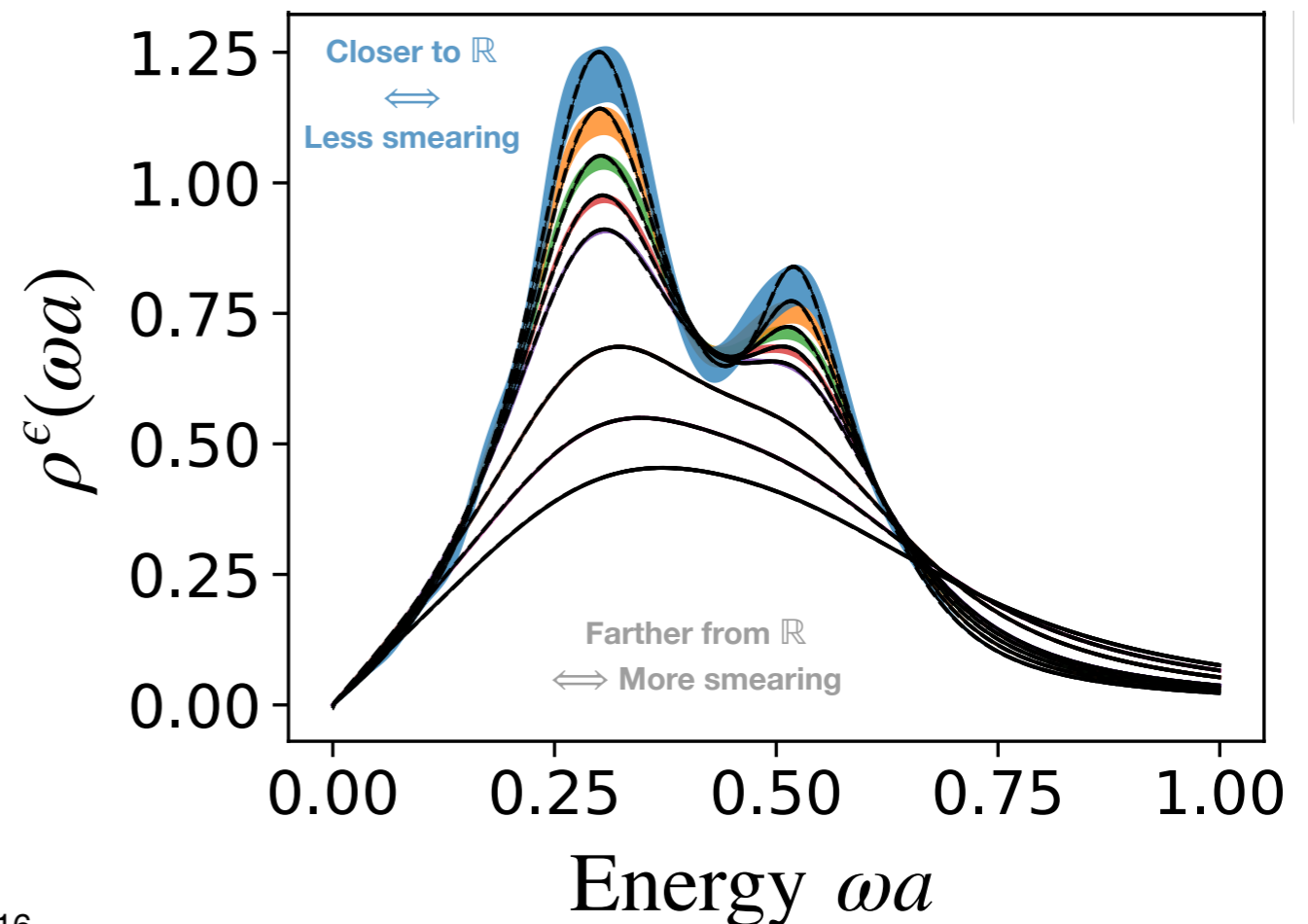
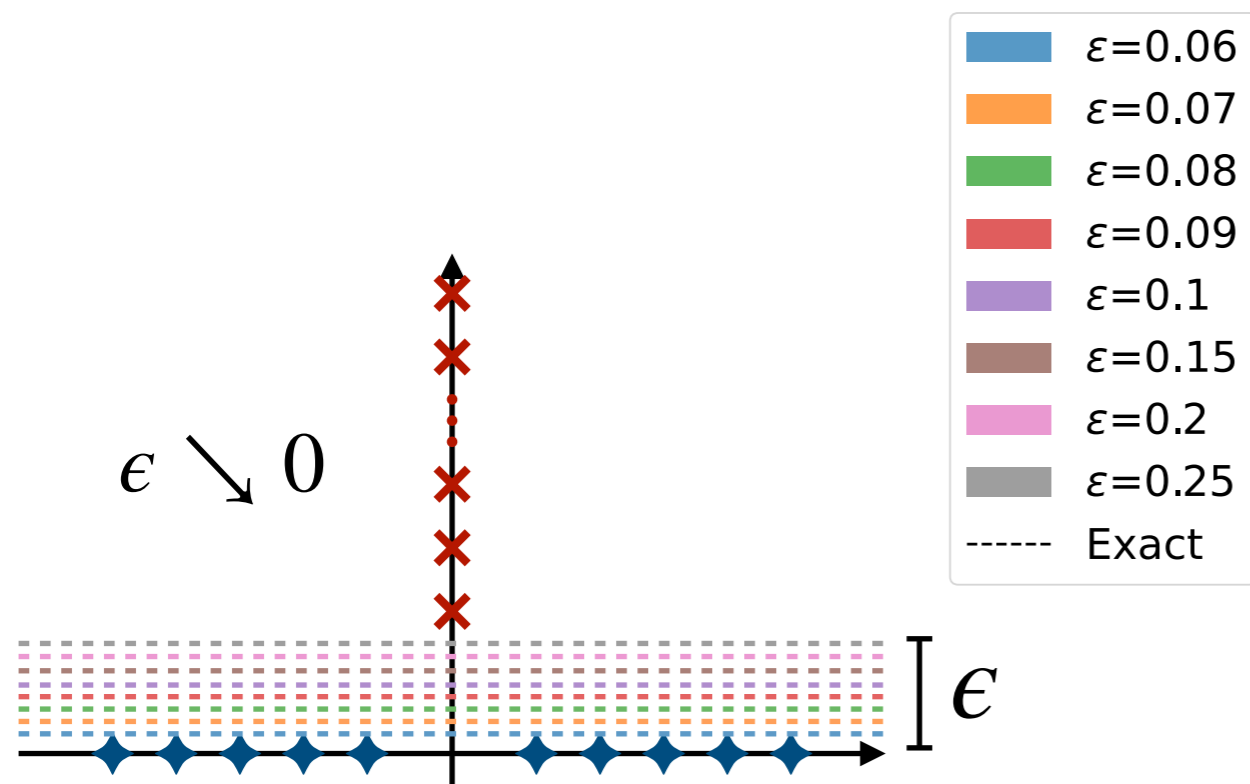
$R(s)$



Numerical Example

The R-ratio – reconstructing a parameterization

- Energies rescaled to line in unit interval \implies lattice units with $a \approx 0.07$ fm, so $am_\rho \approx 0.25$
- Euclidean data generated for $\beta = 96$ total points on the imaginary-energy axis
- ✓ Spectral peaks from $\rho(770)/\omega(782)$ and $\phi(1020)$ clearly visible in reconstructions
- ✓ Exact answer is contained within the bounding envelope of the Wertenorrat



Summary

- Today's spectral reconstruction algorithm:
 - Leverages a century-old theorem due to Nevanlinna
 - Builds off of work in the condensed matter community by Fey, Yeh, and Gull [arXiv:2010.04572]
 - Applies to diagonal correlation functions of bosonic/fermionic operators
 - Bosonic operators (alternative view by Nogaki and Shinaoka [arXiv:2305.03449])
 - Fermionic operators [arXiv:2010.04572]
- Our method:
 - Admits interpretation as a smeared spectral density in the spirit of Hansen, Meyer, and Robaina [arXiv:1704.08993]
 - Bound errors rigorously with Nevanlinna's *Wertevorrat*
 - Wertevorrat \iff Full space of functions consistent with the input data and analyticity
- Next steps: understanding interplay with statistical uncertainties
 - Important existing work in this direction by Huang, Gull, and Lin [arXiv:2210.04187]

Backup

The Pick Criterion

G. Pick
Math. Ann. 77, 7 (1915)

Existence of an Interpolating Function

- Question: When does an interpolating function, analytic on \mathbb{D} , exist for the input data $\{\zeta_\ell\}, \{w_\ell\}$ such that $f(\zeta_\ell) = w_\ell$ for all ℓ ?
- Theorem (Pick, 1915):

Such a function exists if and only if the Pick matrix

$$\begin{bmatrix} 1 - w_i \bar{w}_j \\ 1 - \zeta_i \bar{\zeta}_j \end{bmatrix}$$

is positive semidefinite.

- This condition can fail for noisy data (observed already in arXiv:2010.04572). See also arXiv:2210.04187 for work toward robust reconstructions.

Nevanlinna Coefficients

Concrete Formulae

Convenient notation for manipulating Möbius transformations:

$$\begin{pmatrix} a(\zeta) & b(\zeta) \\ c(\zeta) & d(\zeta) \end{pmatrix} h(\zeta) \equiv \frac{a(\zeta)h(\zeta) + b(\zeta)}{c(\zeta)h(\zeta) + d(\zeta)}$$

The Nevanlinna coefficients:

$$\begin{aligned} f(\zeta) &= U_1(\zeta)U_2(\zeta)\cdots U_N(\zeta)f_N(\zeta) \\ &\equiv \begin{pmatrix} P_N(\zeta)Q_N(\zeta) \\ R_N(\zeta)S_N(\zeta) \end{pmatrix} f_N(\zeta) \end{aligned}$$

Inductive blocks:

$$U_n(\zeta) = \frac{1}{\sqrt{1 - |w_n^{(n-1)}|^2}} \begin{pmatrix} b_{\zeta_n}(\zeta) & w_n^{(n-1)} \\ \bar{w}_n^{(n-1)} b_{\zeta_n}(\zeta) & 1 \end{pmatrix}$$

Blaschke factors:

$$b_a(\zeta) = \frac{|a|}{a} \frac{a - \zeta}{1 - \bar{a}\zeta}$$

“The n th interpolant f_n evaluated at the m th zero”

$$w_m^{(n)} \equiv f_n(\zeta_m)$$

Analytic Continuation

Computing properties of the *Wertevorrat*

- Question: For fixed N and ζ , **what are the possible values that an interpolating function $f(\zeta)$ can take**, by varying possible values of the arbitrary function $f_N(\zeta) \in \mathbb{D}$?
- Answer: For fixed N and ζ , consider the auxiliary function $T : \mathbb{D} \rightarrow \mathbb{D}$

$$T(w) = \frac{P_N w + Q_N}{R_N w + S_N}$$

- We want to evaluate the image of this map.
- First, observe that $T(w)$ is a Möbius transformation, which maps circles to circles. Therefore $T(\mathbb{D})$ must be also be disk $\Delta_N(\zeta)$.
- Since $T(-S_N/R_N) = \infty$, the reflection property of Möbius transformations says that disk is centered at the point c_N . A calculation gives the radius r_N .

$$c_N = \frac{P_N \overline{(-R_N/S_N)} + Q_N}{R_N \overline{(-R_N/S_N)} + S_N} \quad r_N = \frac{|P_N S_N - Q_N R_N|}{|S_N|^2 - |R_N|^2}$$

Bosons and fermions

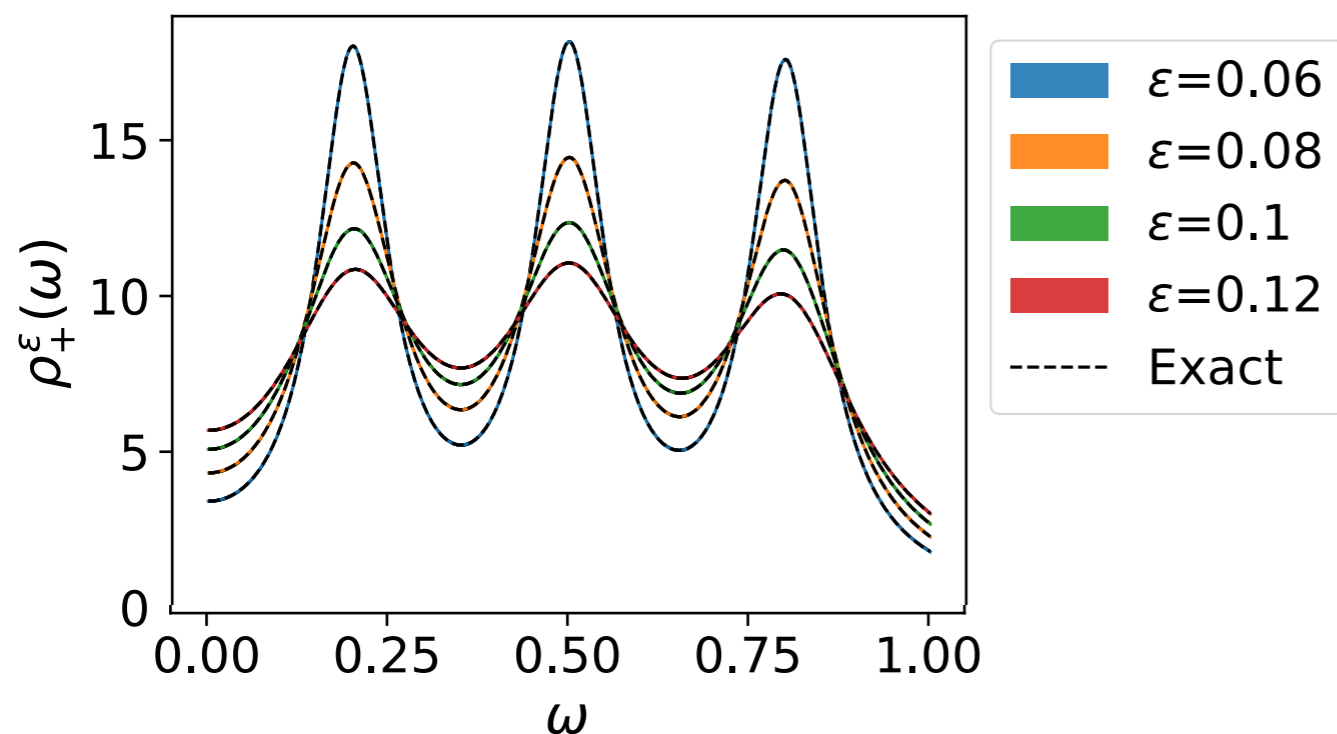
- Can consider two possibilities:
 - Fermionic: $\rho(-\omega) = \rho(\omega)$ [+ from anti-commutator]
 - Bosonic: $\rho(-\omega) = -\rho(\omega)$ [- from commutator]
- Technical detail: slightly different conformal maps used to transform the problem to the unit disk. See our arXiv:2305.16190
- For finite smearing ϵ , $\rho_+^\epsilon(\omega) \neq \rho_-^\epsilon(\omega)$.
- BUT they converge to the same function: $\lim_{\epsilon \rightarrow 0} \rho_\pm^\epsilon(\omega) = \rho(\omega)$, $\omega > 0$

Numerical Examples

Discrete poles

- For $\omega > 0$: $\rho(\omega) = \delta(\omega - 0.2) + \delta(\omega - 0.5) + \delta(\omega - 0.8)$
- Euclidean data generated for $\beta = 64$ total points on the imaginary-energy axis
- The exact answer is always contained with bounding envelope of the Wertevorrat

Fermionic



Bosonic

